

# Baryon Production in A Multi-Phase Transport (AMPT) Model

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# Outline

- Introduction of the AMPT model
- Model results on baryon and electric charge stopping
- Relations to quark stopping
- On net-charge reconstruction
- Summary

*Mostly based on work with PhD student Mason Ross*

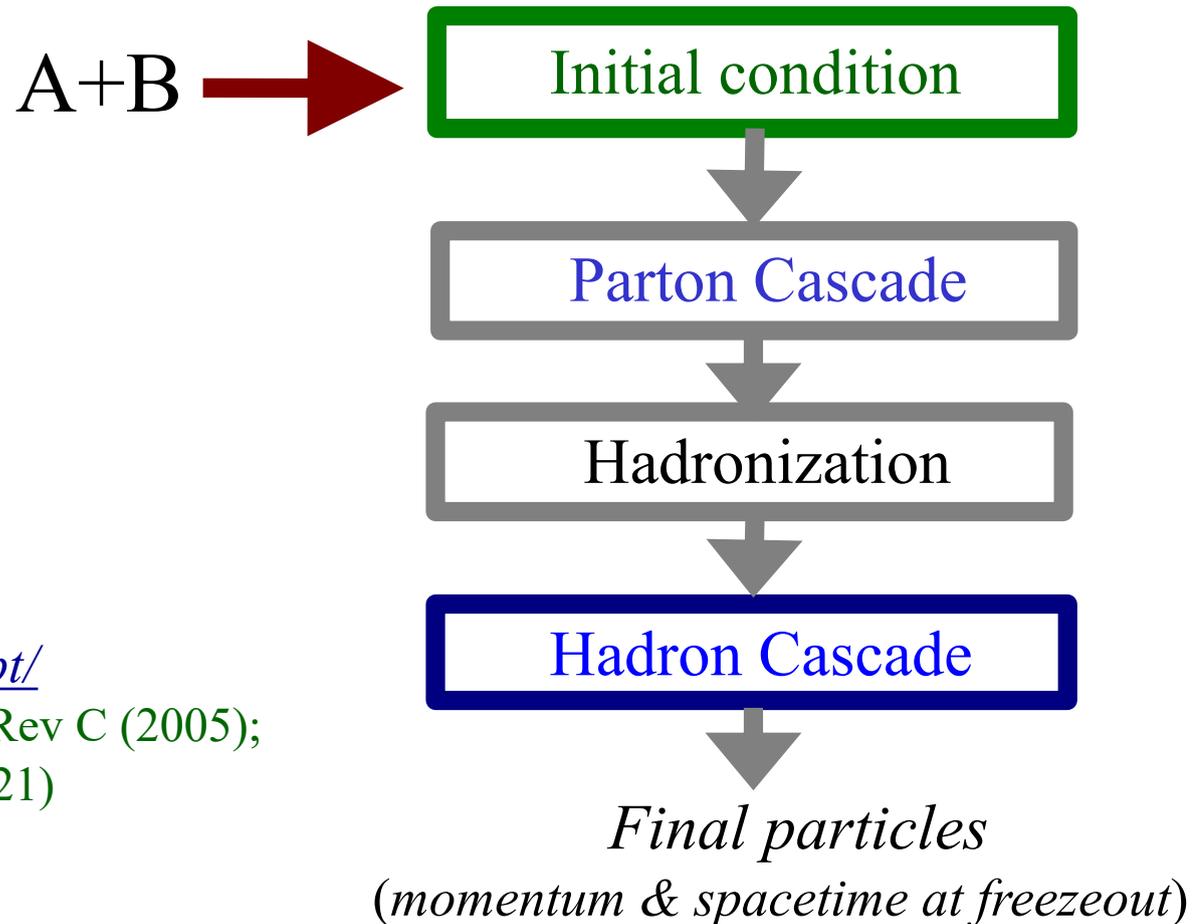


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# A multi-phase transport (AMPT) model

was constructed as a self-contained kinetic description of heavy ion collisions:

- evolves the system from initial condition to final observables;
- has explicit parton stage and hadron stage;
- includes productions of all quark flavors & conserved charges in 3D;
- can address initial condition & dynamics/evolution in non-equilibrium.



