

LIPEI DU

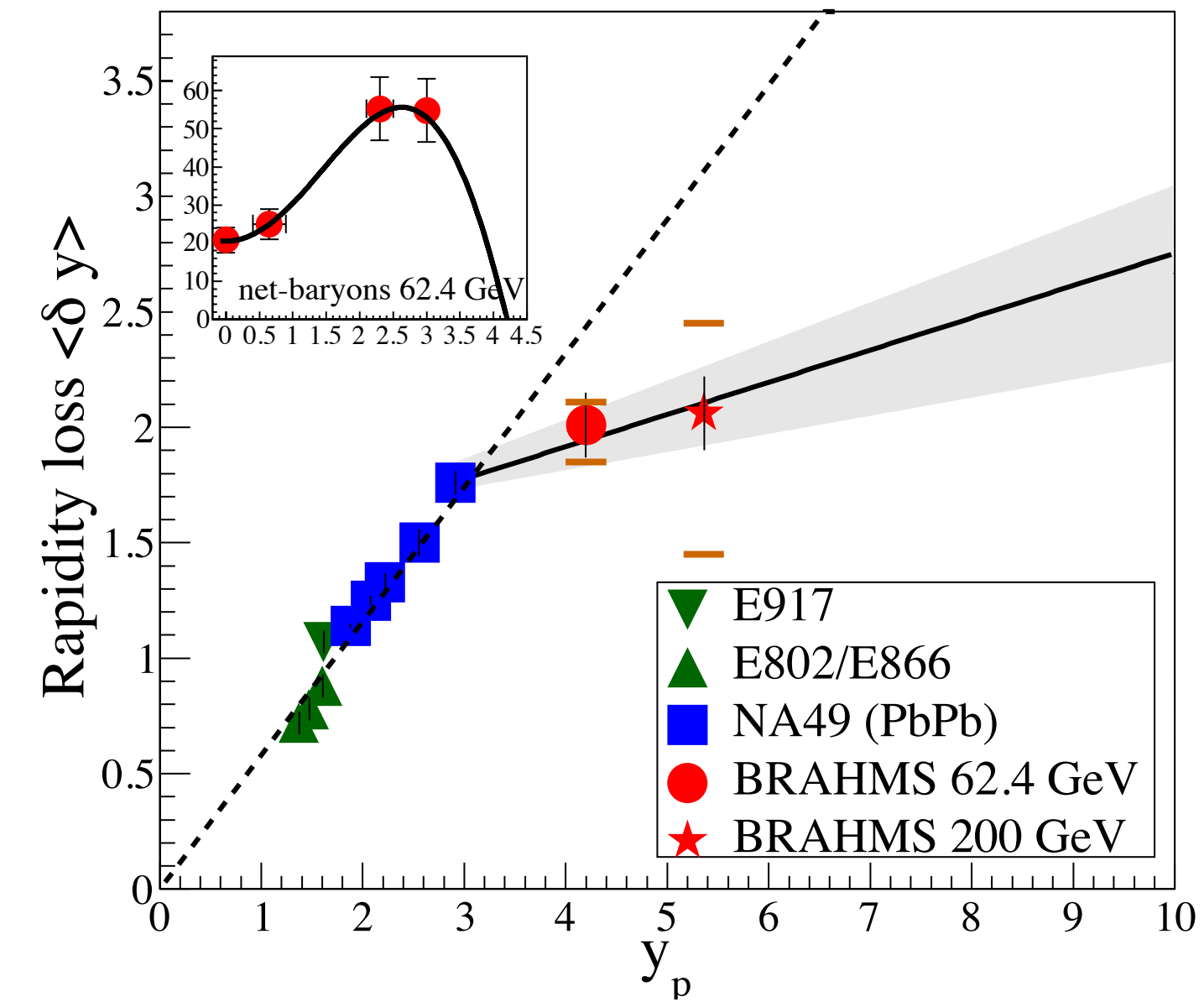
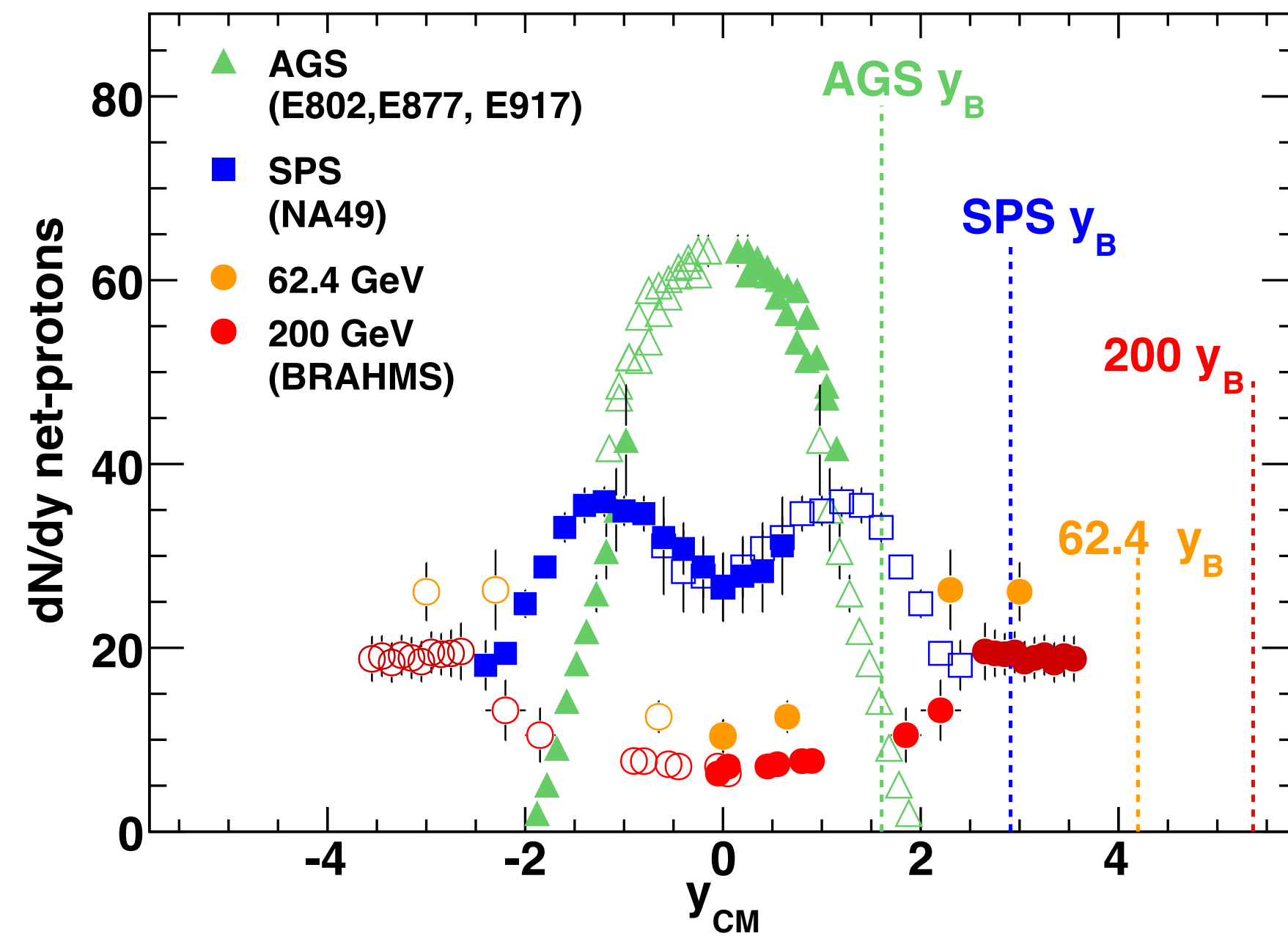
McGILL U. & LBNL

**BARYON STOPPING:
NUCLEON DECELERATION vs STRING JUNCTION BREAKING**

1ST WORKSHOP ON BARYON DYNAMICS FROM RHIC TO EIC

CENTER FOR FRONTIERS IN NUCLEAR SCIENCE, STONY BROOK UNIVERSITY

JANUARY 24, 2024



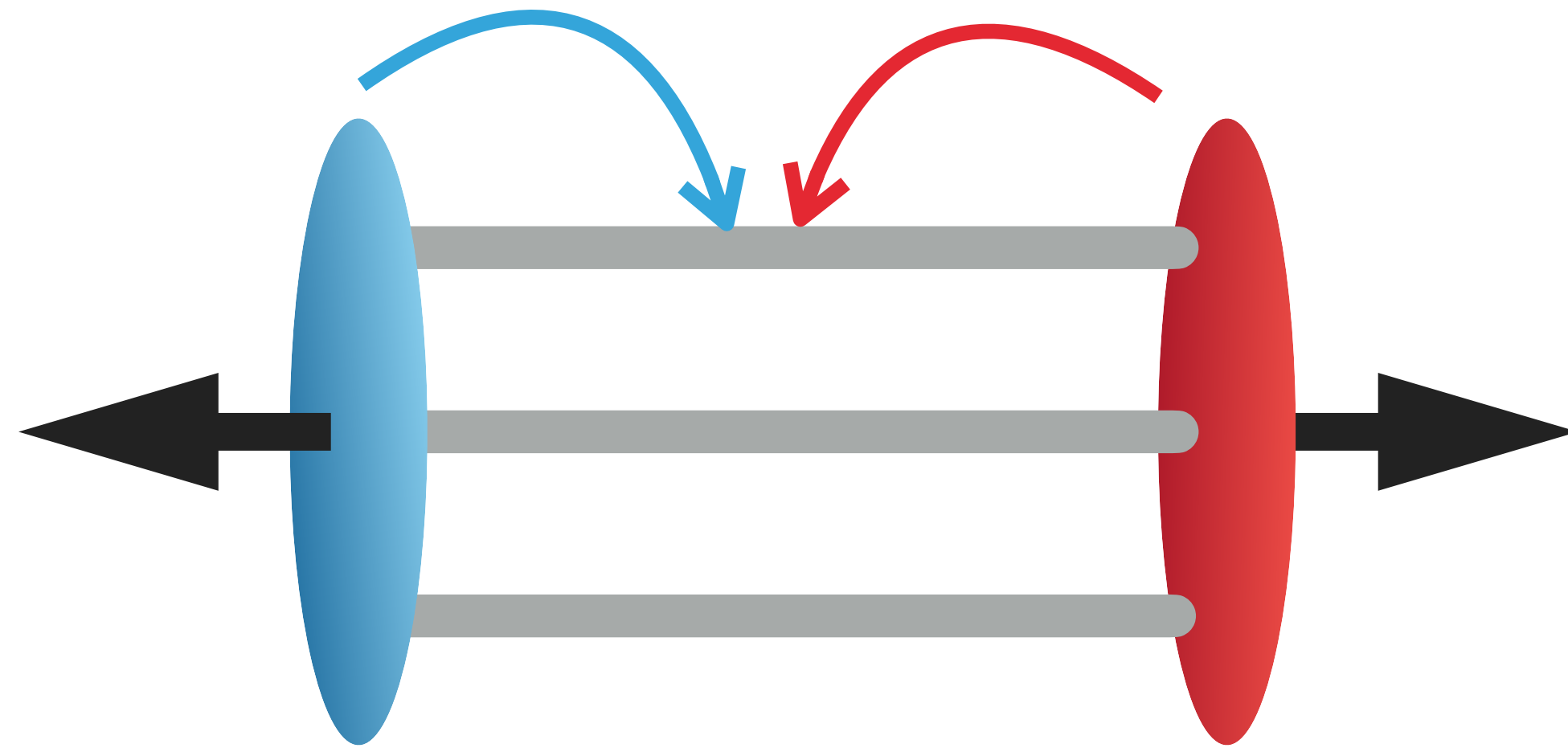
F. Videbaek, Nucl. Phys. A 830 (2009) 43C; BRAHMS, Phys. Lett. B677, 267-271 (2009)

► Baryon rapidity loss

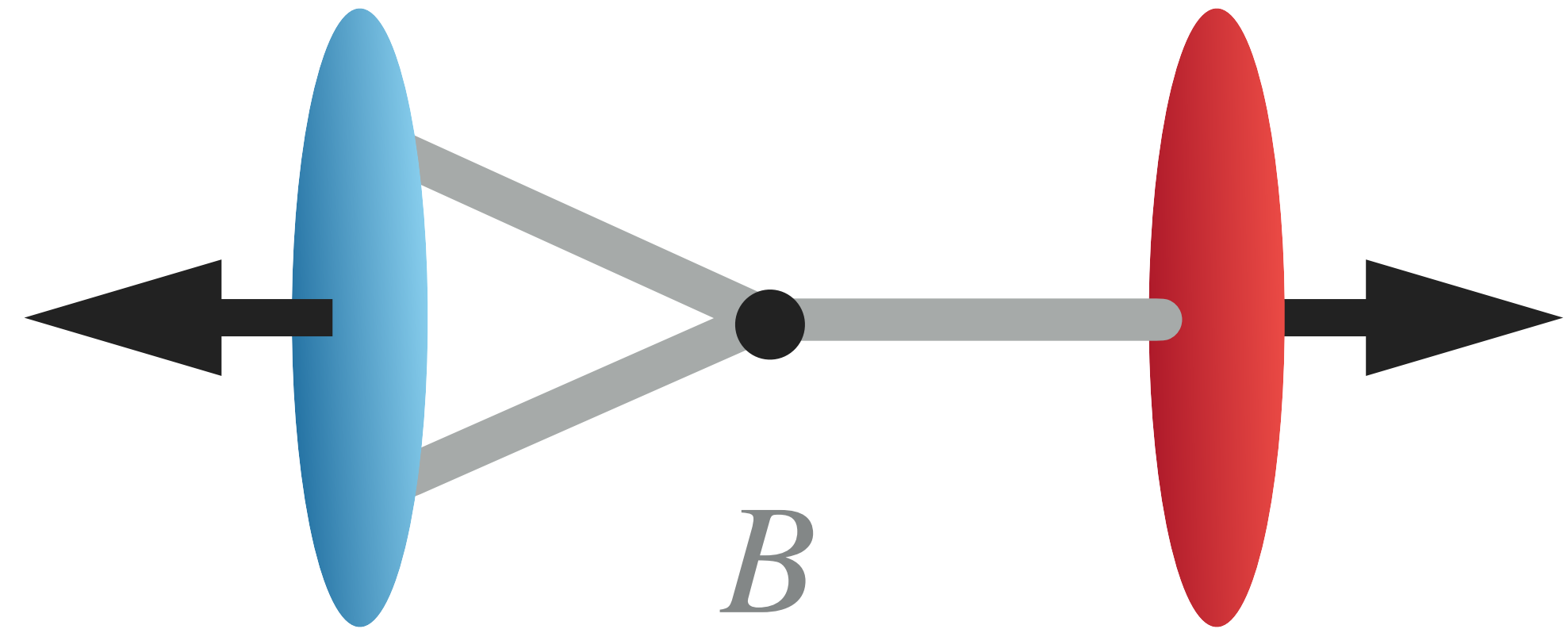
$$\delta y = y_b - \frac{2}{N_{\text{part}}} \int_0^{y_b} y \frac{dN}{dy} dy$$

Lipei Du, arXiv: 2401.00596

BARYON STOPPING & ENERGY DEPOSITION

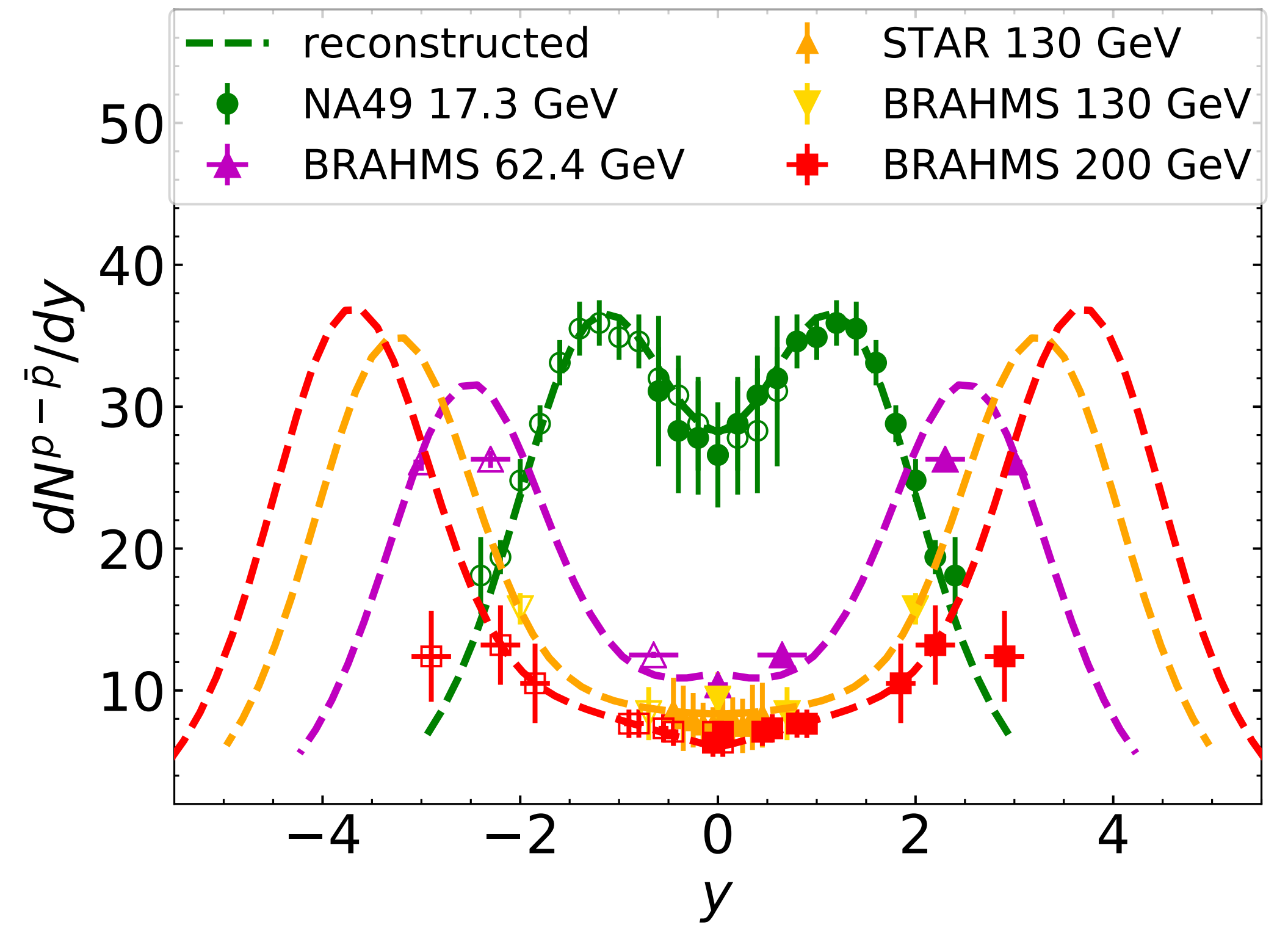
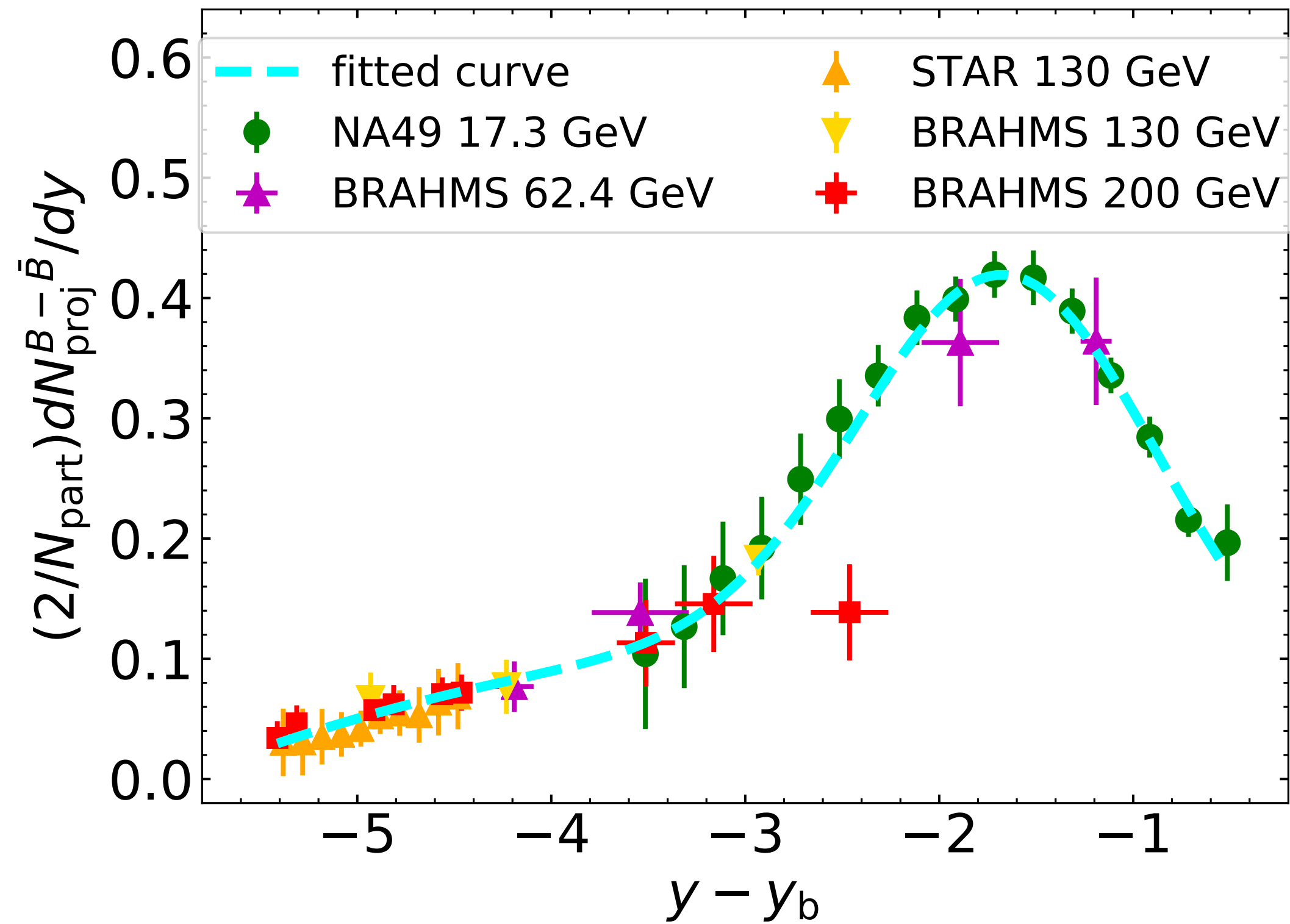


