



Generalized Parton Distributions of low-lying octet baryons

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under the supervision of

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2 Generalized Parton Distributions



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Introduction

Quantum chromodynamics (QCD)

- an underlying theory for the strong interactions.
- describe the internal hadronic structure in terms of their constituent quarks and gluons degree of freedom.



Highly preferable, Lorentz-invariant form.

- four-vectors.
- fundamental concept in special theory of relativity.

So, Special theory of relativity becomes an important concept that must be follow.





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- A generic four Vector x^{μ} in light-cone coordinates is describe as $x^{\mu} = (x^{-}, x^{+}, x_{\perp}).$
- $x^+ = x^0 + x^3$ is called as light-front time.
- $x^- = x^0 x^3$ is called as light-front longitudinal space variable.
- $x^{\perp} = (x^1, x^2)$ is the transverse variable.
- Similarly, we can define the longitudinal momentum $k^+ = k^0 + k^3$ and light-front energy $k^- = k^0 - k^3$.

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What makes Front form dynamics a good choice ?

NO SQUARE ROOT FACTOR in energy-momentum dispersion relation

$$k^{-} = \frac{k_{\perp}^2 + m^2}{k^+},$$

where

- k^- : energy
- *k*⁺: longitudinal momentum
- k_{\perp} : transverse momentum
- Vacuum expectation value is ZERO.
- Seven out of ten fundamental quantities are KINEMATICAL.

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Distribution functions











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Generalized Parton Distributions

$$F_{\lambda\lambda'}^{[\Gamma]} = \frac{1}{2} \int \frac{dz^{-}}{2\pi} e^{ik \cdot z} \left\langle p', \lambda' \middle| \bar{\psi} \left(\frac{-z}{2}\right) \Gamma \mathcal{W} \left(\frac{-z}{2}, \frac{z}{2}\right) \middle| \mathcal{N} \right) \psi \left(\frac{z}{2}\right) \middle| p, \lambda \right\rangle \bigg|_{z^{+} = 0, \mathbf{z}_{\perp} = 0}$$

JHEP 08, 056 (2009)

with

- $\lambda(\lambda')$: Initial (final) state helicity.
- $P = \frac{p+p'}{2}$, average momentum.
- $\Delta = p' p$, momentum transfer.
- Impact parameter space, $\Delta_{\perp} \leftrightarrow b_{\perp}$ via FT.
- Γ basis $[1, \gamma_5, \gamma^{\mu}, \gamma^{\mu}\gamma_5, \sigma^{\mu\nu}].$

 $\Gamma = \gamma^+ : H \text{ and } E$ unpolarized $\Gamma = \gamma^+ \gamma^5$: \tilde{H} and \tilde{E} longitudinal polarized

 $\Gamma = i\sigma^{+i}\gamma^{5} :$ $H_T, E_T, \tilde{H}_T \text{ and } \tilde{E}_T$ transversely polarized

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Generalized Parton Distributions

$$\frac{1}{2\bar{P^+}}\bar{u}(P')\Big[H^q_X(x,0,\Delta_{\perp})\gamma^+ + E^q_X(x,0,\Delta_{\perp})\frac{i\sigma^{+\alpha}(-\Delta_{\alpha})}{2M}\Big]u(P)$$

Eur. Phys. J. A 52(2016)163.



 $\Gamma = \gamma^+ : H \text{ and } E$ unpolarized $\Gamma = \gamma^+ \gamma^5$: \tilde{H} and \tilde{E} longitudinal polarized $\Gamma = i\sigma^{+i}\gamma^{5} :$ $H_T, E_T, \tilde{H}_T \text{ and } \tilde{E}_T$ transversely polarized

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Momentum fraction carried by an active

quark, $x = \frac{k^+}{P^+}$ Skewness: Fraction of longitudinal momentum transfer, $\zeta = -\frac{\Delta^+}{2P^+}$ Momentum transfer squared, $t = (p' - p)^2 = \Delta^2$

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Generalized Parton Distributions

How to access GPDs ?



Phys. Rep. 388, 41-277 (2003)

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Generalized parton distributions (GPDs) For zero skewness

- Reduced to PDFs, $\int f^q(x, 0, t) dt \Rightarrow$ PDFs
- Pathway to FFs $\int f^q(x, 0, t) dx \Rightarrow$ FFs.

$$f^q(x,0,t) \xrightarrow{FT} f^q(x,0,\mathbf{r}_{\perp})$$

 $\int f^q(x,0,\mathbf{r}_{\perp}) \, dx \Rightarrow \text{FT FFs.}$

- 3D hybrid distribution
 - 1D momentum
 - 2D coordinate space



Phys. Rept. 418, 1 - 387(2005).