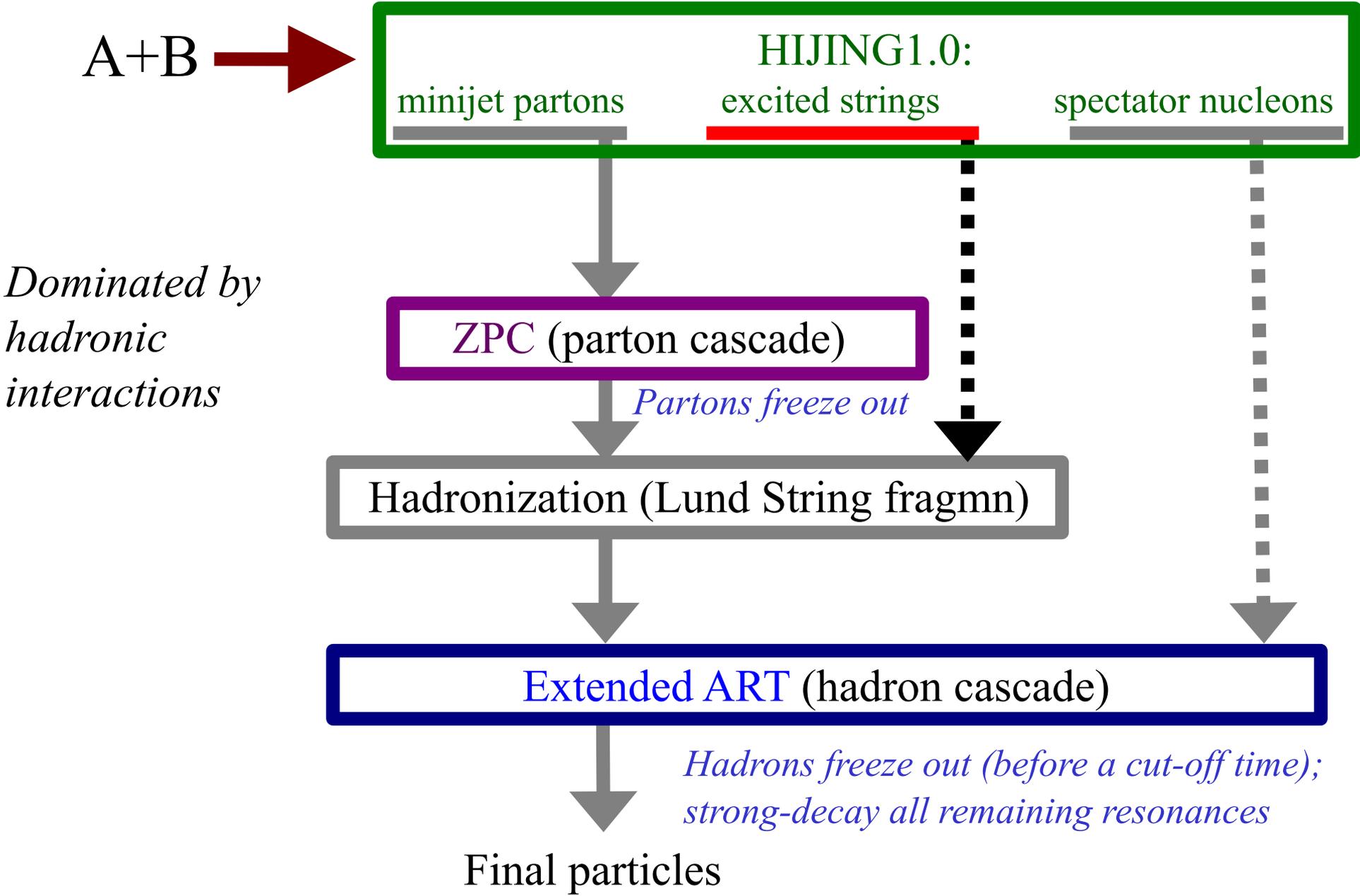
Source codes at

<https://myweb.ecu.edu/linz/ampt/>

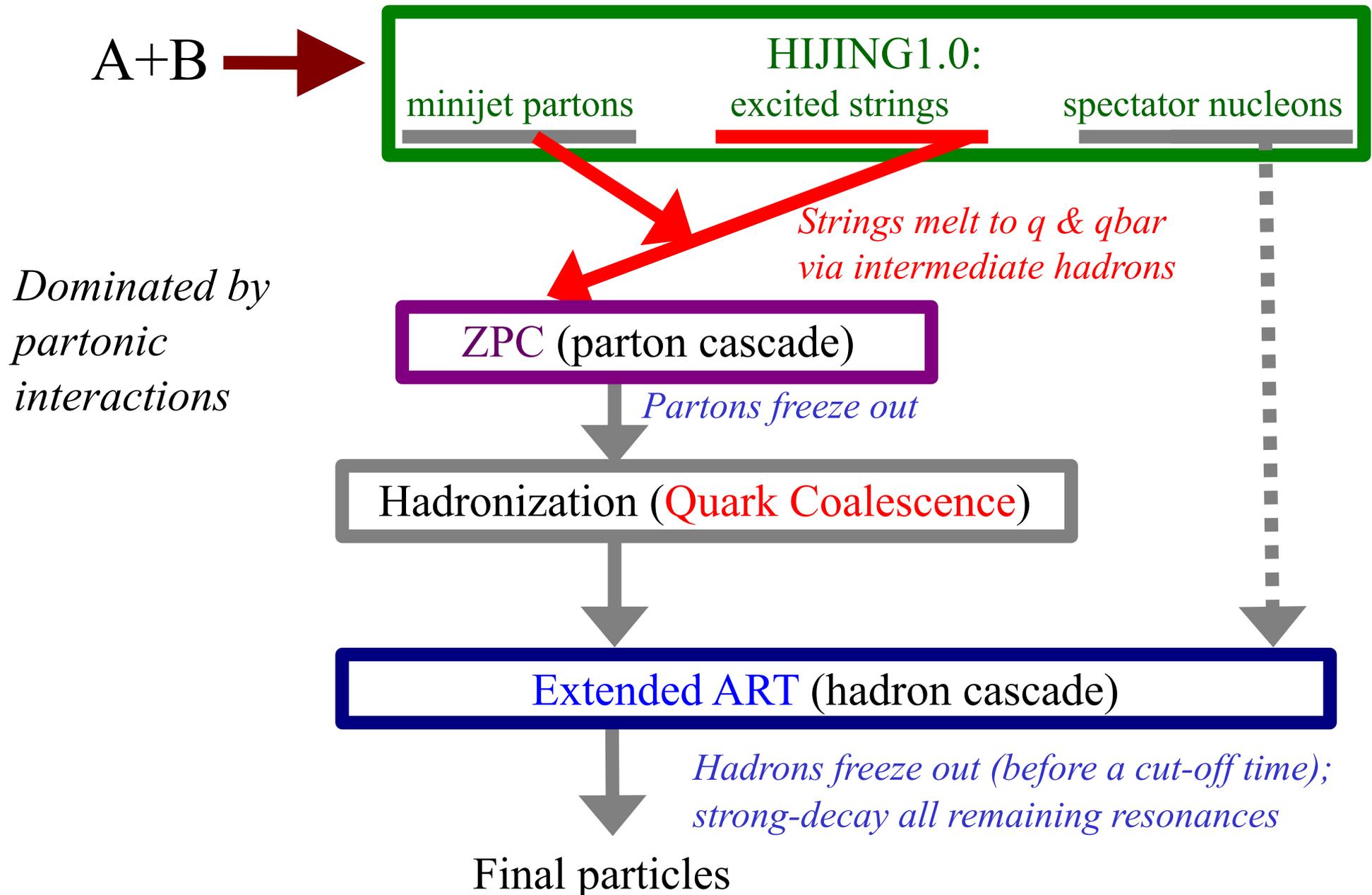
ZWL, Ko, Li, Zhang & Pal, Phys Rev C (2005);

ZWL & Zheng, Nucl Sci Tech (2021)

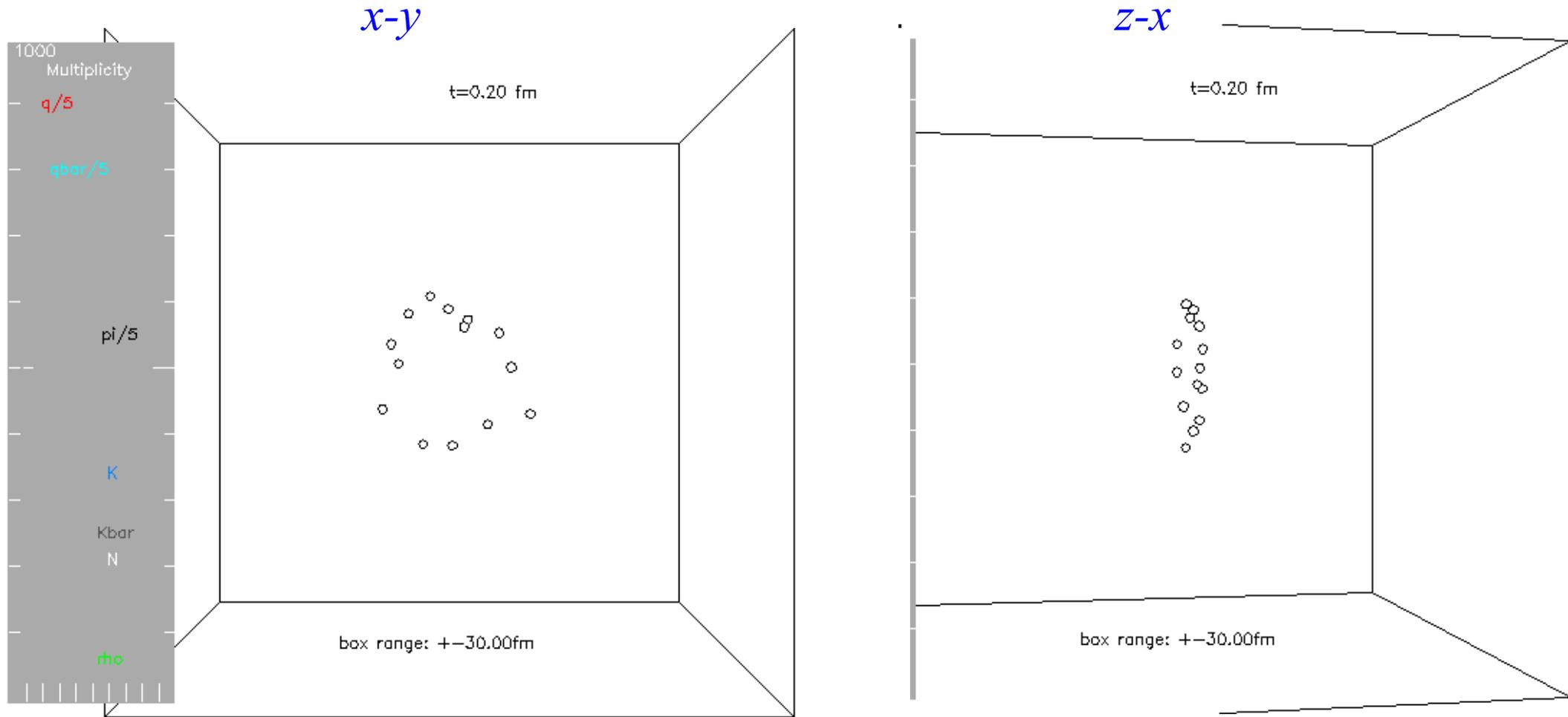
# Structure of the default AMPT model



# Structure of the string melting AMPT model



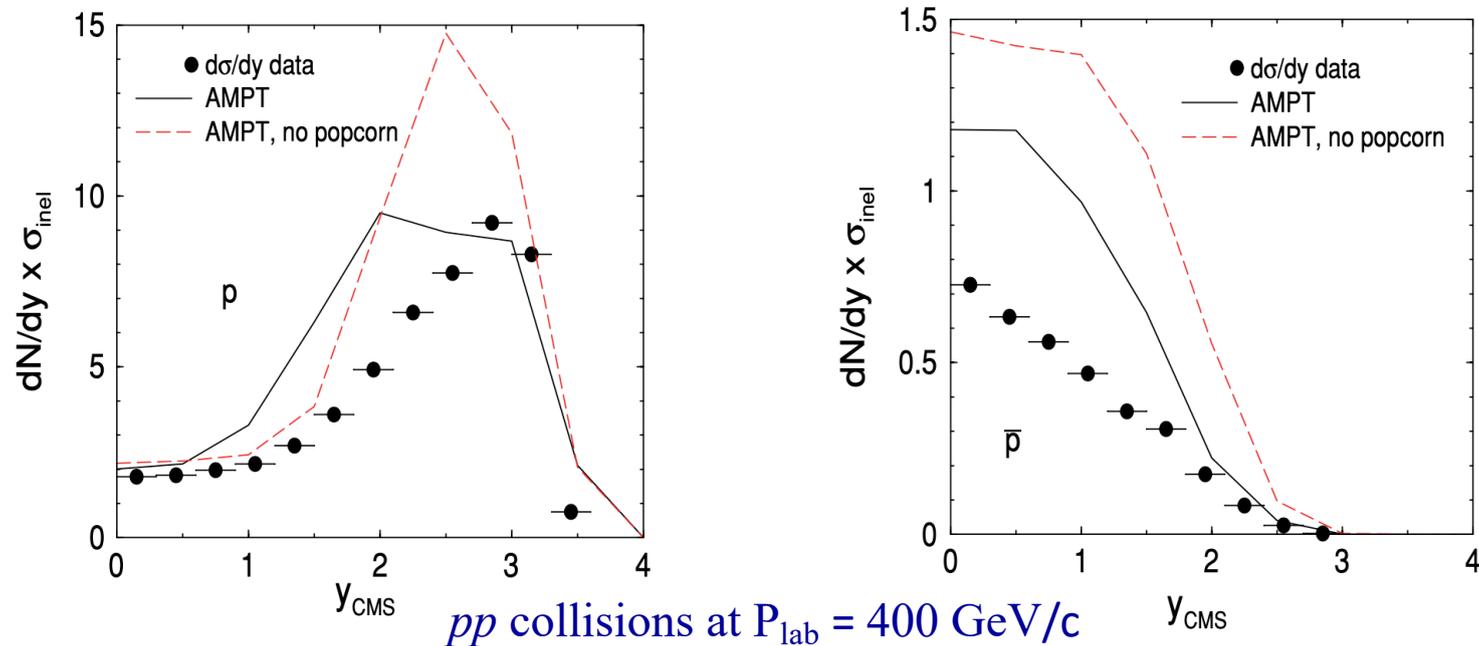
# A central Au+Au event at 200 AGeV from the String Melting AMPT model (*applicable at high energies*):



# Baryon productions in AMPT

In both versions of AMPT, the initial baryon stopping is described by **HIJING1.0** via the Lund string fragmentation in PYTHIA.