Baryons get distributed in rapidity by
deceleration of the incoming nucleons



Baryons get distributed in rapidity
through string junction breaking

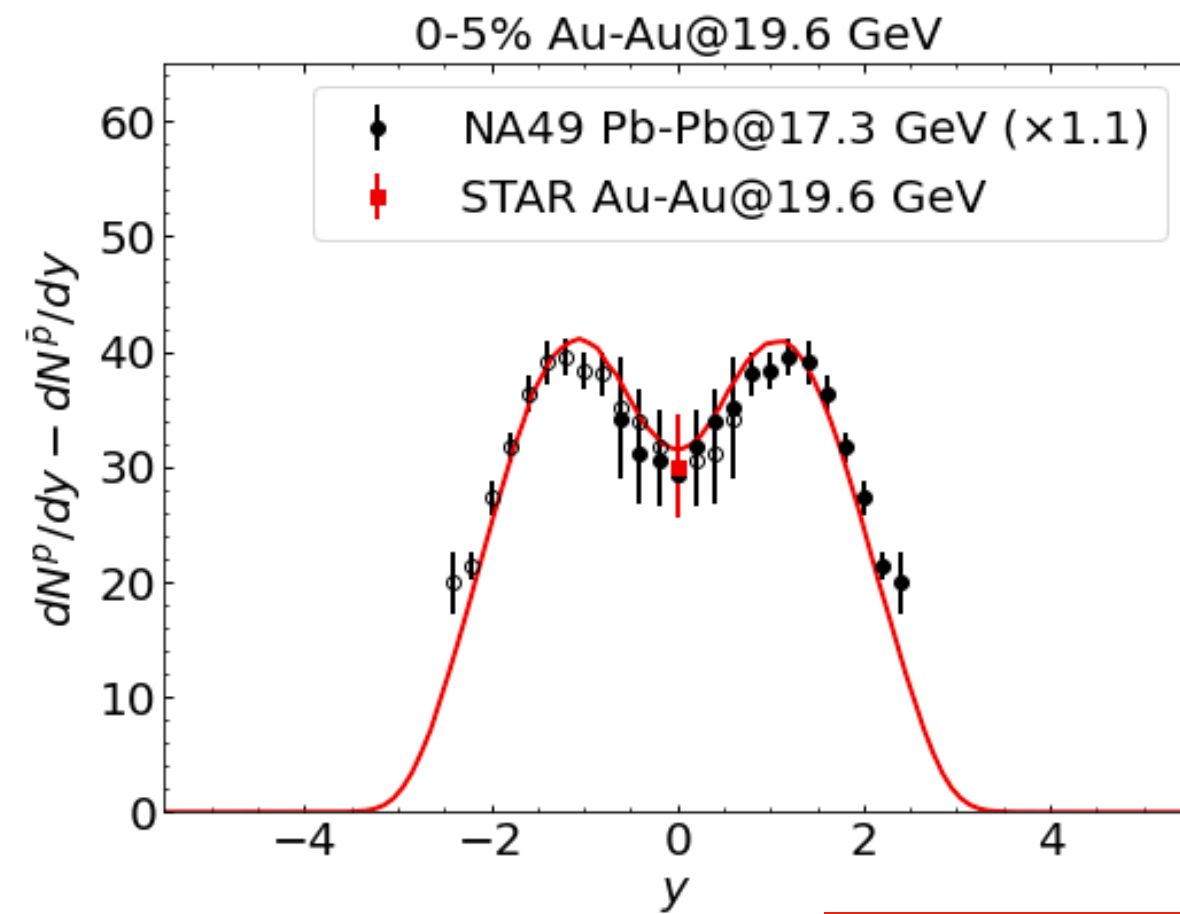
- ▶ “Baryon deceleration” and “string junction breaking” should correspond to largely different energy loss



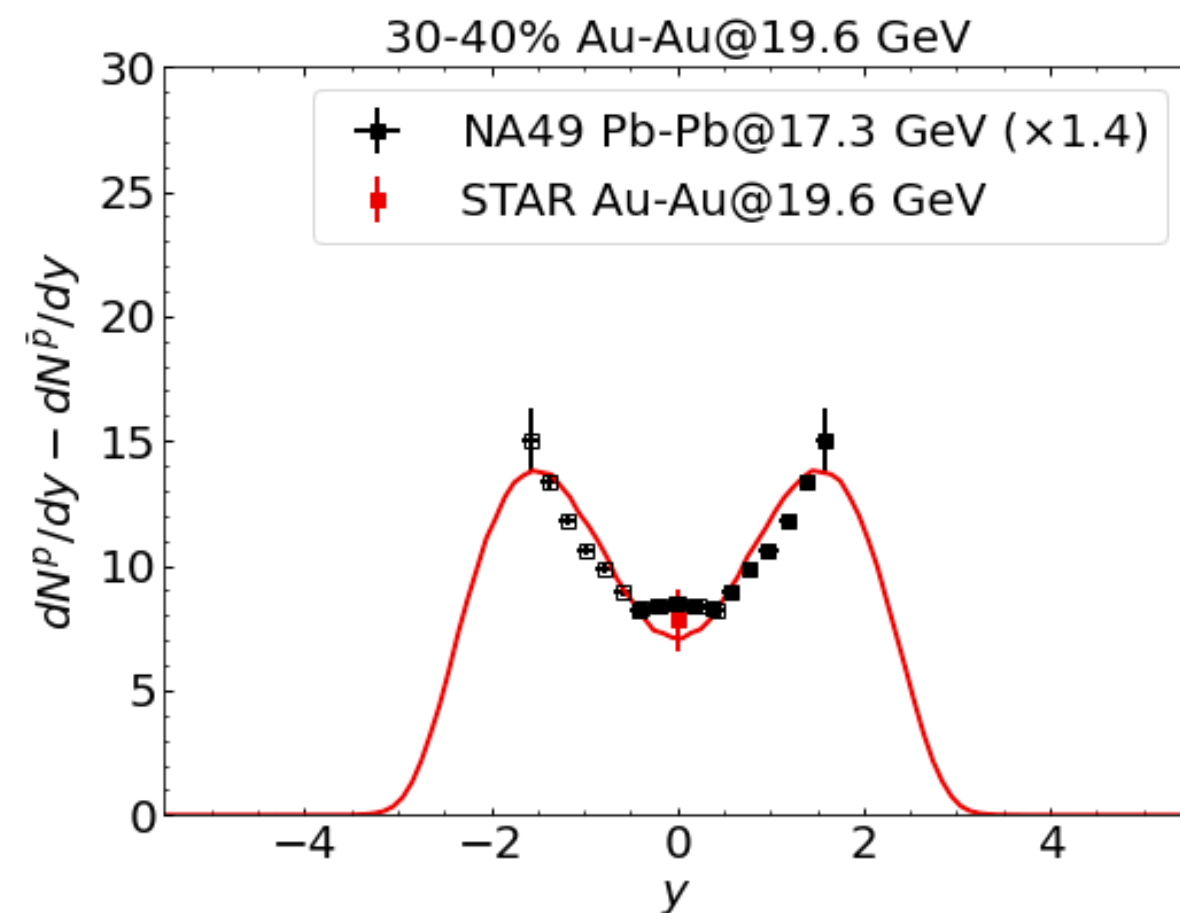
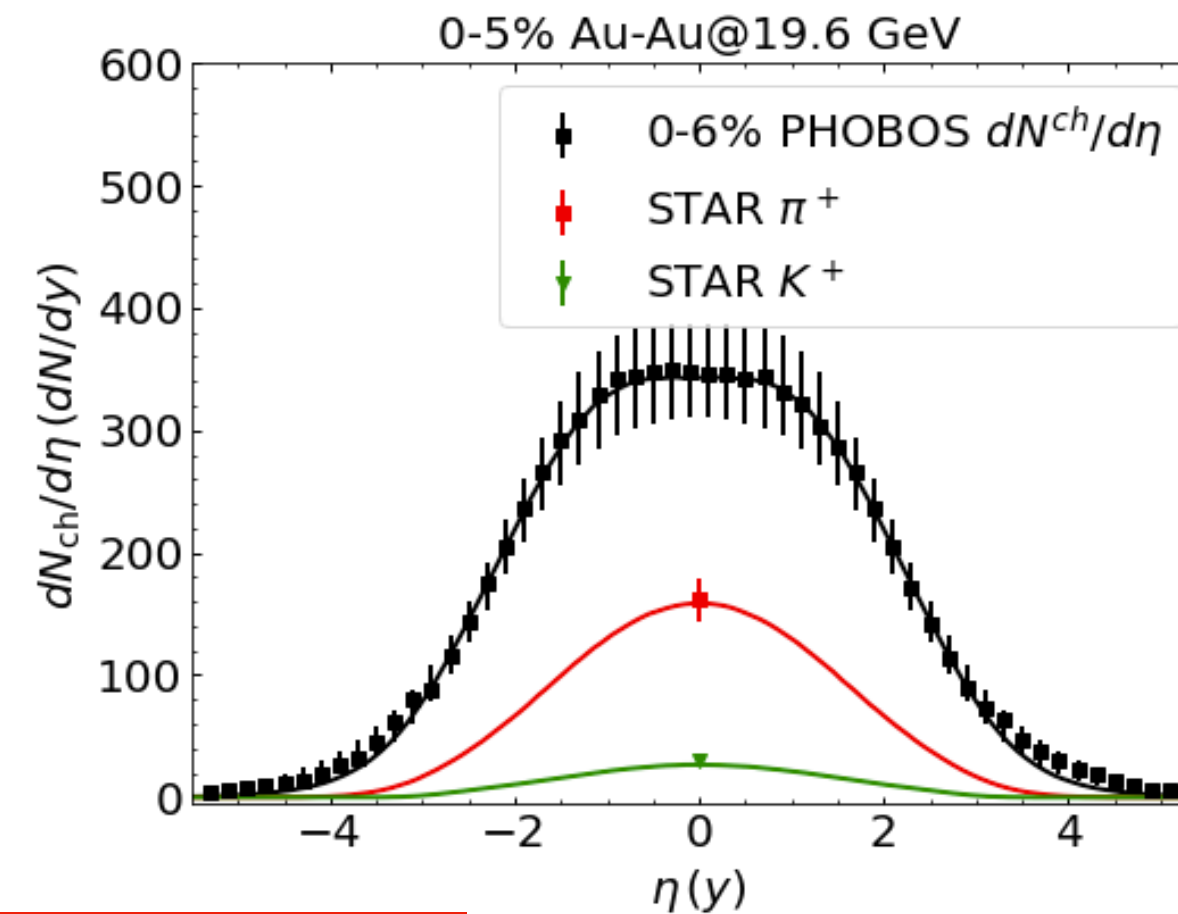
- Data at 130 GeV also fall on the universal curve
- The fitted curve can be used to reconstruct the rapidity density of net protons

ENERGY DEPOSITION

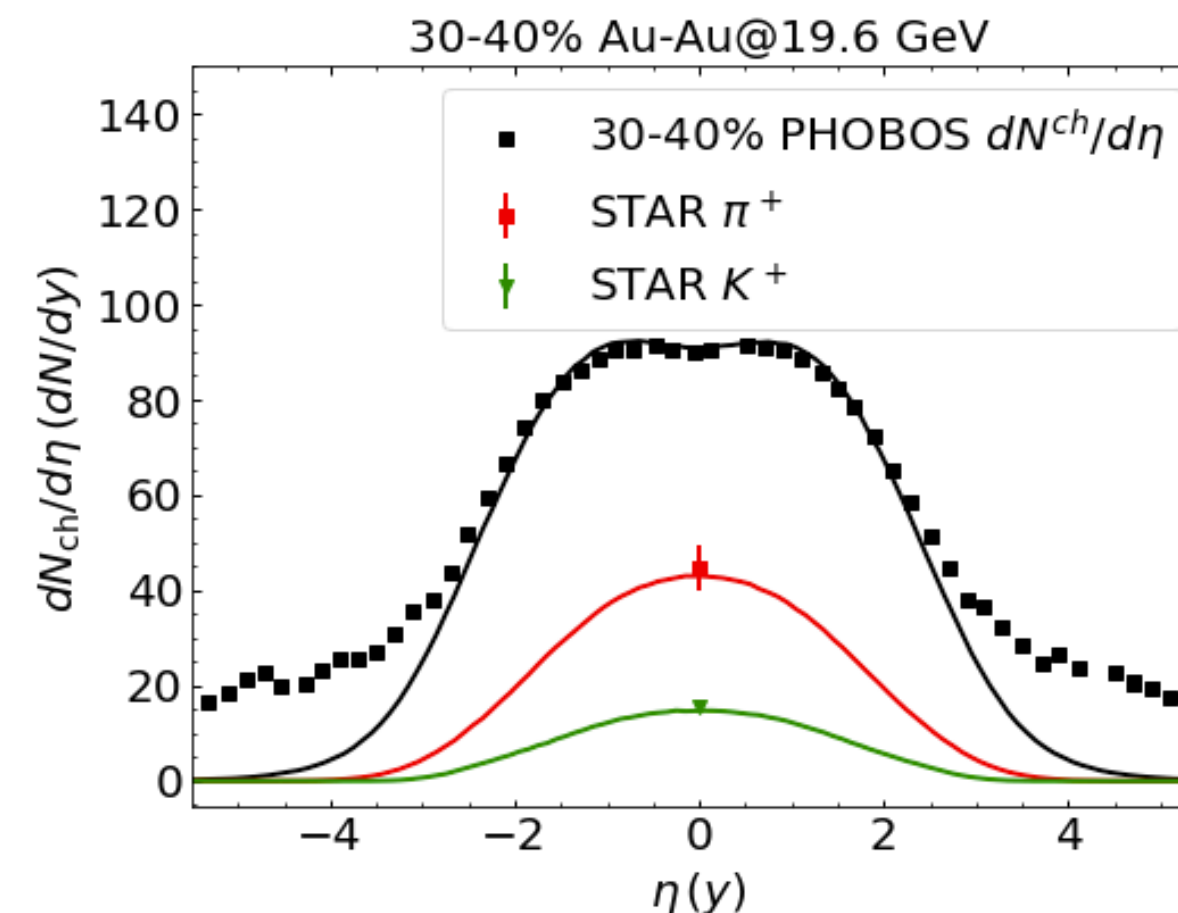
6



0-5% Au-Au @ 19.6 GeV



30-40% Au-Au @ 19.6 GeV



- ▶ Number of participants vs. total net baryon charge:

$$N^{B-\bar{B}} \stackrel{?}{=} N_{\text{part}}$$

- ▶ The energy carried by incoming participant nucleons

$$E_{\text{part}} = N_{\text{part}} \sqrt{s_{\text{NN}}} / 2$$

- ▶ Total energy deposited in the fireball

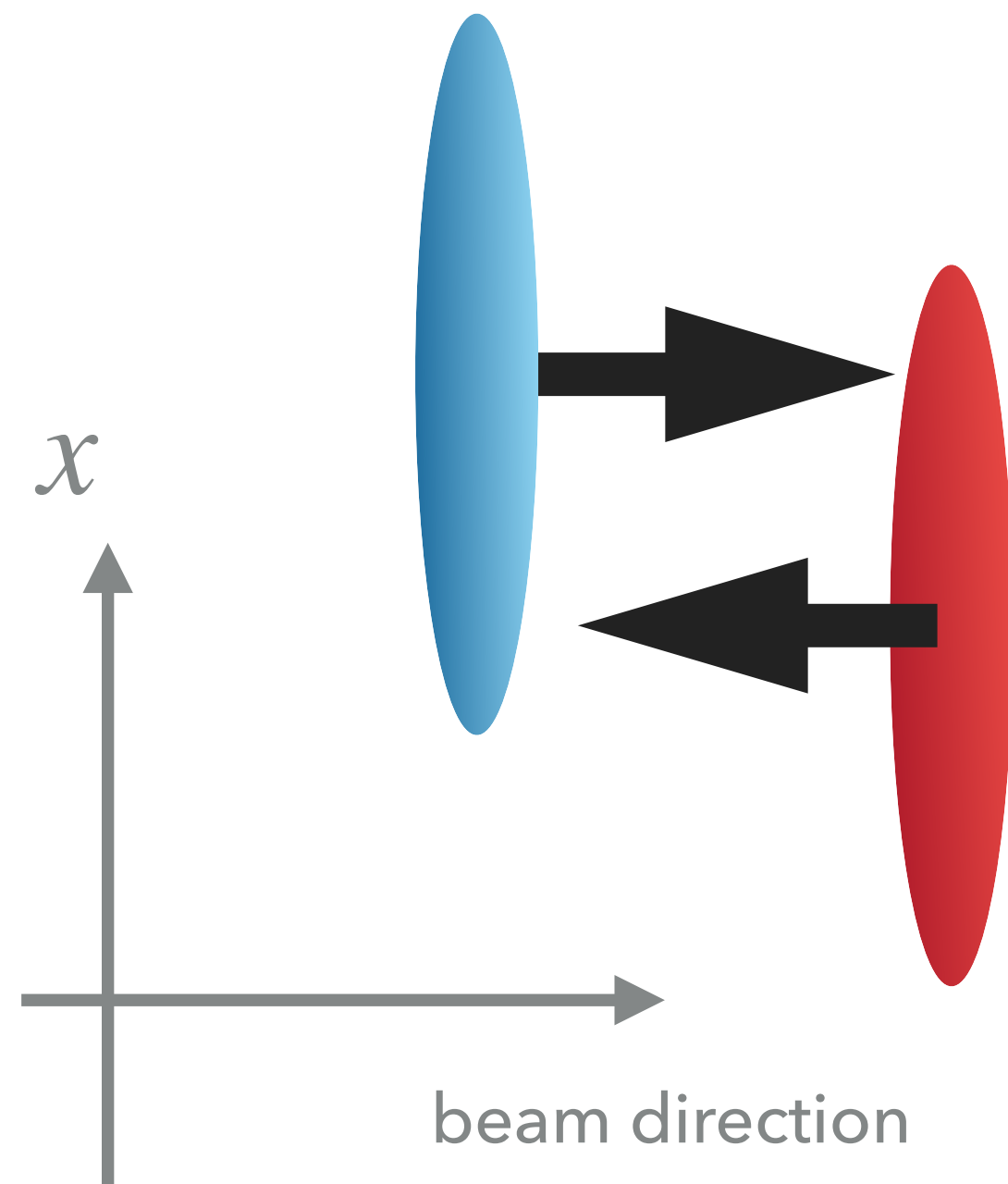
$$E_{\text{hydro}} = \int d^3 \Sigma_{\mu} T^{\mu t}(\mathbf{x}_{\perp}, \eta_s)$$

