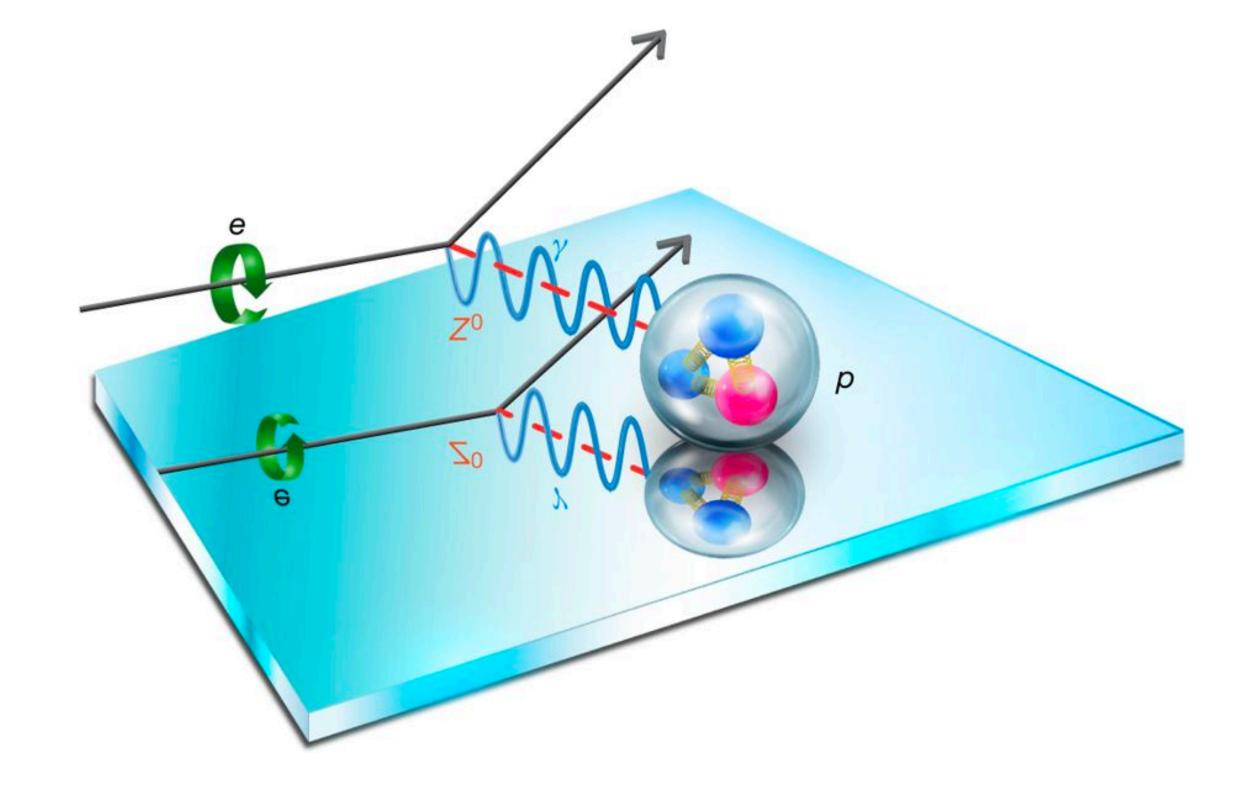
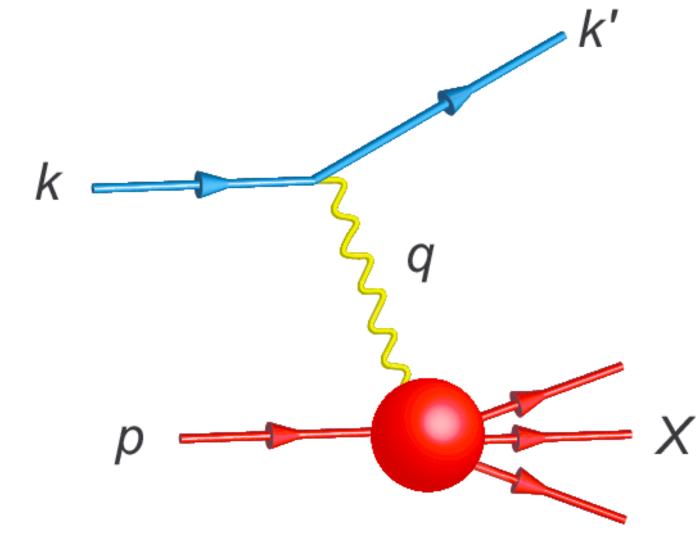
QUADRATIC AND REDUCIBLE TWO LOOP LEVEL FULL ELECTROWEAK e-P SCATTERING WITH COVARIANT APPROACH