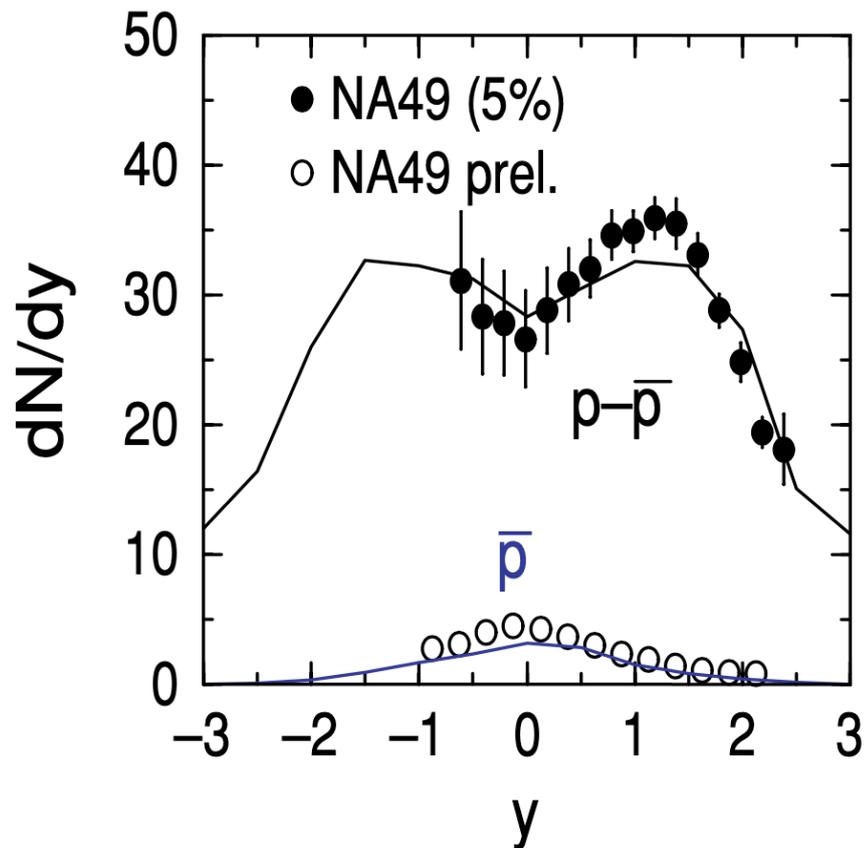
To have more freedom for adjusting the baryon stopping in AMPT, we have included the popcorn mechanism: this introduces two additional baryon-antibaryon pair production channels: the  $B\bar{B}$  and  $BM\bar{B}$  configurations; with a parameter controlling their ratio.



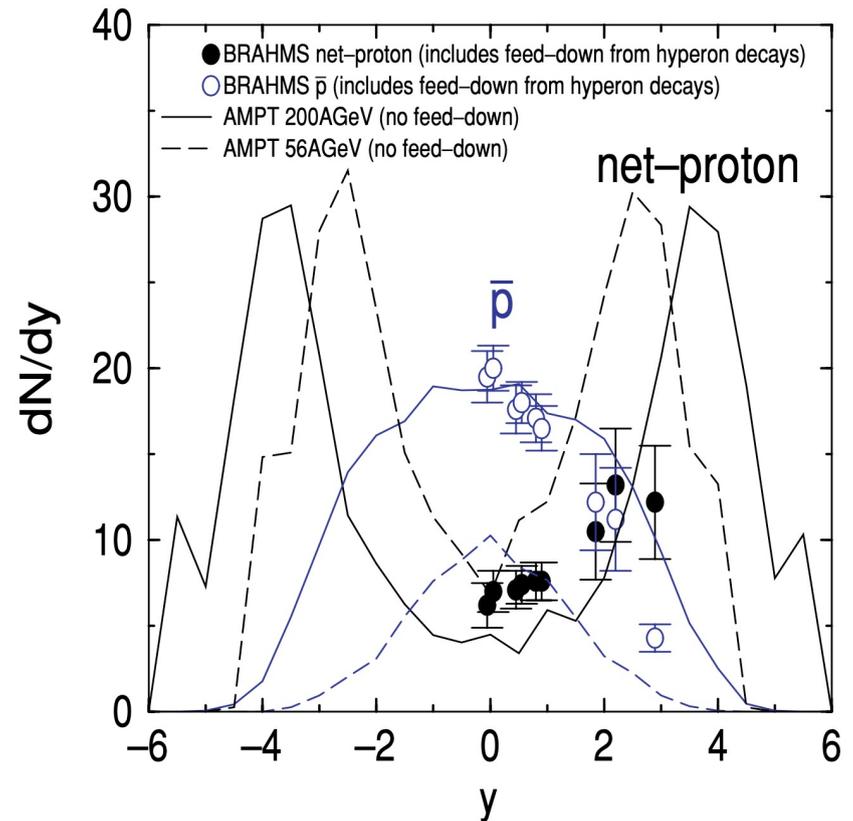
The popcorn mechanism can help but is not enough.

Baryon stopping in the default AMPT model for AA collisions are often reasonable.

Central Pb+Pb collisions  
at top SPS energy (17.3 GeV)



Central Au+Au collisions  
at RHIC energies

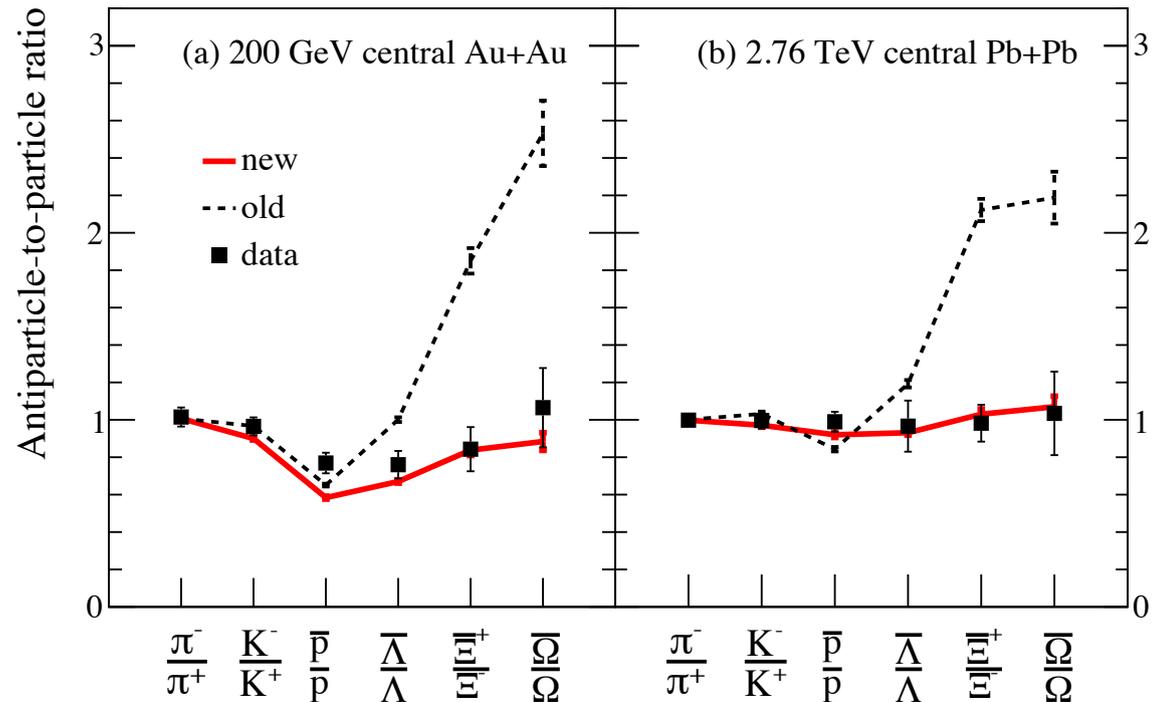


Baryon stopping in the string melting AMPT model also depends on quark coalescence.