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Charge and Magnetization distributions

$$\rho_X^q(x,b_{\perp}) = \int \frac{d^2 \Delta_{\perp}}{(2\pi)^2} e^{-i\Delta_{\perp} \cdot \mathbf{b}_{\perp}} H_X^q(x,0,\Delta_{\perp})$$

$$\tilde{\rho}_M^{X_q}(\mathbf{b}_{\perp}) = \mathbf{b}_{\perp} S in^2(\phi) \int \frac{\Delta_{\perp}^2 d\Delta_{\perp}}{2\pi} J_1(\Delta_{\perp} \mathbf{b}_{\perp}) \int dx \, E_X^q(x, 0, \Delta_{\perp}) \,.$$

Phys. Lett. B **669**, 345-351 (2008) Nucl. Phys. B 969, 115440 (2021) Eur. Phys. J. A **60**, 42 (2024)

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Hadron state; Quark-Scalar diquark on complete basis of Fock-states

$$\begin{split} |\psi_A(P^+,\mathbf{P}_{\perp})\rangle &= \sum_{\mathcal{N}} \prod_{i=1}^{\mathcal{N}} \frac{dx \ d^2 \mathbf{k}_{\perp i}}{2(2\pi)^3 \sqrt{x_i}} \ 16\pi^3 \ \delta \Big(1 - \sum_{i=1}^{\mathcal{N}} x_i\Big) \delta^{(2)} \Big(\sum_{i=1}^{\mathcal{N}} \mathbf{k}_{\perp i}\Big) \\ &\times \psi_{\mathcal{N}}(x_i,\mathbf{k}_{\perp i},\lambda_i) |\mathcal{N}; x_i P^+, x_i \mathbf{P}_{\perp} + \mathbf{k}_{\perp i},\lambda_i\rangle, \end{split}$$

Nucl. Phys. B 596, 99-124 (2001)

k_{$\perp i$}- intrinsic transverse momentum

■ λ_i - helicity

 $J^{z} = +\frac{1}{2}$

$$\psi_{+\frac{1}{2}}^{\uparrow X}(x,\mathbf{k}_{\perp}) = \left(M_{X} + \frac{m_{q}}{x}\right)\varphi_{X},$$

$$\psi_{-\frac{1}{2}}^{\uparrow X}(x,\mathbf{k}_{\perp}) = -\frac{(k^{1} + ik^{2})}{x}\varphi_{X}.$$
 (1)

 $J^z = -\frac{1}{2},$

$$\psi_{\pm\frac{1}{2}}^{\downarrow X}(x,\mathbf{k}_{\perp}) = \frac{(k^{1}-\iota k^{2})}{x}\varphi_{X},$$

$$\psi_{-\frac{1}{2}}^{\downarrow X}(x,\mathbf{k}_{\perp}) = \left(M_{X} + \frac{m_{q}}{x}\right)\varphi_{X}.$$
 (2)

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- **\mathbf{k}_{\perp i}** intrinsic transverse momentum
- λ_i helicity

 $J^z = +\frac{1}{2},$

$$J^z = -\frac{1}{2},$$

 $\psi_{\pm\frac{1}{2}}^{\uparrow X}(x,\mathbf{k}_{\perp}) = \left(M_{X} + \frac{m_{q}}{x}\right)\varphi_{X}, \qquad \qquad \psi_{\pm\frac{1}{2}}^{\downarrow X}(x,\mathbf{k}_{\perp}) = \frac{(k^{1} - \iota k^{2})}{x}\varphi_{X}, \\ \psi_{-\frac{1}{2}}^{\uparrow X}(x,\mathbf{k}_{\perp}) = -\frac{(k^{1} + \iota k^{2})}{x}\varphi_{X}. \quad (1) \qquad \qquad \psi_{-\frac{1}{2}}^{\downarrow X}(x,\mathbf{k}_{\perp}) = \left(M_{X} + \frac{m_{q}}{x}\right)\varphi_{X}. \quad (2)$





2 Generalized Parton Distributions



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GPD $H(x, 0, \Delta_{\perp})$; u - flavorMomentum space



GPD $H(x, 0, \Delta_{\perp})$; *s* – *flavor* Momentum space



GPD $E(x, 0, \Delta_{\perp})$; u - flavorMomentum space



GPD $E(x, 0, \Delta_{\perp})$; *s* – *flavor* Momentum space





"Curiosity is the beginning of knowledge. Action is the beginning of Change." -James Clear