- ▶ How are the energies related

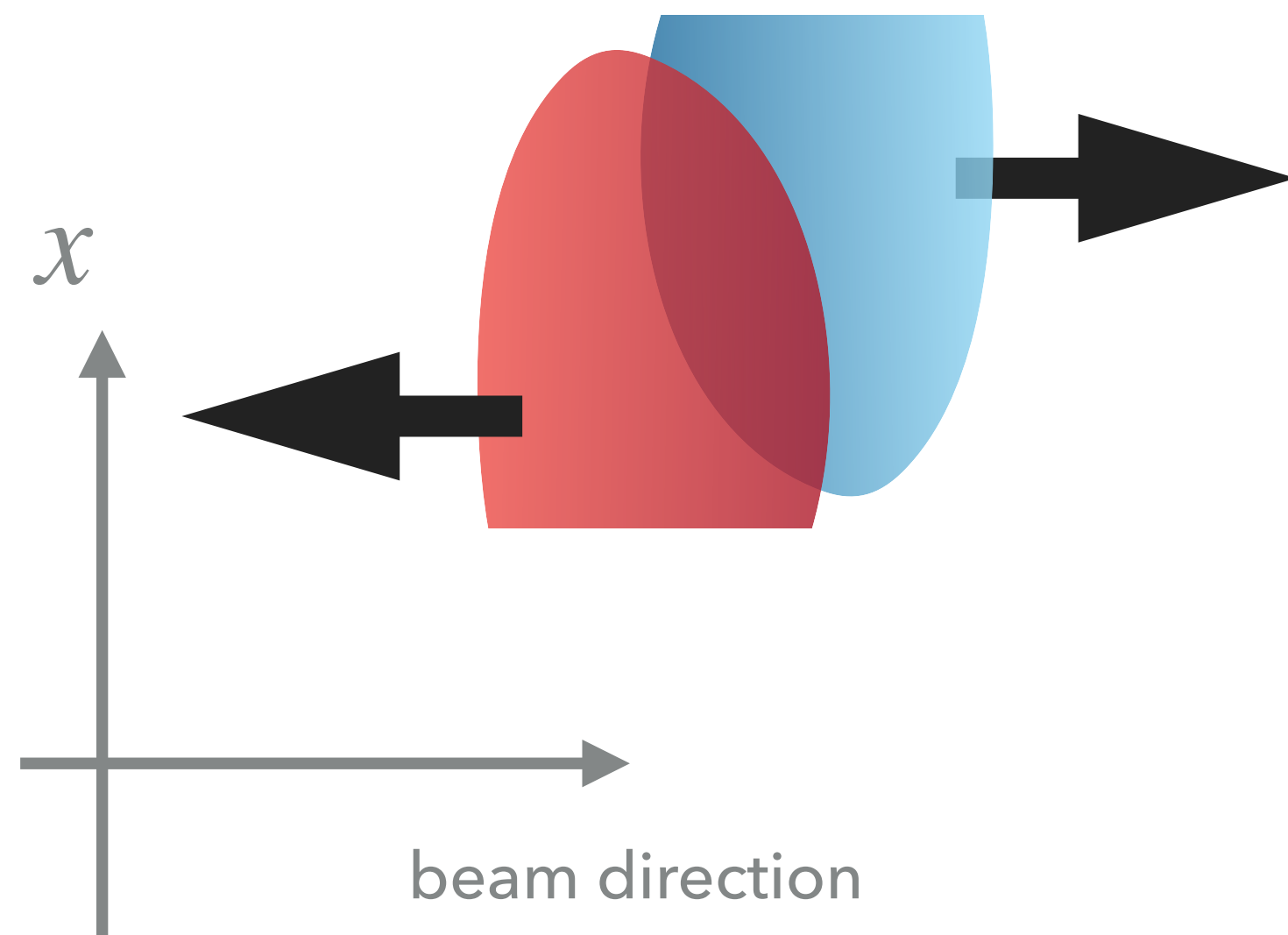
$$E_{\text{hydro}} \stackrel{?}{=} E_{\text{part}}$$

LD, C. Shen, S. Jeon & C. Gale, PRC 108 (2023) L041901

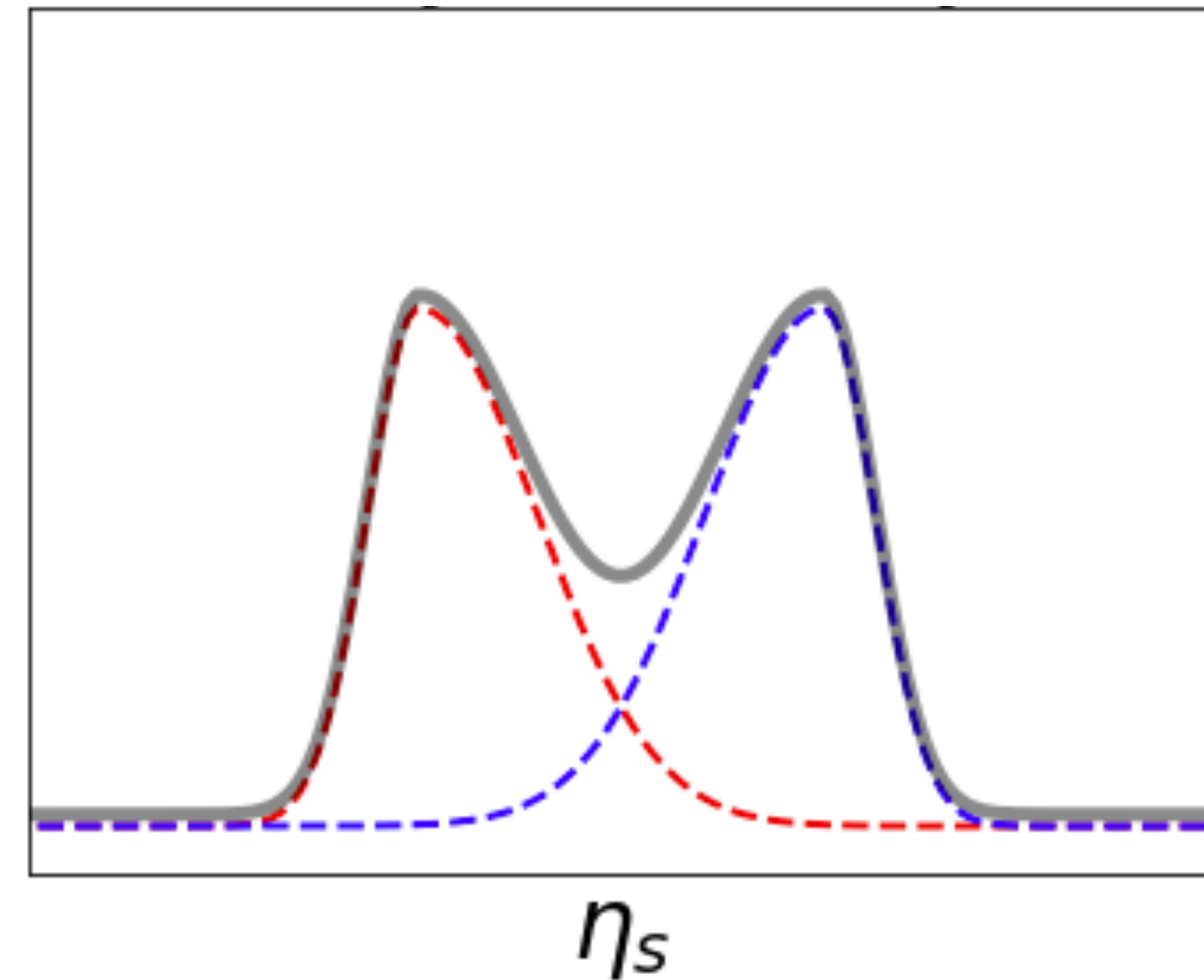
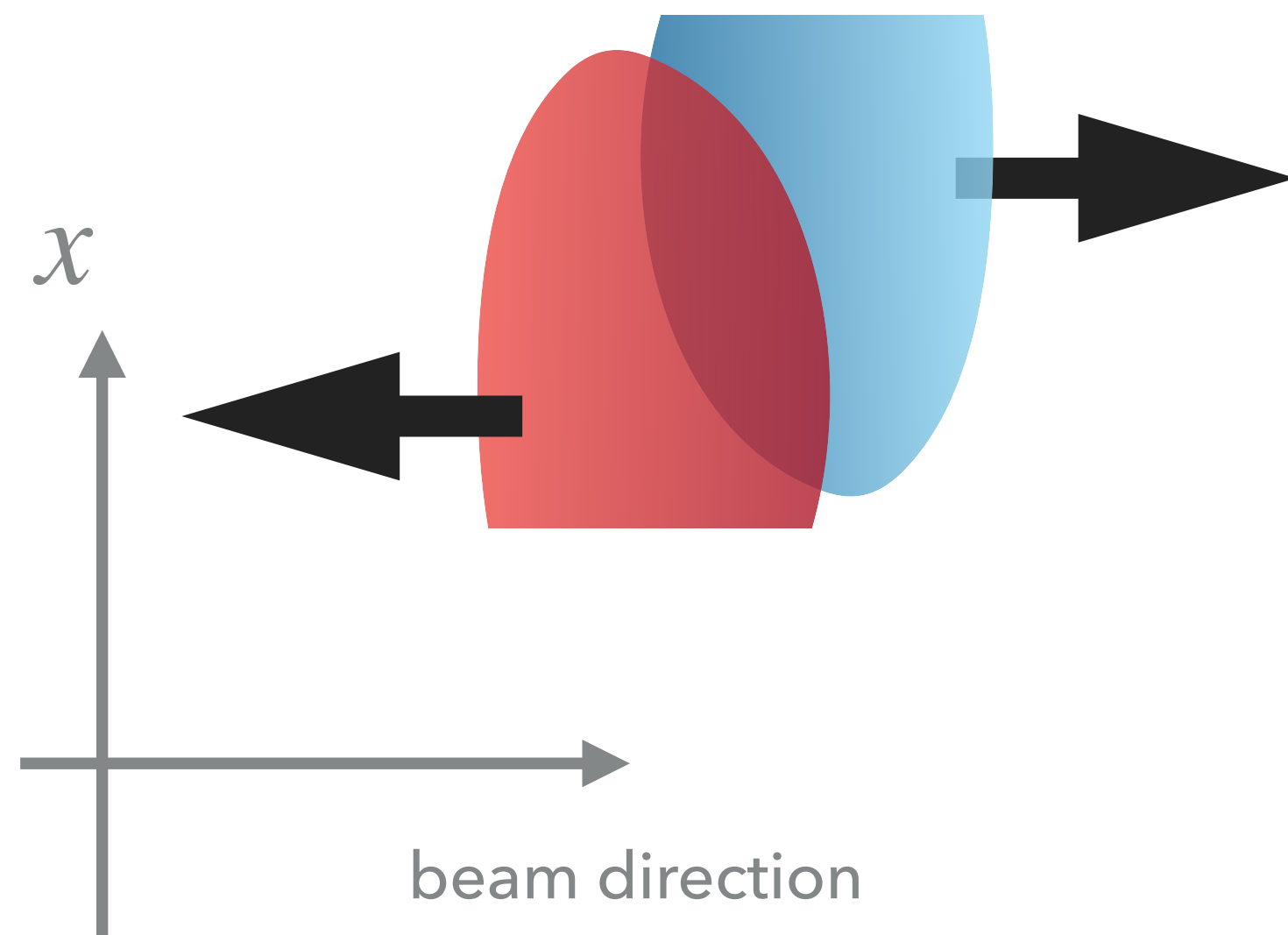
$v_1(y)$ CONSTRAINTS ON BARYON STOPPING



- ▶ From the nucleon deceleration picture, the baryon density profile gets two peaks, naturally giving the double-humped net proton yields in rapidity.

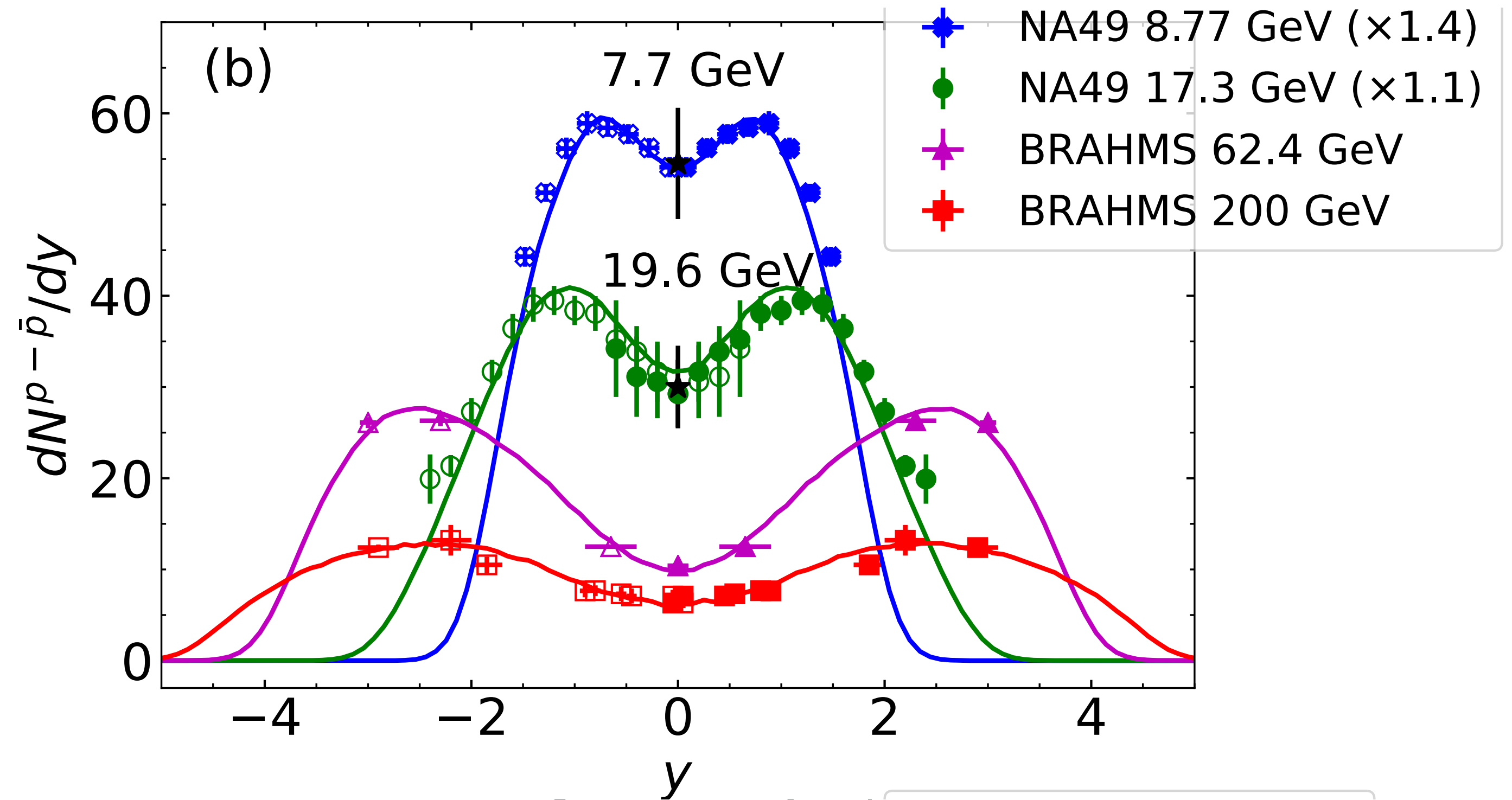
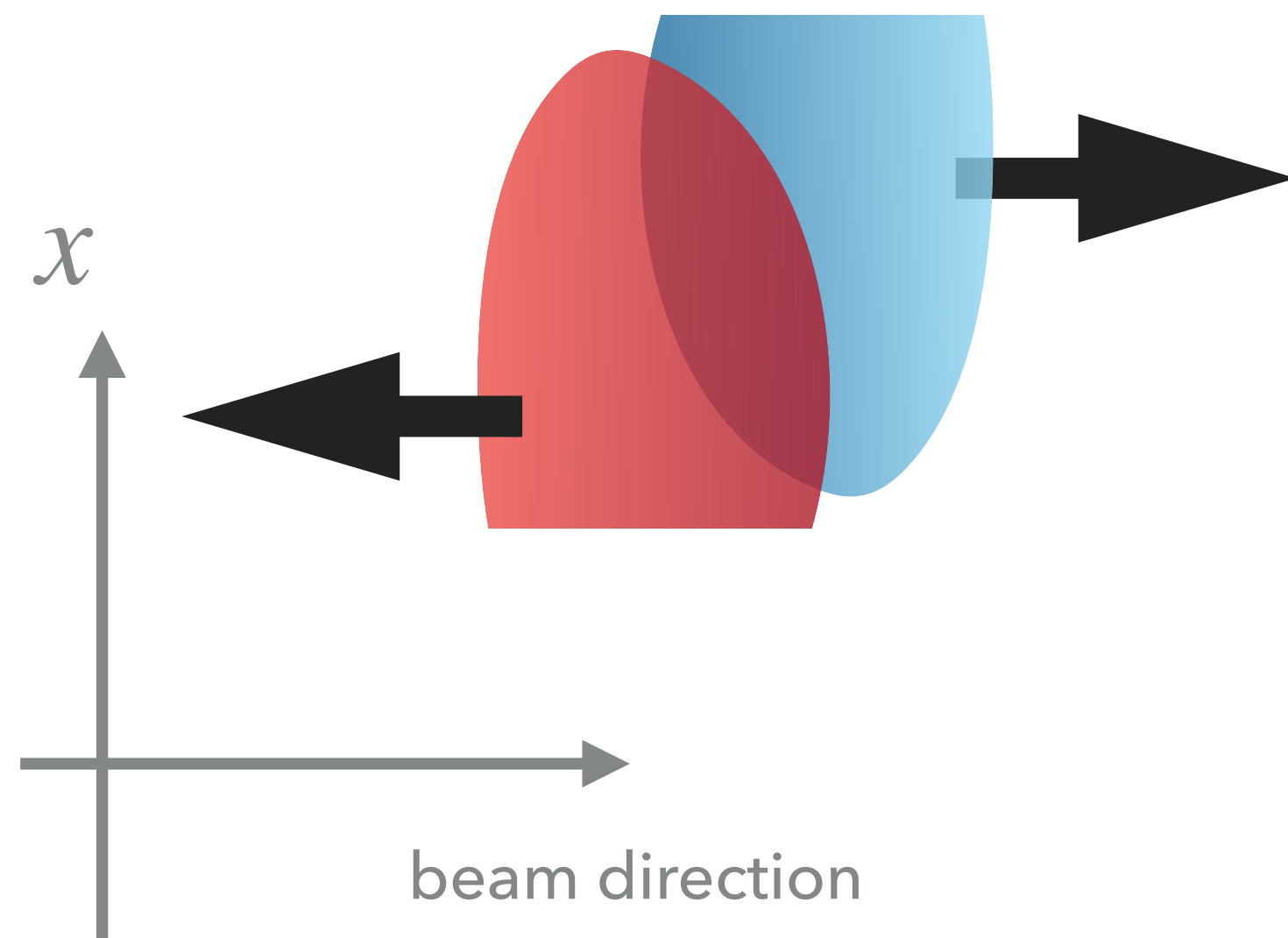