We have improved the quark coalescence model in AMPT, it is more physical and efficient, especially for (anti)baryons.

Y He & ZWL, Phys Rev C (2017)

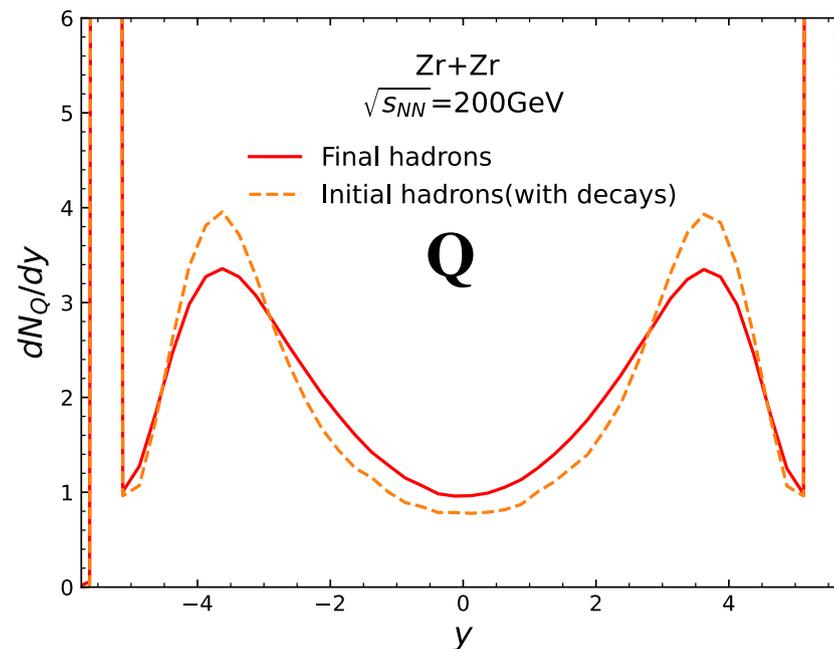
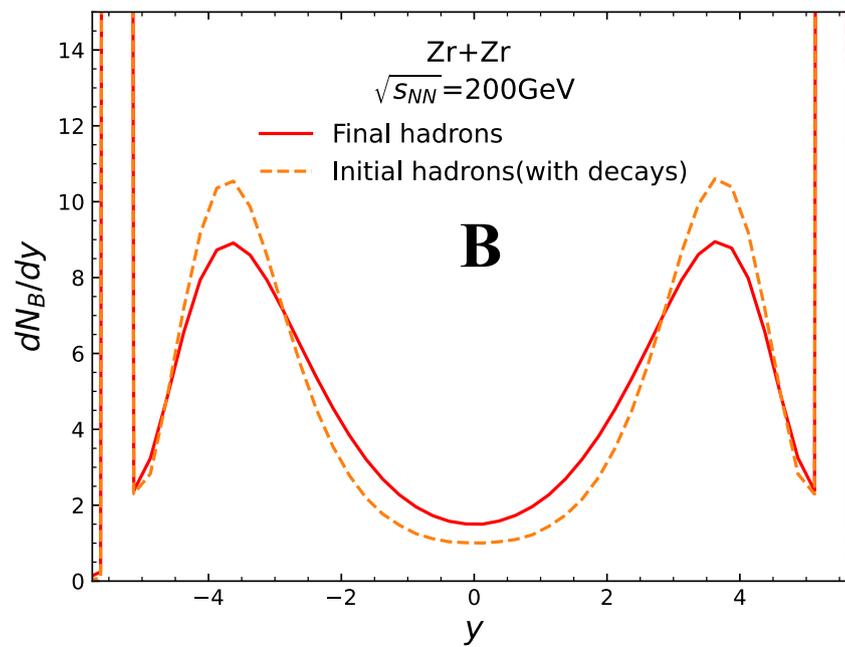
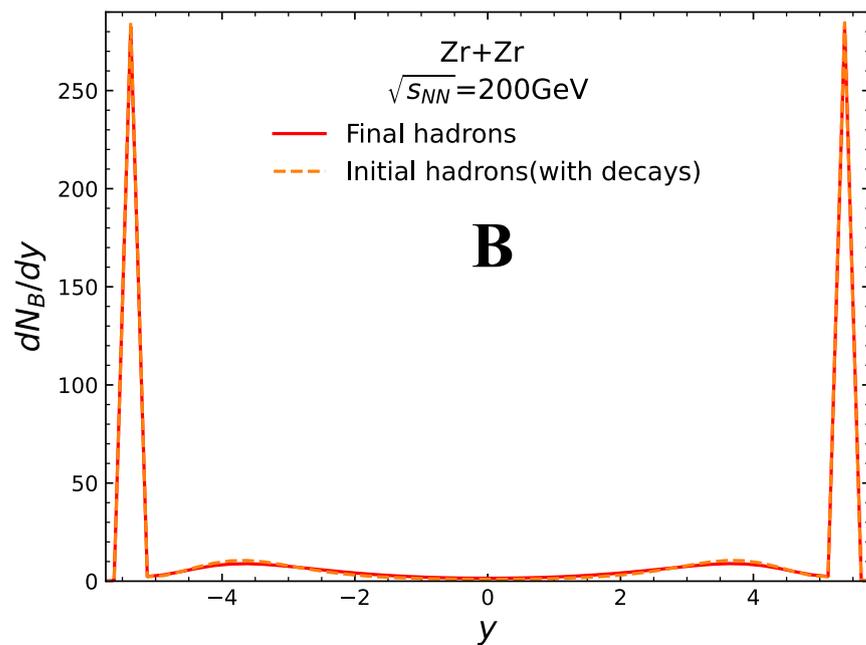
It improves (anti)baryon observables including  $p$  &  $\bar{p}$  yield &  $p_T$ -spectra and multi-strange  $\bar{B}/B$  ratios:



There are still challenges, such as the low  $\bar{p}/p$  ratio at/below 200A GeV.

# Results from string melting AMPT

## Minimum-bias Zr+Zr collisions



$N_B$  or B: net-baryon number.  
 $N_Q$  or Q: net-electric charge number.

Ratio  $B/Q \cdot Z/A$  has been proposed

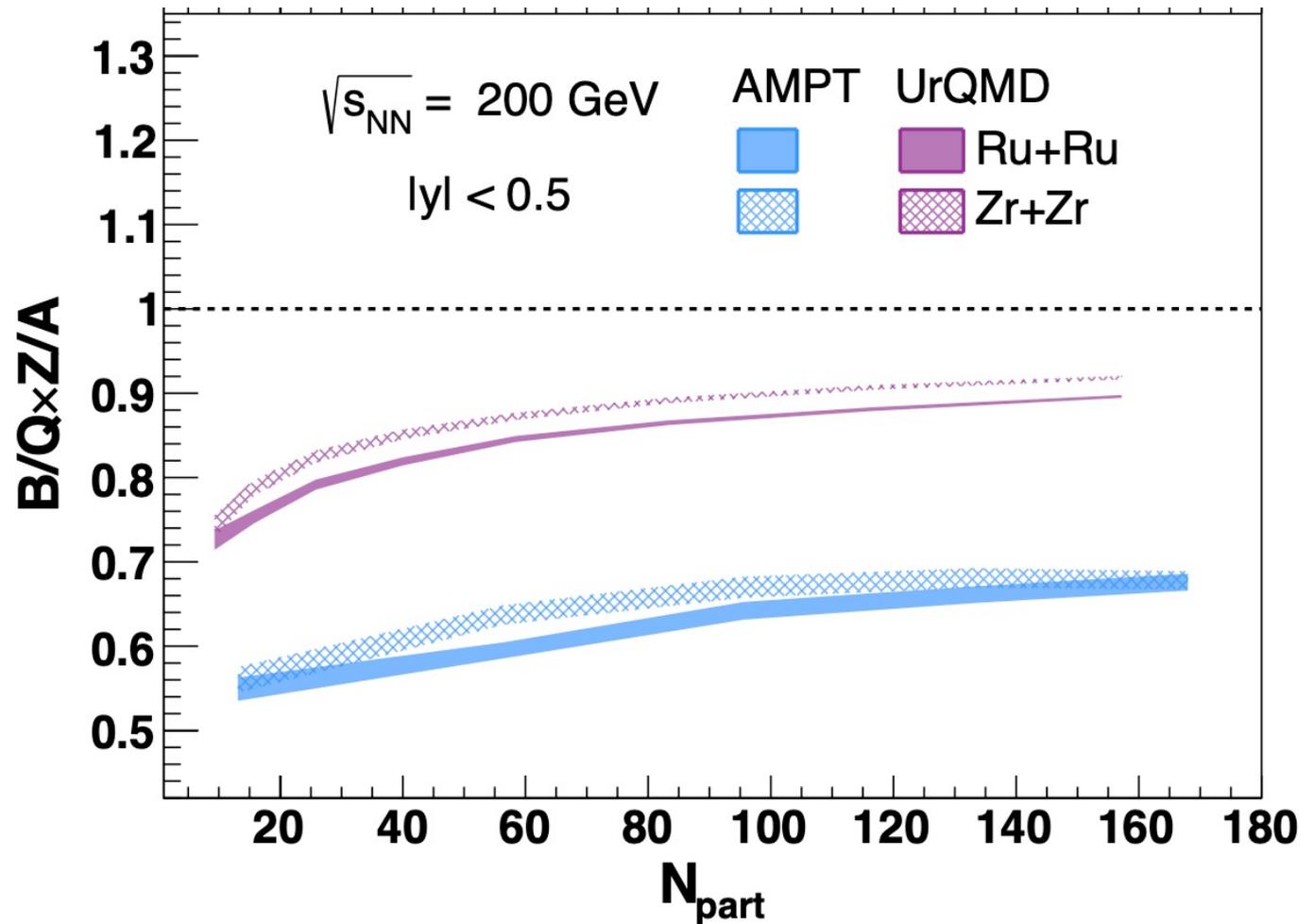
to study baryon stopping.

