Introduction to Small-x Physics

(Part 2)

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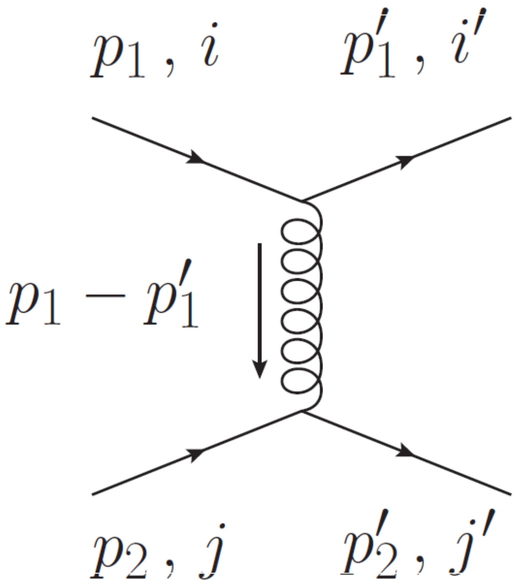
NM STATE

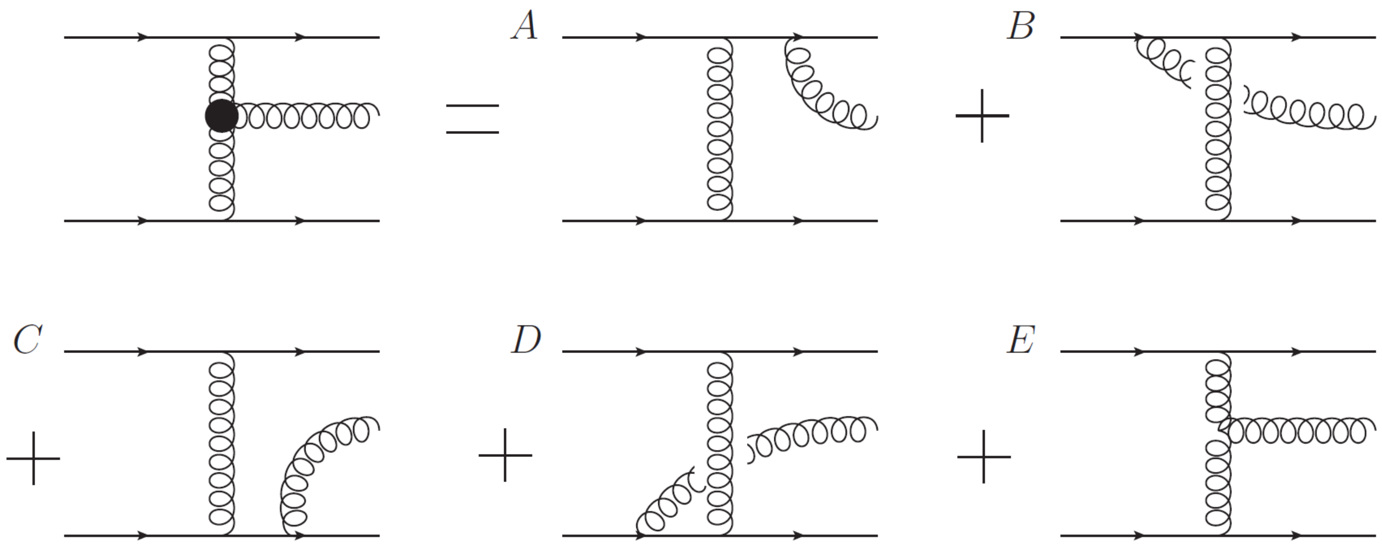
CFNS Summer School 2024 *6/13/2024*



$$i\mathcal{M} \approx 2ig^2 \left(\frac{s}{t}\right)\, (t^a)_{i'i}(t^a)_{j'j}\, \delta_{s_1 s_1'}\, \delta_{s_2 s_2'}$$

$$\frac{d\sigma}{d^2 p_{1\perp}'} = 4\alpha_s^2\, \frac{C_F}{2N_c}\, \frac{1}{p_{1\perp}'^{\,4}}$$

Previously...

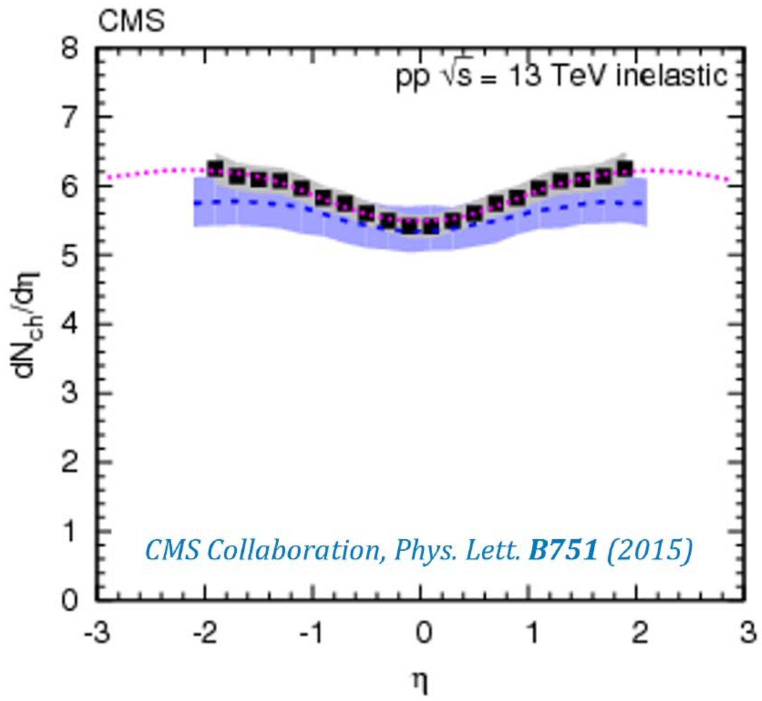


$$iM_A + iM_B = -g^3 f^{abc} (t^c)_{i'i} (t^b)_{j'j} \frac{2s}{q_\perp^2} \frac{(\epsilon_\lambda^*)^+}{k^+} \delta_{s_1 s_1'} \delta_{s_2 s_2'}$$

$$iM = -g^3 f^{abc} (t^c)_{i'i} (t^b)_{j'j} \frac{4s}{q_\perp^2} \; (\epsilon_\lambda^*)_\mu \left(\frac{k_\perp^\mu}{k_\perp^2} - \frac{(k+q)_\perp^\mu}{(k+q)_\perp^2} \right) \delta_{s_1 s_1'} \delta_{s_2 s_2'}$$

$$\frac{d\sigma}{d^2k_\perp \, d^2q_\perp \, dy} = \frac{2\alpha_s^3 C_F}{\pi^2} \frac{1}{q_\perp^2 \, k_\perp^2 \, (k+q)_\perp^2}$$

Previously...



$$\frac{dN^g}{d^2k_\perp\, dy} \equiv \frac{1}{\sigma_{el}} \frac{d\sigma}{d^2k_\perp\, dy} = \frac{2\alpha_s N_c}{\pi^2} \frac{m^4}{k_\perp^4} \ln \frac{k_\perp^2}{m^2}$$

$$\frac{dN^g}{dy} = \frac{2\alpha_s N_c}{\pi^2}$$

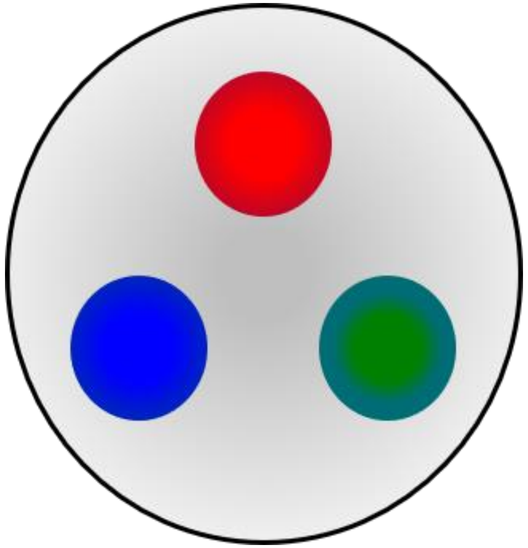
$$N^g_{tot} = \frac{2\alpha_s N_c}{\pi^2} \ln \frac{s}{m^2}$$

Deep Inelastic Scattering:

Accessing Proton Structure

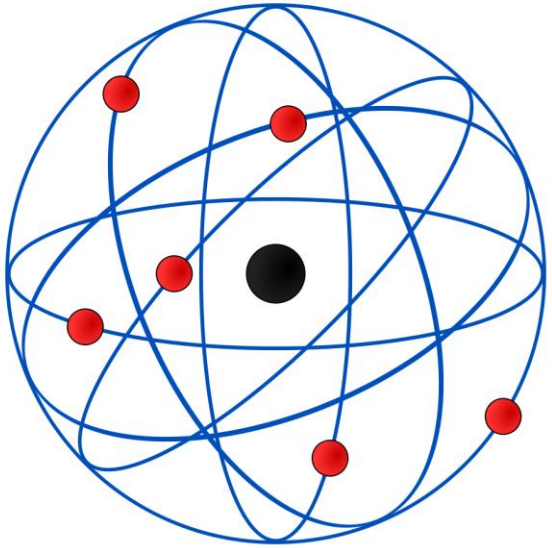
A Tale of Two Systems

The Proton



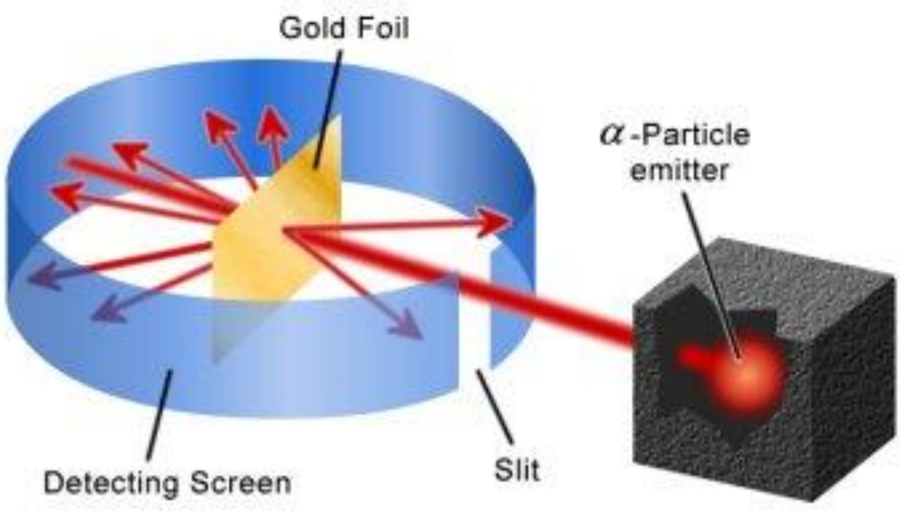
Bound together by
the **strong nuclear force**

The Atom



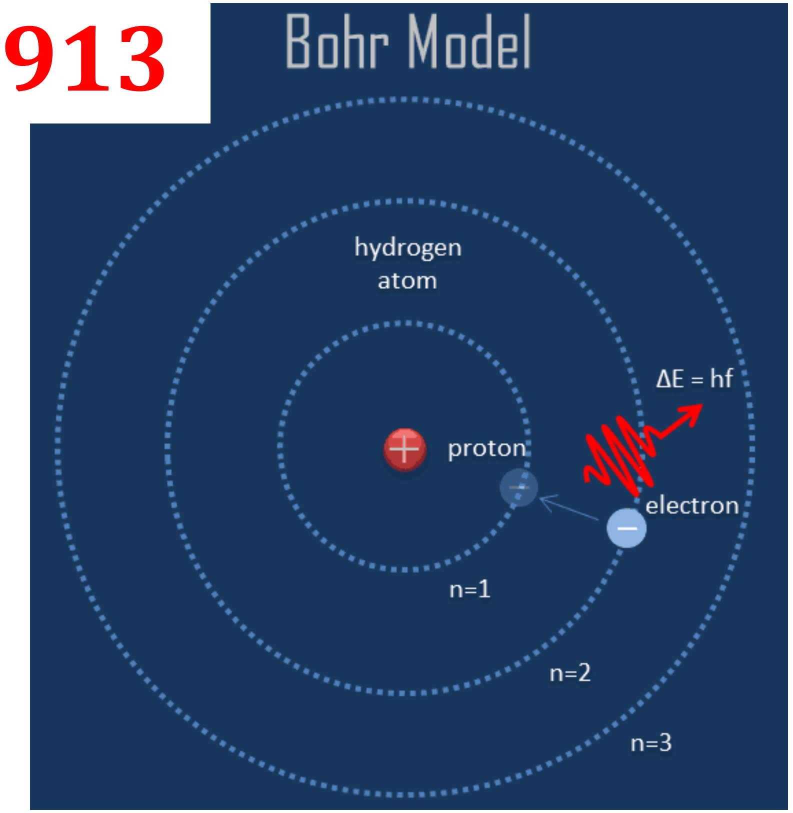
Bound together by
the **electromagnetic force**

1908 - 1917



Rutherford Gold Foil Experiment

1913

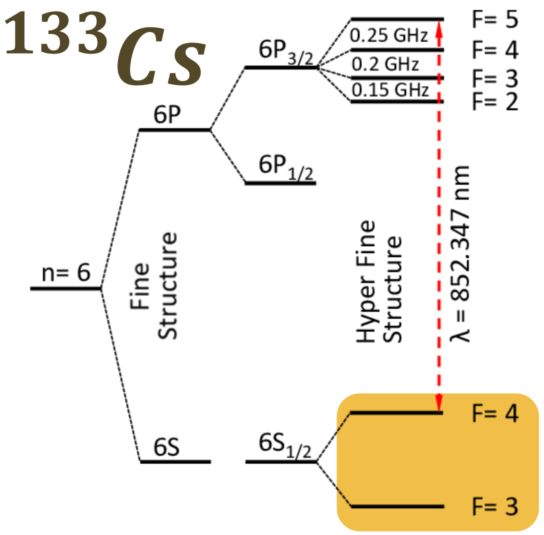


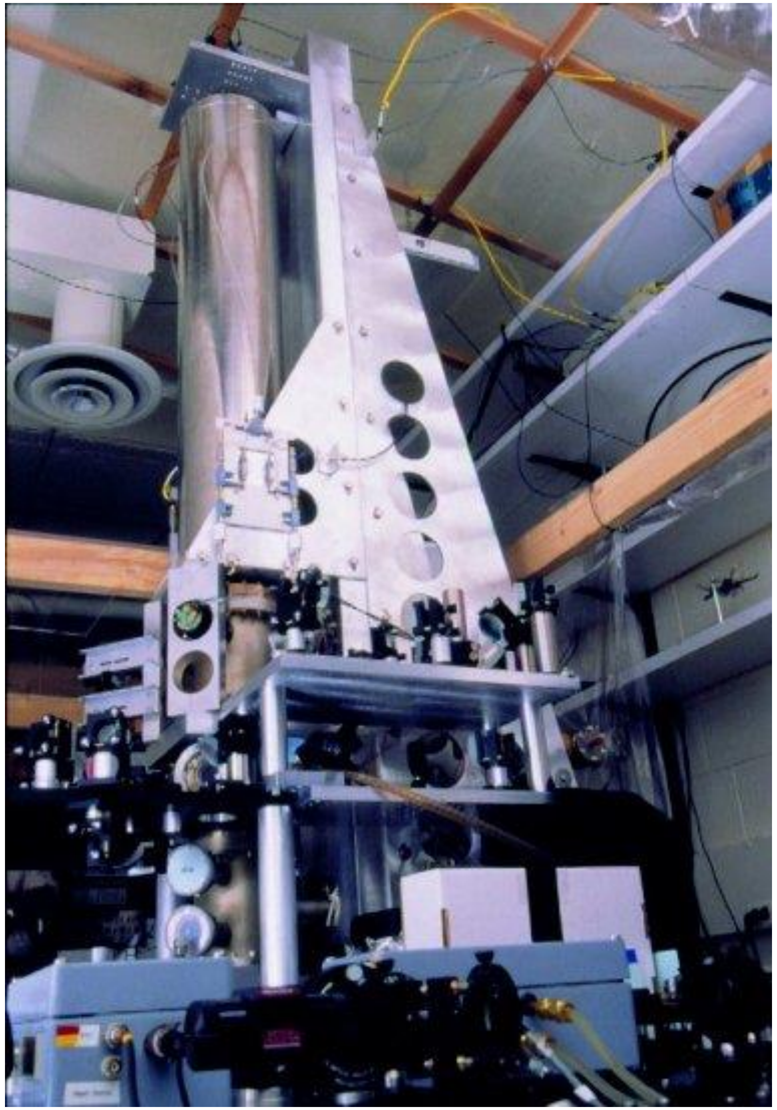
Bohr Model of the Atom

- The atom is used as the **definition of the second**

Combined Fractional Uncertainty: $0.52\text{x}10^{-15}$ **!!!**

Mar. 2016, NIST-F1 BIPM Report





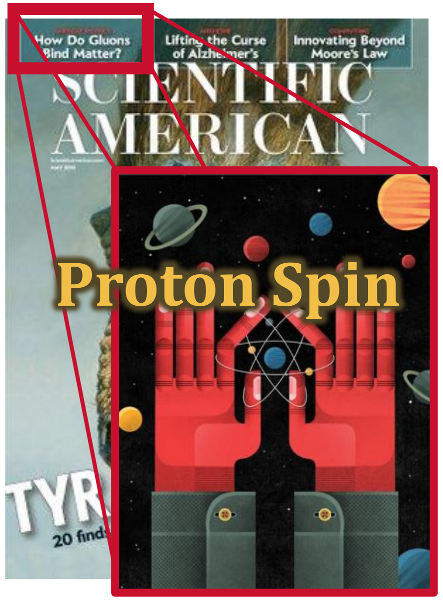
- Calculations to **extreme precision**

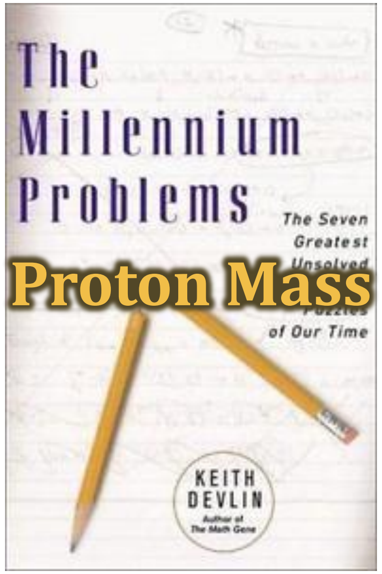
I have evaluated up to 1100 digits of precision the contribution of the 891 4-loop Feynman diagrams contributing to the electron g-2 in QED. The total mass-independent 4-loop contribution is