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longitudinal baryon profile

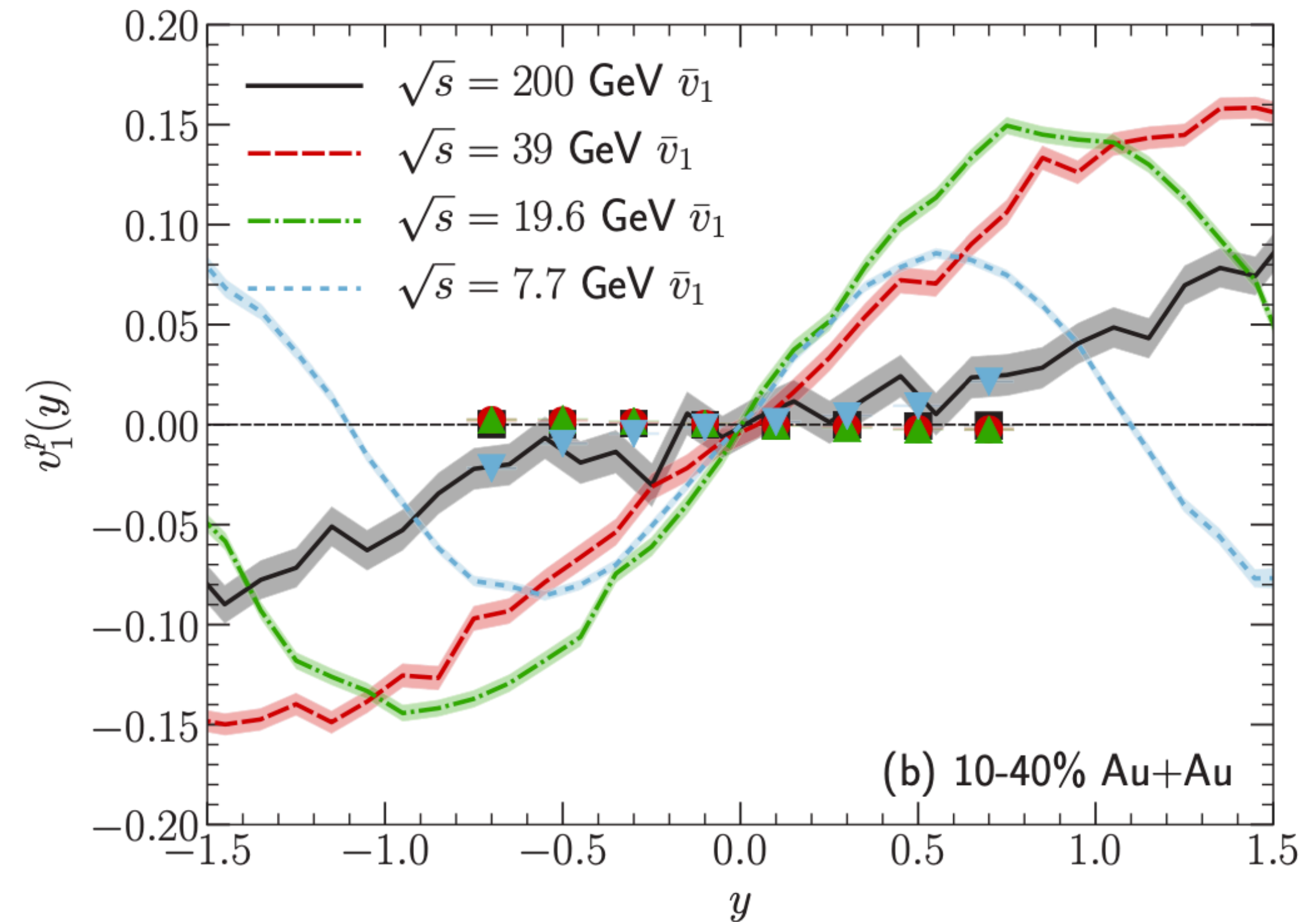
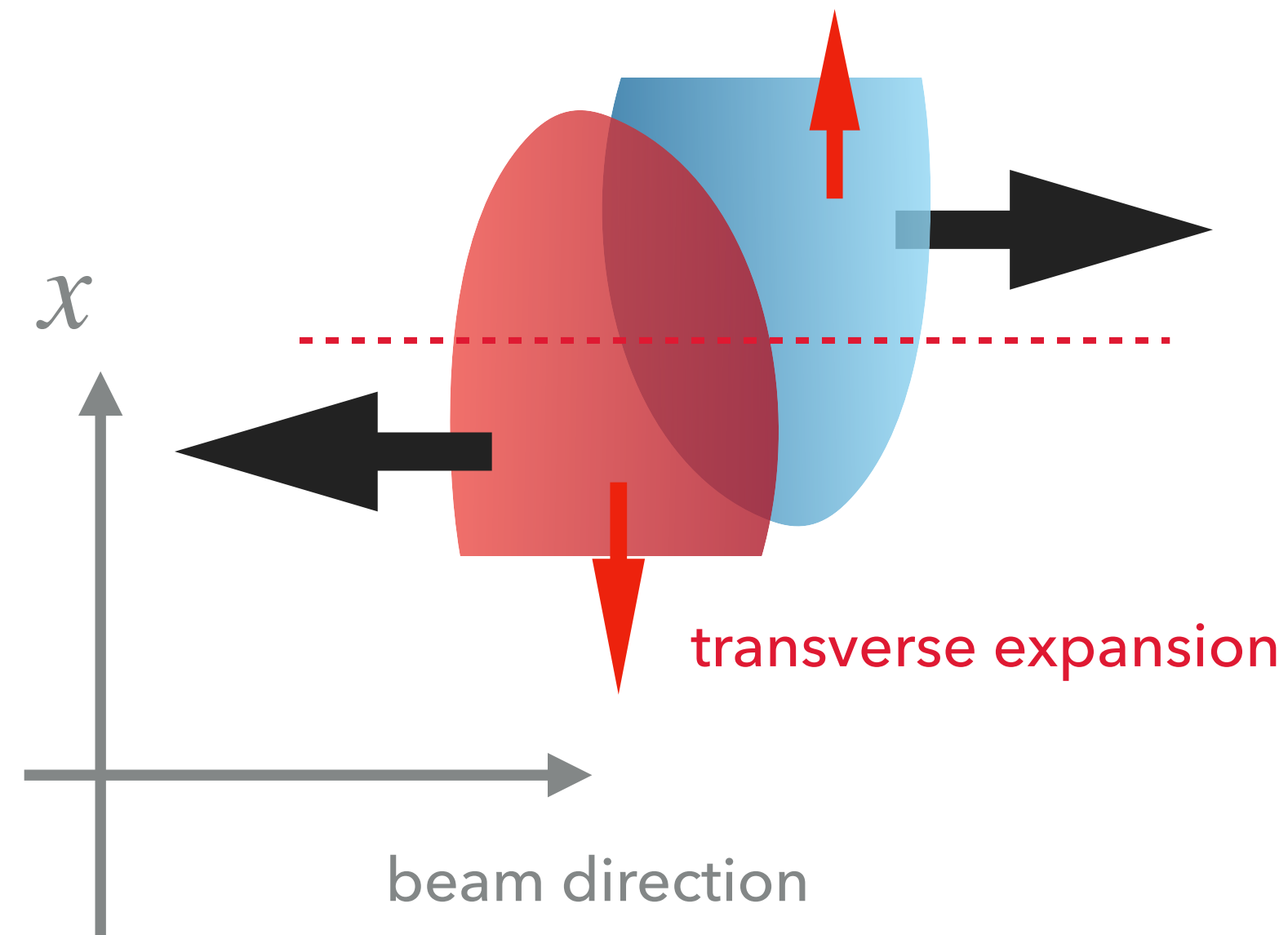
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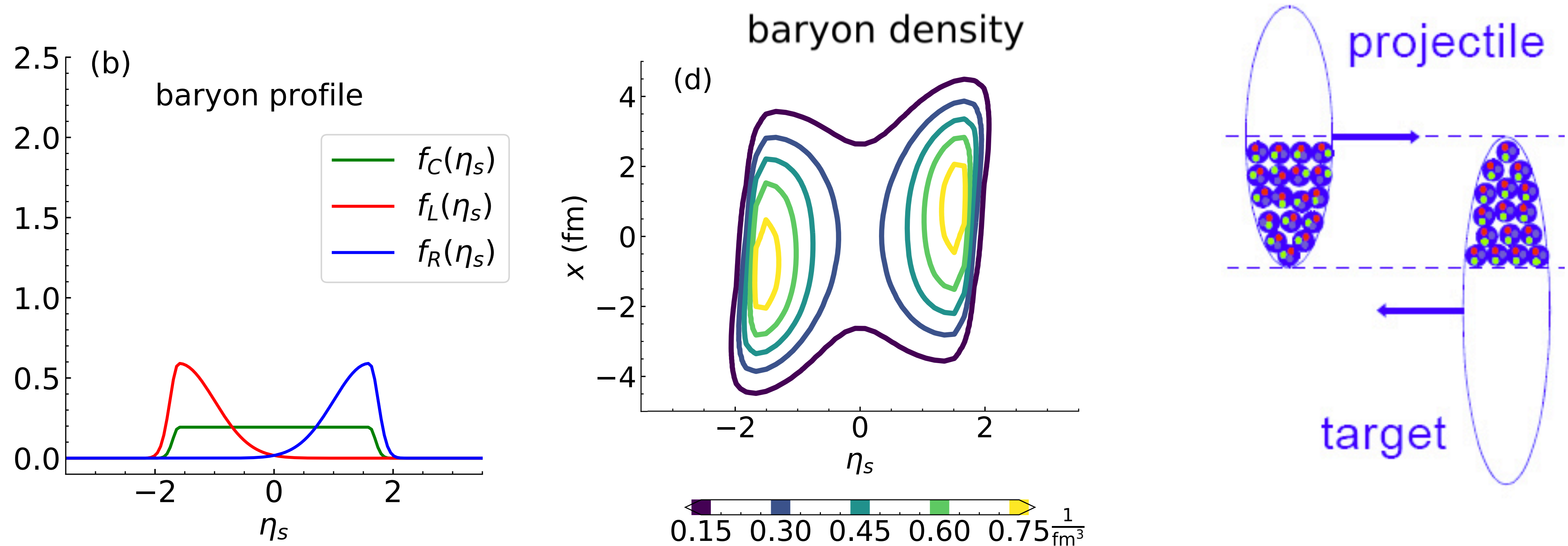
DIRECTED FLOW $v_1(y)$ OF PROTONS

9

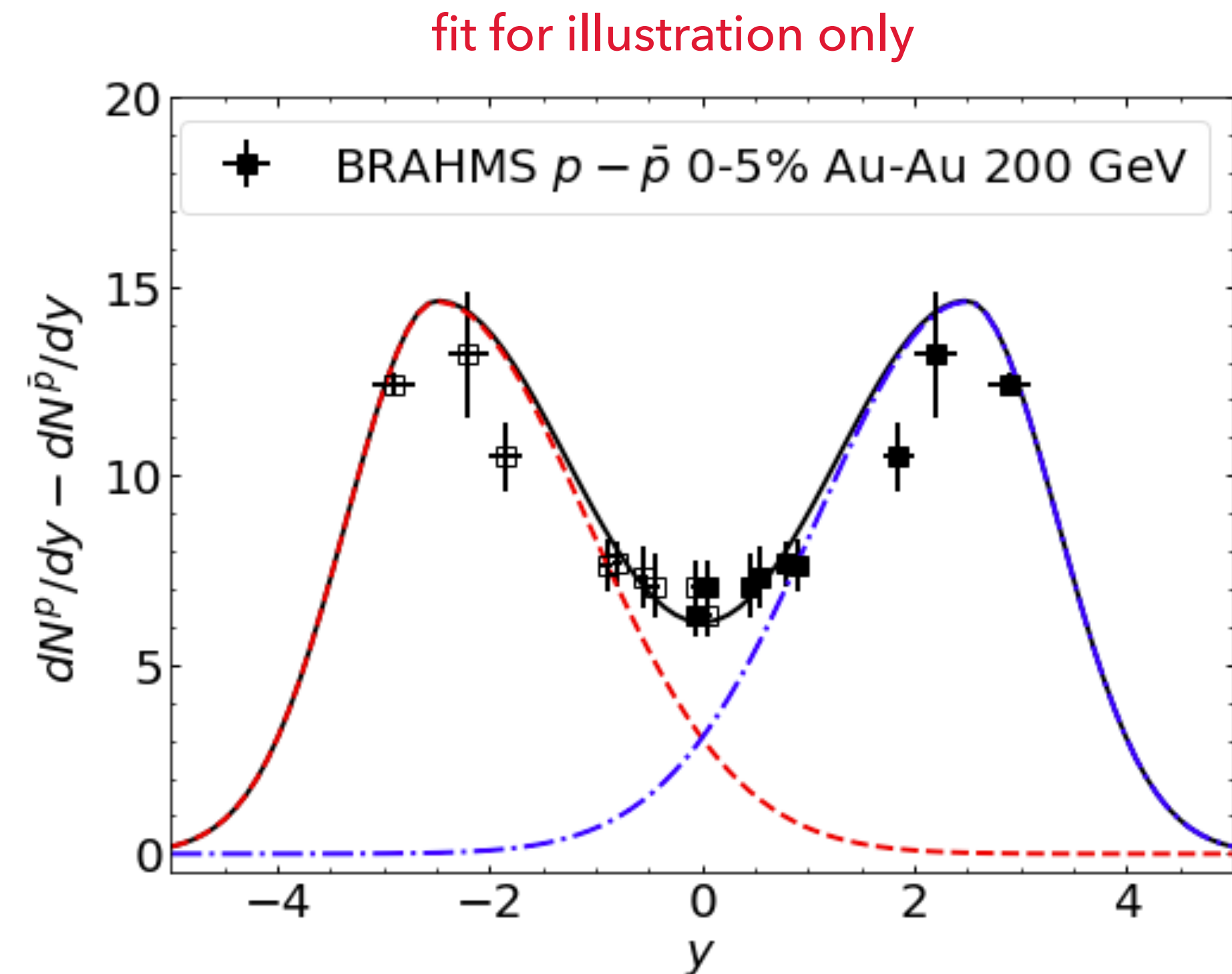
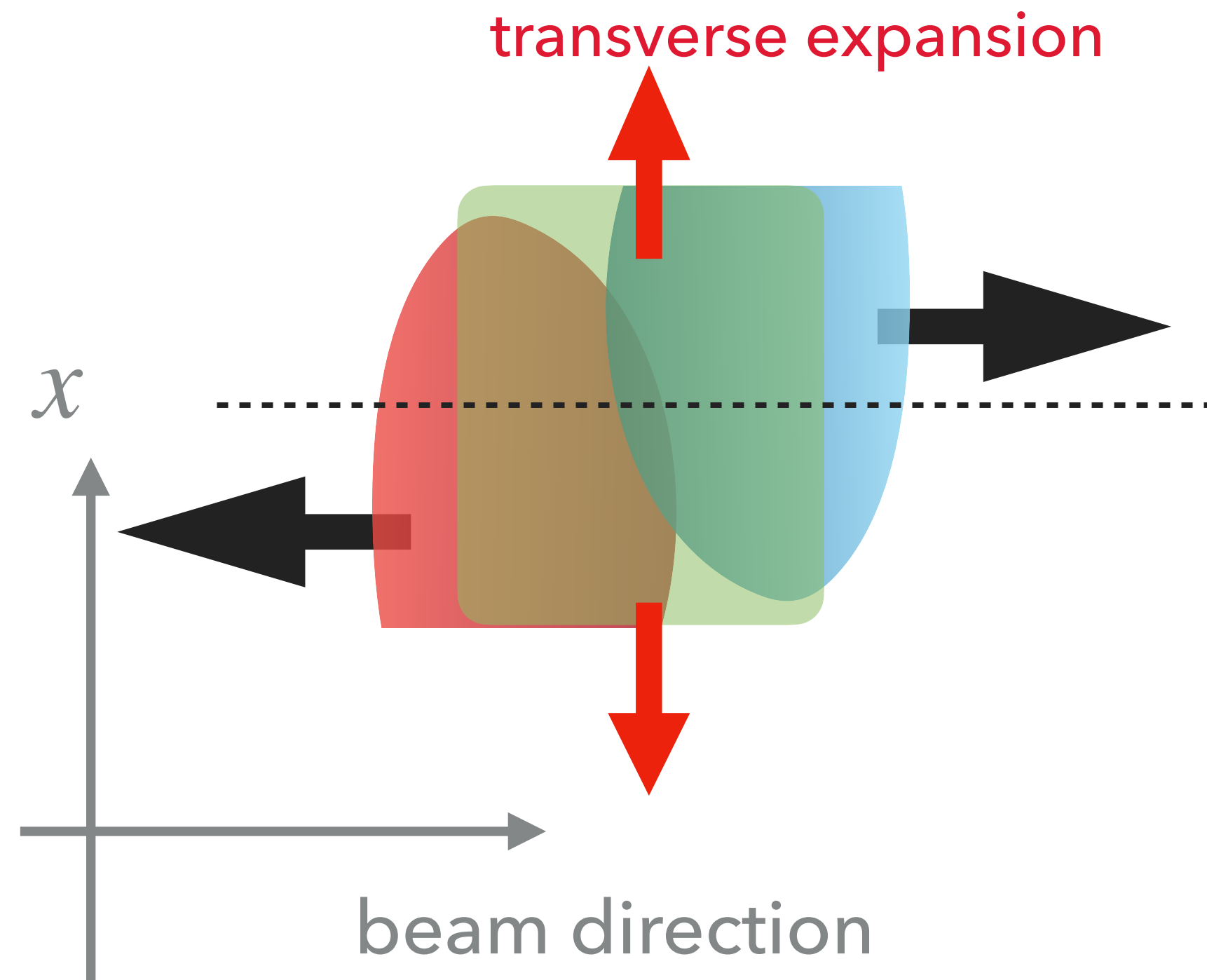


Shen and Alzhani, PRC102, 014909 (2020)

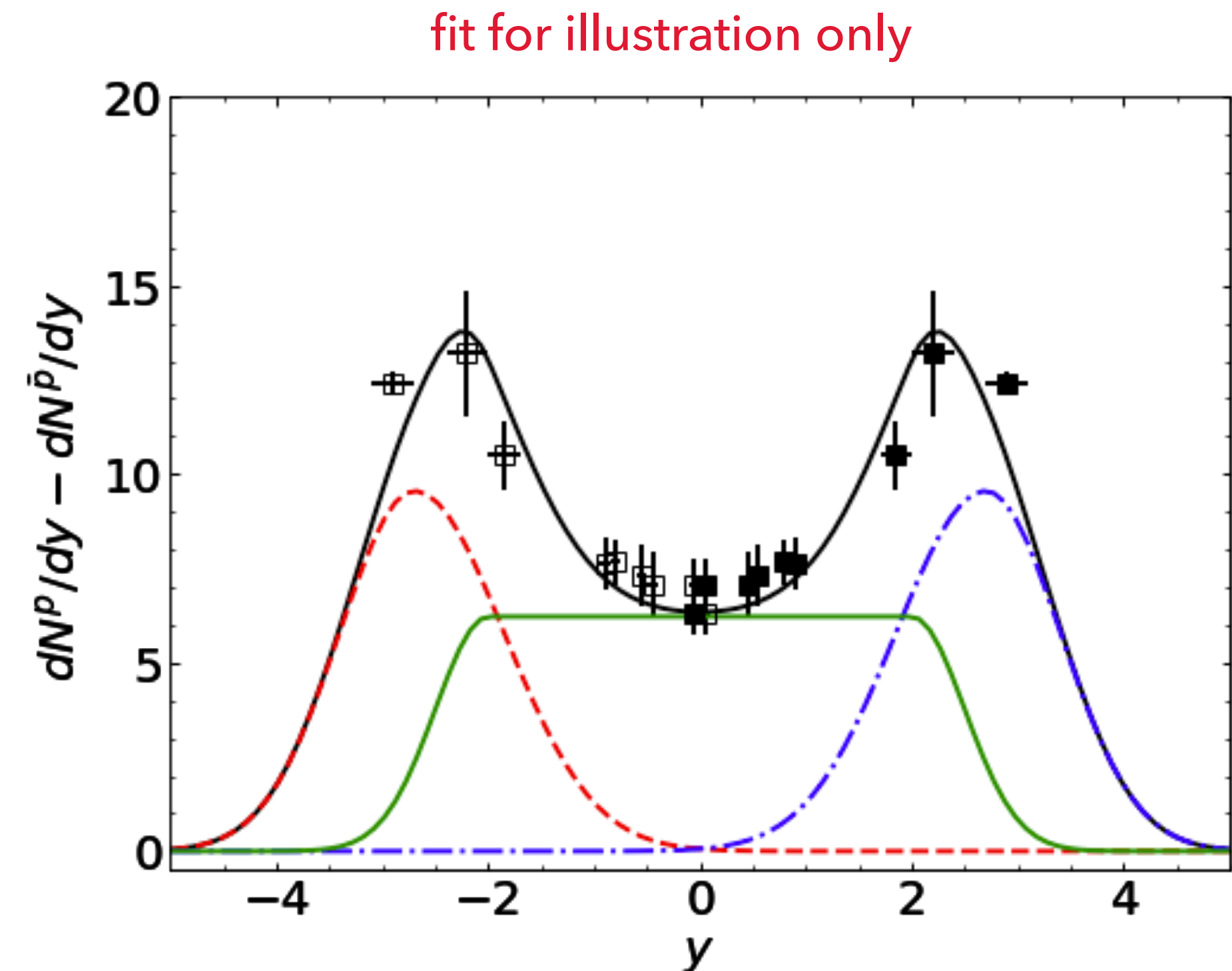
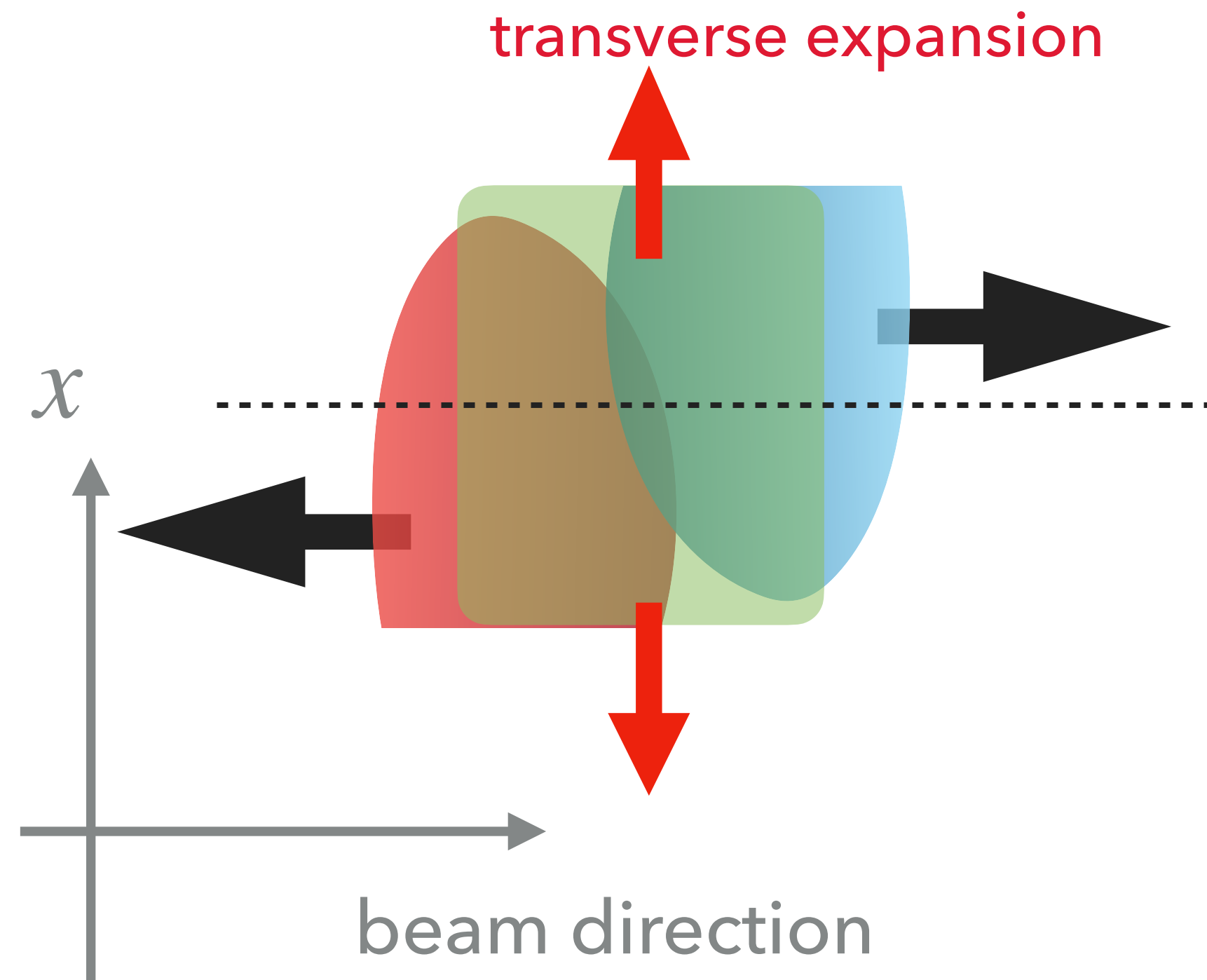
- ▶ $v_1(y)$ of baryons is mainly driven by the asymmetric distribution of baryon density with respect to beam axis + transverse expansion;
- ▶ The widely used baryon-stopping picture results in $v_1(y)$ strongly overshooting the experimental measurements for protons at all beam energies.