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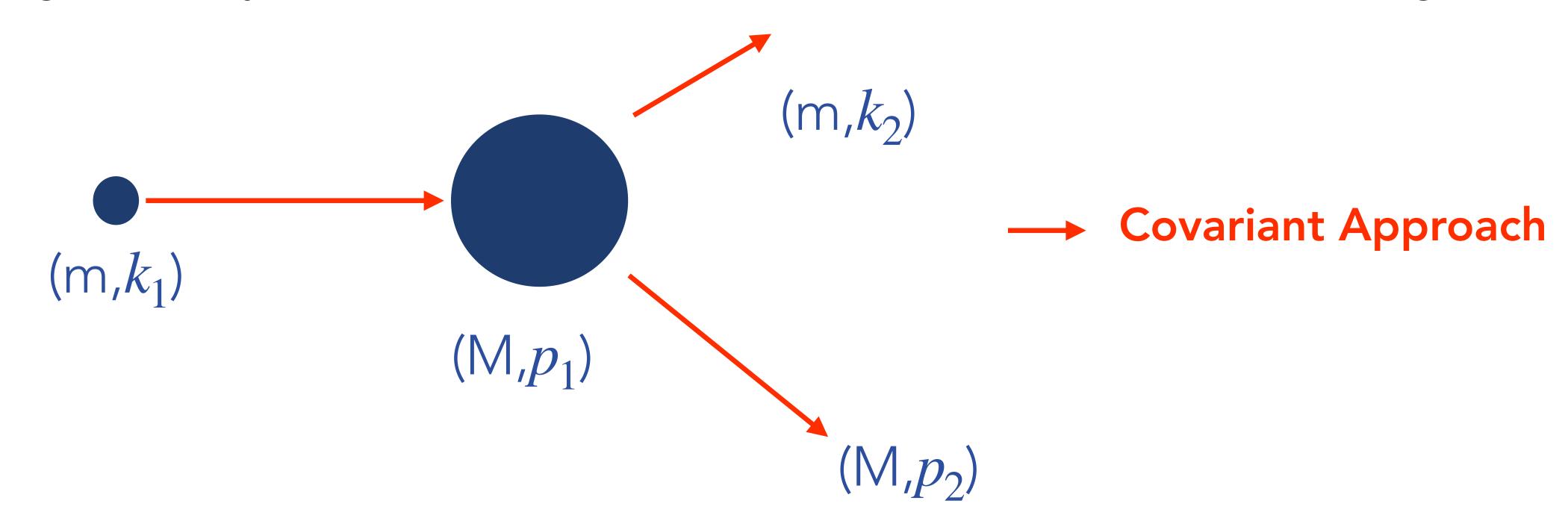
### MOTIVATION

- The theory of Standard Model (SM)  $\rightarrow$  unifies Electromagnetic, Weak and Strong interactions  $\rightarrow$  can make predictions that match experiments to one part in ten billion.
- SM limitations → don't include gravity, dark matter/dark energy existence, hierarchies of scale related to Higgs boson etc.
- Theoretical door open for Beyond the SM (BSM) physics to be observed at TeV scale → but till date
  no concrete evidence of BSM at 13 TeV centre-of-mass energy at LHC.
- Low energy precision physics becomes important → provides a way to reach mass scales not directly accessible at existing high-energy colliders.
- We are doing precision physics with full electroweak Parity Violating Asymmetry  $(A_{PV}) \rightarrow$  achieve by calculating the higher order corrections up to NNLO ( $\alpha^4$ ) using Covariant/ leptonic tensor approach.

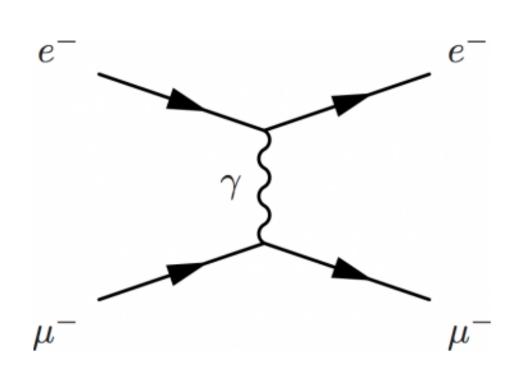
June 14th, 2024 Mahumm Ghaffar 2/20 CFNS Summer School 2024

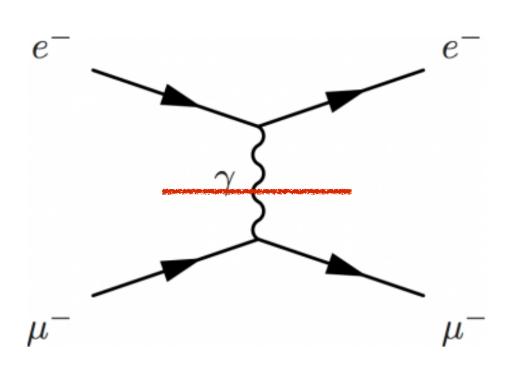
## FULL ELECTROWEAK $e^-p$ SCATTERING

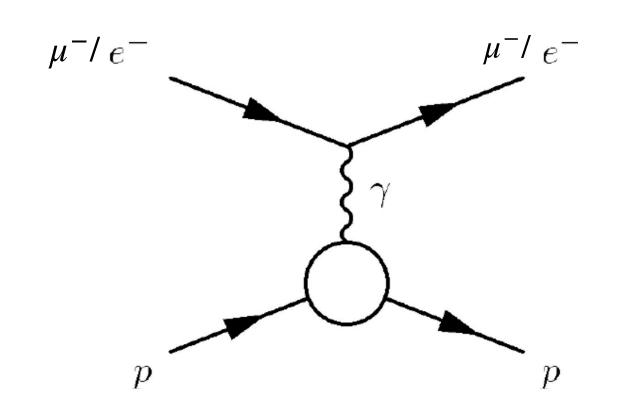
- Elastic  $e^-p$  scattering is studied up to the NNLO level considering all SM particles in the loop.
- ullet A longitudinally polarized  $e^-$  scatters off an unpolarized proton target