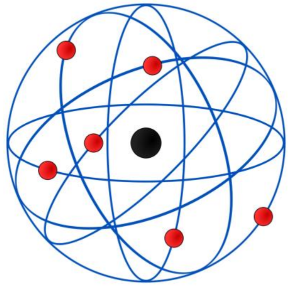
$$a_e = -1.912245764926445574152647167439830054060873390658725345\ldots\left(\frac{\alpha}{\pi}\right)^4 .$$

S. Laporta, Phys. Lett. ***B772*** *(2017) 232*

- Even **fundamental questions** about the **structure of the proton** remain **unanswered**



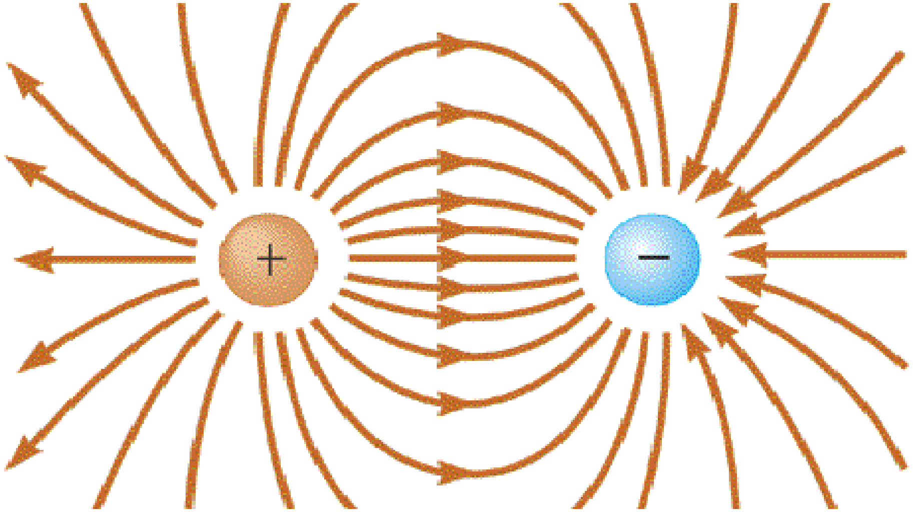




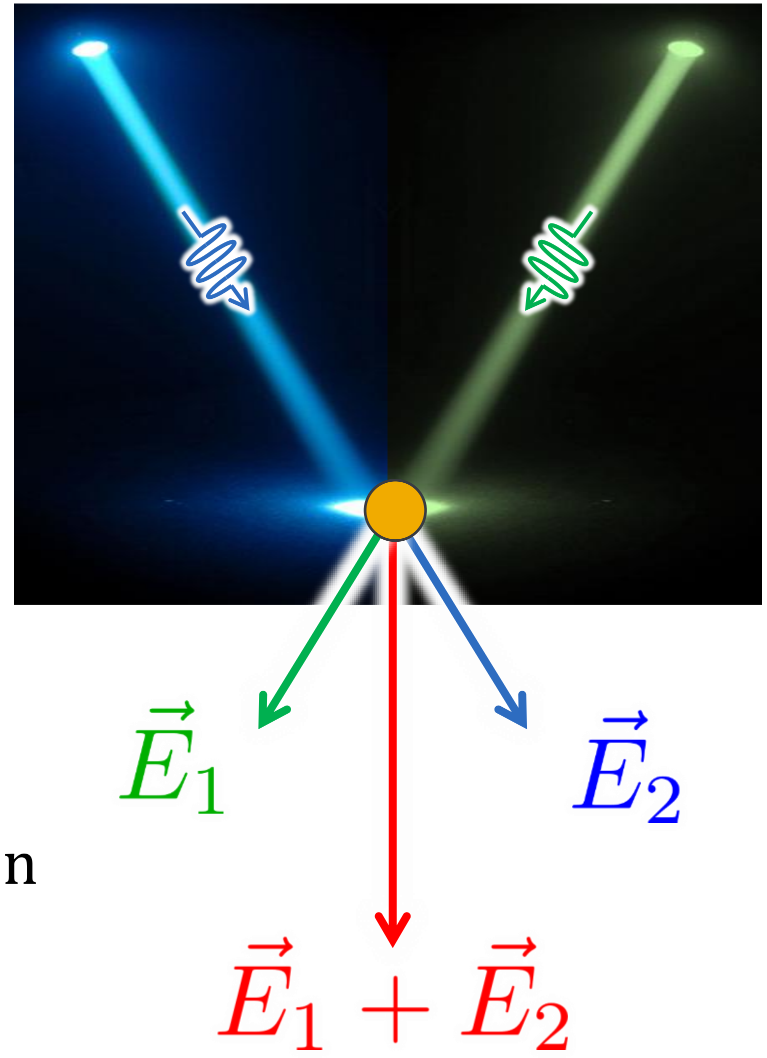
Atom: Electrodynamics

electrons

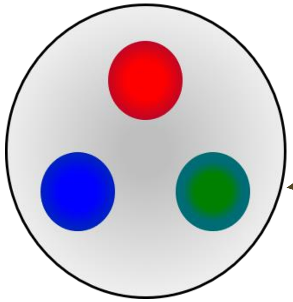
- **Charges** (electrons) radiate **fields** (photons)
- Electric charge is a **scalar** **(+/−)**



Linear:
Superposition
Principle



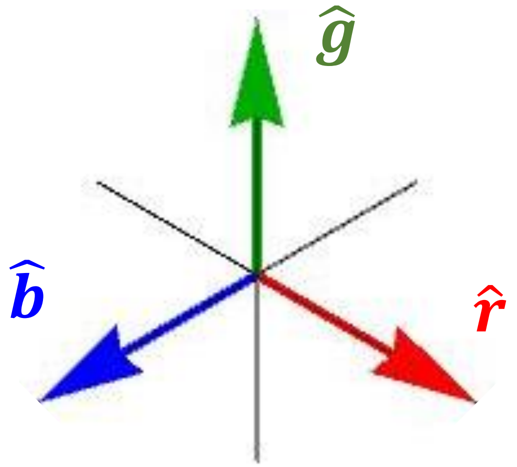
Chromo-Dynamics: One Crucial Difference

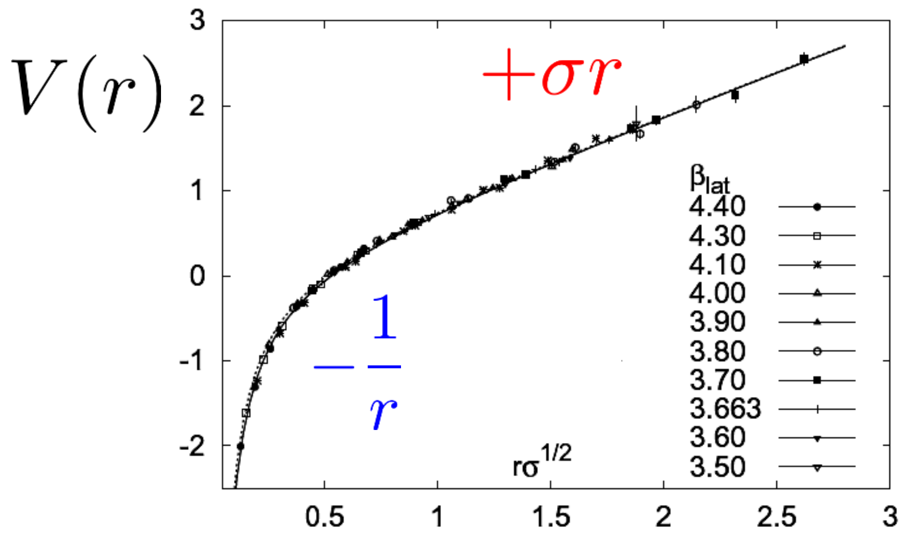


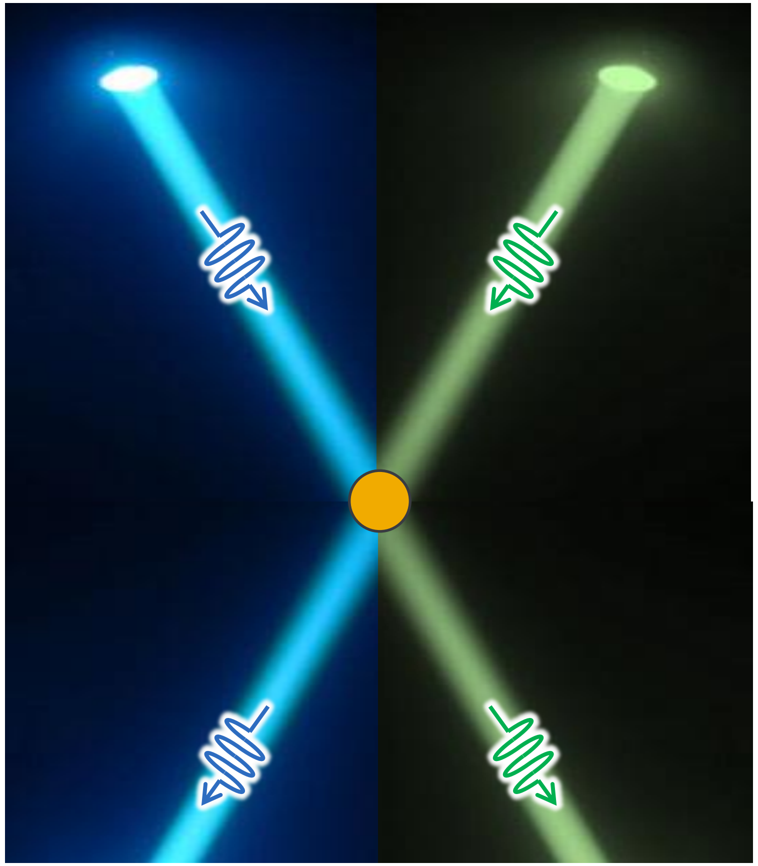
Proton: Chromodynamics

"quarks"

- **Charges** (quarks) radiate **fields** (gluons)
- Color charge is a **vector**



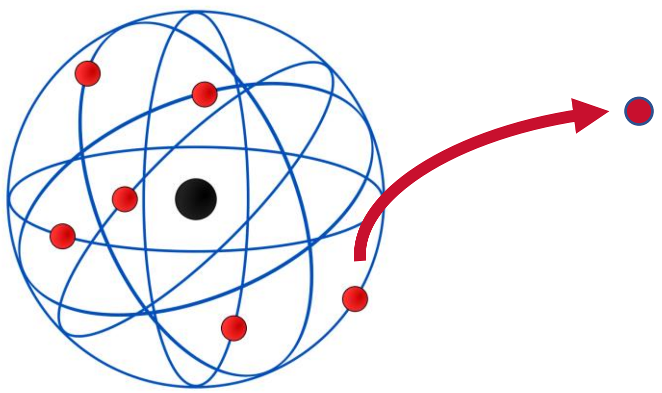


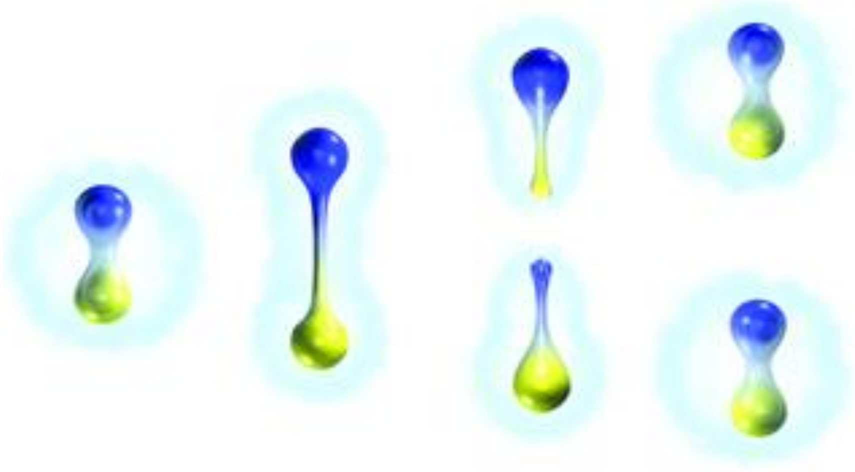


Non-Linear:
Self-interactions of fields

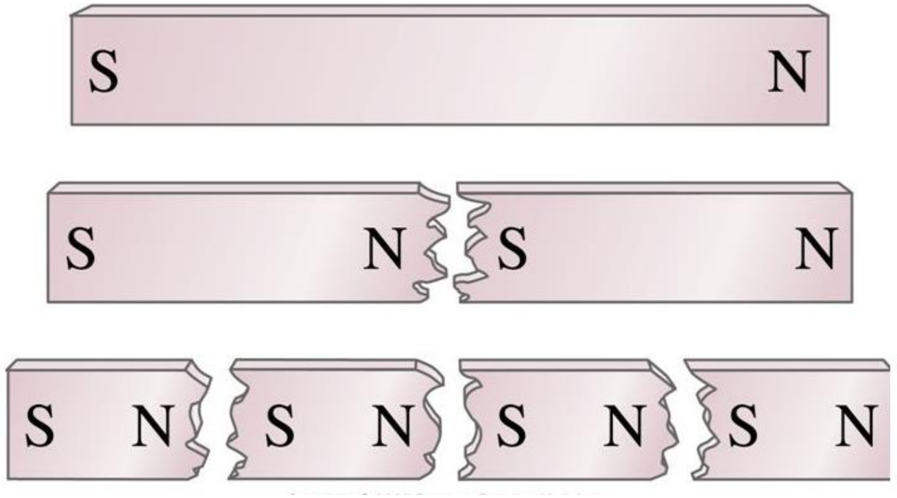
Dissecting the Proton is Hard…

- You **can't** just **"ionize" a quark** out of the proton
  - If you **hit it hard enough** to knock out a quark…
  - You create a **shower of new particles** instead!