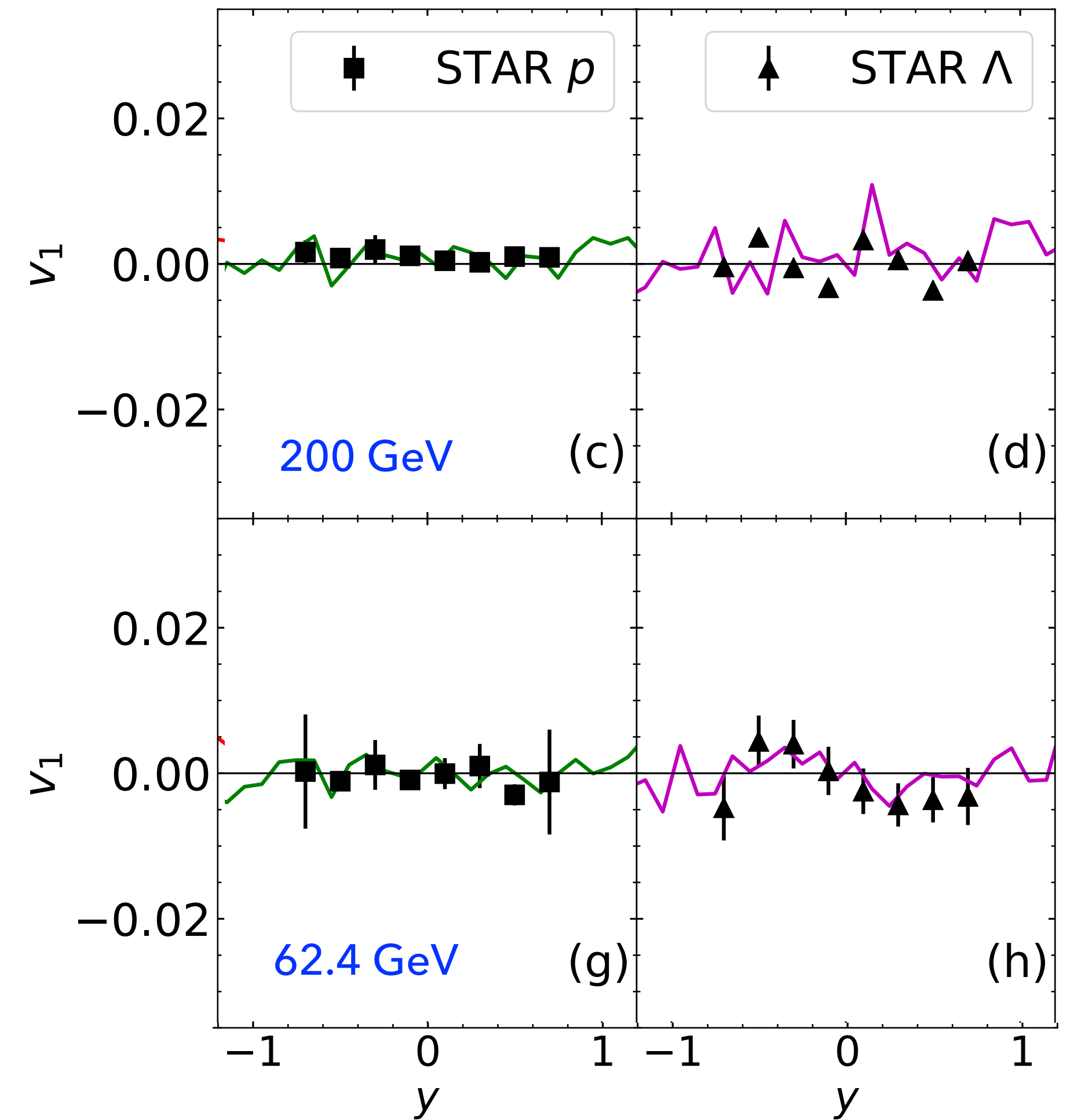
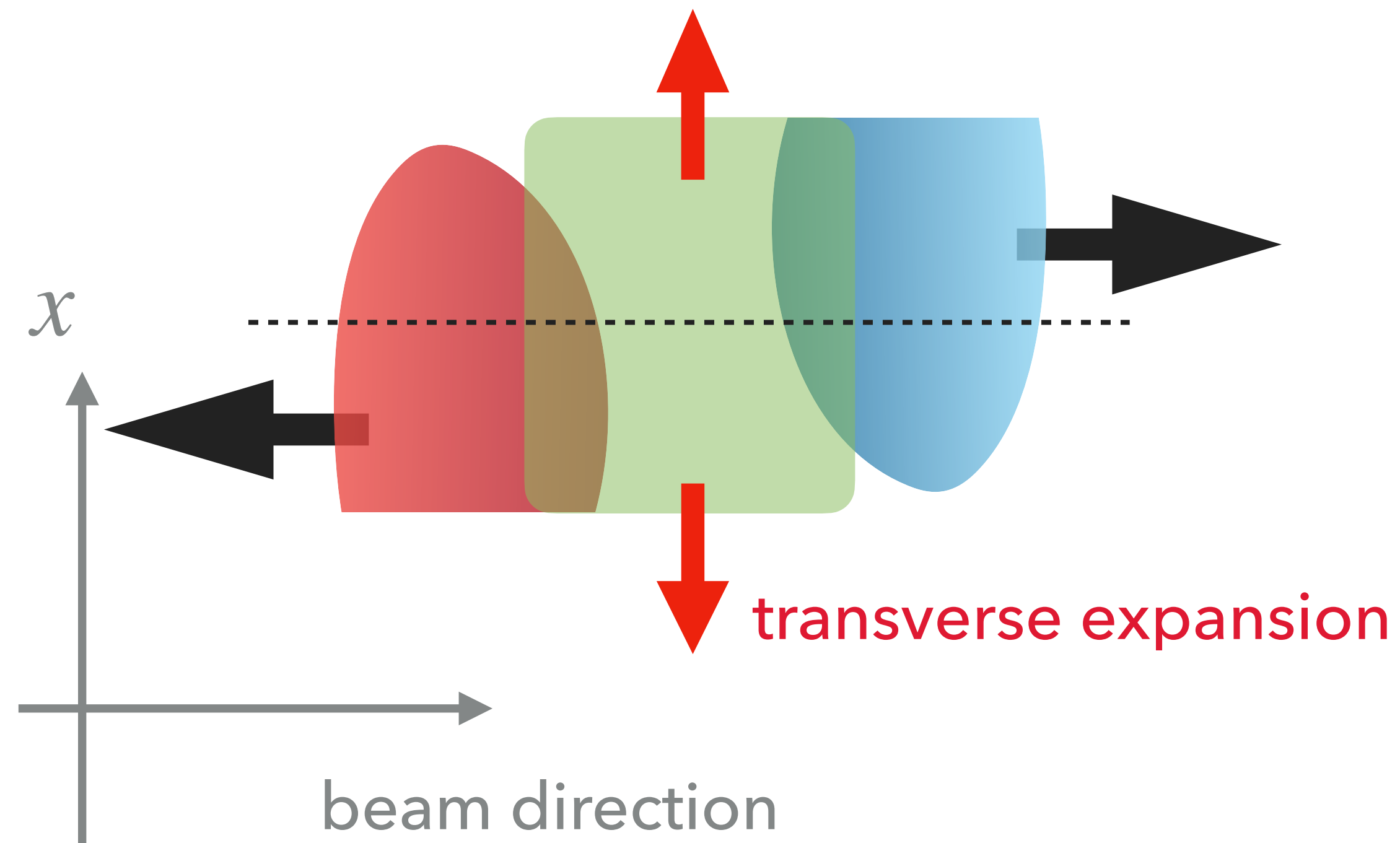
- ▶ A rapidity-independent “plateau” component in initial baryon profile & tilted baryon peaks describing the varying baryon stopping in the transverse plane



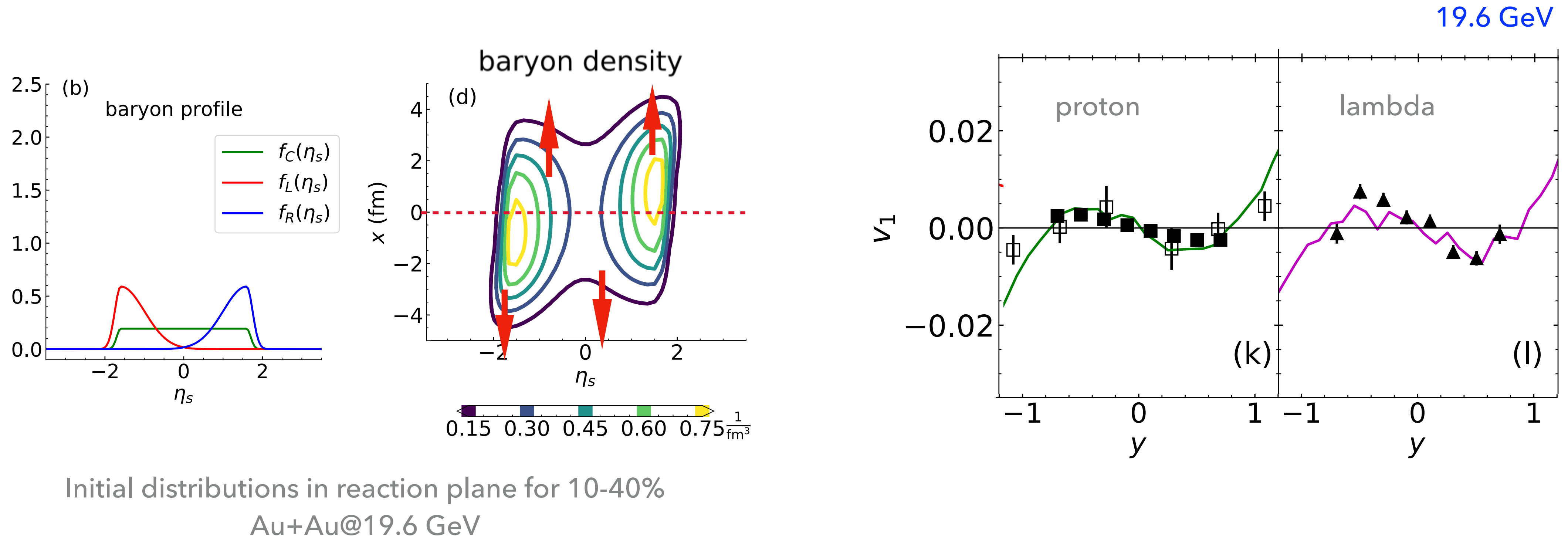
- ▶ To explain the rapidity distributions of net proton yields and baryons' directed flows simultaneously, the plateau is favored;
- ▶ It helps to reduce baryons' $v_1(y)$ while giving enough net proton yields around midrapidity.



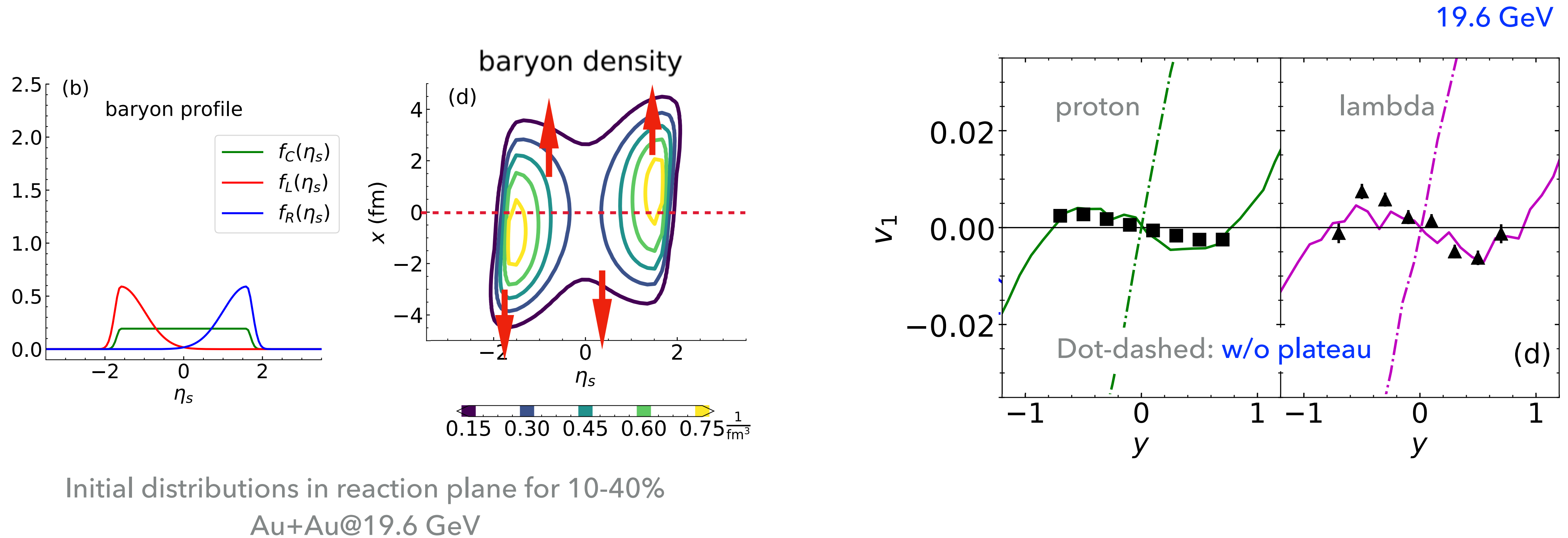
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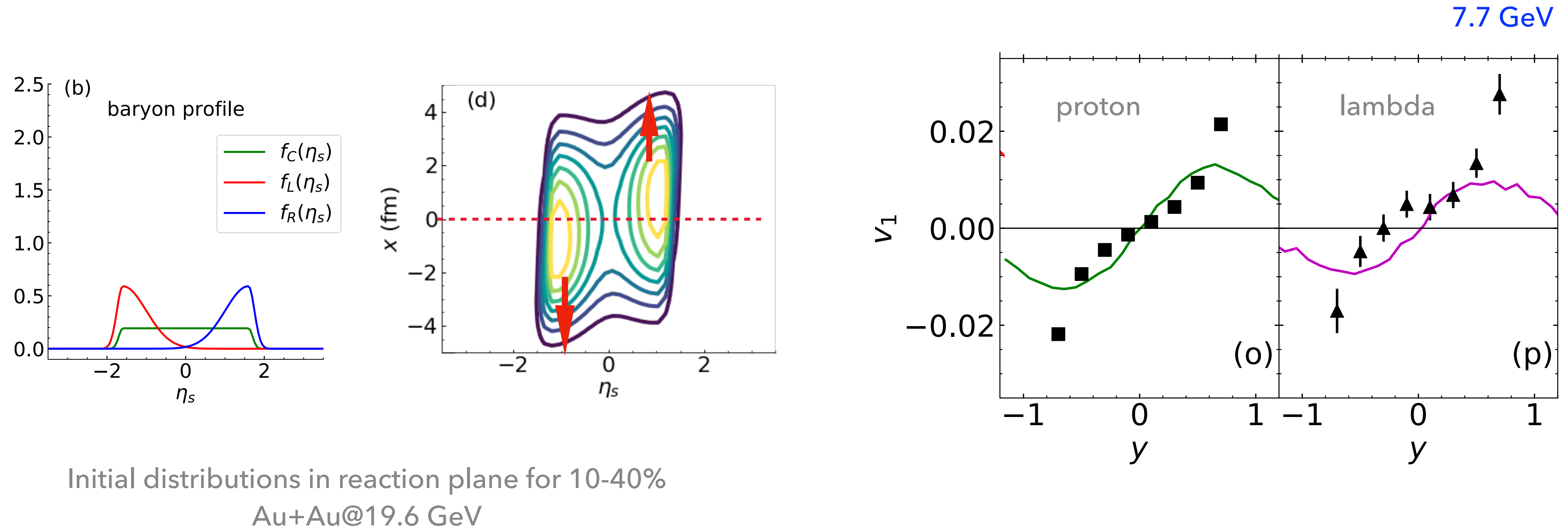
- At high beam energies with large beam rapidities, the plateau dominates and strongly reduces baryons' $v_1(y)$.



- ▶ Transverse expansion + asymmetric distribution of baryon density along $x \Rightarrow$ double sign change in the slope of $v_1(y)$ for baryons at 19.6 GeV, and positive slope at 7.7 GeV
- ▶ The sign change of $dv_1(y)/dy|_{y=0}$ is naturally reproduced without a 1st-order phase transition.



