

Spin-orbit correlations in QCD

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[2404.04208](#); [2404.04209](#) with Shohini Bhattacharya, Renaud Boussarie,

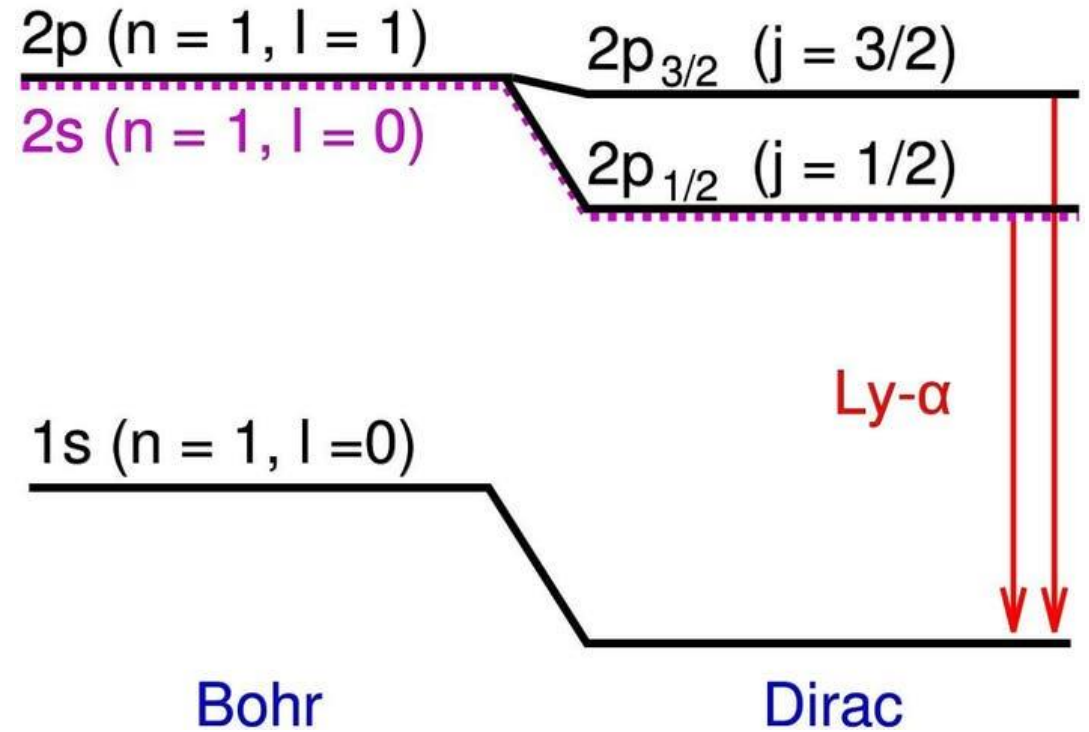
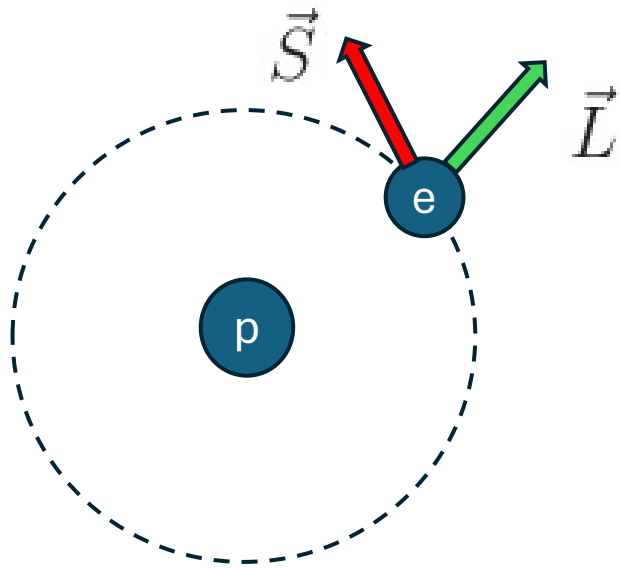
[2404.18872](#) with Jakob Schoenleber

Contents

- Spin-orbit correlation in QCD
- New momentum sum rule
- Small- x , quantum entanglement
- Connection to experimental observable