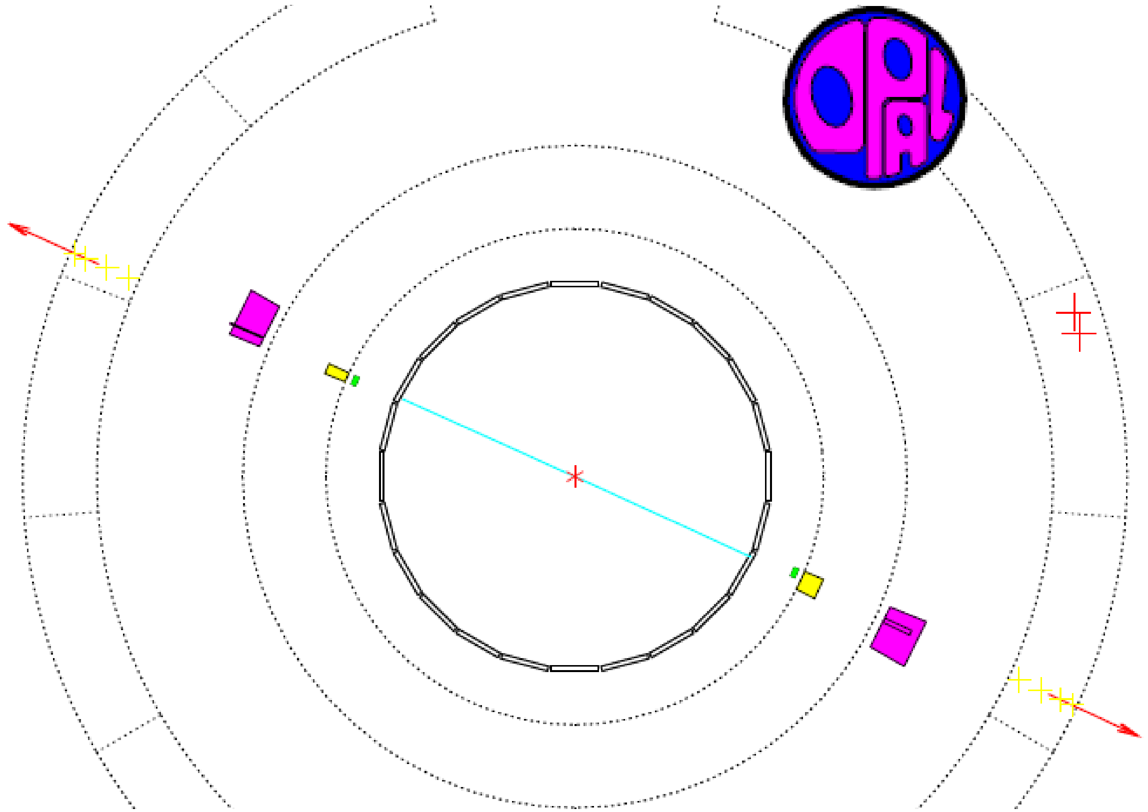


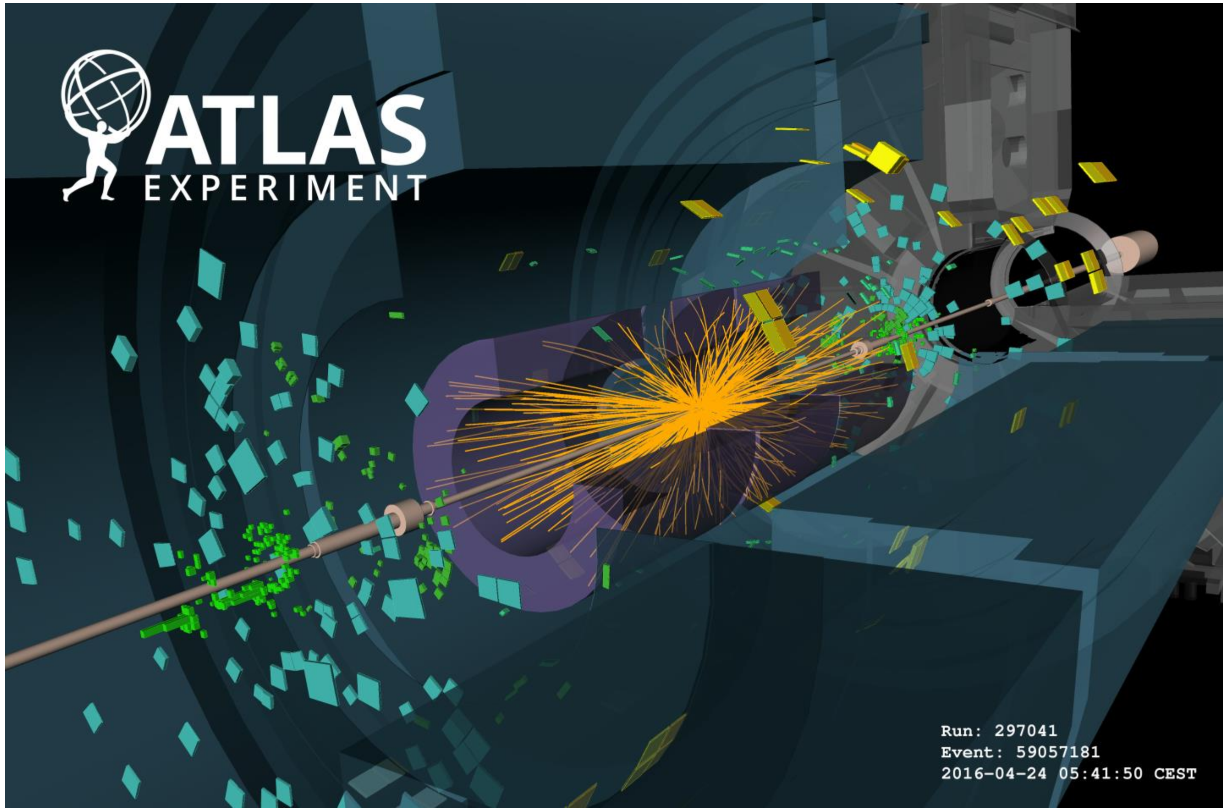
No Isolated Quarks



No Magnetic Monopoles

...and Messy!





E&M: Very **clean**

QCD: **TONS** of radiation

How Do We Measure Proton Structure?

- An **optical microscope** is limited by the **wavelength of visible light**:

  $E \sim 2\ eV$ $\Delta x \geq 100\ nm$

- A **scanning electron microscope** uses a thermally produced **beam of electrons**:

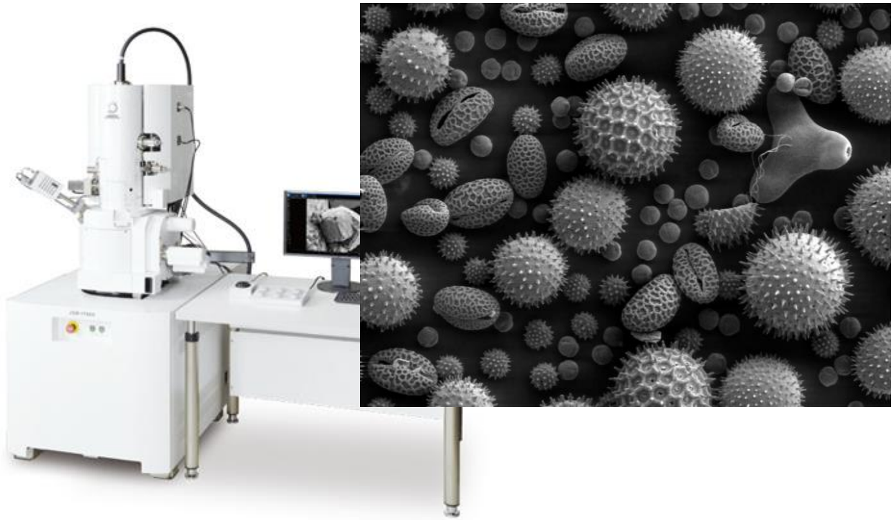
  $E \sim 10\ keV$ $\Delta x \geq 1\ nm$

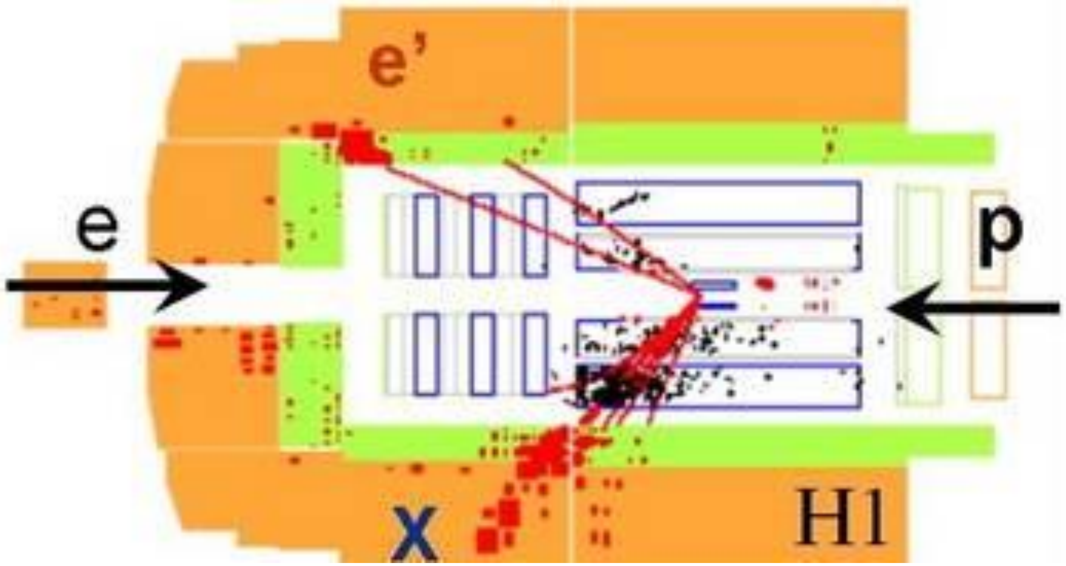
- An **electron-proton collider** can use the same principles at **top collider energy**.

$$e + p \;\rightarrow e' + \boldsymbol{X}$$

$E \gg 1\ GeV$ $\Delta x \ll 1\ fm$

D *eep*
I *nelastic*
S *cattering*





Evolution of the DIS Femtoscope

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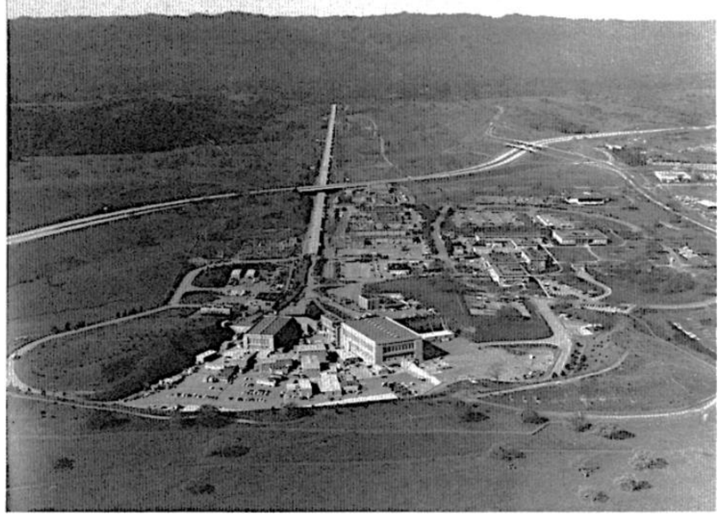
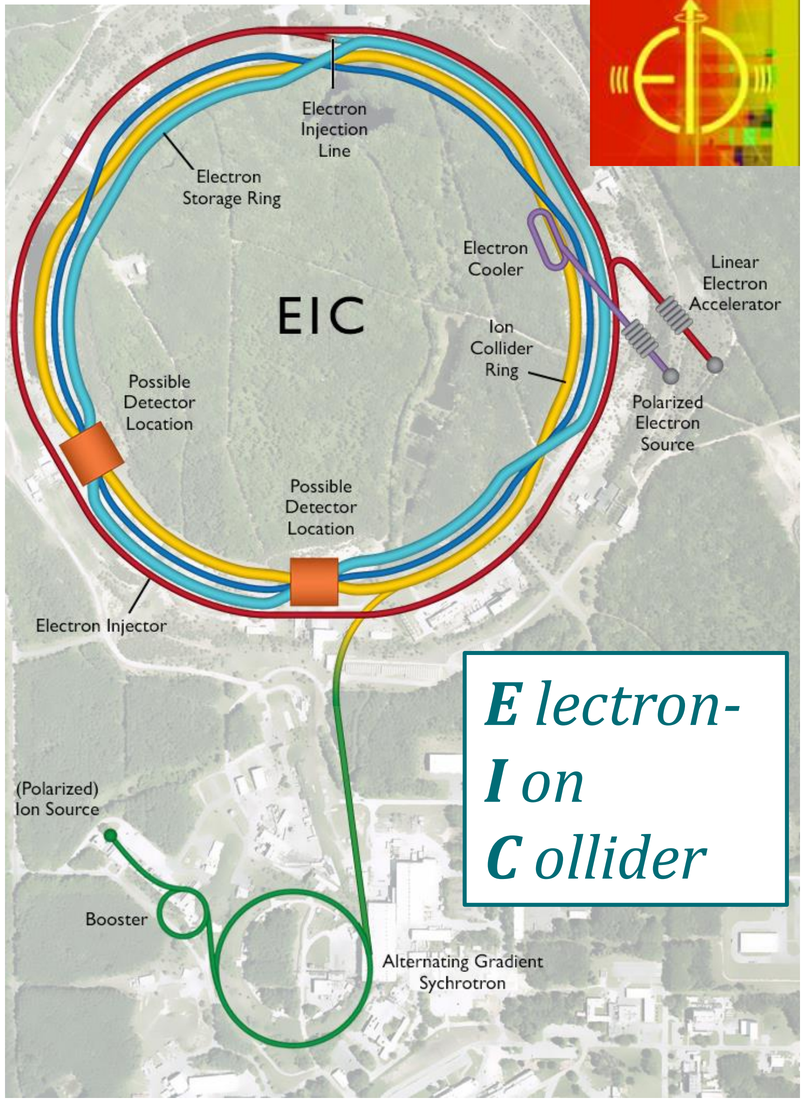




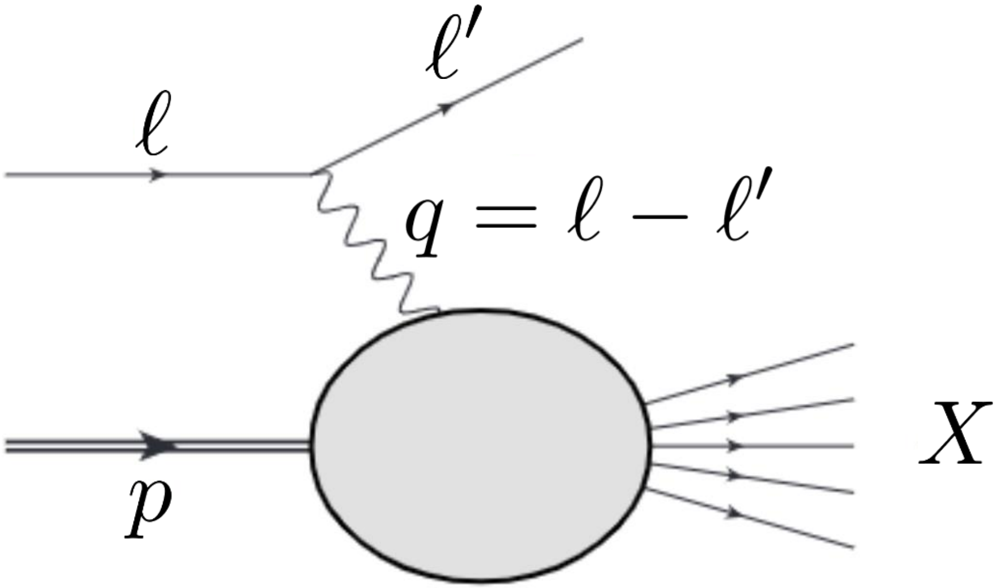
***E** lectron-*
***I** on*
***C** ollider*

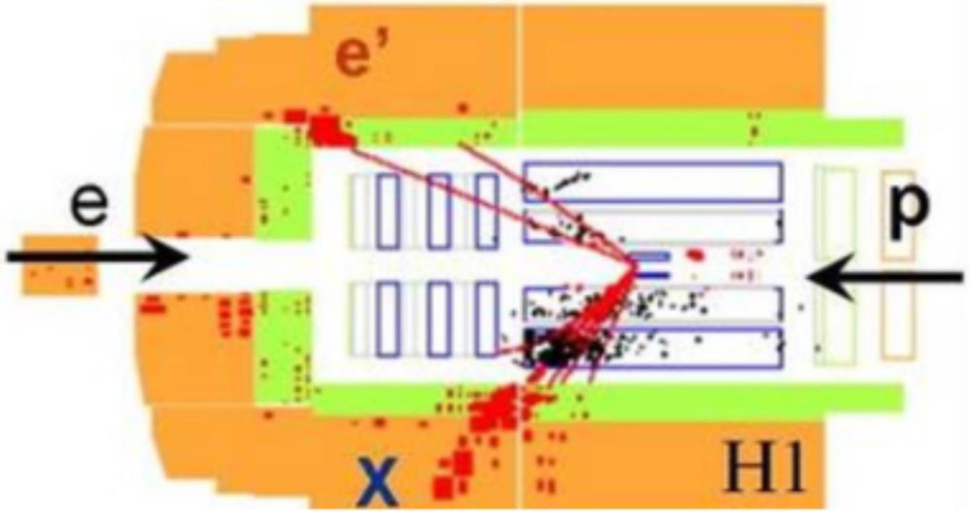
- 1968: **SLAC-MIT Experiment** — Fixed target, $E_{beam} \sim 20\ GeV$
- 1992 – 2007: **HERA** — $E_{CM} = 318\ GeV$
- 2030+ **Electron-Ion Collider** — $E_{CM} = 140\ GeV$ + polarization + ions

Theory of DIS

(What does "x" mean, anyway?)