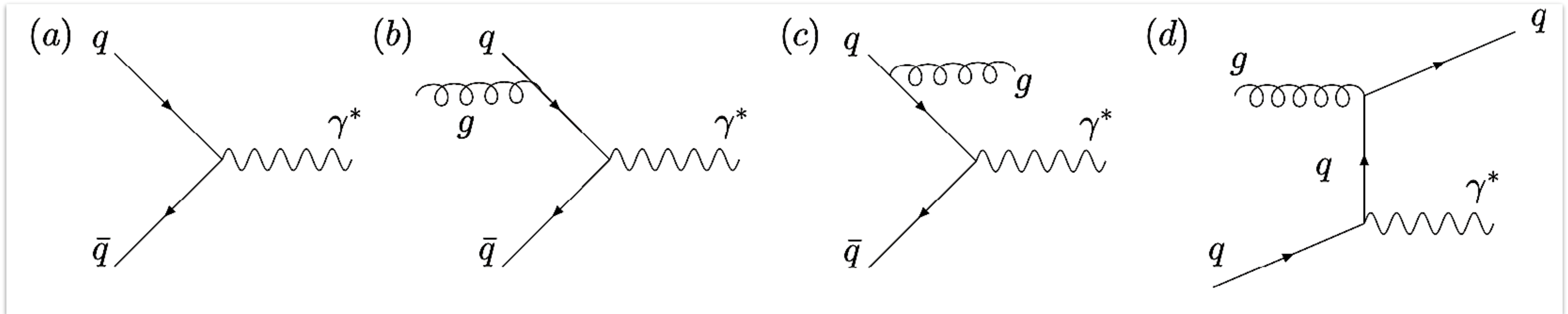
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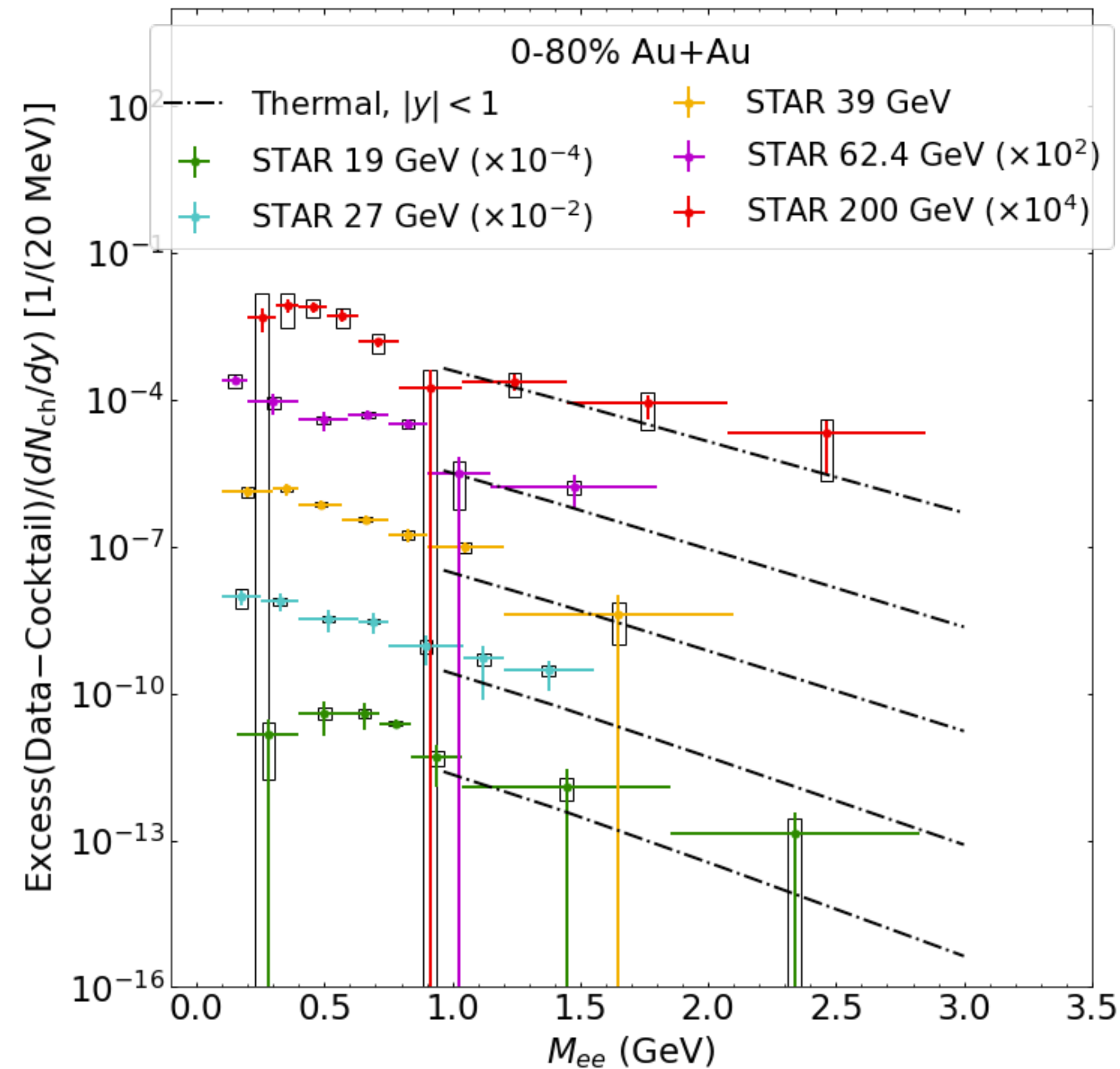
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J. Churchill, LD, C. Gale, G. Jackson & S. Jeon, arXiv: 2311.06675, arXiv: 2311.06951

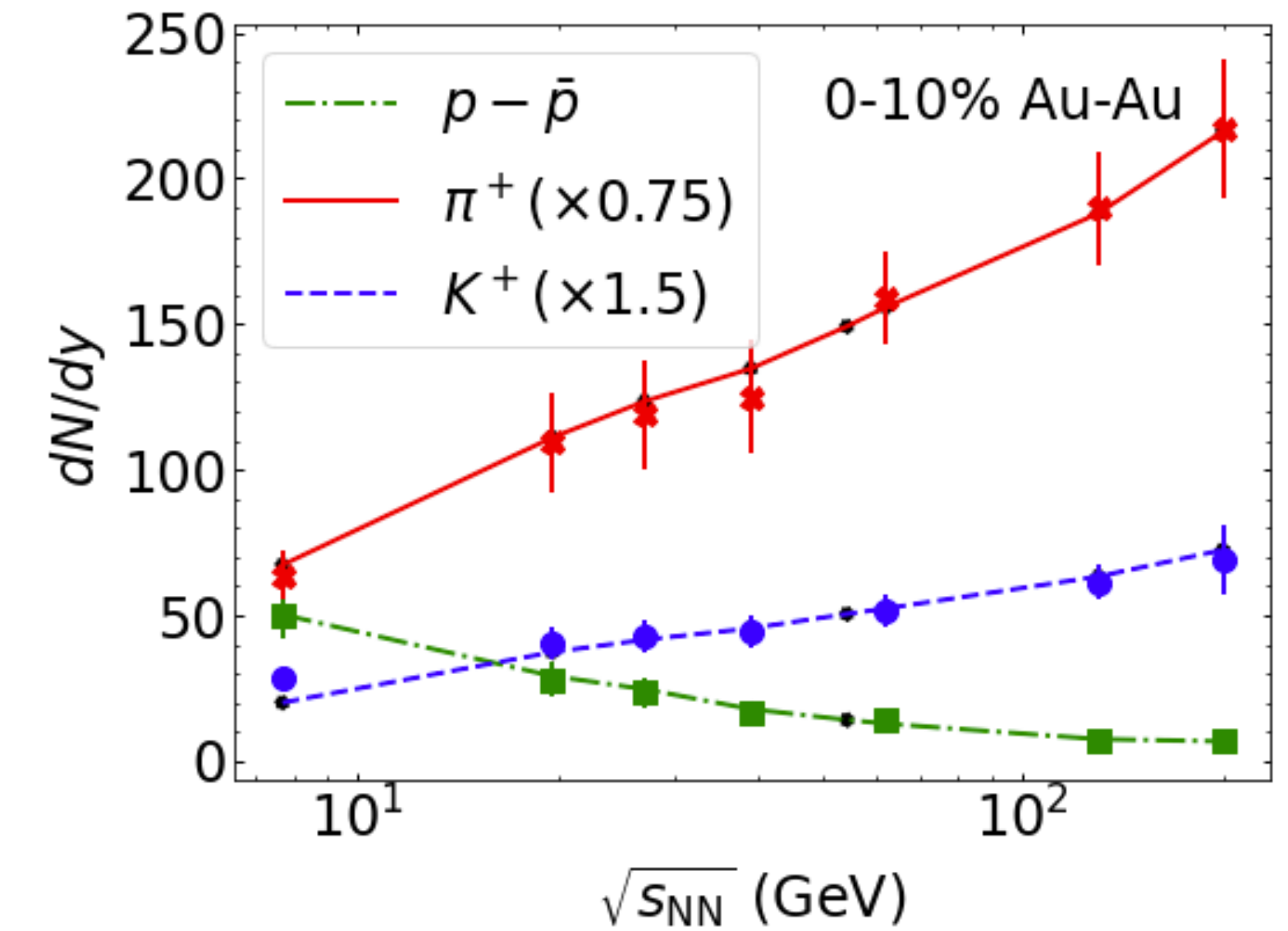
THERMAL DILEPTON AT $\mu_B \neq 0$



- ▶ Baryons: both quarks and string junctions could be the possible source
- ▶ Dileptons and photons production: quarks and anti-quarks are needed
- ▶ Electromagnetic probes of quarks and anti-quarks
 - ▶ μ_B -dependent emission rates assume $\mu_B = 3\mu_q$ (quarks carry baryon number)



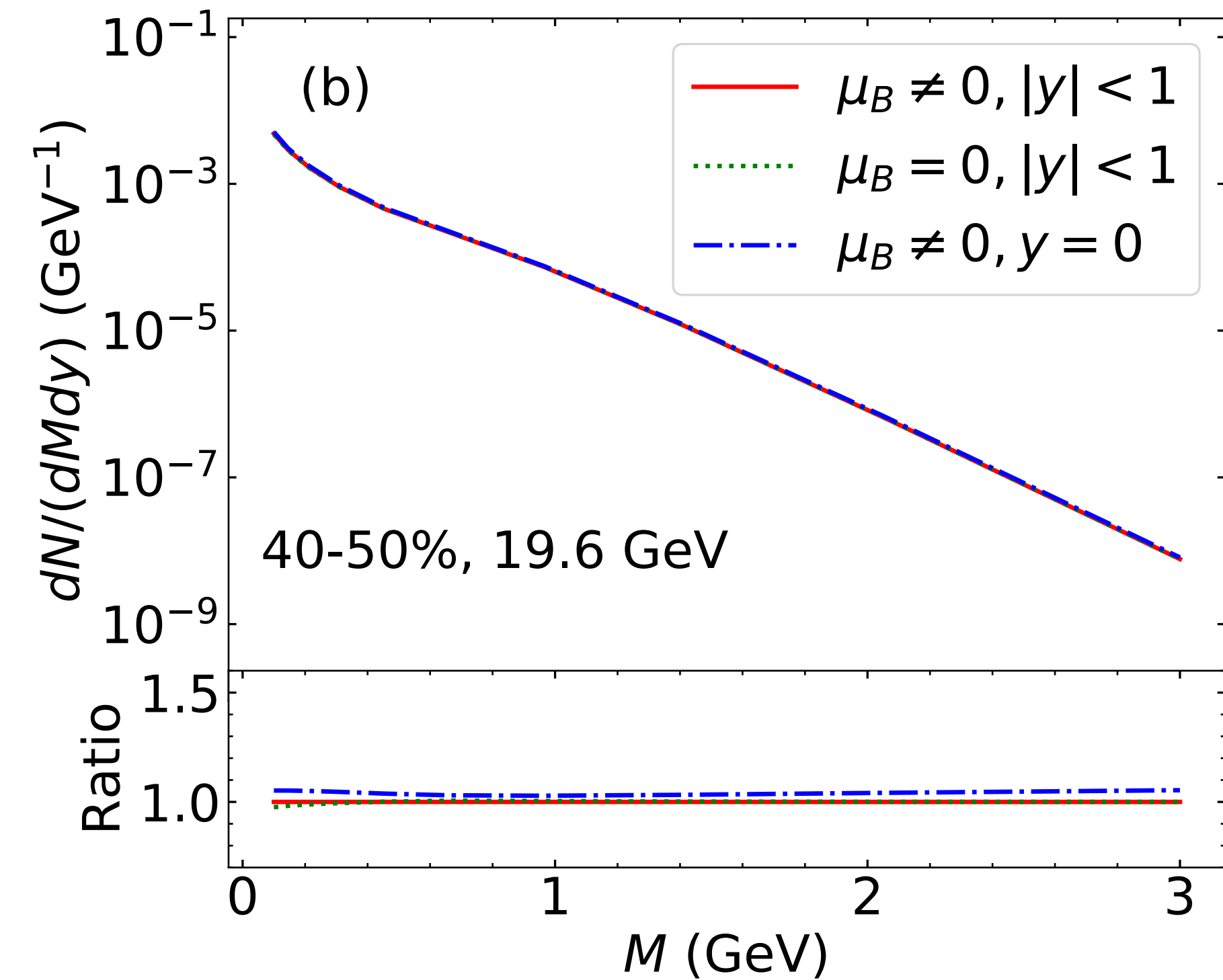
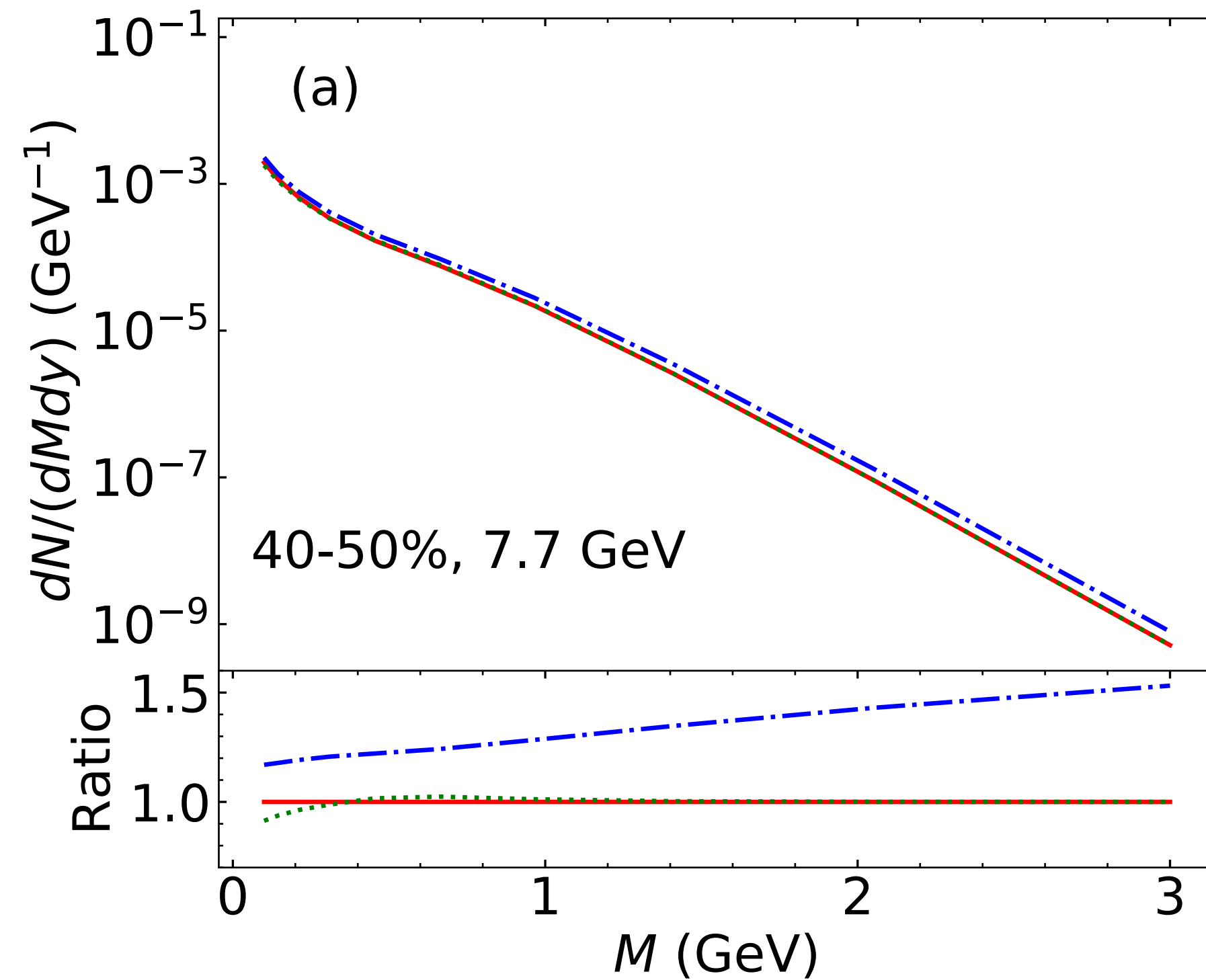
Data: STAR, PRL 113 (2014) 022301, PLB 750 (2015) 64-71, PRC 107 (2023) L061901



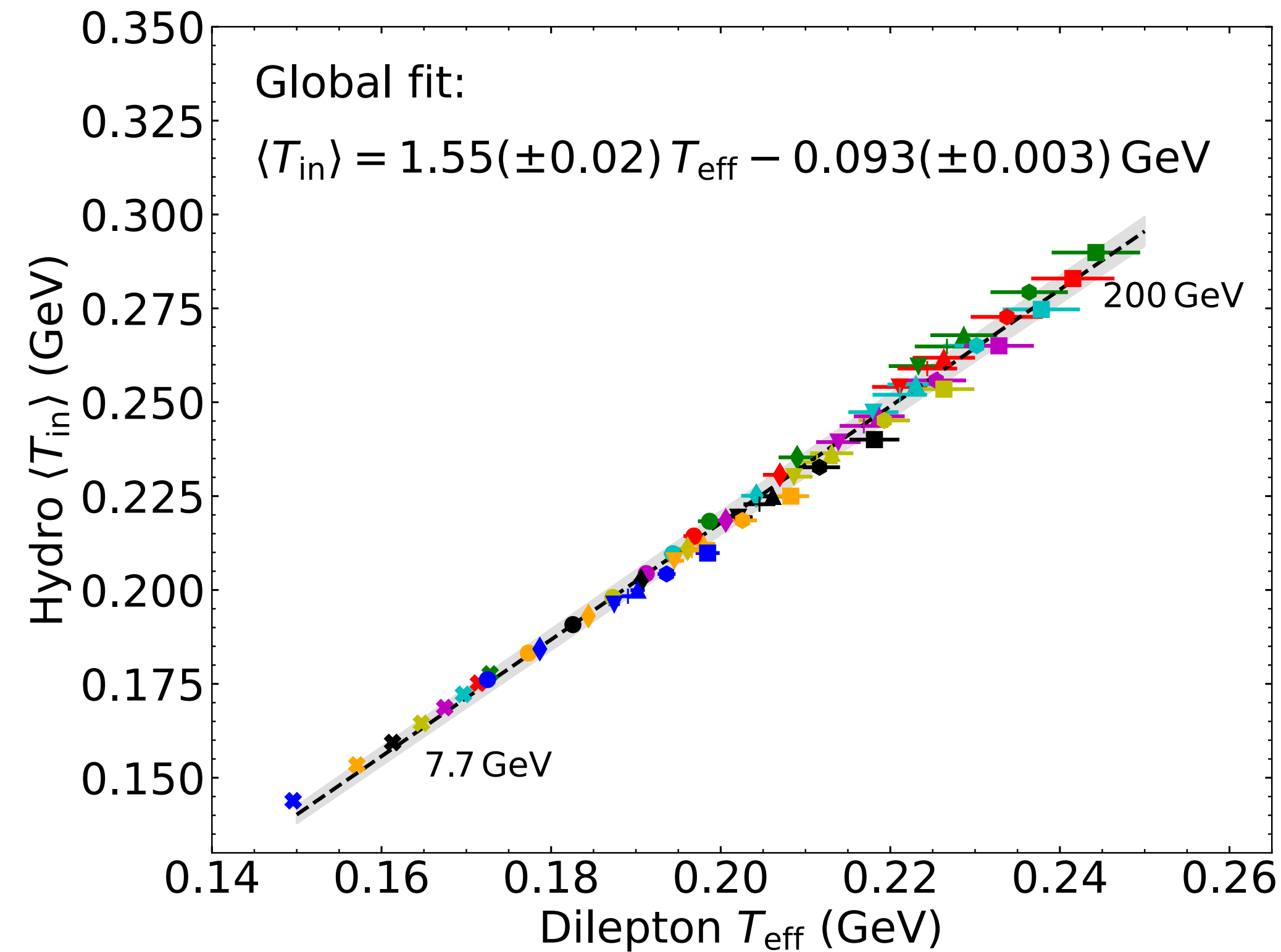
<https://github.com/LipeiDu/DileptonEmission>

J. Churchill, LD, C. Gale, G. Jackson & S. Jeon, arXiv: 2311.06675, 2311.06951

- First estimate of NLO dilepton emission at $\mu_B \neq 0$ with hydrodynamics



- ▶ μ_B -dependence is not significant for 7.7 GeV and higher beam energies
- ▶ Sensitivity of real photons to μ_B (excess of q to $\bar{q}$) might be stronger



- ▶ Extract the initial QGP temperature from the effective temperature derived from the dilepton spectrum
- ▶ Not about μ_B extraction, but very useful!