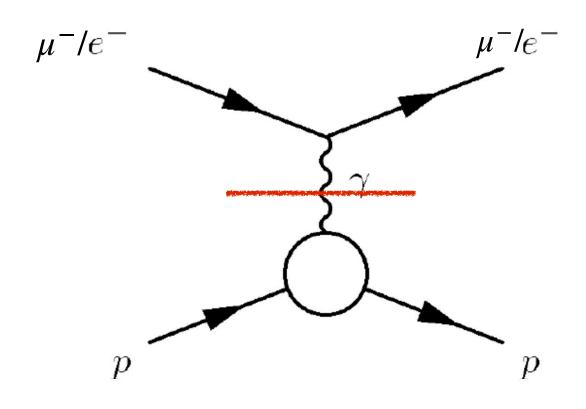


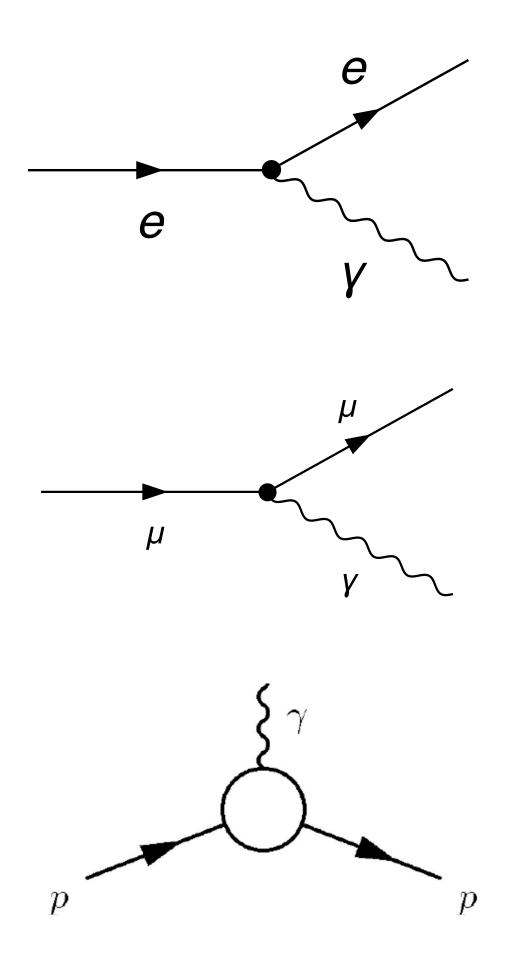
## WHAT IS A COVARIANT APPROACH?











## COVARIANT APPROACH WITH LEPTONIC-HADRONIC TENSORS

• The differential cross section of general lepton-lepton/hadron scattering can be obtained by:

$$d\sigma \sim L^{\mu\nu} L_{\mu\nu}$$
 or  $d\sigma \sim L^{\mu\nu} W_{\mu\nu}$ 

• where  $W_{\mu\nu}$  is the hadronic tensor which in case of elastic  $e^-p$  scattering:

$$W_{\mu\nu} = H_1 g_{\mu\nu} + H_2 p_{1\mu} p_{1\nu} + H_3 p_{2\mu} p_{2\nu} + H_4 p_{1\mu} p_{2\nu} + H_5 p_{2\mu} p_{1\nu} + H_6 \epsilon_{\mu,\nu,p_1,p_2}$$

where  $p_1$  and  $p_2$  are incoming and outgoing protons momenta.  $H_1$ ,  $H_2$ ,  $H_3$ ,  $H_4$ ,  $H_5$  and  $H_6$  are the hadronic structure functions which can be extracted from experimental data.

## QED AND ELECTROWEAK HADRONIC COUPLINGS WITH FORM FACTORS

$$\Gamma^{\mu}_{\gamma-p}(q^2) = ieCnp2\left(f2p\gamma^{\mu} + gp\gamma_L\gamma^{\mu}\omega_- + gp\gamma_R\gamma^{\mu}\omega_+ - \frac{f2p(p_1^{\mu} + p_2^{\mu})}{2m_p}\right)$$

$$\Gamma^{\mu}_{Z-p}(q^2) = -ieCnp2\left(F2W\gamma^{\mu} + gpz_{L}\gamma^{\mu}\omega_{-} + gpz_{R}\gamma^{\mu}\omega_{+} - \frac{F2W(p_1^{\mu} + p_2^{\mu})}{2m_p}\right)$$

$$F2W = \frac{F2Vp - 4\sin^2\theta_W f2p}{4\cos\theta_W \sin\theta_W} \rightarrow \text{EW form factor} \qquad gpz_{(L,R)} = \frac{F1Vp - 4\sin^2\theta_W f1p \pm G1p}{4\sin\theta_W \cos\theta_W}$$

$$gpz_{(L,R)} = \frac{F1Vp - 4\sin^2\theta_W f1p \pm G1p}{4\sin\theta_W \cos\theta_W}$$