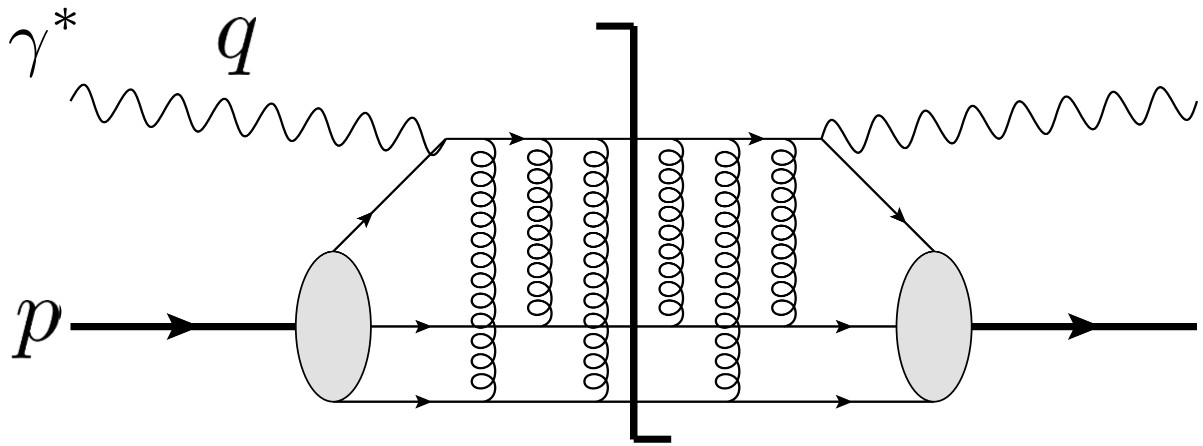
Using QED to Measure QCD





- **Inclusive** electron/proton scattering: $e + p \rightarrow e' + \boldsymbol{X}$
  - ➢ Electron acts as a **source of virtual photons** via well-controlled **QED vertex**
  - ➢ Couples to **electromagnetic currents** $\sim A_\mu \langle X|\, j^\mu |p\rangle$ **inside the proton**

Field-Theoretic Description of DIS

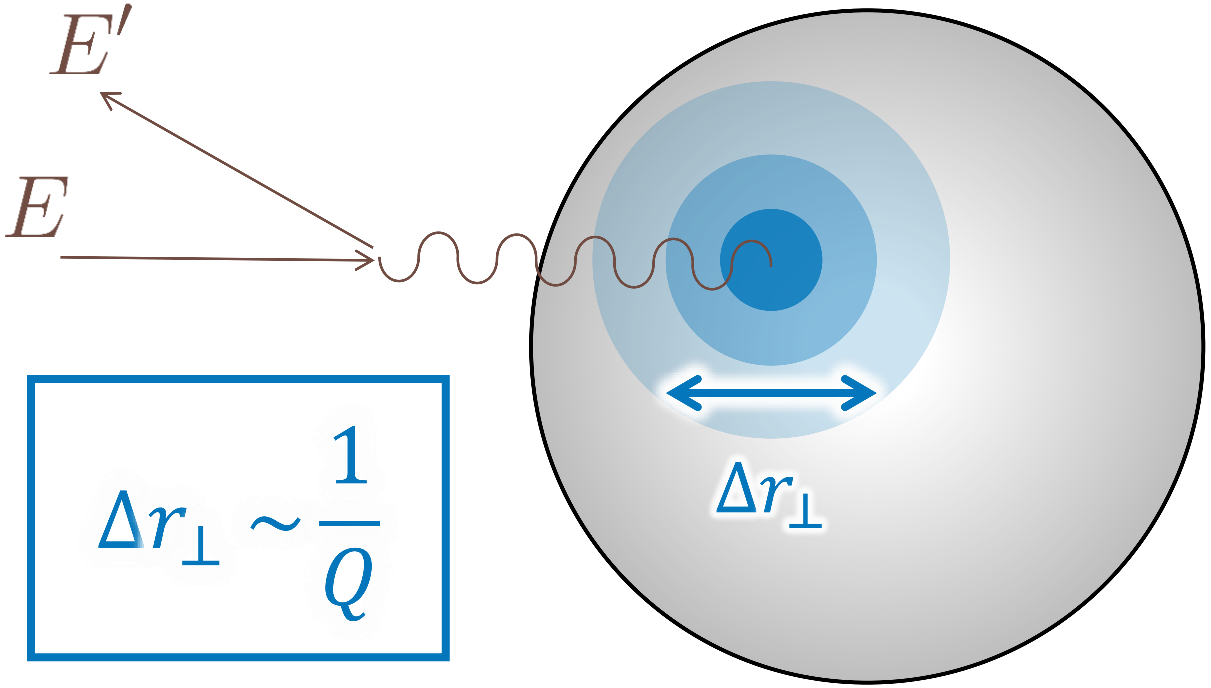


$$E'_\ell \frac{d\sigma}{d^3\ell'} = \frac{\alpha_{EM}^2}{E_\ell Q^4} L_{\mu\nu} W^{\mu\nu}$$

- Trivial **QED part** can be calculated and **removed**: $L_{\mu\nu} = \frac{1}{2} \sum_{ss'} [\bar{u}_{p's'} \gamma_\mu u_{ps}][\bar{u}_{ps} \gamma_\nu u_{p's'}]$
- Nontrivial QCD info from $\boldsymbol{\gamma^* p \to X}$ **subprocess** : $W^{\mu\nu} = \frac{1}{4\pi m} \int d^4x \, e^{iq\cdot x} \langle p| \, j^\mu(x) j^\nu(0) \, |p\rangle$
  - ➢ **Tensor-valued** function of **two momenta**: p^μ, q^μ
  - ➢ **Decompose** $W^{\mu\nu}$ into invariant **structure functions** depending on **invariants**

Image Resolution:

$$Q^2 \equiv -q_\mu q^\mu = 4EE' \sin^2 \frac{\theta}{2}$$



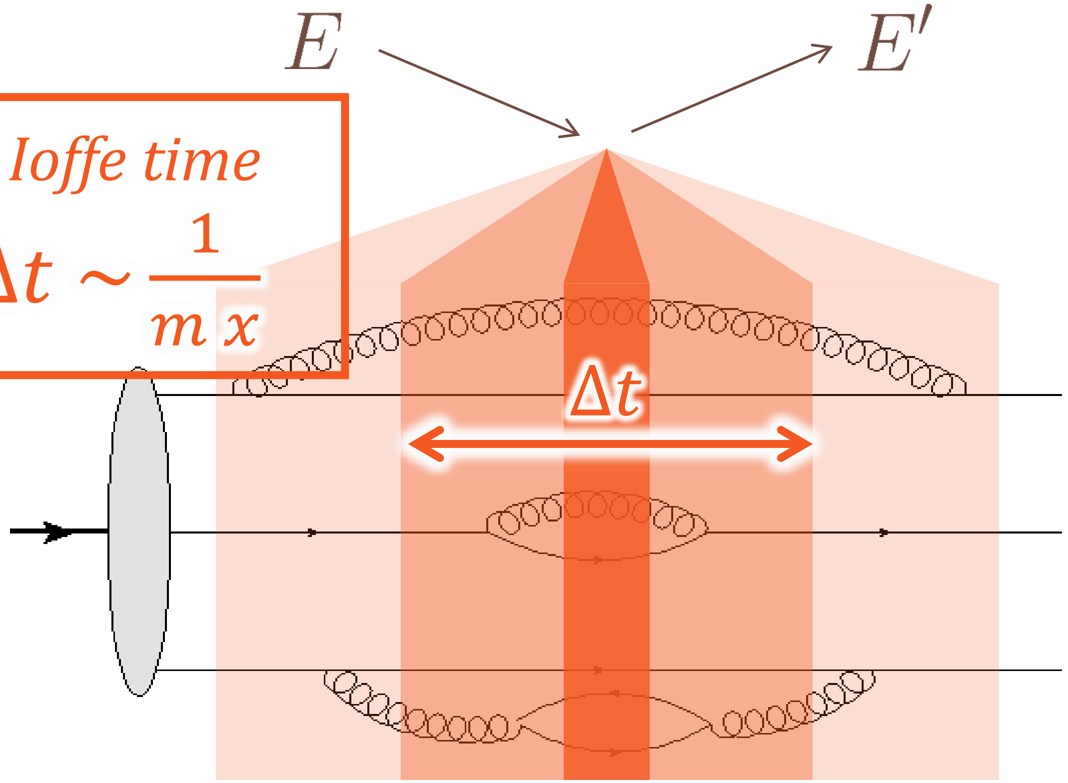
$$\boxed{\Delta r_\perp \sim \frac{1}{Q}}$$

Exposure Time:

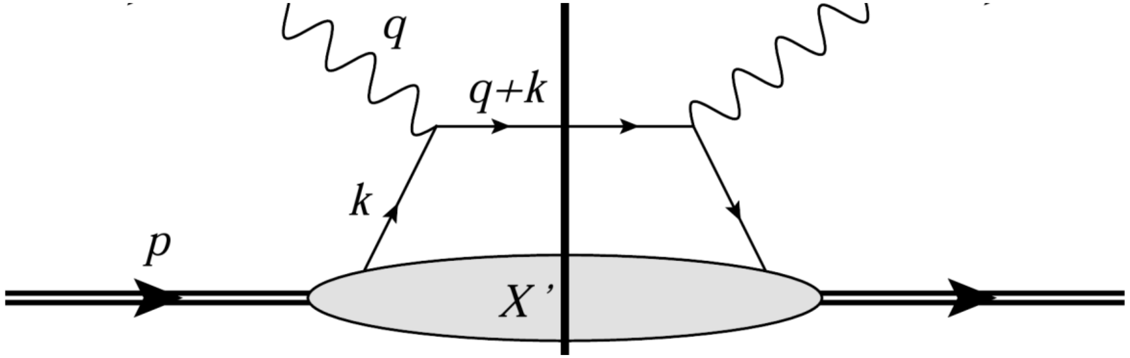
$$x_B \equiv \frac{Q^2}{2p \cdot q} \overset{R.F}{=} \frac{Q^2}{2m(E - E')}$$

Ioffe time

$$\Delta t \sim \frac{1}{m\,x}$$



Feynman x Versus Bjorken x



$$d\sigma \sim \int \frac{d^3k}{(2\pi)^3 2E_{q+k}} \langle |M(k)|^2 \rangle$$

$$\sim \int \frac{d^4k}{(2\pi)^4} \langle |M(k)|^2 \rangle \boxed{\delta\left((q+k)^2\right)} \theta(q^0 + k^0)$$

$$\delta\left((q+k)^2\right) \approx \delta\left(2(x_F - x_B)p^+ q^-\right)$$

$$= \frac{1}{s} \boxed{\delta(x_F - x_B)}$$

$$x_B \equiv \frac{Q^2}{2p \cdot q} \qquad x_F \equiv \frac{k^+}{p^+}$$

- Breit frame:

$$p^\mu = \left[p^+ \, , \, 0^- \, , \, \vec{0}_\perp\right]^\mu$$

$$q^\mu = \left[-x_B p^+ \, , \, q^- \, , \, \vec{0}_\perp\right]^\mu$$

$$k^\mu = \left[x_F p^+ \, , \, 0^- \, , \, \vec{k}_\perp\right]^\mu$$

$$(q+k)^\mu = \left[(x_F - x_B)p^+ \, , \, q^- \, , \, -\vec{k}_\perp\right]^\mu$$

➢ DIS is a “snapshot” of the proton wave function (squared)

Quarks are For Real

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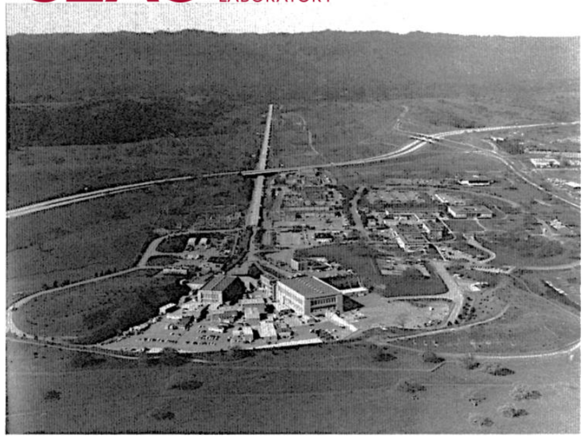
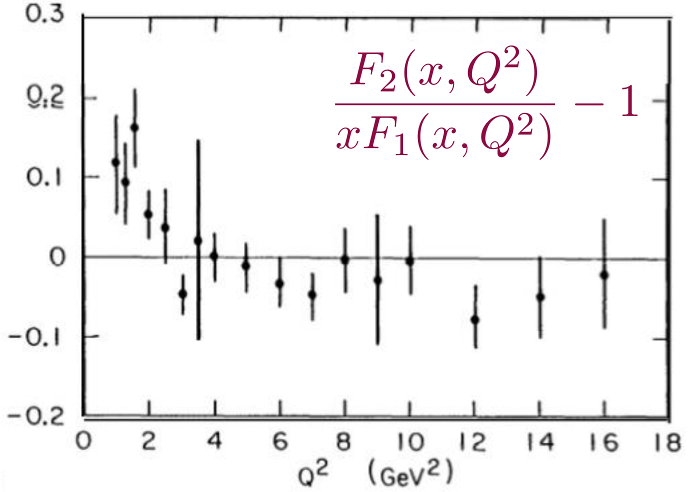
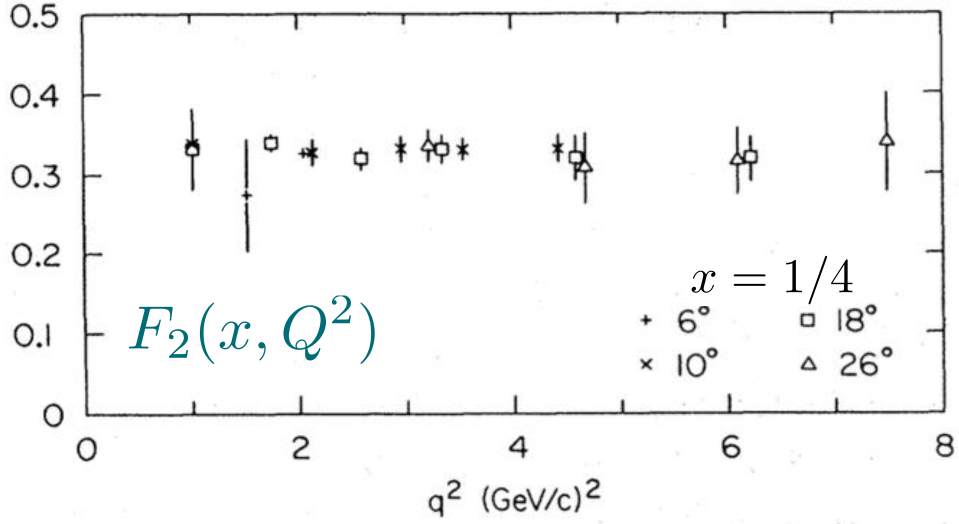


Callan-Gross Relation



Bjorken Scaling



- Proton described by two **structure functions:**

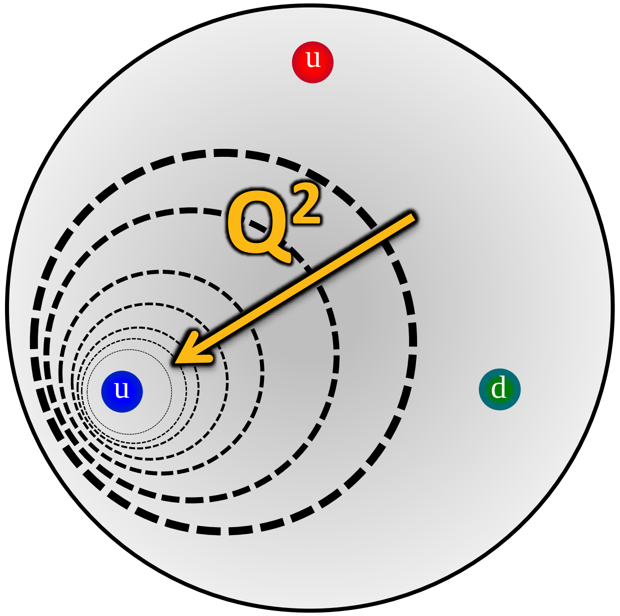
$$\frac{d\sigma}{d\Omega\, dE'} = \sigma_{point} \left[\frac{2mx}{Q^2} F_2(x,Q^2) + \frac{1}{m} F_1(x,Q^2) \tan^2\frac{\theta}{2} \right]$$

- ➢ Proportional to scattering on **point-like particles**
- ➢ F_1 and F_2 are not independent: **Spin ½ fermions (Quarks!!)**
- ➢ F_1 and F_2 are independent of Q^2: **Point-like at any resolution (???)**