SUMMARY

I. Rapidity distributions

- ▶ Nucleon deceleration and string junction breaking correspond to different energy losses.
- ▶ Probing the **total energy deposited** in the fireball helps differentiate these two mechanisms.

II. Baryonic $v_1(y)$

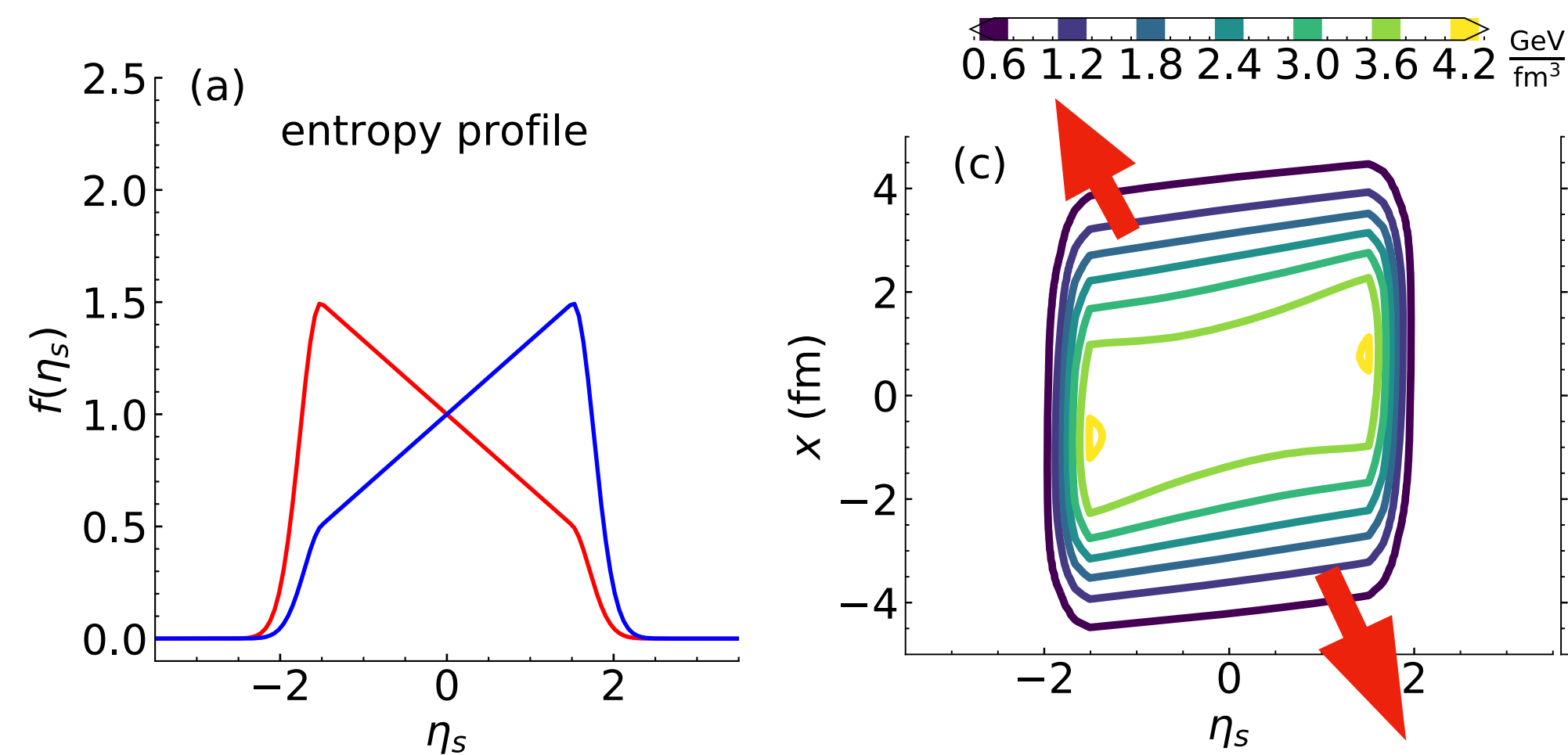
- ▶ We introduced a **central plateau component** in the initial baryon distribution, essential for explaining characteristic features of $v_1(y)$ at various beam energies.
- ▶ The plateau component may originate from **string junctions**.

III. Thermal dilepton

- ▶ **Electromagnetic probes** help to measure quarks and anti-quarks.

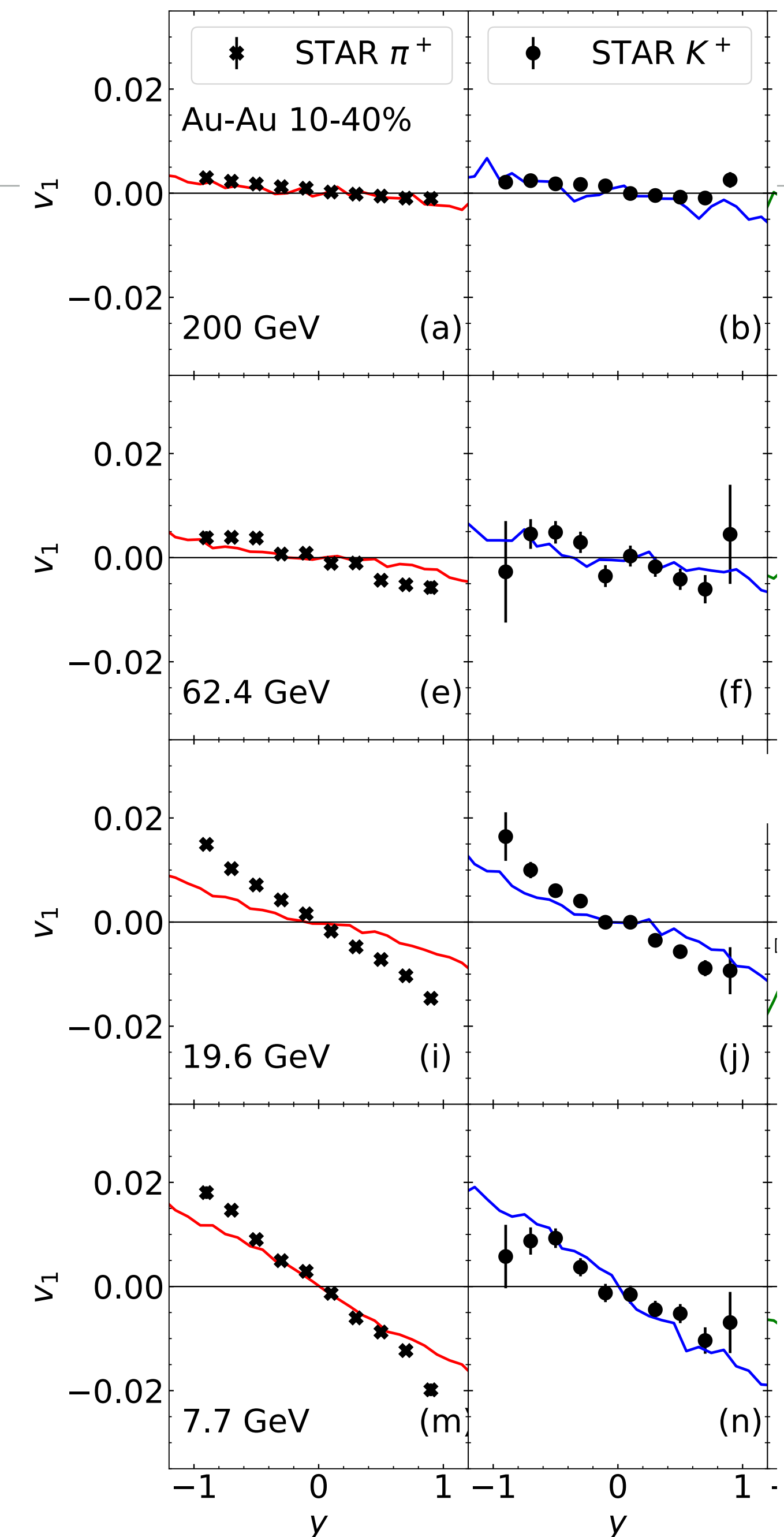
THANK YOU FOR YOUR ATTENTION!

DIRECTED FLOW OF MESONS



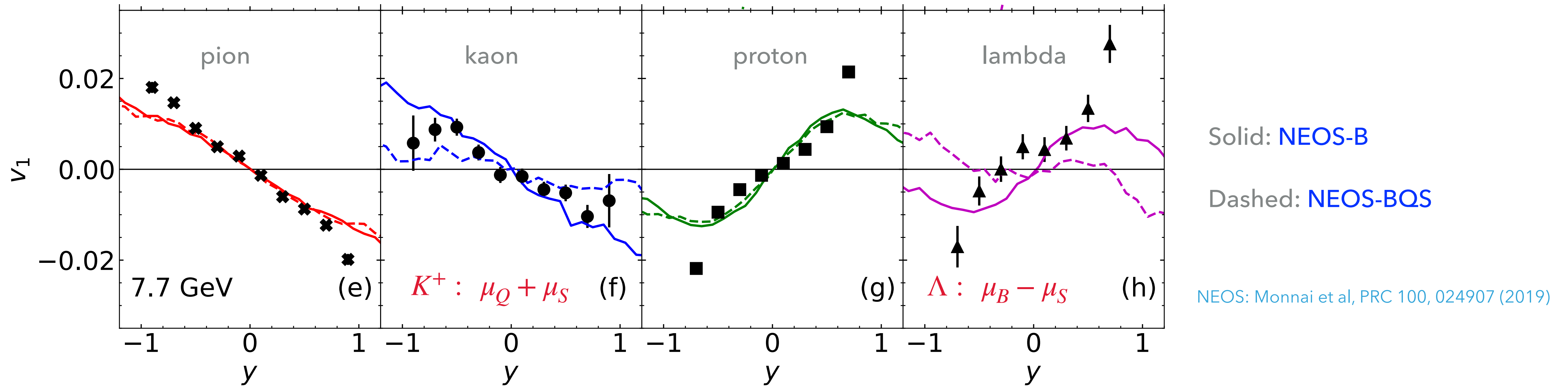
- ▶ $v_1(y)$ of mesons with **negative slope** at all beam energies; slope increases when beam energy decreases
- ▶ Slight shift of energy density along x (tilted structure) generates **sideward pressure gradient**

see also: Bożek & Wykiel Phys. Rev. C 81, 054902 (2010)



Data: STAR, PRL 112, 162301 (2014); PRL 120, 062301 (2018)

PROBING EOS AT FINITE CHEMICAL POTENTIALS



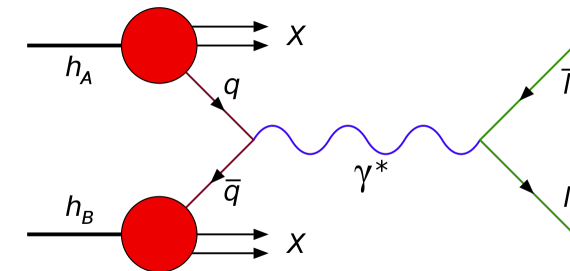
- ▶ Two limits of EoS: NEOS-B, $\mu_S = \mu_Q = 0$ and NEOS-BQS, $n_S = 0, n_Q = 0.4n_B$ (2D projection of a 4D EOS)
- ▶ Local strangeness neutrality suppresses the $v_1(y)$ of K^+ and Λ around midrapidity, and even alters the sign of Λ 's $v_1(y)$ beyond $|y| \gtrsim 0.6$
- ▶ The $v_1(y)$ of identified particles can be used to probe EoS at finite chemical potentials.

DILEPTON SPECTRA

STAR, PLB 750 (2015) 64-71

High-mass region (HMR)

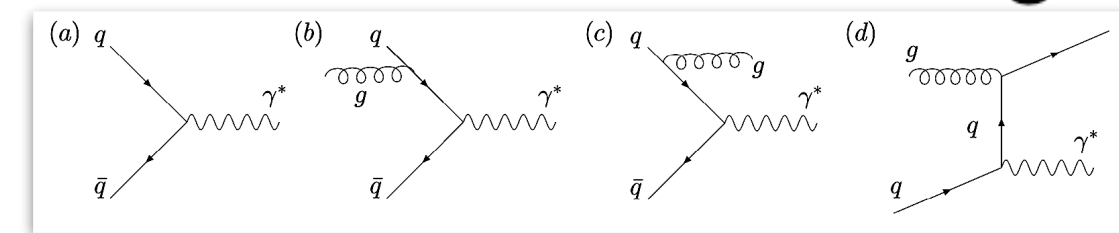
Drell-Yan process



decays of J/ψ

Intermediate-mass region (IMR)

thermal emission from QGP



semileptonic decays from open heavy flavor-antiflavor pair, e.g. $D/\bar{D}$

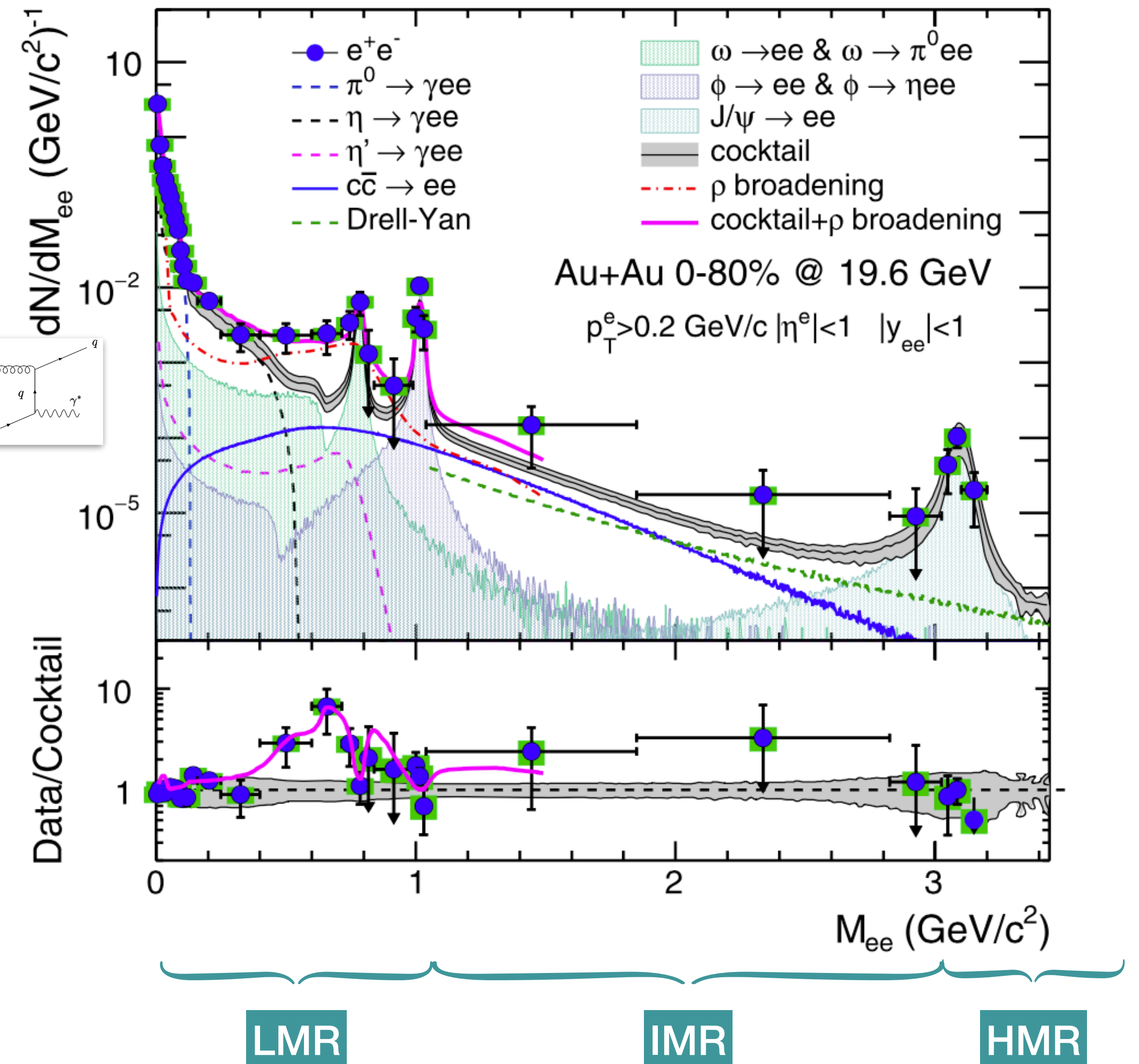
Low-mass region (LMR)