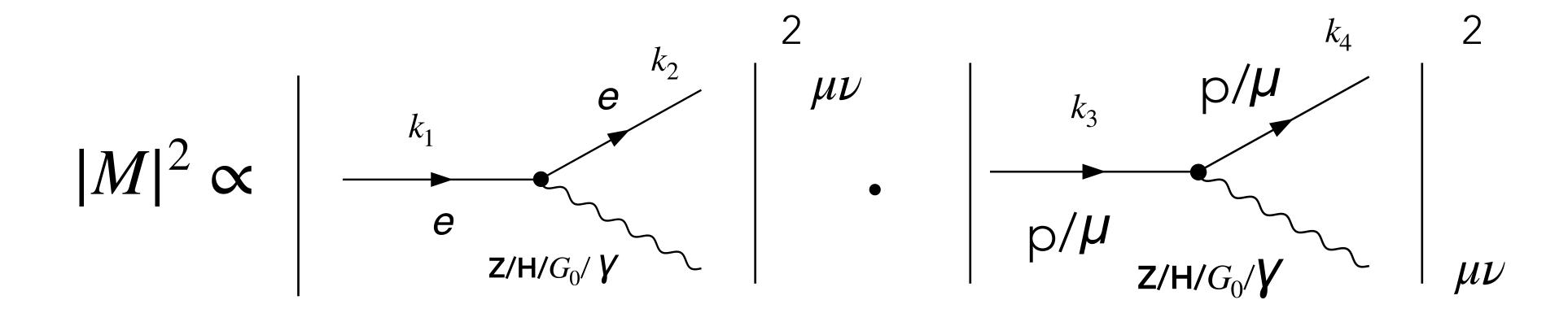
$$gp\gamma_L = gp\gamma_R = f1p(0) \rightarrow$$
 Electric form factor

$$Cnp2 = \left(\frac{\Lambda^2}{\Lambda^2 - t}\right)^2, \quad \Lambda = \sqrt{0.83 \ m_p^2}$$

$$G1p = 1.267 \rightarrow Axial form factor$$

$$F(1,2)Vp = f(1,2)p - f(1,2)n$$

### TREE-LEVEL LEPTONIC TENSOR ( $\alpha$ -ORDER)

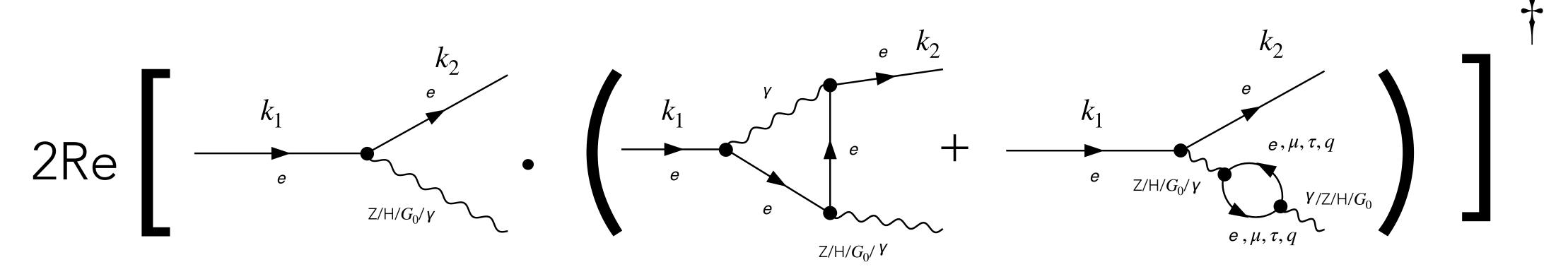


• For tree-level upper part of the diagram (say  $e^-p$  scattering), one can calculate leptonic tensor which is:

$$L_{\mu\nu}^{0} \propto 4\pi\alpha((l_{1})g_{\mu\nu} + (l_{2})k_{2\mu}k_{1\nu} + (l_{3})k_{1\mu}k_{2\nu} + \dots)$$

where  $k_1$ ,  $k_2$  are incoming and outgoing  $e^-$  momenta and  $l_{1,2..}$  are tree level leptonic tensor structure functions.

# NEXT TO THE LEADING ORDER (NLO) LEPTONIC TENSOR ( $\alpha^2$ -ORDER)



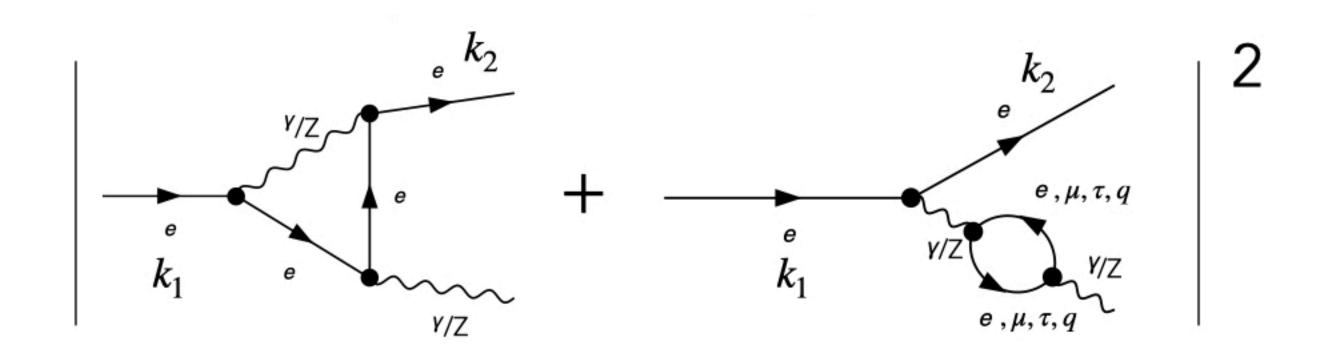
• The NLO leptonic tensor can be obtained by multiplying tree-level upper diagram with the sum of one-loop level self energy (SE) and triangular diagrams.