The Missing Link: Asymptotic Freedom

Accurate taxonomy of hadrons…

Bjorken scaling…



Pauli exclusion problem in Ω^- …

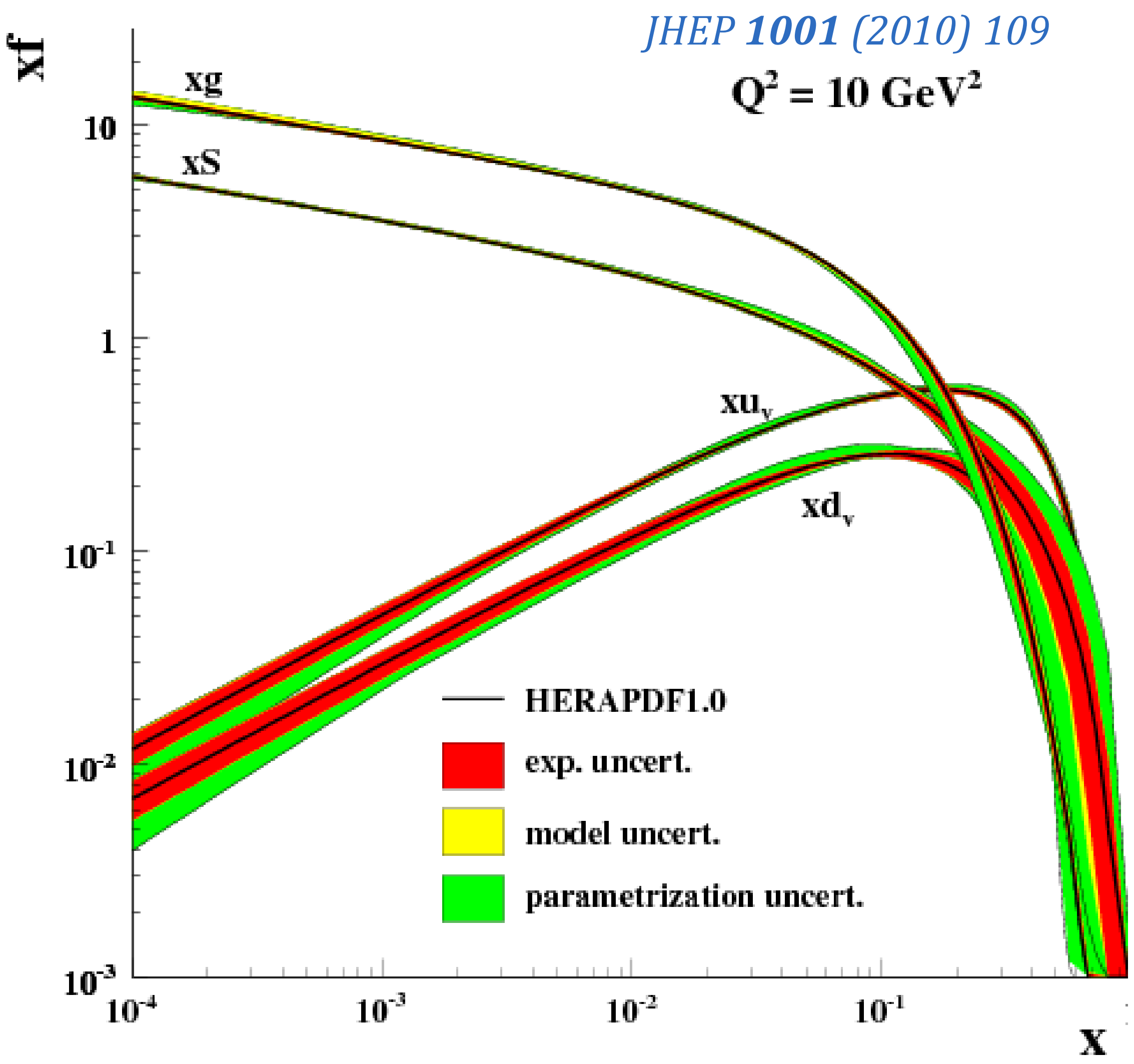
Free quarks never found…

- **Bjorken scaling: No change** in proton structure with **increased resolution**

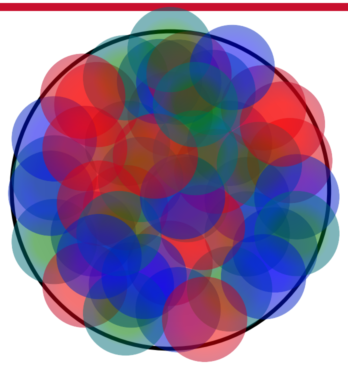
- **Uncertainty principle:** the **largest** quantum fluctuations happen over **short times and distances**…

  - ➢ … Unless the **interactions** of the nuclear force **go to zero** at short distances!
  - ❖ Bjorken scaling demands **asymptotic freedom** – and QCD

The 1D Picture of the Proton



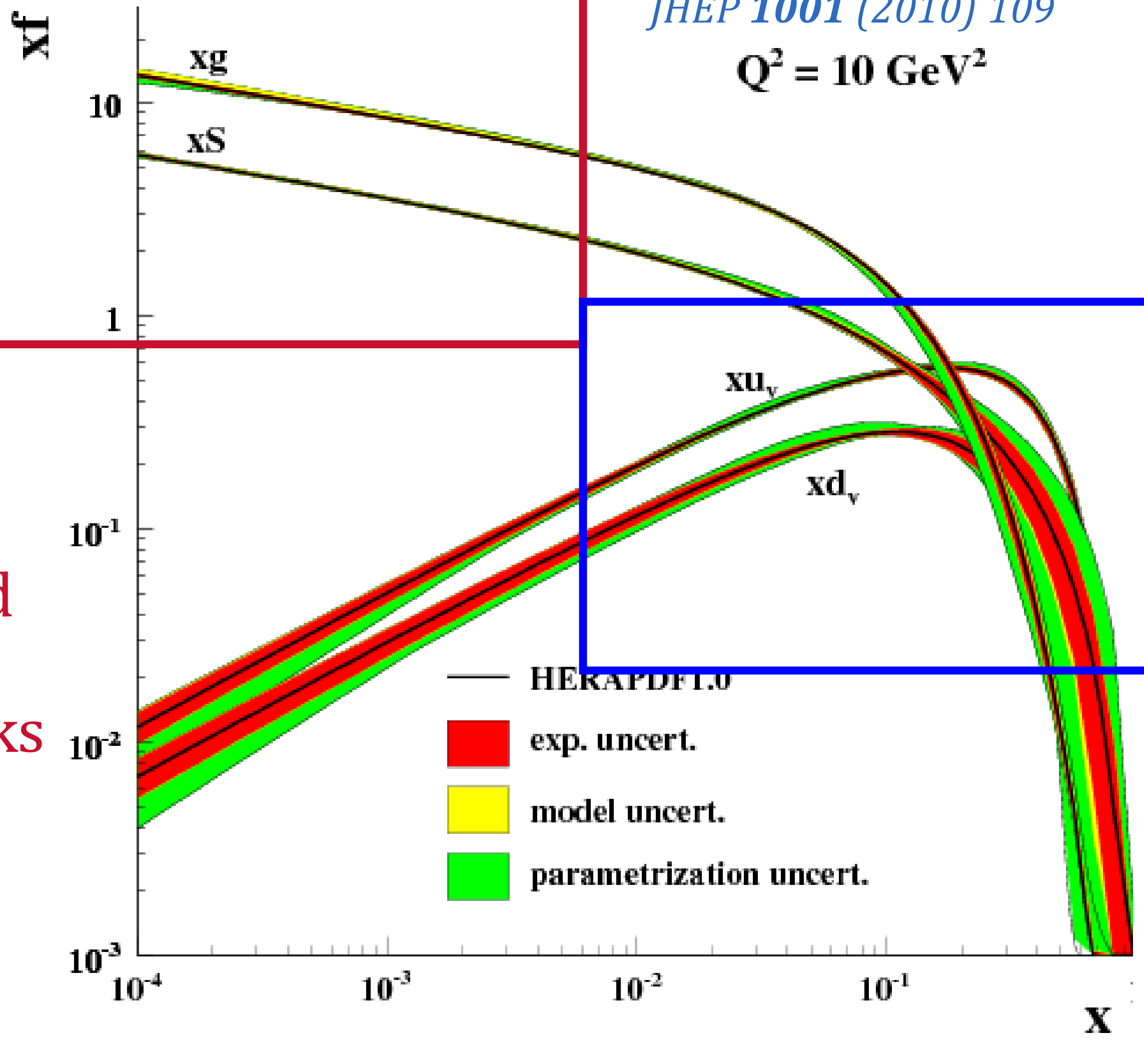
Two Very Different Regimes of Proton Structure



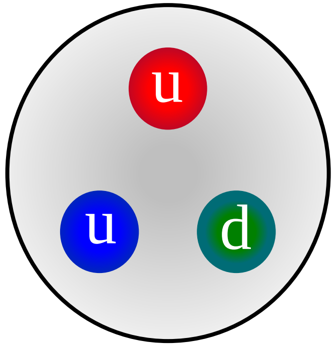
Dense Proton

- Small $x \ll 1$
- Radiation-dominated
- Gluons and sea quarks

JHEP ***1001*** *(2010) 109*

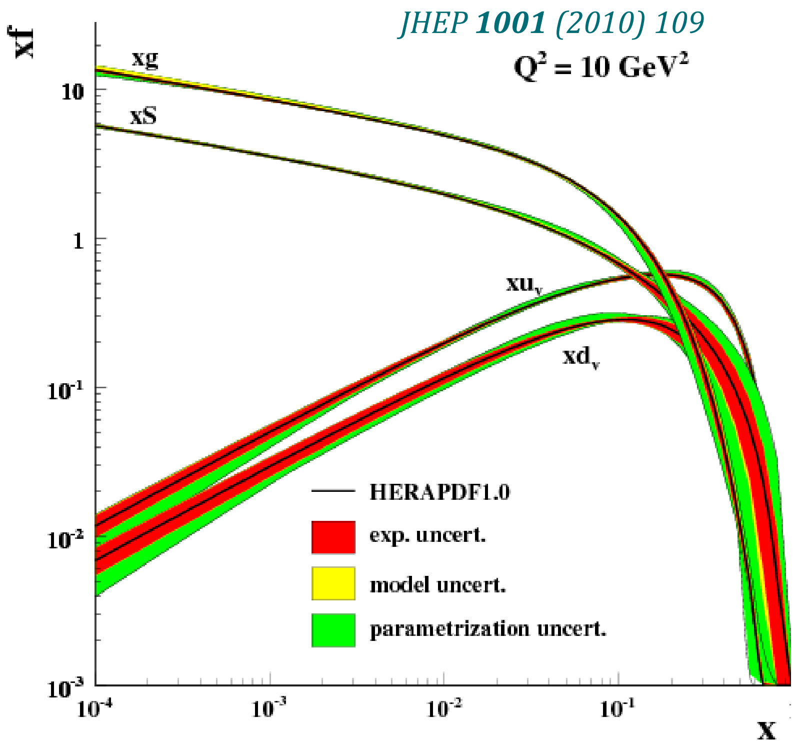


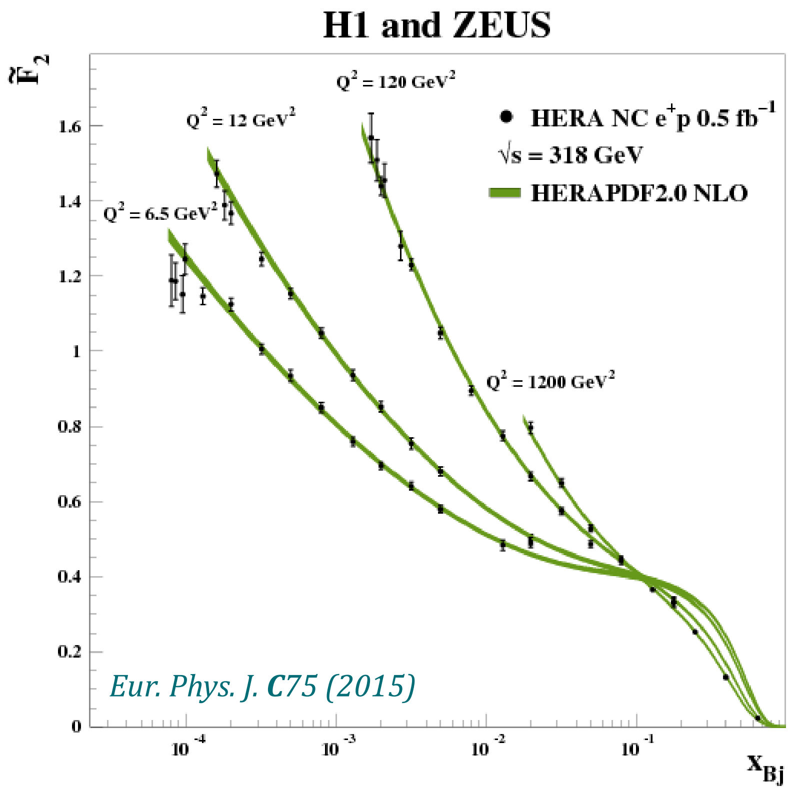
Dilute Proton

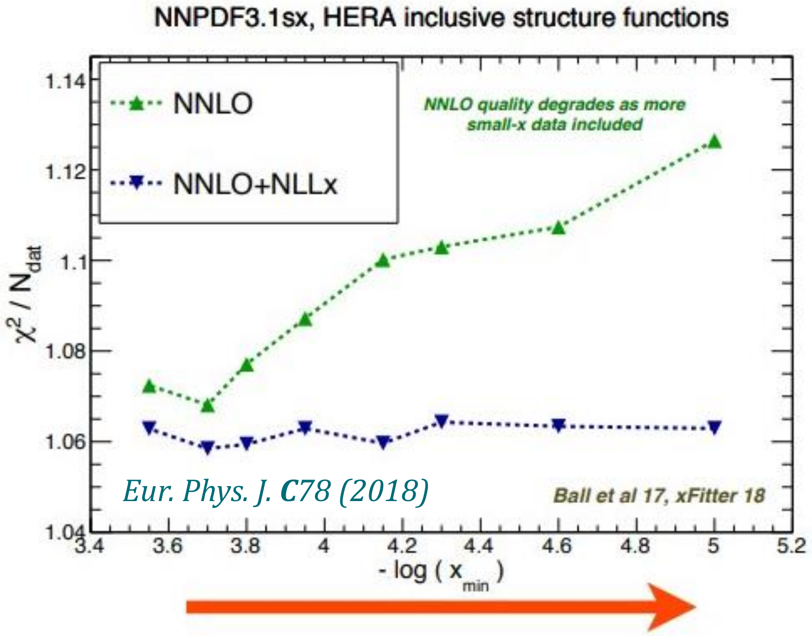


- Large $x \sim \mathcal{O}(1)$
- Valence region

Something Changes at Small x







- The **dense proton** at **small x** reflects an **explosion of soft gluon** bremsstrahlung:

Systematic tension in the HERA fits at **small x**...

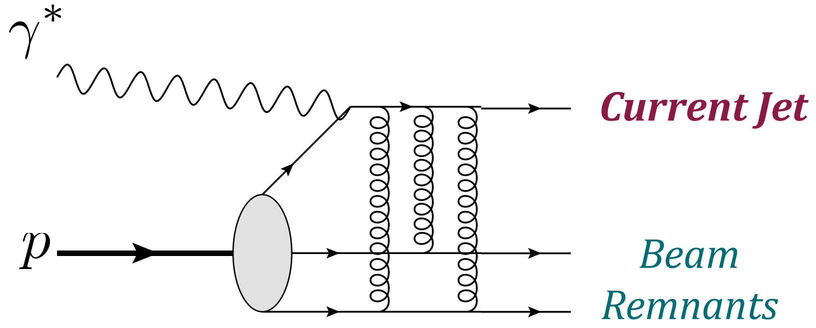
...reflect the **onset of new degrees of freedom**

DIS at Small x

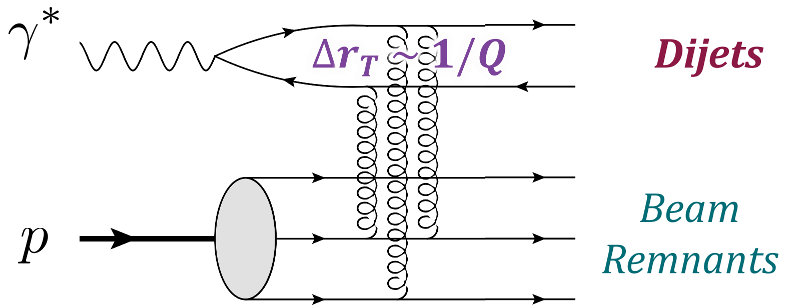
The Intersection of Structure and Hadronic Scattering

Onset of Small-x Degrees of Freedom in DIS

"Knockout" DIS at Large x:



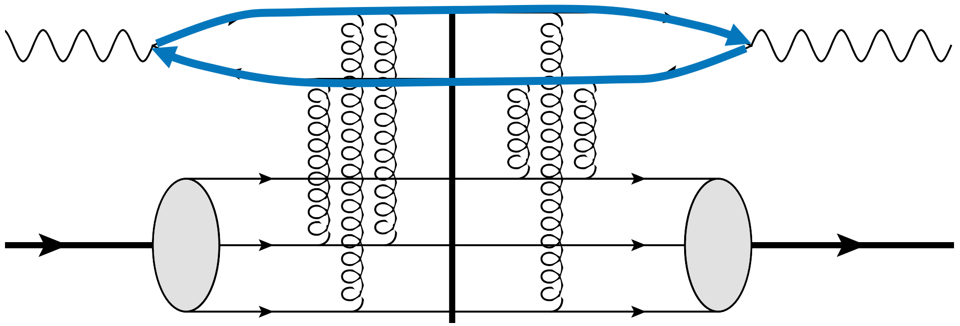
"Dipole" DIS at Small x:



- At **small x**, DIS resembles a **hadronic scattering** process
  ("baby" pp collision)
  - γ^* fluctuates into a **$q\bar{q}$ dipole** with **variable size**