Z. Xu, RBRC workshop on Physics  
Opportunities from isobar run (2022)

Ratio at mid-rapidity:

$B/Q \cdot Z/A < 1$

in both models, for every centrality.



N. Lewis et al.,  
2205.05685 [hep-ph]

More UrQMD results in  
W. Lv et al., 2309.06445 [nucl-th]

## Naive expectation for $B/Q \cdot Z/A$

When we include all particles at all rapidities,  
conservation laws give  $B = 2A$  and  $Q = 2Z$ ,

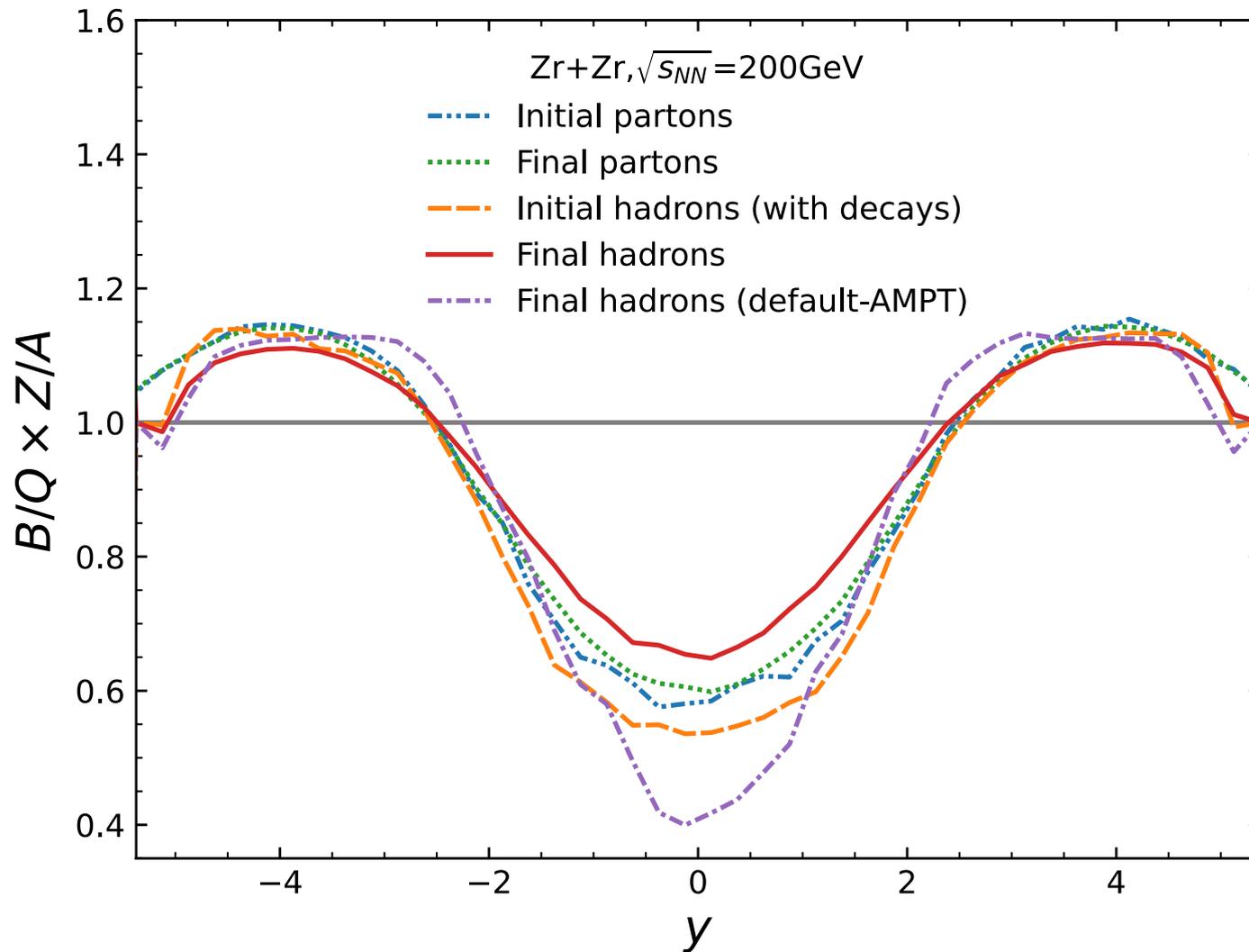
so  $B/Q \cdot Z/A = 1$ .

This can be called the **naive expectation**.

It's tempting to expect this to be true at any rapidity  $y$ ,  
but we already see strong deviations.

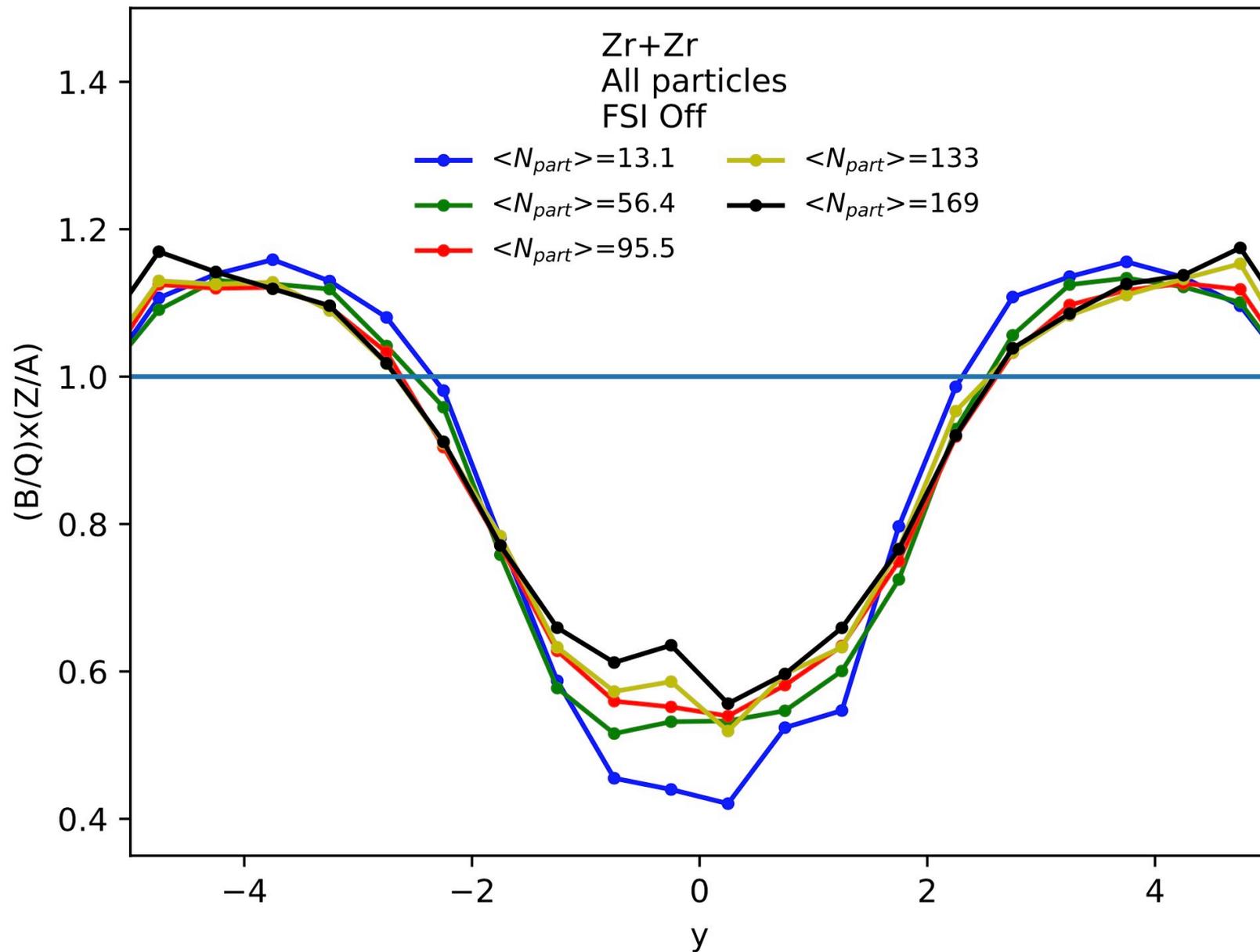
## Ratio versus rapidity (for minimum bias events):