$$L_{\mu\nu}^{NLO} = (m_1)g_{\mu\nu} + (m_2)k_{1\nu}k_{2\mu} + (m_3)k_{1\mu}k_{2\nu} + (m_4)k_{1\mu}k_{1\nu} + (m_5)k_{2\mu}k_{2\nu} + \dots$$

Where  $m_{1,2,3...}$  are leptonic structure functions which depend on the momentum transfer "t" and written in terms of Passarino-Veltman integral functions. We used LoopTools Mathematica package to calculate them.

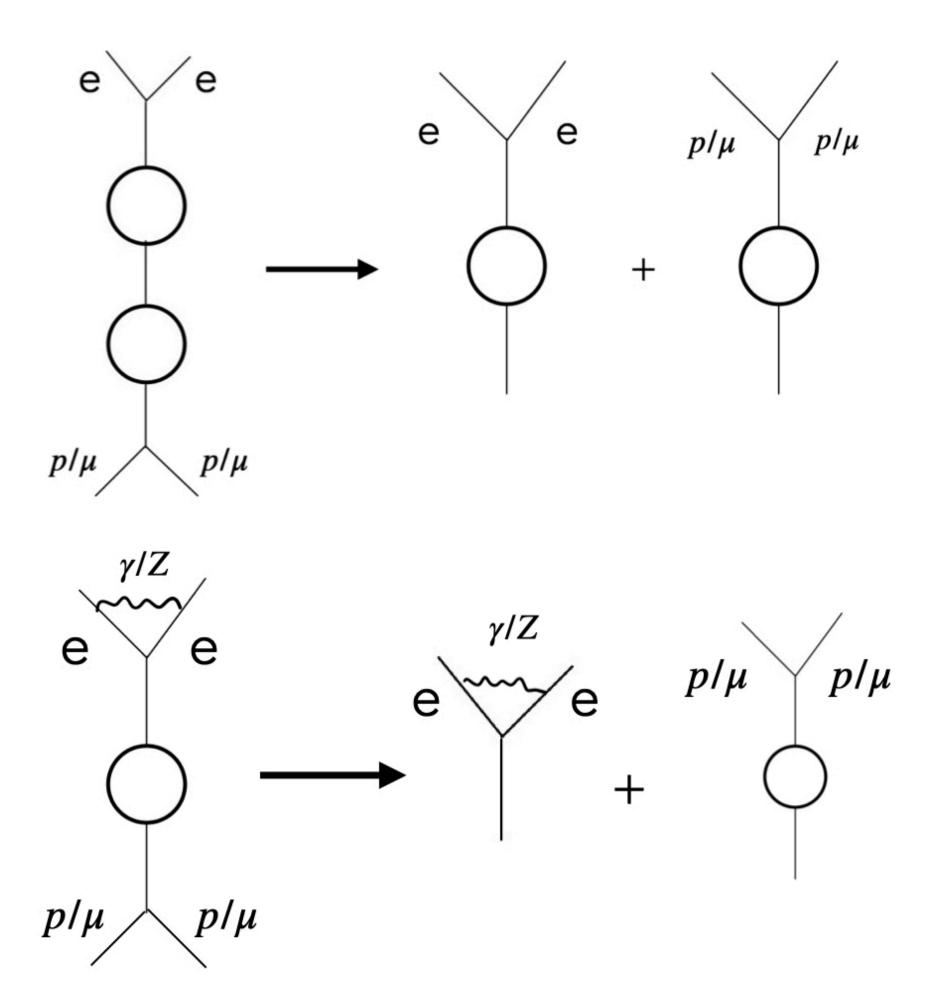
In total 307 graphs SE and triangular graphs.