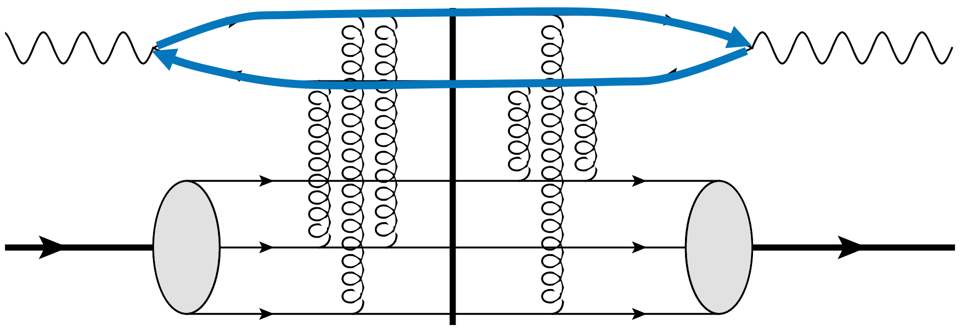
$$V_{x_\perp} = \mathcal{P} \exp\left[ig \int_{-\infty}^{\infty} dz^- \, A^{+a}(0^+, z^-, \vec{x}_\perp) t^a \right]$$

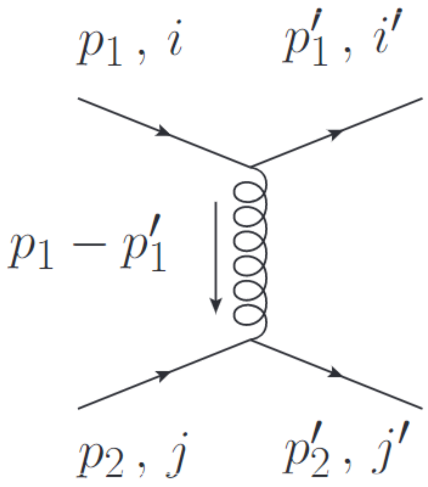
$$\langle \mathrm{tr} \left[V_{x_\perp} V^\dagger_{y_\perp} \right] \rangle$$



- New degrees of freedom: **Wilson lines**
  - **Boosted, coherent** scattering in a **background field**

$$\langle \mathrm{tr}\left[V_{x_\perp} V^\dagger_{y_\perp}\right]\rangle$$

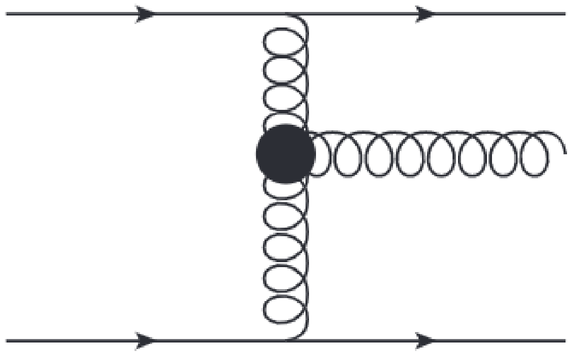




- ❖ Elastic q/q scattering (LO)
  - ➢ Unsuppressed
  - ➢ Wilson line!

$$V_{x_\perp} = \mathcal{P}\exp\left[ig\int_{-\infty}^{\infty} dz^- \; A^{+a}(0^+, z^-, \vec{x}_\perp)t^a\right]$$

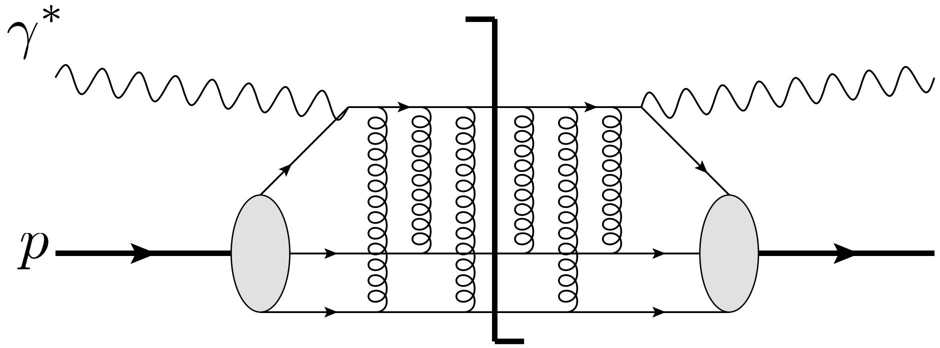
- ❖ Naturally suited to (light-front) "time"-ordered perturbation theory
  - ➢ Wave functions × operators (Wilson lines)



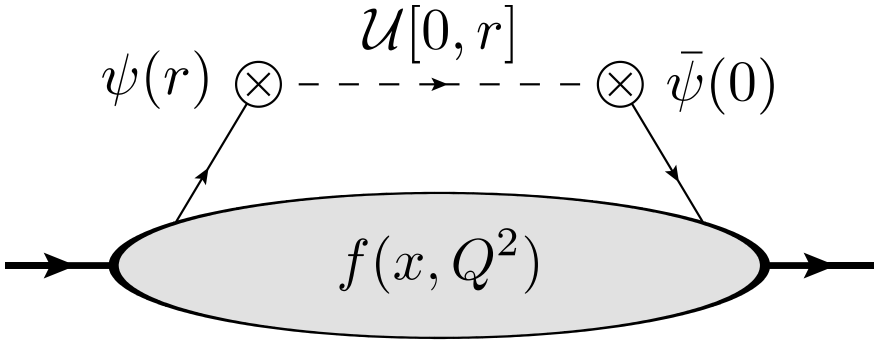
- ❖ Plus soft gluon radiation (NLO)
  - ➢ Large logs…
  - ➢ Resummation…?

- **Factorization** provides a one-to-one correspondence between the DIS **cross section** and hadronic structure: **PDFs**

$$\frac{Q^2}{4\pi^2\alpha_{EM}}\underline{\frac{d\sigma^{(\gamma^* p)}}{dx\, dQ^2}} = F_2(x,Q^2) \overset{L.O.}{=} \sum_f e_f^2\, x\underline{q_f(x,Q^2)}$$



$$\underline{q(x,Q^2)} = \int \frac{dr^-}{2\pi}\, e^{ixp^+r^-}\, \langle p|\, \bar\psi(0)\, \mathcal{U}[0,r^-]\, \frac{\gamma^+}{2}\, \psi(r^-)\, |p\rangle$$



- At small x, the DIS cross section, and therefore the PDFs themselves, are expressed in terms of **dipole scattering amplitudes**

$$x q_f(x, Q^2) = \frac{Q^2 N_c}{4\pi^2 \alpha_{EM}} \int \frac{d^2 x_{10}\, dz}{4\pi z(1-z)} \sum_{L,T} \left|\Psi(x_{10}^2, z)\right|^2 \int d^2 b_{10} \left[2 - \frac{1}{N_c} \left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right] \right\rangle_{(zs)} - \frac{1}{N_c} \left\langle \mathrm{tr}\left[V_1 V_0^\dagger\right] \right\rangle_{(zs)} \right]$$

Photon splitting wave functions

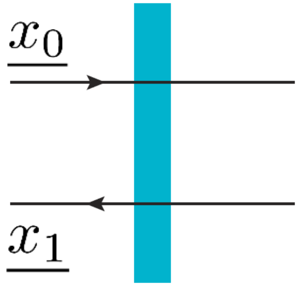
Non-interacting terms

Dipole amplitudes

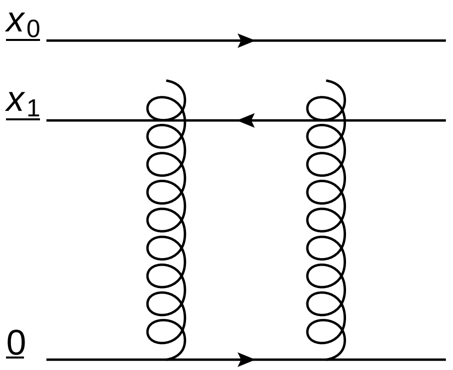
Initial Conditions

- The **x-dependence of the PDF** (TMD) is governed by the **energy dependence of the dipole amplitude**

$$\frac{1}{N_c}\left\langle \mathrm{tr}\left[V_{\underline{0}} V_{\underline{1}}^\dagger\right]\right\rangle_{(zs)}$$

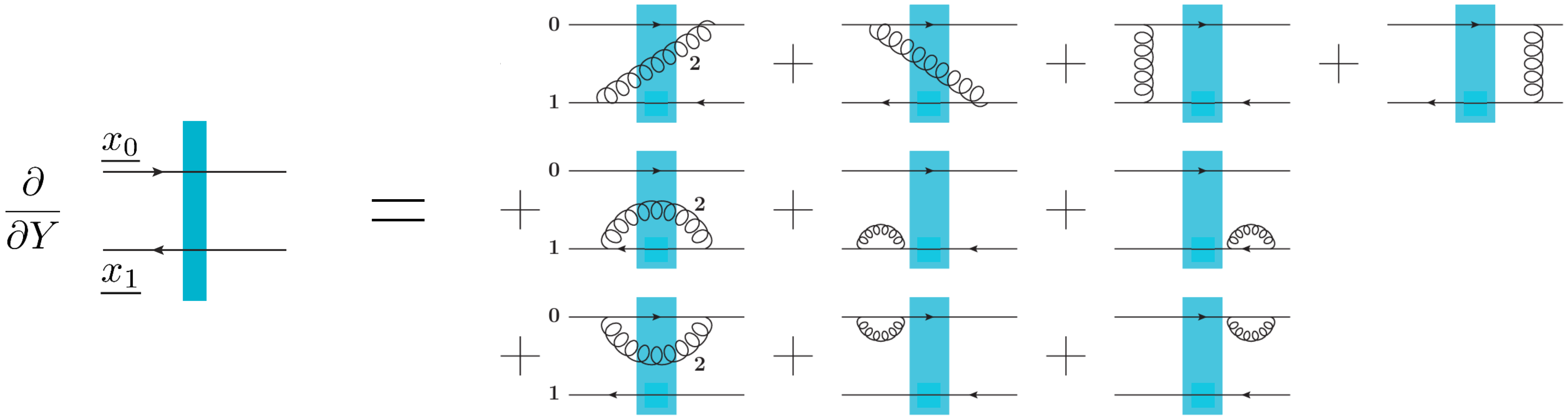


- Arises from the **phase-space enhanced quantum corrections** which are **resummed**



- The **initial conditions** can be taken from PDF fits at large x or, e.g.) the quark target model

$$\frac{1}{N_c}\left\langle \mathrm{tr}\left[V_{\underline{0}} V_{\underline{1}}^\dagger\right]\right\rangle_{(zs)}^{(0)} = \frac{2\alpha_s^2 C_F}{N_c} \ln^2 \frac{x_{0T}}{x_{1T}}$$



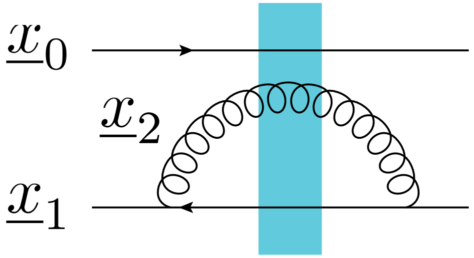
$$\frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(zs)} = \frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(zs)}^{(0)}$$

$$+\frac{\alpha_s N_c}{2\pi^2}\int\limits_{\frac{\Lambda^2}{s}}^{z}\frac{dz'}{z'}\int d^2x_2\,\frac{x_{10}^2}{x_{20}^2\,x_{21}^2}\left[\frac{1}{N_c^2}\left\langle \mathrm{tr}\left[V_2 V_1^\dagger\right]\,\mathrm{tr}\left[V_0 V_2^\dagger\right]\right\rangle_{(z's)} - \frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(z's)}\right]$$

- **"Real" gluon emissions** propagate through the **gauge field** of the proton

$$\frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 t^a V_1^\dagger t^b\right] U_2^{ba}\right\rangle_{(z's)}$$

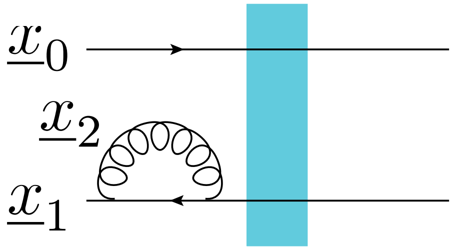
Adjoint Wilson line (gluon)



$$(x_2^-)_i < 0^- < (x_2^-)_f$$

- **"Virtual" gluon emissions** propagate through the **vacuum**, before or after hitting the proton.

$$-\frac{C_F}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(z's)}$$



+ "c.c."

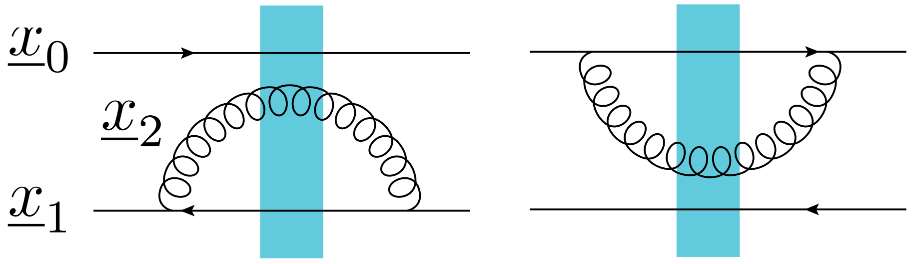
$$(x_2^-)_i\,,\,(x_2^-)_f < 0^-$$

$$0^- < (x_2^-)_i\,,\,(x_2^-)_f$$

Radiative Corrections: "Ladder" and "Non-Ladder"

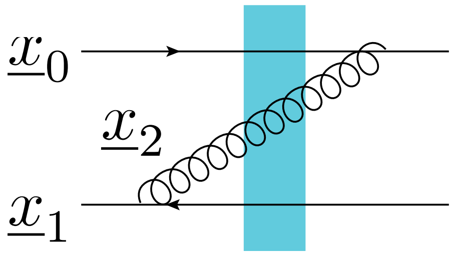
- **"Ladder" emissions** are emitted and absorbed by the same parton

$$\frac{\alpha_s}{\pi^2} \int\limits_{\frac{\Lambda^2}{s}}^{z} \frac{dz'}{z'} \int d^2 x_2 \left(\frac{1}{x_{21}^2} + \frac{1}{x_{20}^2} \right) \times \left[\text{operator} \right]$$

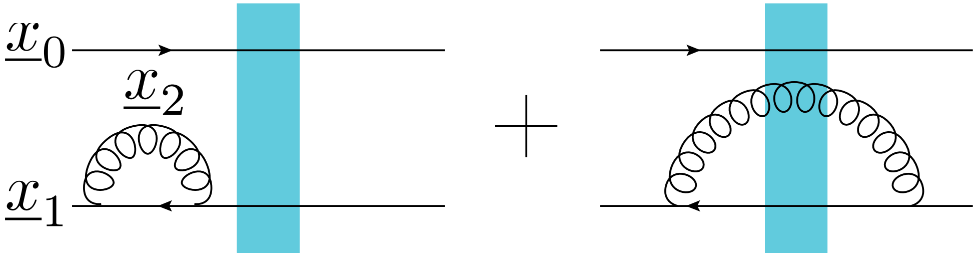


- **"Non-ladder" emissions** are emitted and absorbed by different partons

$$\frac{\alpha_s}{\pi^2} \int\limits_{\frac{\Lambda^2}{s}}^{z} \frac{dz'}{z'} \int d^2 x_2 \left(-2 \frac{\underline{x}_{21} \cdot \underline{x}_{20}}{x_{21}^2 \, x_{20}^2} \right) \times \left[\text{operator} \right]$$



What Happens to the Transverse Integral?



+ “c.c.”

$$\frac{\alpha_s N_c}{2\pi^2} \int\limits_{\frac{\Lambda^2}{s}}^{z} \frac{dz'}{z'} \int \frac{d^2 x_2}{x_{21}^2} \times \left[\frac{1}{N_c^2} \left\langle \mathrm{tr}\left[V_2 V_1^\dagger \right] \, \mathrm{tr}\left[V_0 V_2^\dagger \right] \right\rangle_{(z's)} - \frac{1}{N_c} \left\langle \mathrm{tr}\left[V_0 V_1^\dagger \right] \right\rangle_{(z's)} \right]$$

- “Ladder” emissions of **small-sized fluctuations are enhanced**
- **Transverse integral** appears divergent… a **second logarithm?**
- **No**: **Cancellation** of real + virtual diagrams due to **color transparency**

Full Structure: The Balitsky Operator Hierarchy

$$\frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(zs)} = \frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(zs)}^{(0)}$$

$$+\frac{\alpha_s N_c}{2\pi^2}\int\limits_{\frac{\Lambda^2}{s}}^{z}\frac{dz'}{z'}\int d^2x_2\,\frac{x_{10}^2}{x_{20}^2\,x_{21}^2}\left[\frac{1}{N_c^2}\left\langle \mathrm{tr}\left[V_2 V_1^\dagger\right]\,\mathrm{tr}\left[V_0 V_2^\dagger\right]\right\rangle_{(z's)} - \frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(z's)}\right]$$

Rapidity Logarithm

BFKL Kernel

New, more complex operator

- The dipole evolves into **increasingly complex operators**….