thermal emission from hadronic matter

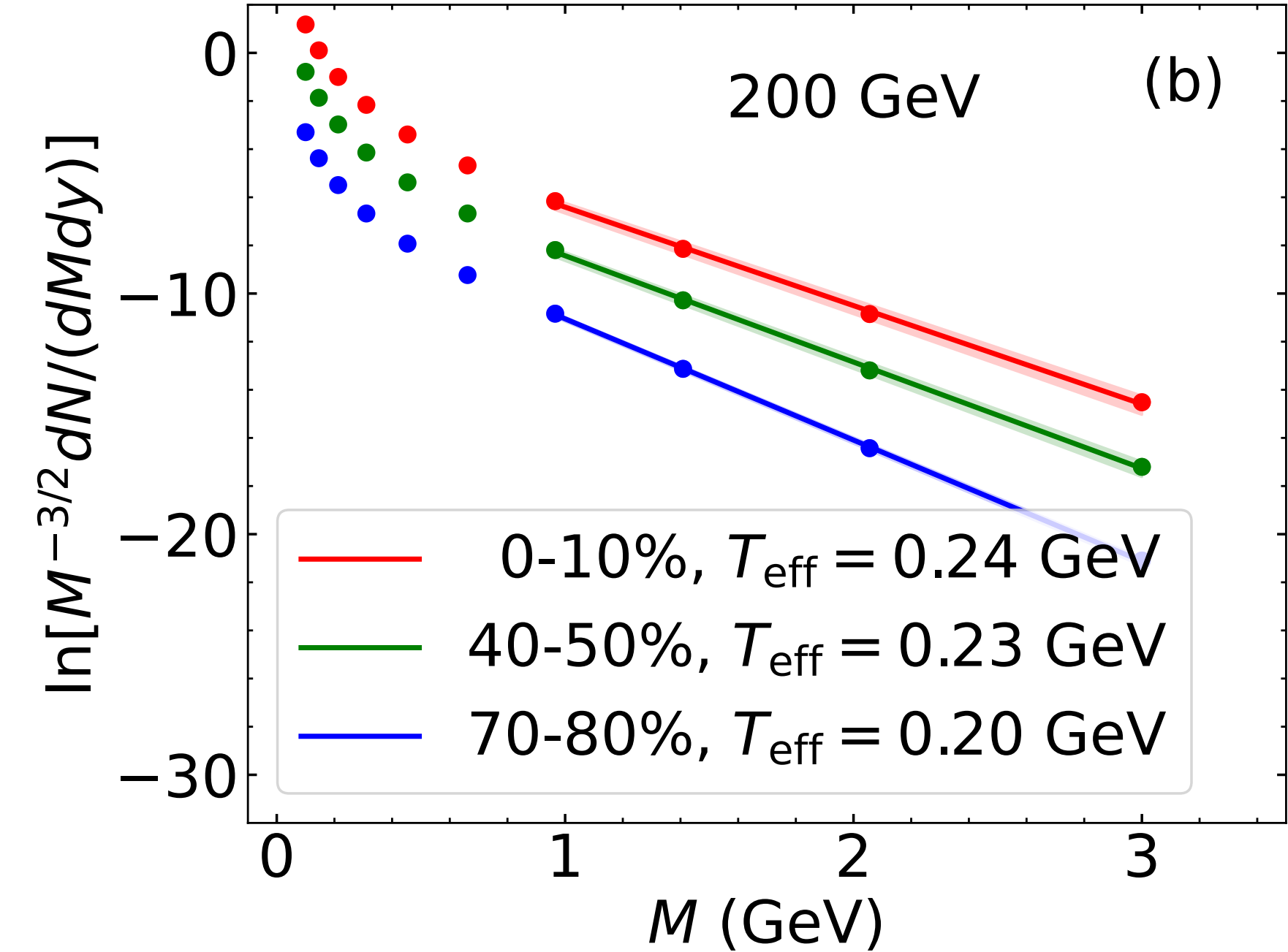
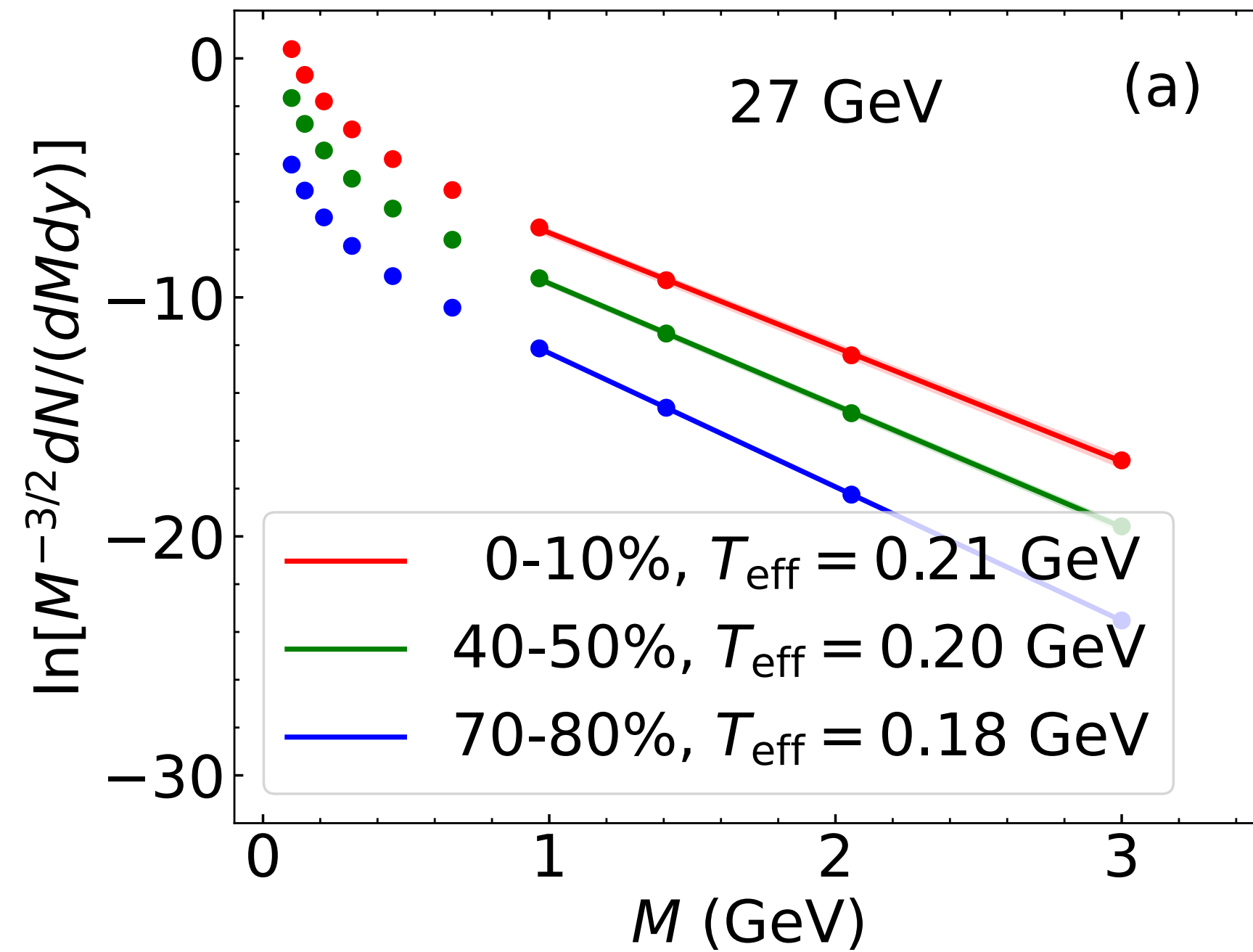
direct decays of $\rho/\omega/\phi$ (ρ is short-lived)

Dalitz (three body) decays of $\pi^0/\eta/\eta'$

Dilepton cocktail: late decays of hadrons

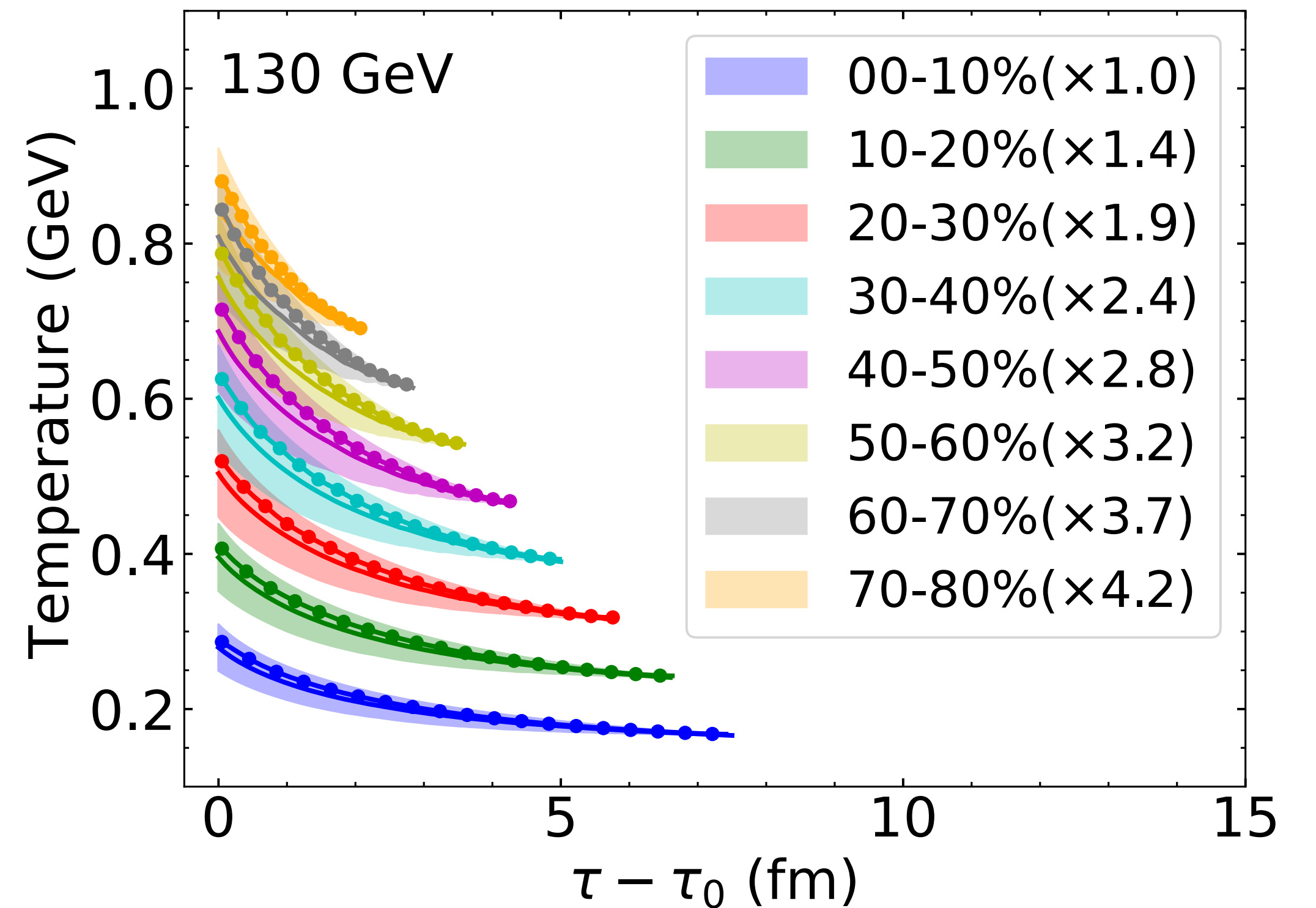
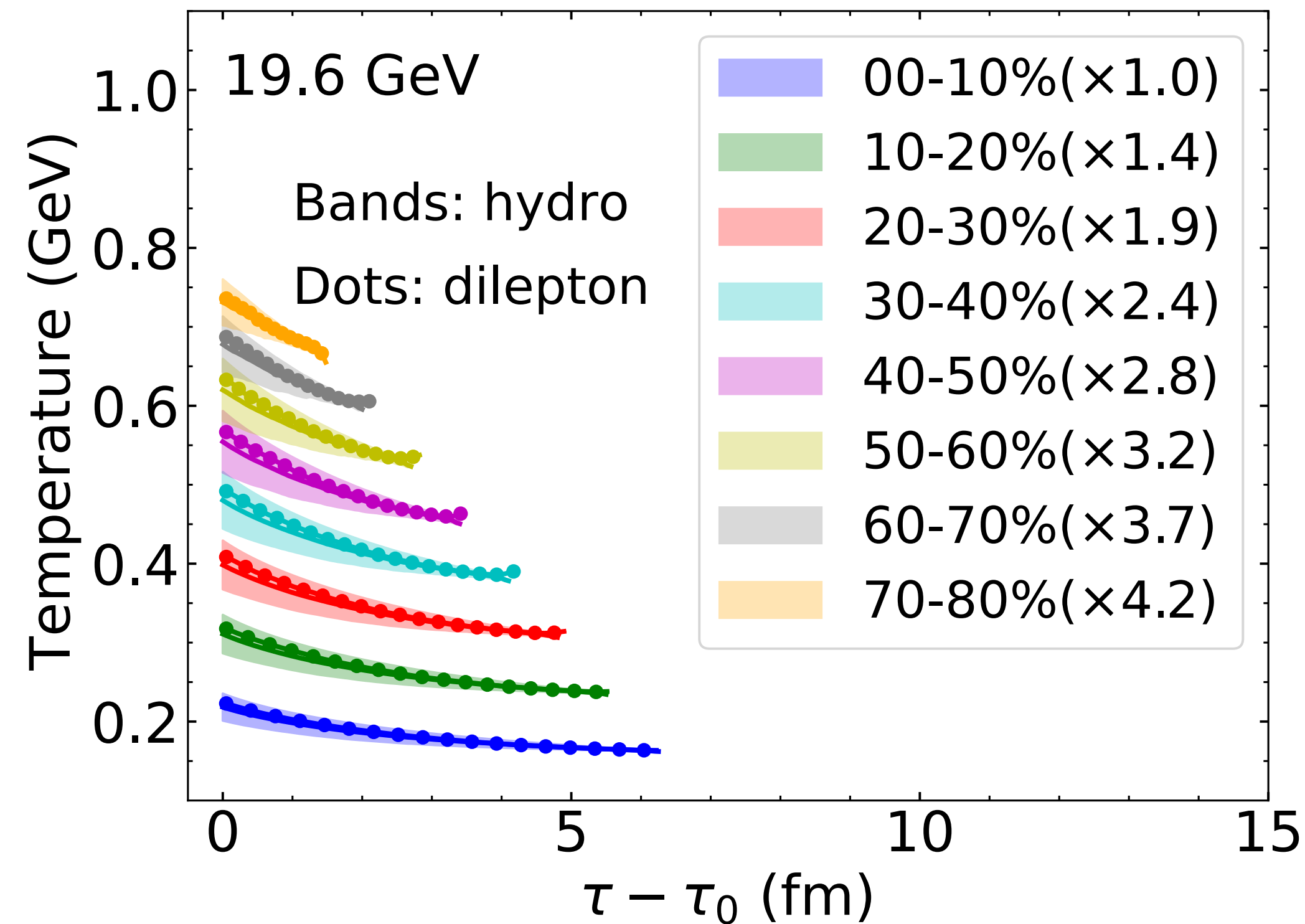


TEMPERATURE EXTRACTION



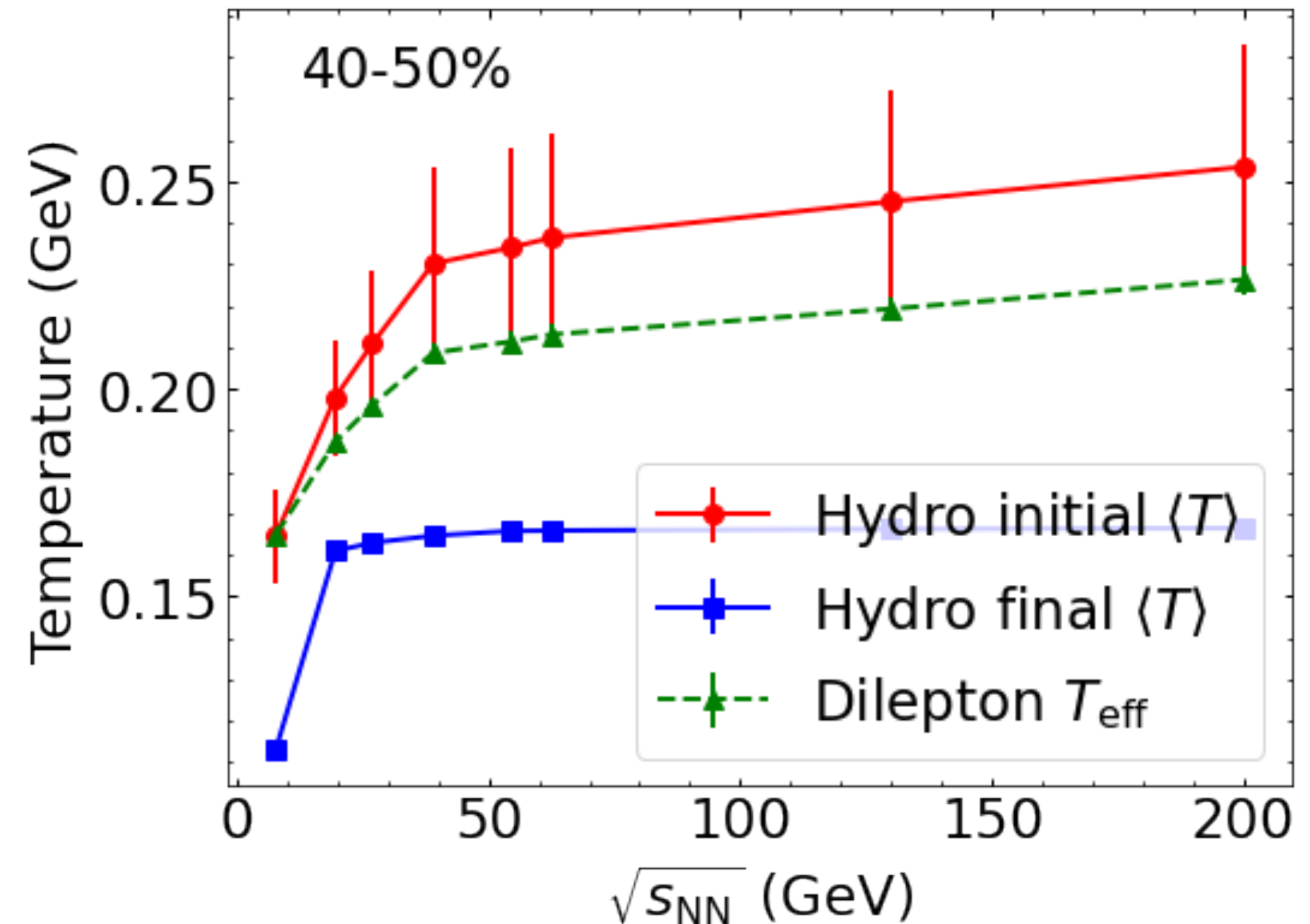
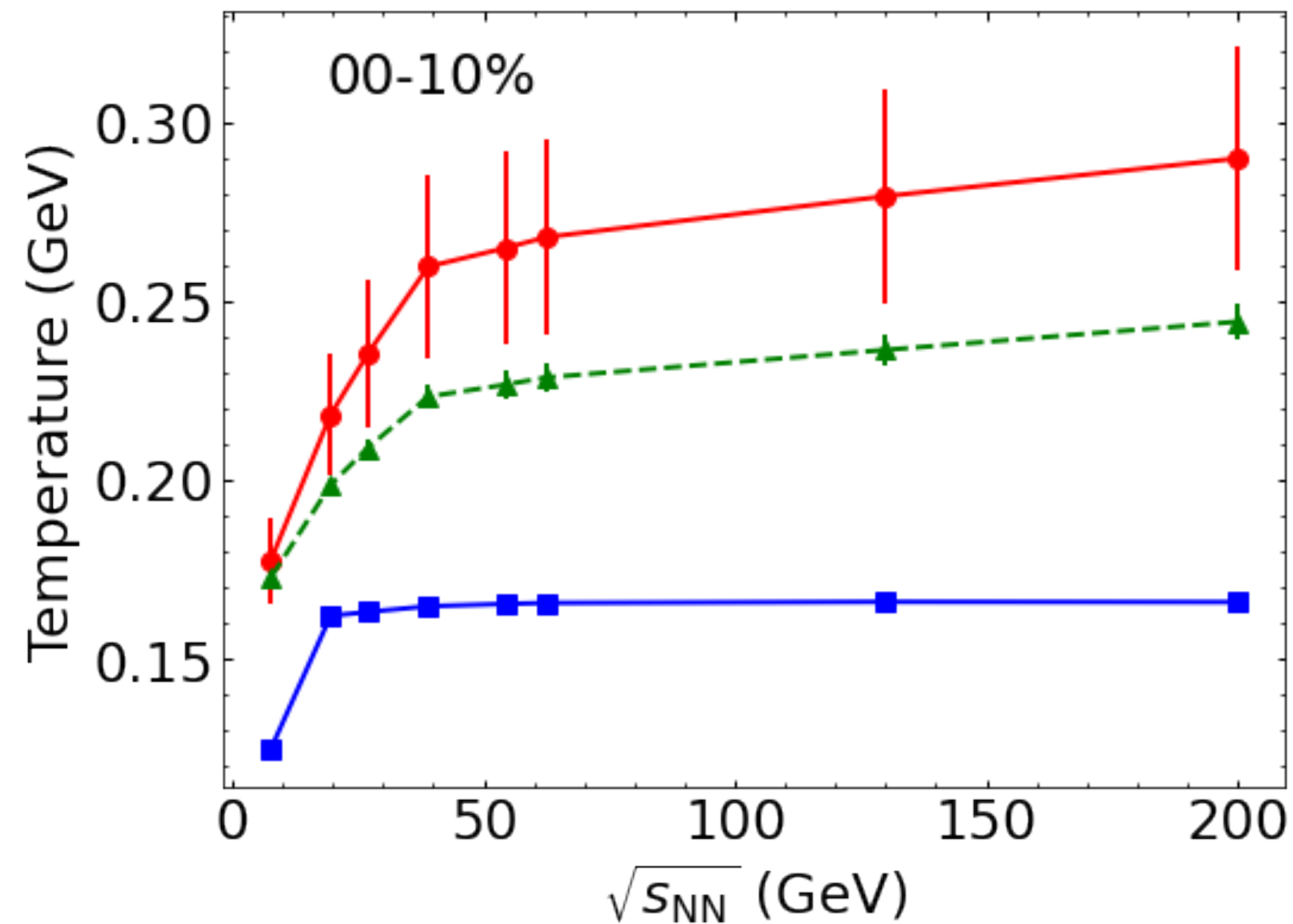
- ▶ In non-relativistic approximation, the emission rate $\frac{dR}{dM} \propto (MT)^{3/2} e^{-M/T}$
 - ▶ Use thermal dileptons within $1\text{ GeV} \lesssim M \lesssim 3\text{ GeV}$ for temperature extraction
- $$\ln \left(\frac{dN}{dM} M^{-3/2} \right) \propto \ln(T^{3/2} e^{-M/T}) = -\frac{M}{T} + \dots$$

TEMPERATURE EXTRACTION IN τ



- ▶ To obtain hydrodynamic temperature distribution, the temperature of each fluid cell is weighted by its energy density. curve: mean temperature; band: standard deviation
- ▶ Dots: effective temperature extracted from dilepton spectra

THERMAL DILEPTON AS THERMOMETER



- ▶ The final temperature is approximately a constant for $\sqrt{s_{NN}} \gtrsim 19.6$ GeV and much smaller at $\sqrt{s_{NN}} = 7.7$ GeV because of the baryon chemical potential
- ▶ A correlation between the hydro initial temperature and effective temperature can be observed