# NEW RESULTS: QED AND ELECTROWEAK NNLO LEVEL LEPTONIC TENSOR ( $\alpha^3$ -ORDER)

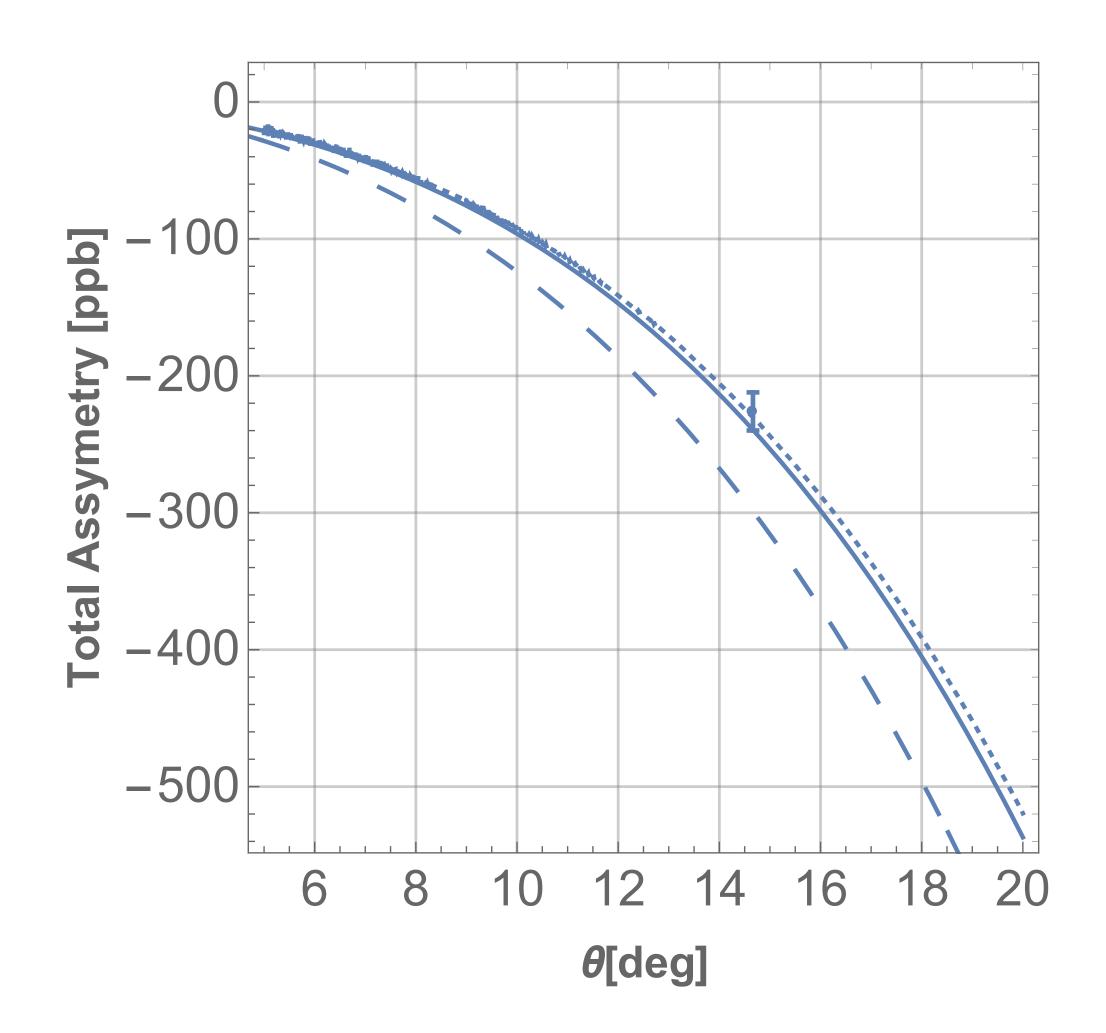


$$L_{\mu\nu}^{NNLO} = (n_1)g_{\mu\nu} + (n_2)k_{1\nu}k_{2\mu} + (n_3)k_{1\mu}k_{2\nu} + (n_4)k_{1\mu}k_{1\nu} + (n_5)k_{2\mu}k_{2\nu} + \dots$$

- FEYNARTS and FORMCALC as base languages to calculate leptonic tensor structure functions.
- Kept the mass of electron.



# Tree level, NLO and NNLO level $A_{PV}$ for $e^-p$ scattering versus $\theta_{CM}$ using QWEAK kinematics



 $(e^-p)$  Tree level, NLO and NNLO level  $A_{PV}$  versus  $heta_{CM}$ 

$$(\theta = 14.6^{0})$$
---- Tree  $A_{PV} \sim -298.9 \; ppb$ 
---- NLO  $A_{PV} \sim -239 \; ppb$ 
..... NNLO  $A_{PV} \sim -230 \; ppb$ 