I. Balitsky, Nucl. Phys. ***B463*** *(1996) 99*
I. Balitsky, Phys. Rev. ***D60*** *(1999) 014020*

- Equivalent to a **functional differential equation**….

(JIMWLK)

Jalilian-Marian et al., Phys. Rev. ***D59*** *(1998) 014015*
Jalilian-Marian et al., Phys. Rev. ***D59*** *(1998) 014014*
Iancu et al., Phys. Lett. ***B510*** *(2001) 133*
Iancu et al., Nucl. Phys. ***A692*** *(2001) 583*

Linear BFKL Evolution at Small x

- The equations **linearize** in the **dilute limit (BFKL)**

*Kuraev, et al., Sov. Phys. JETP **45** (1977) 199*
*Balitsky and Lipatov, Sov. J. Nucl. Phys. **28** (1978) 822*

$$S_{01} = 1 - \frac{1}{N_c} \left\langle \mathrm{tr} \left[V_0 V_1^\dagger \right] \right\rangle \ll 1$$

$$\frac{1}{N_c} \left\langle \mathrm{tr} \left[V_0 V_1^\dagger \right] \right\rangle_{(zs)} = \frac{1}{N_c} \left\langle \mathrm{tr} \left[V_0 V_1^\dagger \right] \right\rangle_{(zs)}^{(0)}$$

$$+ \frac{\alpha_s N_c}{2\pi^2} \int\limits_{\frac{\Lambda^2}{s}}^{z} \frac{dz'}{z'} \int d^2 x_2 \, \frac{x_{10}^2}{x_{20}^2 \, x_{21}^2} \left\langle \frac{1}{N_c} \mathrm{tr} \left[V_2 V_1^\dagger \right] + \frac{1}{N_c} \mathrm{tr} \left[V_0 V_2^\dagger \right] - \frac{1}{N_c} \mathrm{tr} \left[V_0 V_1^\dagger \right] - 1 \right\rangle_{(z's)}$$

➢ Solution by **Laplace-Mellin transform** leads to **power-law growth** at small x

"Pomeron Intercept"

$$xq(x, Q^2) \sim xG(x, Q^2) \sim \left(\frac{1}{x} \right)^{\frac{4\alpha_s N_c}{\pi} \ln 2}$$

$$\alpha_P - 1 = \frac{4\alpha_s N_c}{\pi} \ln 2$$

Gluon Saturation

What It Is, and Why We Need to Find It

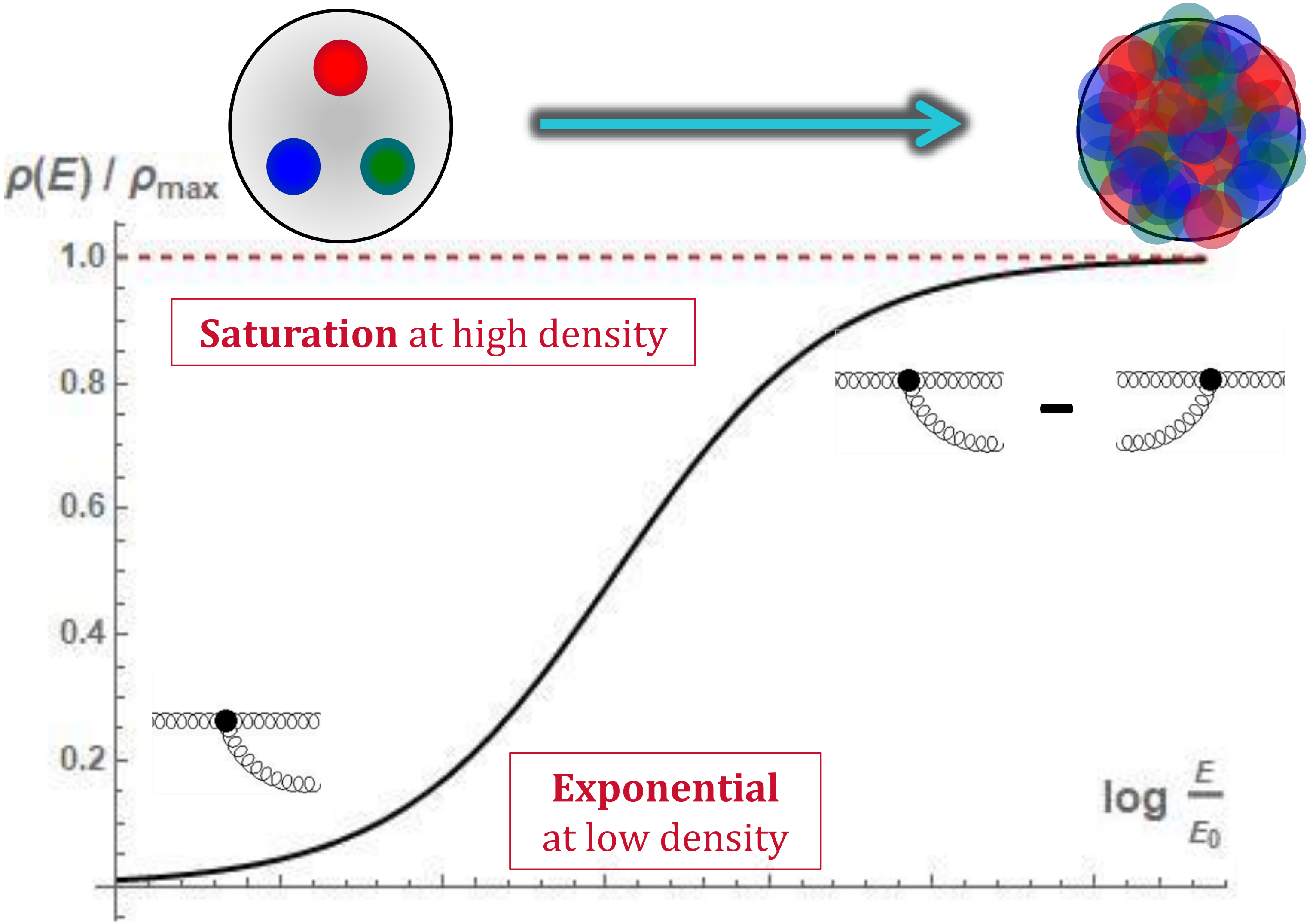
The Drive to Gluon Saturation

- QCD radiates an abundance of **soft gluons at small x**

➢ **Exponential growth** of charge density

- Also provides a mechanism to **regulate the growth**

➢ Gluon **recombination** can balance radiation, leading to a **saturation** of the density



Nonlinear BK Evolution at Large x

- The full nonlinear operator **hierarchy closes** under a **mean-field approximation**.

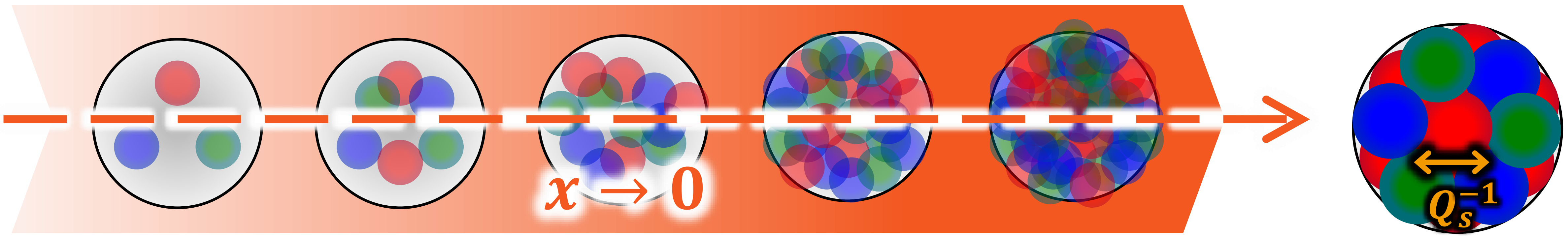
Balitsky, Nucl. Phys. ***B463*** *(1996) 99*
Balitsky, Phys. Rev. ***D60*** *(1999) 014020*
Kovchegov, Phys. Rev. ***D60*** *(1999) 034008*
Kovchegov, Phys. Rev. ***D61*** *(2000) 074018*

➢ **Large-Nc limit!**

$$\frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(zs)} = \frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(zs)}^{(0)}$$

$$+ \frac{\alpha_s N_c}{2\pi^2}\int\limits_{\frac{\Lambda^2}{s}}^{z}\frac{dz'}{z'}\int d^2x_2\,\frac{x_{10}^2}{x_{20}^2\,x_{21}^2}\left[\frac{1}{N_c^2}\left\langle \mathrm{tr}\left[V_2 V_1^\dagger\right]\right\rangle_{(z's)}\times\left\langle \mathrm{tr}\left[V_0 V_2^\dagger\right]\right\rangle_{(z's)} - \frac{1}{N_c}\left\langle \mathrm{tr}\left[V_0 V_1^\dagger\right]\right\rangle_{(z's)}\right]$$

❖ **Mean field** ❖ **Gluon-dominated** ❖ **Fully nonlinear** ❖ **Solvable!**

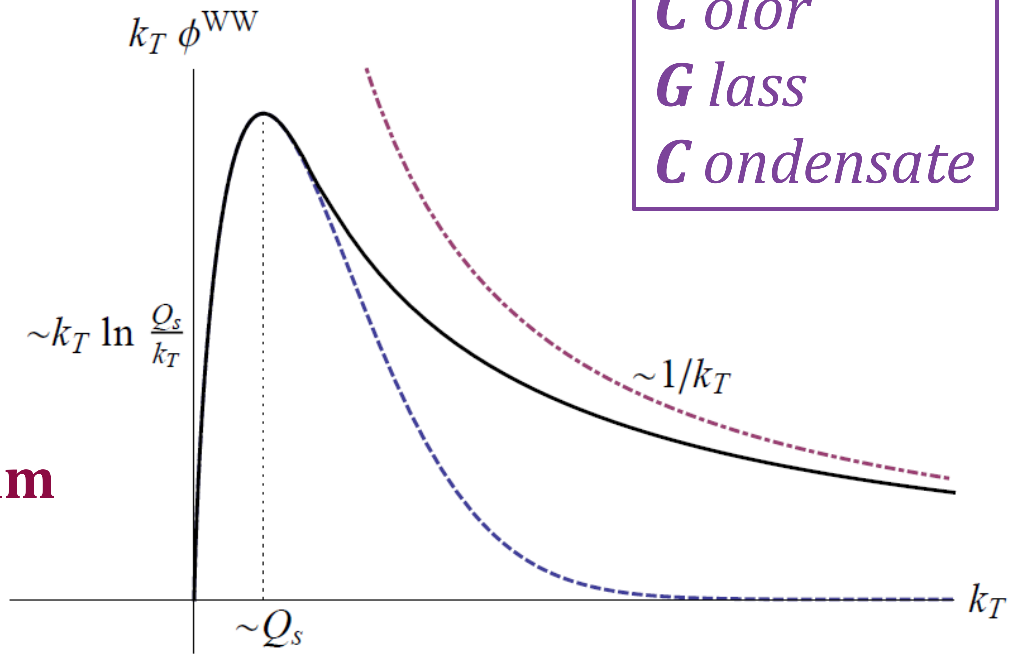


L. McLerran, R. Venugopalan, *Phys. Rev.* **D49** (1994), **D50** (1994)

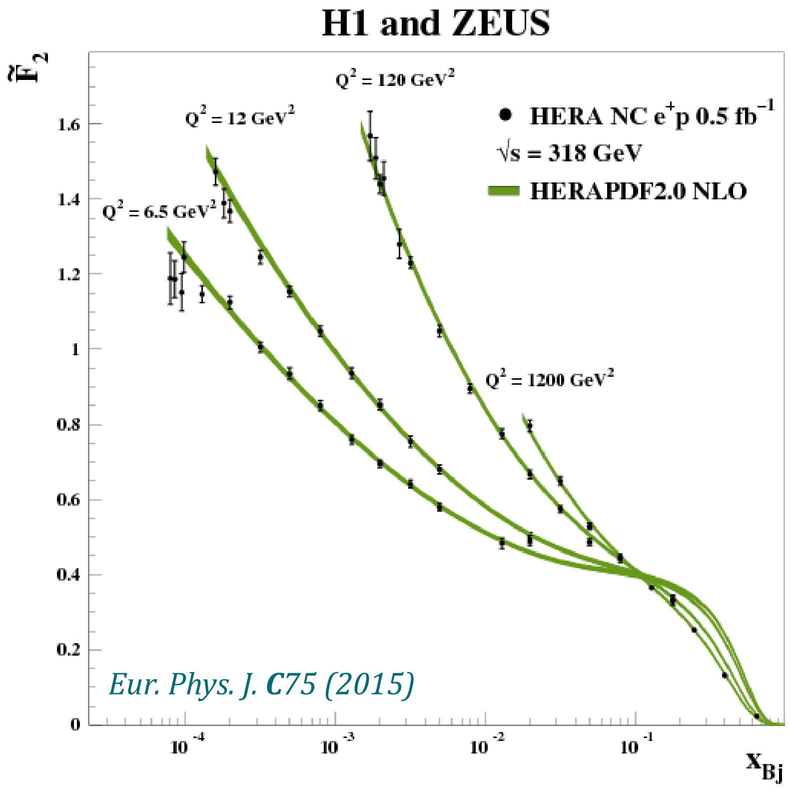
C *olor*
G *lass*
C *ondensate*

- At small x, we approach an **emergent**, **universal**, **high-density state** of QCD characterized by:
  - Resummation of **coherent multiple scattering**
  - **Strong gluon fields**, semi-classical QCD
  - **Color domains** with emergent **saturation momentum**

$$Q_s^2(x) \approx 3.70\,\alpha_s(Q_s^2)\,\rho_G(x, Q_s^2)$$

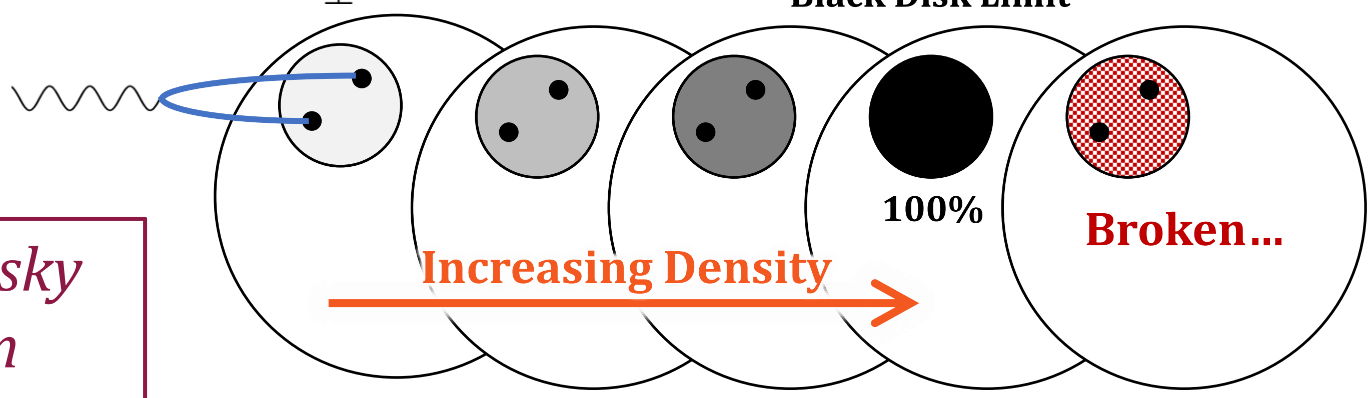


The Status Quo (BFKL) Is Unsustainable



$$S_{dip}(x_\perp) = \frac{d\sigma^{dip}(x_\perp)}{d^2 b_\perp}$$

"Little Unitarity Problem"