- $B/Q \times Z/A$  changes strongly with  $y$
- Rescatterings and hadronization only modestly change the ratio
- Default AMPT and string melting AMPT are qualitatively similar



Ratio versus rapidity:

$B/Q \cdot Z/A$  changes only modestly with centrality

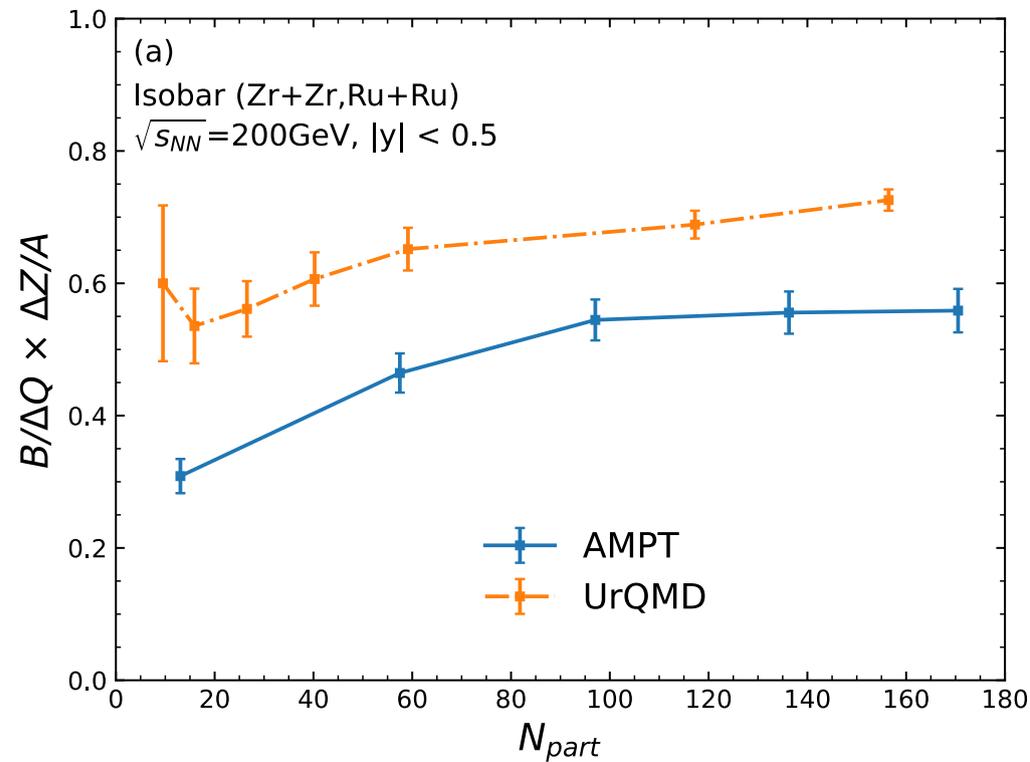
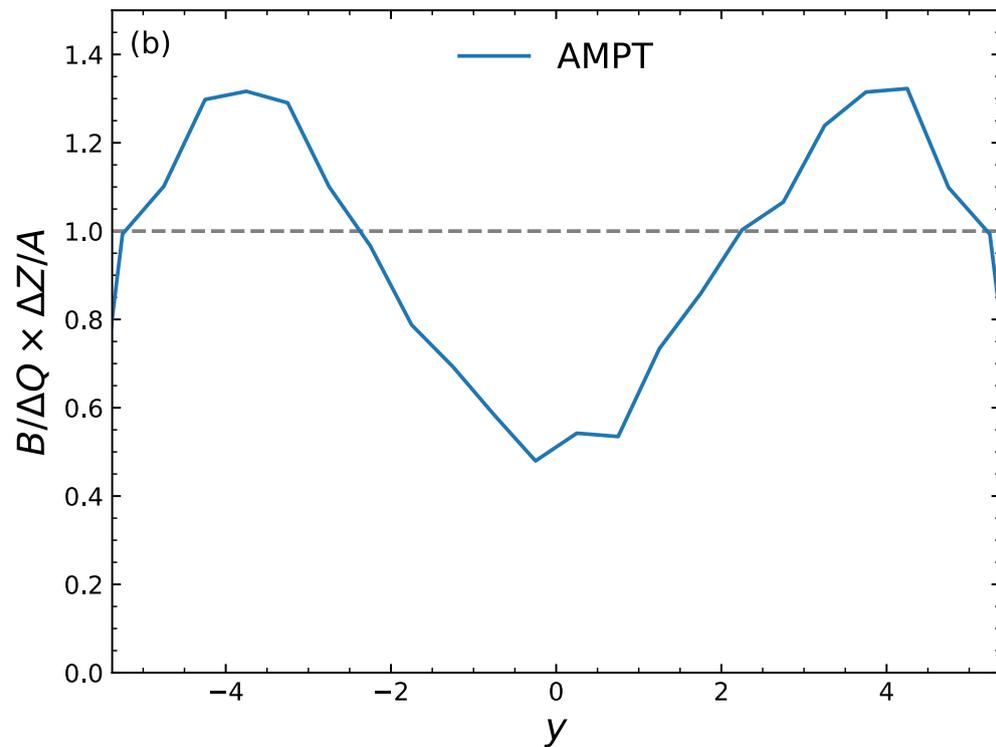


For the difference between Zr and Ru isobars:

$$B/\Delta Q * \Delta Z/A$$

is qualitatively similar to

$$B/Q * Z/A$$



UrQMD results from

N. Lewis et al., 2205.05685 [hep-ph]

## Relate B/Q\*Z/A to quark stopping

We consider a quark matter under the conventional scenario:  
baryon and electric charge are carried by quarks and antiquarks  
while neglecting the baryon junction.

Let  $f_i \equiv f_i(y) \equiv dN/dy$  of quark flavor  $i$ , we then have

$$f_B \equiv \frac{dN_B}{dy} = \frac{1}{3}(f_u - f_{\bar{u}} + f_d - f_{\bar{d}} + f_s - f_{\bar{s}}),$$
$$f_Q \equiv \frac{dN_Q}{dy} = \frac{1}{3}(2f_u - 2f_{\bar{u}} - f_d + f_{\bar{d}} - f_s + f_{\bar{s}}).$$

Let's separate the  $dN/dy$  shape from the magnitude by writing

$$N_i = \int f_i(y) dy, \text{ and } s_i(y) \equiv f_i(y)/N_i,$$

$s_i(y)$  is the normalized  $dN/dy$   
& represents the stopping of quark flavor  $i$

Consider “central” collisions, conservation laws give

$$\rightarrow N_u = 2A + 2Z + N_{\bar{u}}, \quad N_d = 4A - 2Z + N_{\bar{d}}.$$

Due to isospin symmetry of the strong interaction (& ignore neutron skin), we assume that the relative stopping of  $u$  is the same as  $d$  :

$$s_u(y) = s_d(y) \equiv s_q, \quad \text{similarly } s_{\bar{u}}(y) = s_{\bar{d}}(y) \equiv s_{\bar{q}}.$$

Then we have

$$3f_B = 6As_q + (N_{\bar{u}} + N_{\bar{d}})(s_q - s_{\bar{q}}) + N_s(s_s - s_{\bar{s}}),$$

$$3f_Q = 6Zs_q + (2N_{\bar{u}} - N_{\bar{d}})(s_q - s_{\bar{q}}) - N_s(s_s - s_{\bar{s}}).$$

*Note:  $N_u$  can still be different from  $N_d$ .*

We consider two simple limits below.