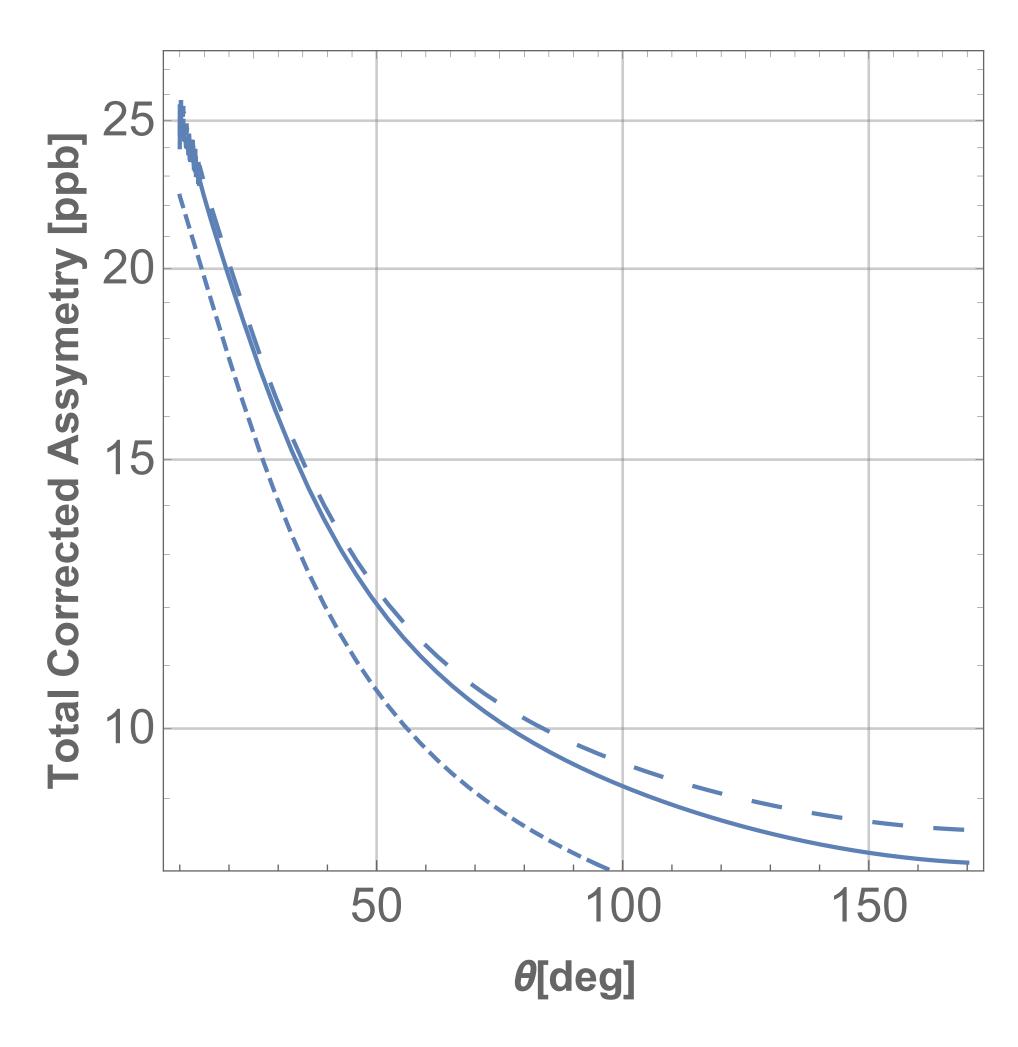
#### **QWEAK Measured**

 $\sim -226.5 \pm 7.3 (statistical) \pm 5.8 (systematic) ppb$ 

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$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$$

# NLO and NNLO level Corrected $A_{PV}$ Asymmetry with Correction factors in percentage for QWEAK kinematics



 $(\theta=14.6^0)$  ...... NLO level Corrected  $A_{PV}=19.9~\%$  ...... Quadratic level Corrected  $A_{PV}=22.5~\%$ 

Corrected 
$$A_{PV}\% = \left(\frac{Tree\ level\ A_{PV} - (NLO, NNLO)A_{PV}}{Tree\ level\ A_{PV}}\right) \times 100$$

- - - Total Corrected  $A_{PV}=22.9~\%$ 

NLO and NNLO level Corrected  $A_{PV}$  versus  $heta_{CM}$ 

### RESULTS:

- For completeness, work in progress to include soft and hard photon bremsstrahlung cross sections in the results.
- Next goal is to consider the polarized proton and study the  ${\cal A}_{PV}$  effects in electron-proton scattering.
- We make predictions for the  $e^-p$  NNLO level radiative corrections. These theoretical predictions will be important for many experimental programs such as QWEAK, P2, EIC and MOLLER (background studies) searching for physics beyond the Standard Model at the precision frontier.

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## Backup Slides

### PARITY VIOLATING ASYMMETRY

- Formula:  $A_{PV}=\frac{\sigma_R-\sigma_L}{\sigma_R+\sigma_L}$ , where:  $\sigma_R\propto |M_R|^2$  and  $\sigma_L\propto |M_L|^2$
- For QED,  $|M_{\gamma R}| = |M_{\gamma L}|$ , numerator contains just weak+electroweak cross terms.
- Denominator contains just QED terms as  $m_Z$  (90 GeV)>  $m_e$ -(0.5 MeV)

$$A_{PV} = \frac{|M_{ZZ}|_R^2 - |M_{ZZ}|_L^2 + |M_{\gamma Z}|_R^2 - |M_{\gamma Z}|_L^2}{|M_{\gamma \gamma}|_R^2 + |M_{\gamma \gamma}|_L^2}$$

• Using covariant approach  $\rightarrow$  calculated tree level, one-loop level and quadratic level scattering amplitudes squared along with the parity violating asymmetry.

#### QED

$$i\mathcal{M} = \bar{u}(k_2)(-ie\gamma_{\mu})u(k_1)\frac{-i}{t}\bar{u}(p_2)(-ie\Gamma^{\mu}_{\gamma-p})u(p_1)$$

$$\Gamma^{\mu}_{\gamma-p} = F_1^p(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}F_2^p(t)$$

 $F_1(t)$  and  $F_2(t) \to {\sf Dirac}$  and Pauli form factors depending on momentum transfer "t".