***B** alitsky*
***F** adin*
***K** uraev*
***L** ipatov*

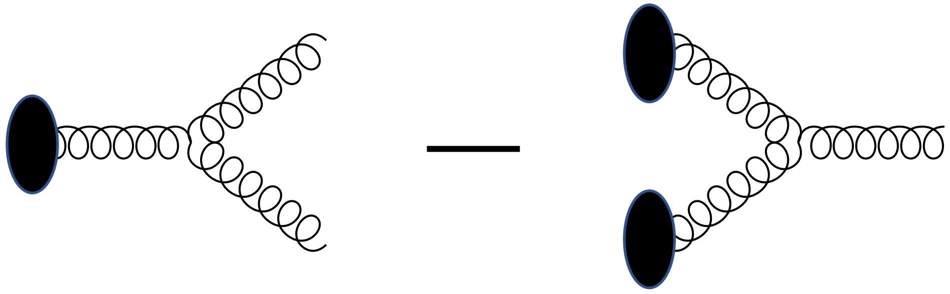
*Kuraev, Lipatov, and Fadin, Sov. Phys. JETP **45** (1977) 199*
*Balitsky and Lipatov, Sov. J. Nucl. Phys. **28** (1978) 822*

$$xG(x,Q^2) \sim \left(\frac{1}{x}\right)^{2.65\,\alpha_s(Q^2)}$$

$$\sigma^{q\bar{q}} \approx 3.29 \frac{\alpha_s(Q^2)}{Q^2}\, xG(x,Q^2)$$

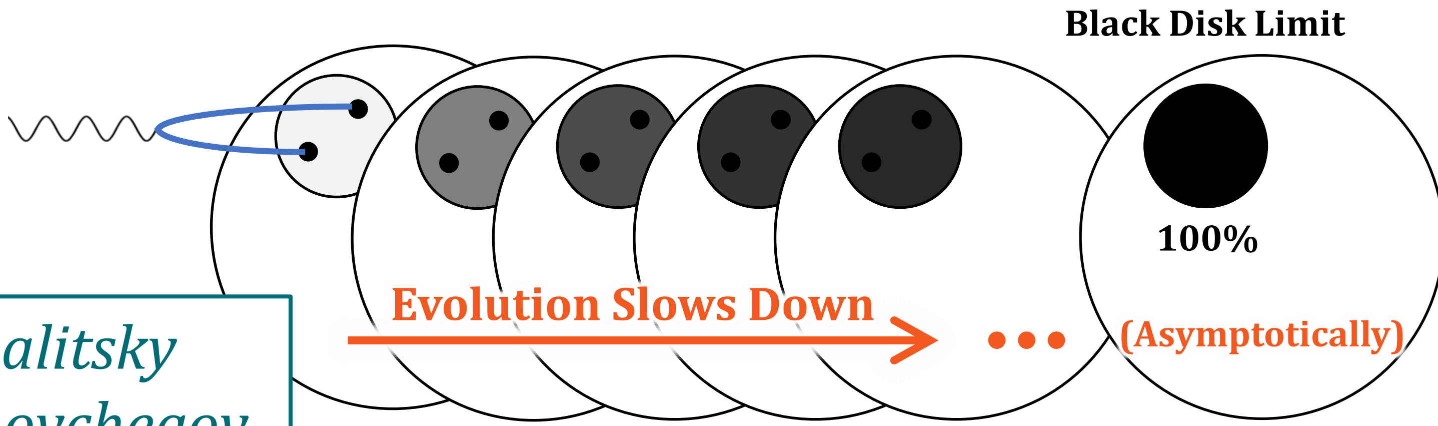
- **Resumming multi-emission** leads to a **power law growth** of the gluon density at small x.
  - **Breaks unitarity** locally: $\boxed{S_{dip}(x_\perp) > 1}$

Saturation: Necessary but not Sufficient for Unitarity



I. Balitsky, Nucl. Phys. **B463** (1996)
Y. Kovchegov, Phys. Rev. **D60** (1999)
J. Jalilian-Marian et al., Nucl. Phys. ***B504*** *(1997), Phys. Rev.* ***D59*** *(1998)*
A. Kovner, J. G. Milhano, H. Weigert, Phys. Rev. ***D62*** *(2000)*
E. Iancu, A. Leonidov, L. McLerran, Nucl. Phys. ***A692*** *(2001)*

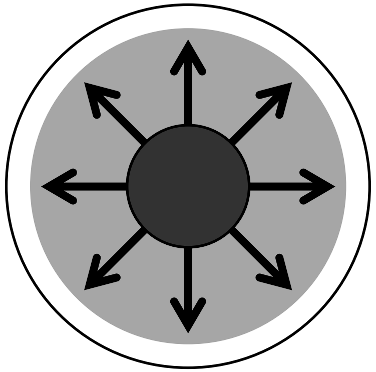
***B** alitsky*
***K** ovchegov*



$$\rho_G \sim 1 - c\, e^{-0.556\alpha_s^2 Y^2}$$

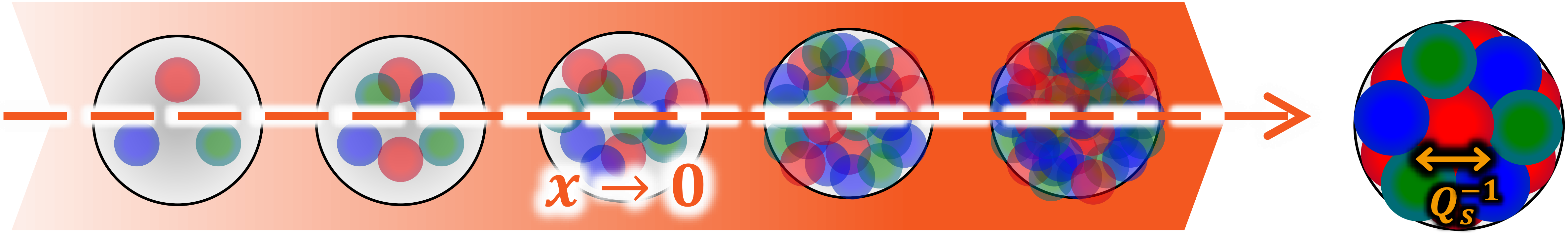
- At high density, **nonlinear gluon fusion slows down** the gluon cascade
  - ➢ Asymptotically **saturates** to the **black disk limit**
  - ➢ **Solves** the **"little unitarity problem"** (perturbatively!)
- But does **not** explicitly unitarize **total** cross sections due to the **diffusion** of the **black disk**. (nonperturbative)

$$R_{black} \sim e^{Y}$$



E. Ferreiro et al., Nucl. Phys. **A710** (2002)
A. Kovner, U. A. Wiedemann,
Phys. Lett. **B551** (2003)

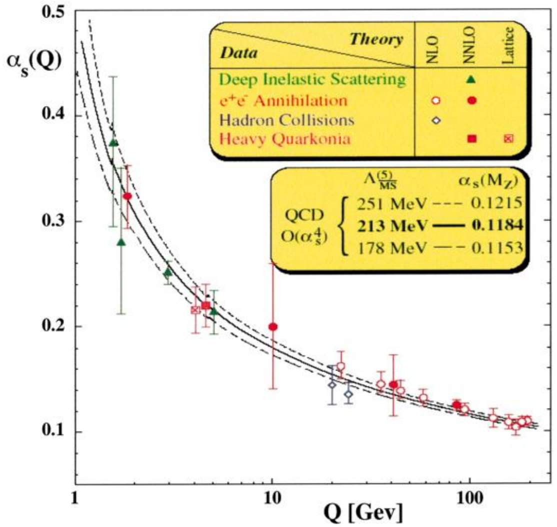
Why Should We Care About Saturation?



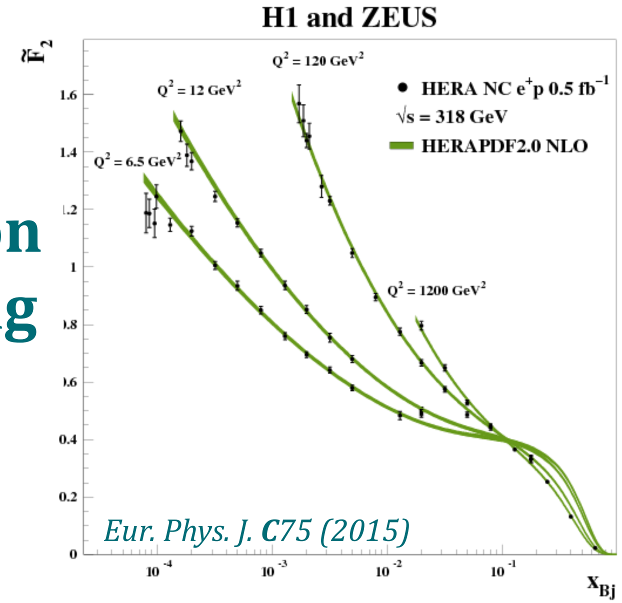
- Exotic state of nuclear matter characterized by (maximally!) intense gluon fields
- Initial conditions for heavy-ion collisions
- Perturbatively controllable regime of QCD with different dynamics (dense proton)
- ...Yes, and...

Why Should We Care About Saturation?

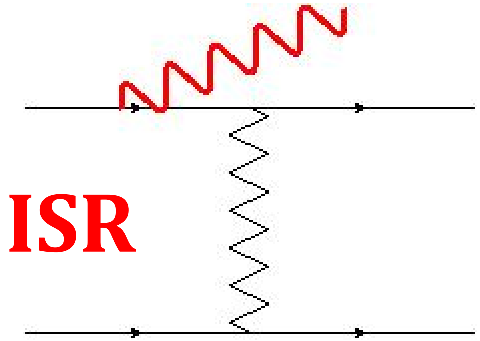
Any theory which is **asymptotically free** like **QCD**…



…will have an **explosion** of **soft bremsstrahlung** at **small x**.



- Radiation patterns at high energy $s \to \infty \quad (x \to 0)$



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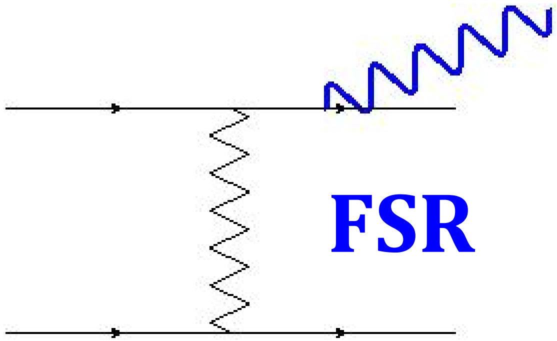
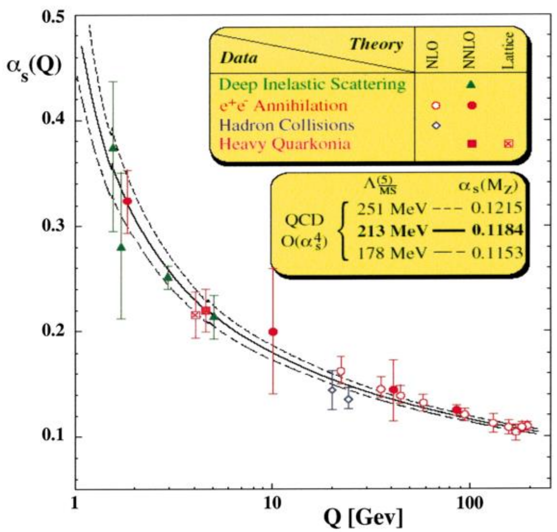
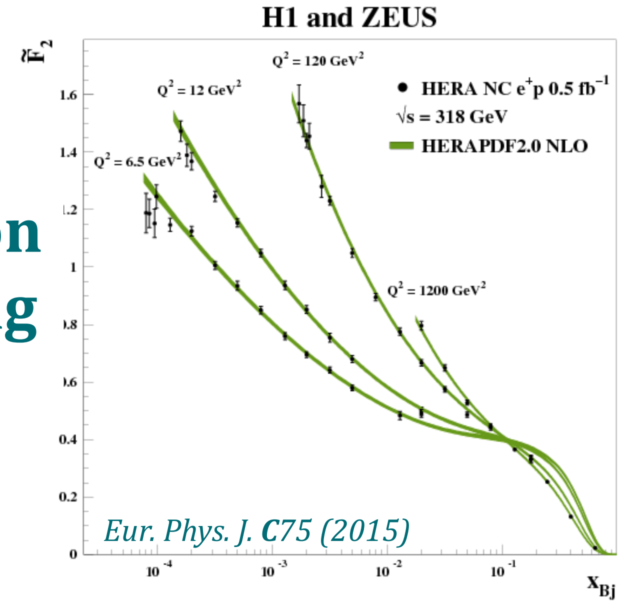


$$\begin{cases} \dfrac{dN}{dy} = 0 & \textbf{Abelian (QED)} \\[2ex] \dfrac{dN}{dy} = \alpha_s \times (\text{const}) & \textbf{Non-Abelian (QCD)} \end{cases}$$

Any theory which is **asymptotically free** like **QCD**…

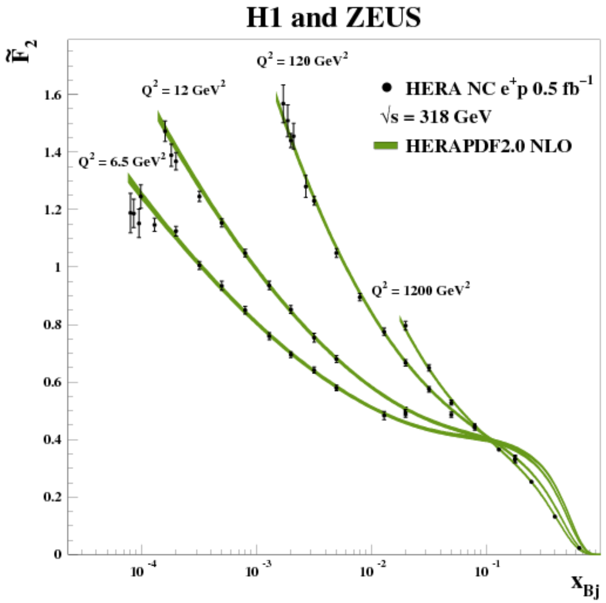


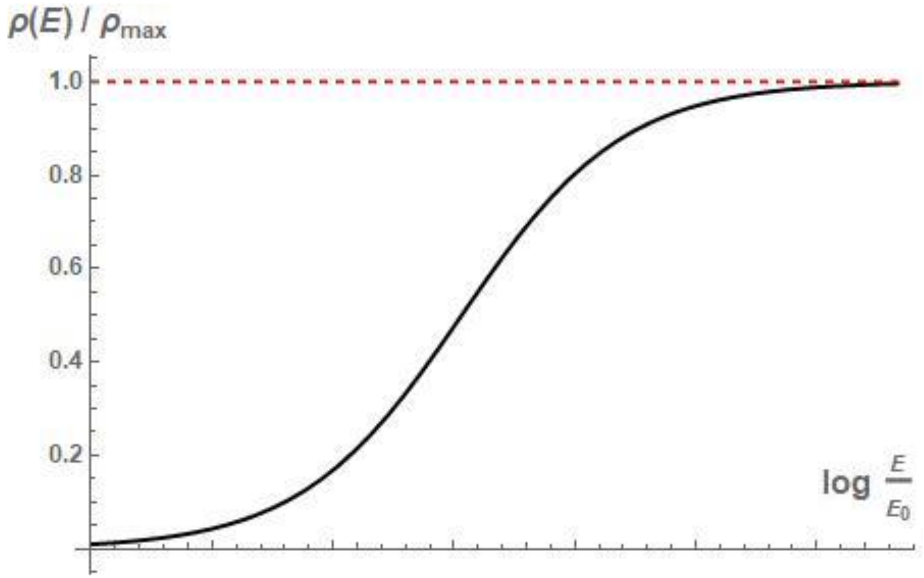
…will have an **explosion** of **soft bremsstrahlung** at **small x**.

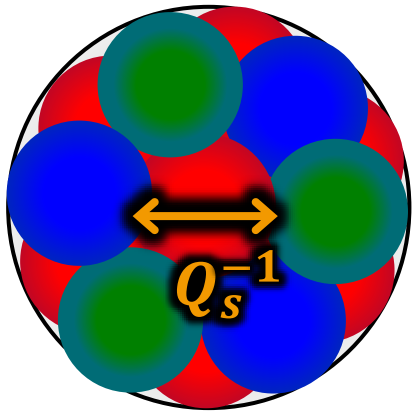


- The very feature that **makes QCD work** as an asymptotically free, UV-complete theory **threatens to break it** at small x.
  - *"Saturation is a non-negotiable consequence of QCD."* - *N. Armesto*
- If saturation **does not exist**, something is **irredeemably broken**
  - Not just for QCD, but for **any non-Abelian quantum field theory**

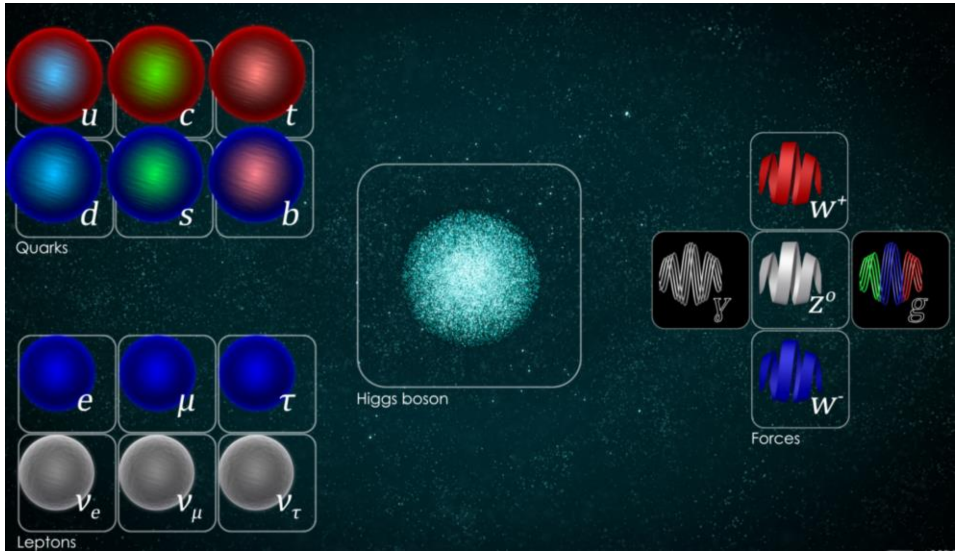
My Claim: Saturation is as Essential as the Higgs Boson







- Saturation, like the Higgs, is necessary for the **self-consistency of the QFT**
- Nothing guarantees we could find it…
- …But if it **does not exist**, we **need to know!**



Thank You!

Enjoy the Rest of the Summer School!