A)

If we can neglect sea-quarks and strange quarks (e.g., at very low energies):

$$3f_B = 6As_q + \cancel{(N_{\bar{u}} + N_{\bar{d}})(s_q - s_{\bar{q}})} + N_s(s_s - s_{\bar{s}}),$$

$$3f_Q = 6Zs_q + \cancel{(2N_{\bar{u}} - N_{\bar{d}})(s_q - s_{\bar{q}})} - N_s(s_s - s_{\bar{s}}).$$

we get the naive expectation (at any  $y$ ):

$$\frac{f_B}{f_Q} \frac{Z}{A} \equiv B/Q \times Z/A = 1$$

## B)

For heavy ion collisions with  $A=2Z$  (like isobar  $^{96}\text{Cd}+\text{Cd}$ ), isospin symmetry gives  $N_{\bar{u}} = N_{\bar{d}} \equiv N_{\bar{q}}$ ,

so we have

$$3f_B = 12Zs_q + 2N_{\bar{q}}(s_q - s_{\bar{q}}) + \underline{N_s(s_s - s_{\bar{s}})},$$
$$3f_Q = 6Zs_q + N_{\bar{q}}(s_q - s_{\bar{q}}) - \underline{N_s(s_s - s_{\bar{s}})}.$$

- Without the strangeness terms, we have (at any  $y$ )

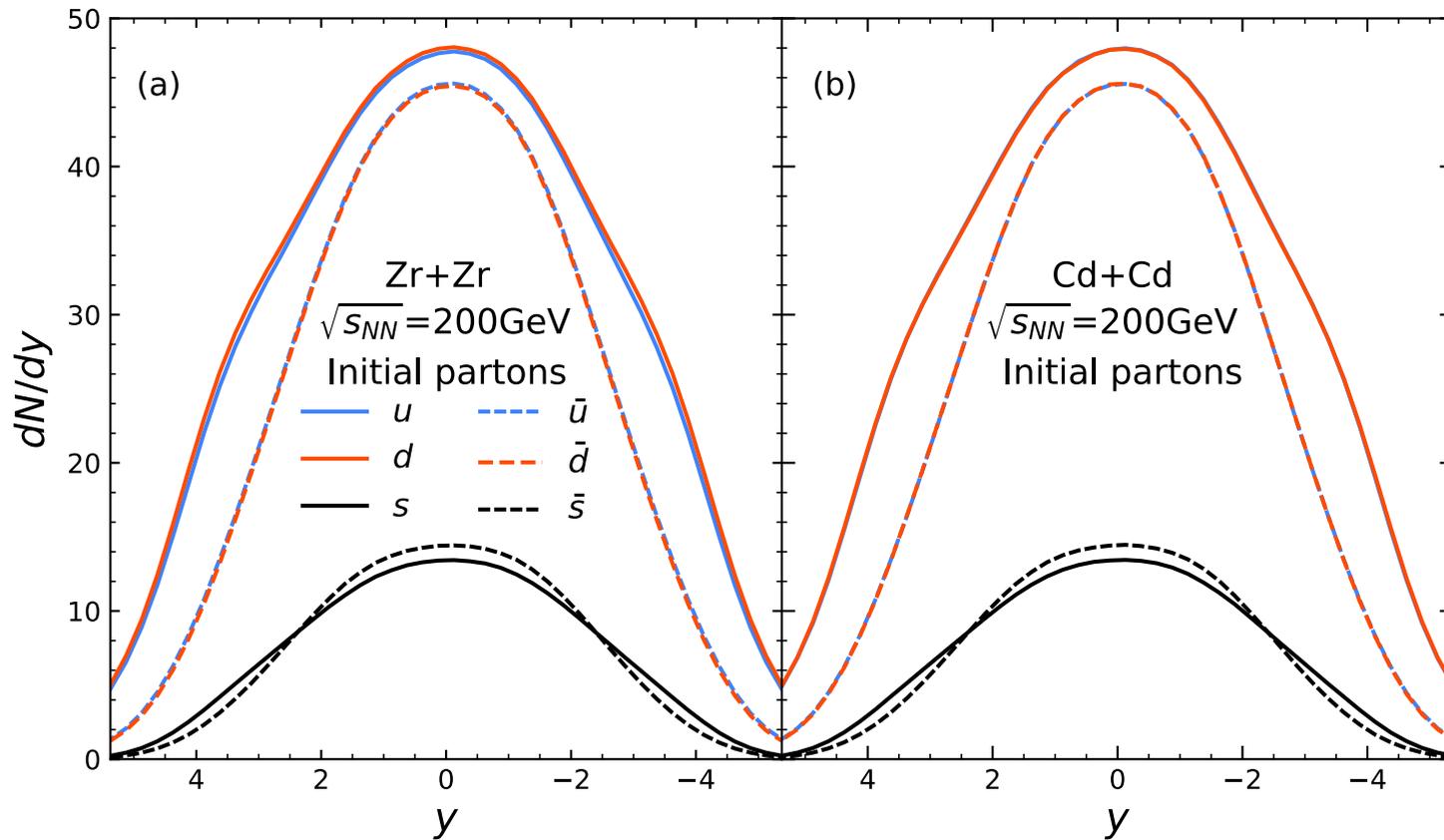
$$f_B/f_Q = 2, \text{ or } B/Q \times Z/A = 1$$

- $B/Q * Z/A \neq 1$  (for isospin symmetric nucleus)  
is due to (anti)strange quarks  
and it means  $s_s \neq s_{\bar{s}}$

*Note: same conclusion for non-central events (when ignoring neutron skin)*

# Quark results from string melting AMPT

for minimum-bias  $^{96}\text{Zr}+\text{Zr}$  &  $^{96}\text{Cd}+\text{Cd}$



- We see  $s_u = s_d$ ,  $s_{\bar{u}} = s_{\bar{d}}$  (previous assumptions)
- $s_s \neq s_{\bar{s}}$  even for isospin symmetric  $^{96}\text{Cd}+\text{Cd}$
- More  $\bar{s}$  than  $s$  at mid-rapidity  
 ---> this decreases B, increases Q & causes  $B/Q \cdot Z/A < 1$

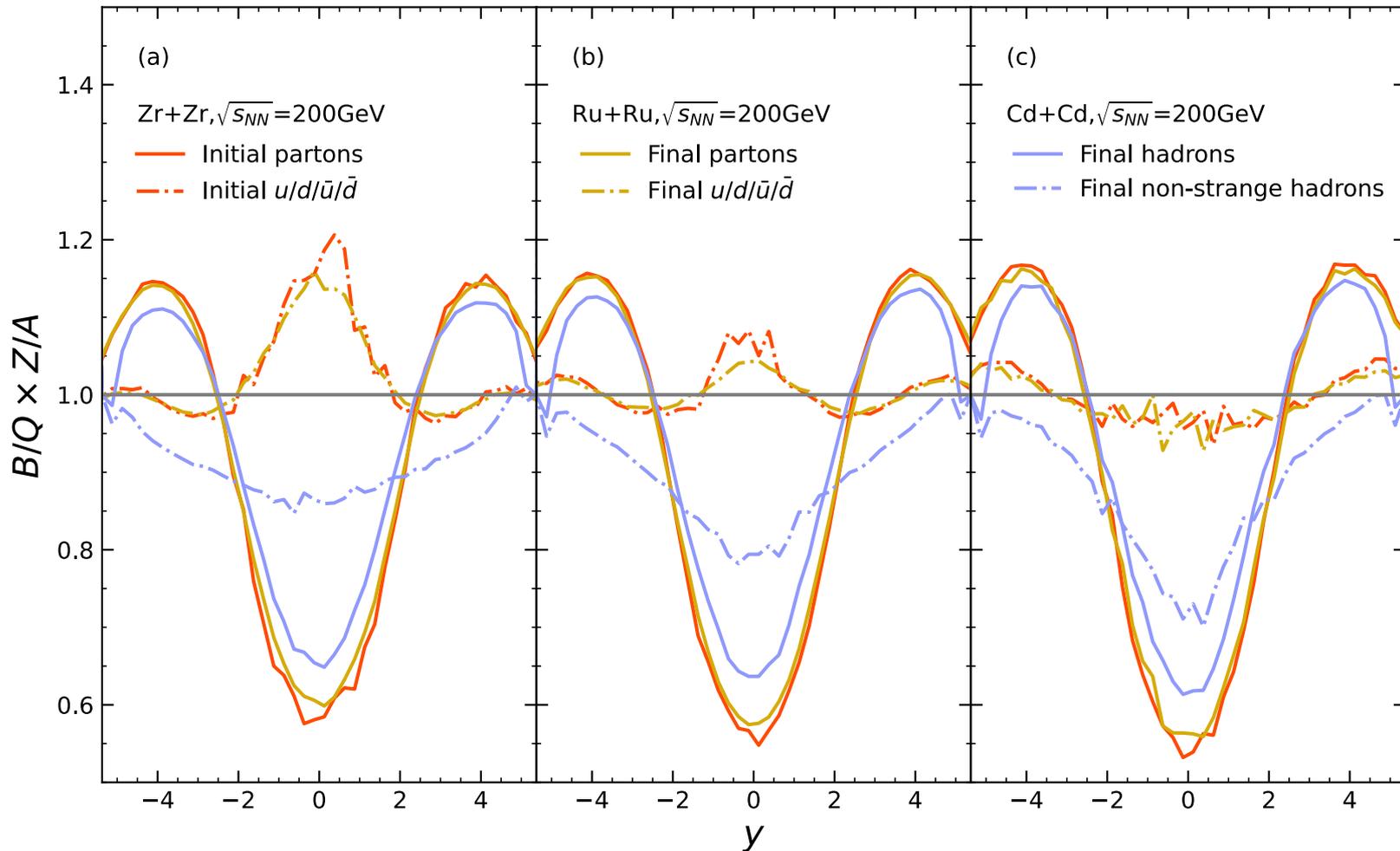
# Results from string melting AMPT

for MB

$^{96}\text{Zr}+\text{Zr}$

$^{96}\text{Ru}+\text{Ru}$

$^{96}\text{Cd}+\text{Cd}$



When excluding  $\bar{s}$  and  $s$  quarks

- $^{96}\text{Cd}+\text{Cd}$  gives parton  $B/Q * Z/A \approx 1.0$  (as expected)
- Parton  $B/Q * Z/A$  is closer to 1.0 for  $\text{Ru}+\text{Ru}$  than  $\text{Zr}+\text{Zr}$  (as expected)

$\rightarrow B/Q * Z/A \neq 1$  is mainly due to  $s_s \neq s_{\bar{s}}$

# On reconstruction of net charges

Using AMPT results, we check the experimental reconstruction method by comparing “Reconstructed” charge with the direct model result.