### WEAK

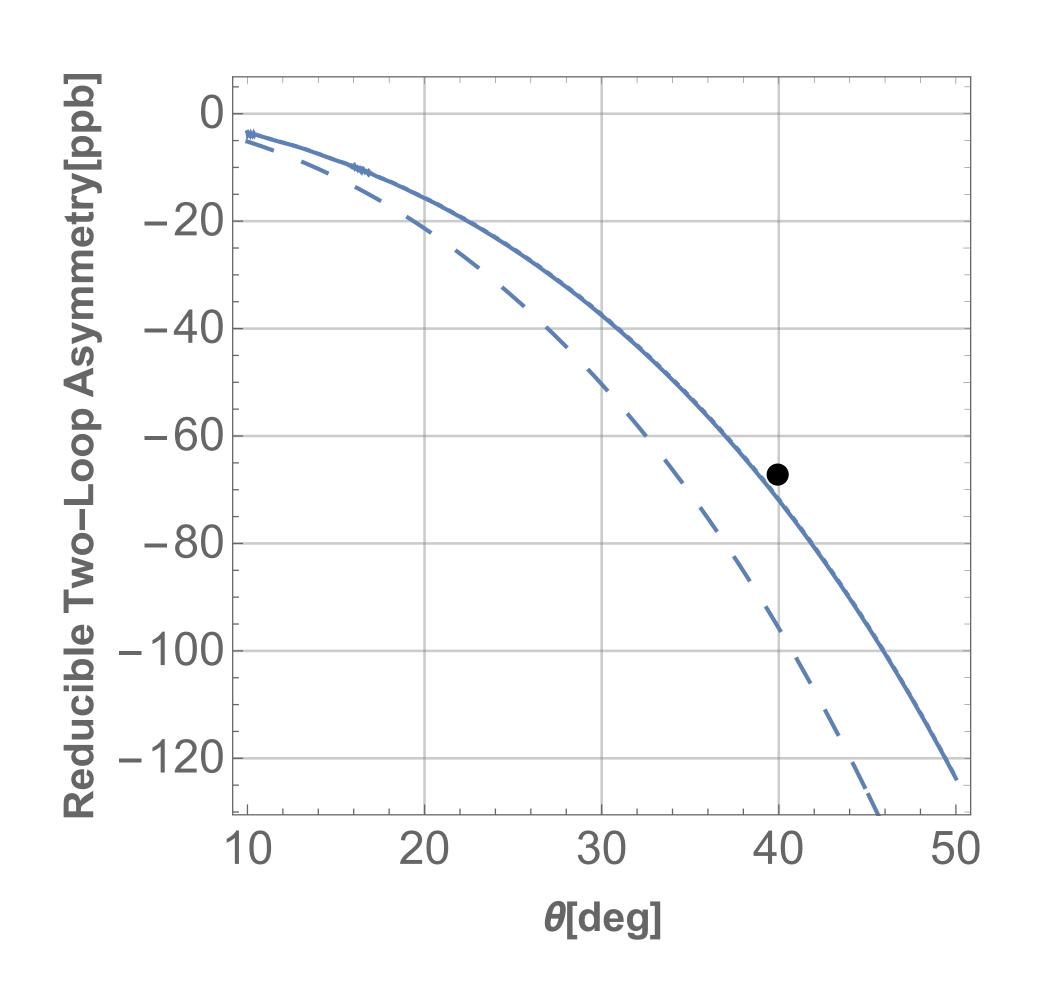
$$i\mathcal{M} = \bar{u}(k_2)(-ie(a_v\gamma_{\mu} + a_p\gamma_{\mu}\gamma_5))u(k_1)\frac{-i}{t - M_Z^2}\bar{u}(p_2)(-ie\Gamma_{Z-p}^{\mu})u(p_1)$$

$$\Gamma_{Z-p}^{\mu} = f_1^p(t)\gamma^{\mu} + \frac{i\sigma^{\mu\nu}q_{\nu}}{2M}f_2^p(t) + g_1^p(t)\gamma^{\mu}\gamma_5$$

 $f_1^p(t)$ ,  $f_2^p(t)$  and  $g_1^p(t) \rightarrow$  weak electric, magnetic and axial vector form factors.

$$a_{v} = rac{I^{3} - 2\sin^{2}\theta_{W}Q_{f}}{2\sin\theta_{W}\cos\theta_{W}}, \ a_{p} = rac{I^{3}}{2\sin\theta_{W}\cos\theta_{W}}, \ \text{where}$$
  $Q_{f} = -1(e^{-}) \ \text{and} \ I_{3} = -rac{1}{2}.$ 

# Tree level, NLO and NNLO level $A_{PV}$ for $e^-p$ scattering versus $\theta_{CM}$ using P2 kinematics



$$(\theta=39.97^0)$$
 ---- Tree  $A_{PV}\sim-95.6~ppb$  ---- NLO  $A_{PV}\sim-71.81~ppb$  ..... NNLO (Reducible Two loop)  $A_{PV}\sim-71.6~ppb$ 

P2 Proposed  $A_{PV} \sim -67.34 \ ppb$ 

$$A_{PV} = \frac{\sigma_R - \sigma_L}{\sigma_R + \sigma_L}$$

 $(e^-p)$  Tree level, NLO and NNLO level  $A_{PV}$  versus  $\theta_{CM}$ 

### REFERENCES FOR BOX DIAGRAMS

[1] M. Gorchtein, Phys. Rev. C **73**, 055201 (2006)

[2] Peter G. Blunden et al., Physical Review Letters 91(14)