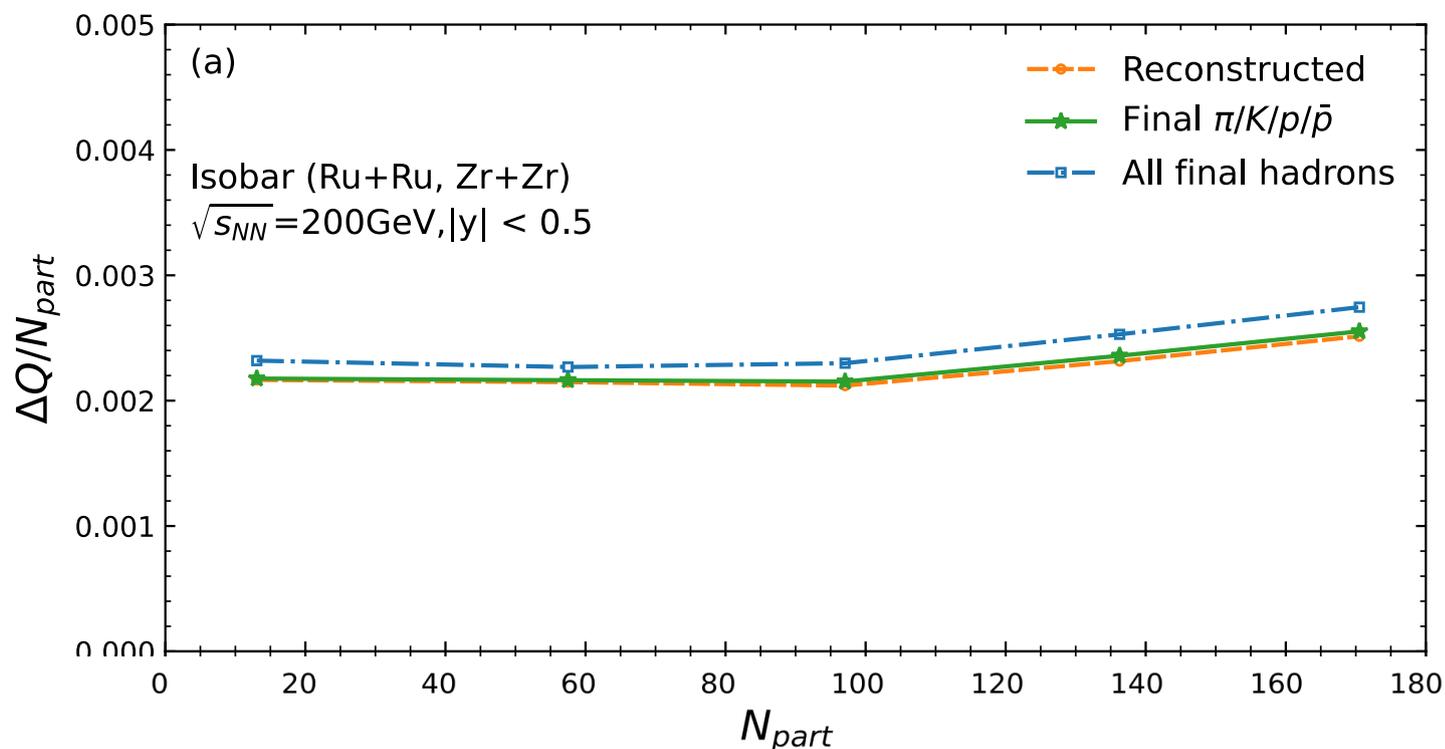
## Net-electric charge reconstruction:

$$\Delta Q = N_{\pi}[(R2_{\pi} - 1) + \frac{N_K}{N_{\pi}}(R2_K - 1) + \frac{N_p}{N_{\pi}}(R2_p - 1)].$$

$$R2_{\pi} = (N_{\pi^+}^{\text{Ru}} / N_{\pi^-}^{\text{Ru}}) / (N_{\pi^+}^{\text{Zr}} / N_{\pi^-}^{\text{Zr}})$$

CY Tsang (for STAR)  
at GHP 2023 &  
Quark Matter 2023

→ Reconstructed  
 $\Delta Q$  is accurate  
(from  $\pi/K/p/\bar{p}$ )



## Net-baryon reconstruction:

$$B_{net} = (N_p - N_{\bar{p}}) + (N_n - N_{\bar{n}})$$

$$\approx (N_p - N_{\bar{p}}) + \bar{p} \sqrt{\frac{d}{\bar{d}}} - p \sqrt{\frac{\bar{d}}{d}}$$

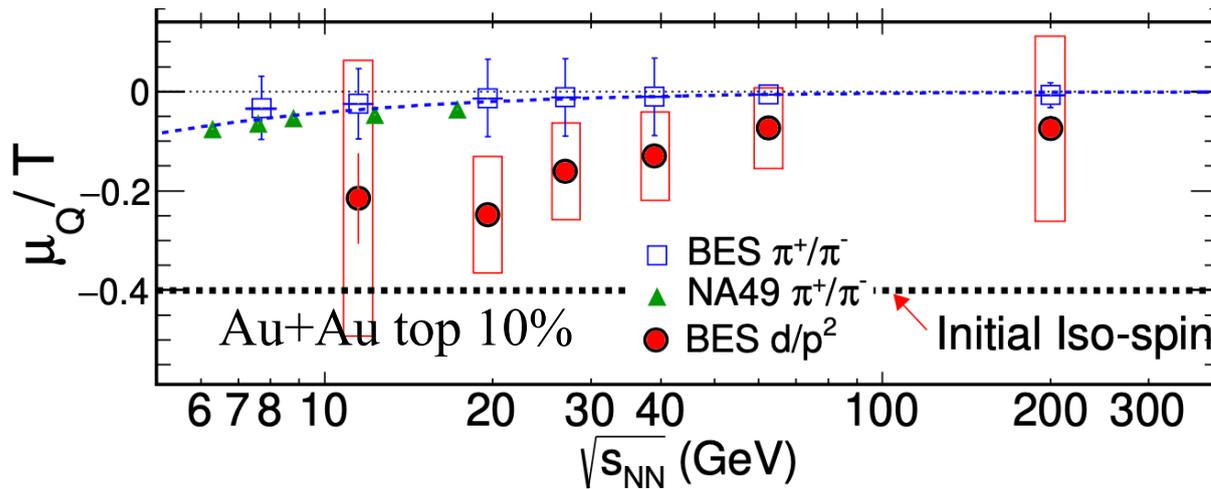
CY Tsang (for STAR)  
at GHP 2023 &  
Quark Matter 2023

This is based on  
the thermal model:

$$\frac{\mu_Q}{T} = \frac{1}{2} \ln \left( \frac{\pi^+}{\pi^-} \right),$$

$$\frac{\mu_Q}{T} = \frac{1}{2} \ln \left( \frac{\bar{d}/\bar{p}^2}{d/p^2} \right).$$

STAR Collaboration,  
Phys Rev C 99, 064905 (2019)



However,

“ The  $\mu_Q/T$  values extracted from  $\pi^+/\pi^-$  are systematically larger than those from  $(\bar{d}/\bar{p}^2)/(d/p^2)$  at small  $\sqrt{s_{NN}}$  ”

→ Using  $d$  &  $\bar{d}$  data reconstructs more net-neutrons than using  $\pi^+$  &  $\pi^-$  data due to the uncertainty of  $\mu_Q/T$  values

# Summary

A Multi-Phase Transport (AMPT) model  
enables us to study baryon production at both parton and hadron phases

Model results for isobars show

$B/Q * Z/A$  is strongly  $y$ -dependent &  $< 1$  at mid-rapidity

Simple analysis shows

$s$ -stopping  $\neq \bar{s}$ -stopping

is the main reason why parton  $B/Q * Z/A \neq 1$ ,

is the only reason for isospin symmetric A+A collisions

Reconstruction of the net-baryon number suffers from  
the uncertainty of  $\mu_Q/T$  values extracted from different observables

*Thanks for your attention!*