Spin-orbit coupling in atoms

$$V = -\frac{\mu_B e}{mc^2 r^3} \vec{S} \cdot \vec{L}$$



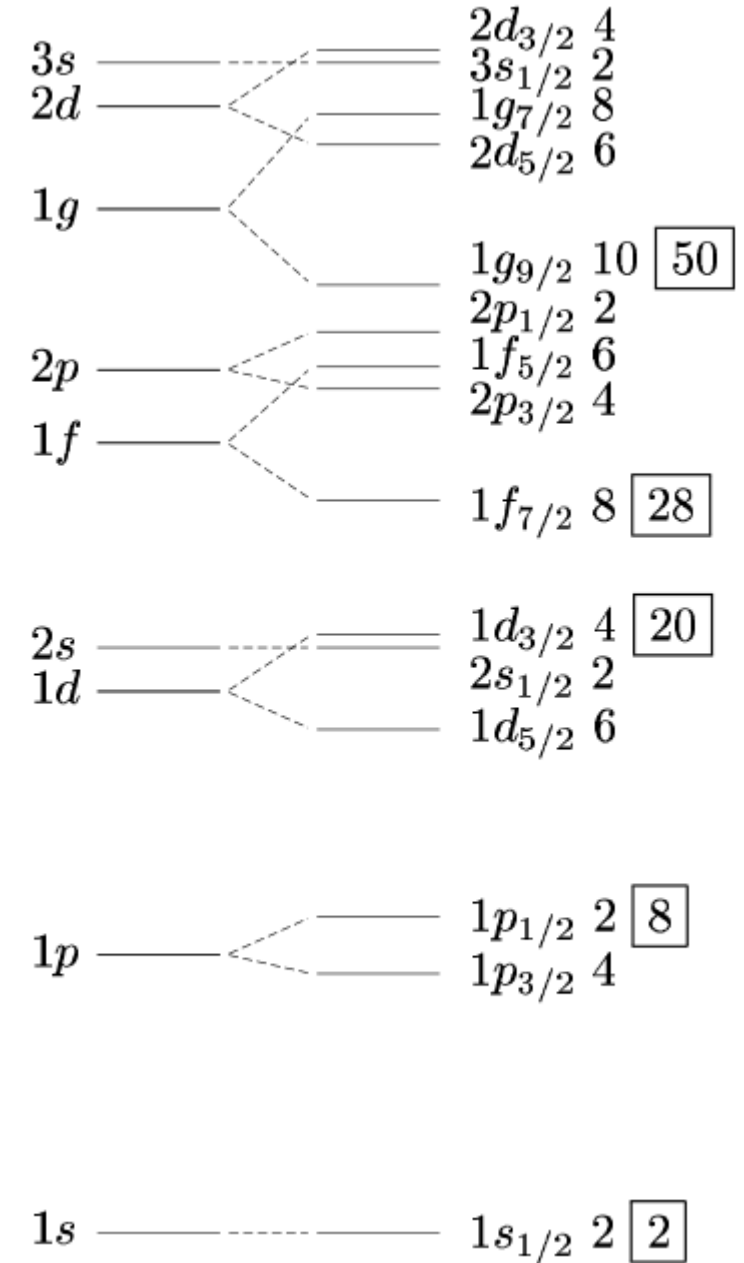
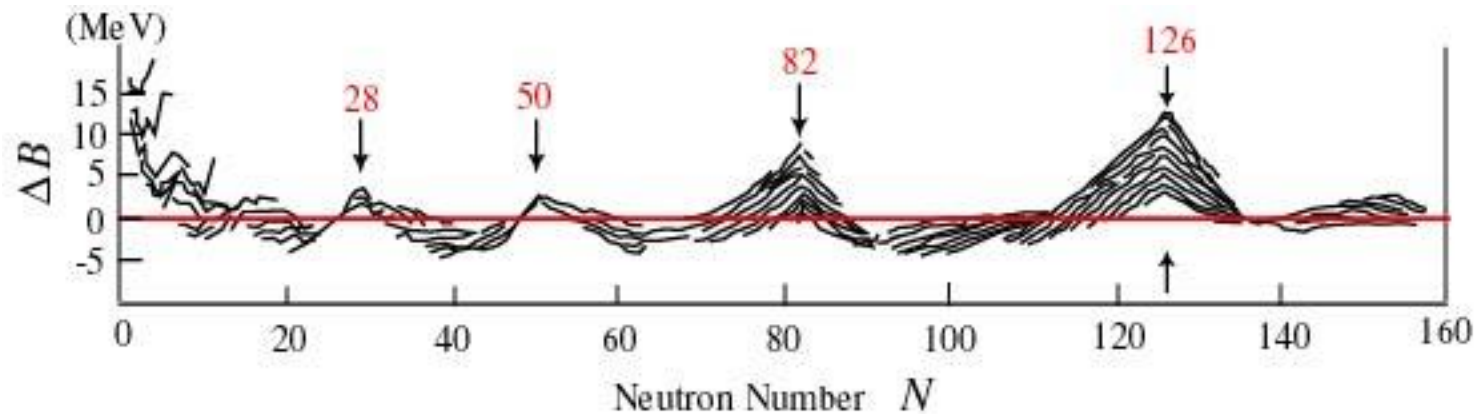
$\vec{\mu} \cdot \vec{B}$ in the electron rest frame + relativistic effects contributes to the **fine structure** of atoms

Spin-orbit coupling in nuclei

In the nuclear shell model, nucleons orbiting inside a nucleus feel a spin-orbit force

Strong spin-orbit coupling → **magic numbers**

Mayer & Jensen Nobel prize (1963)



Spin-orbit coupling in nucleons?

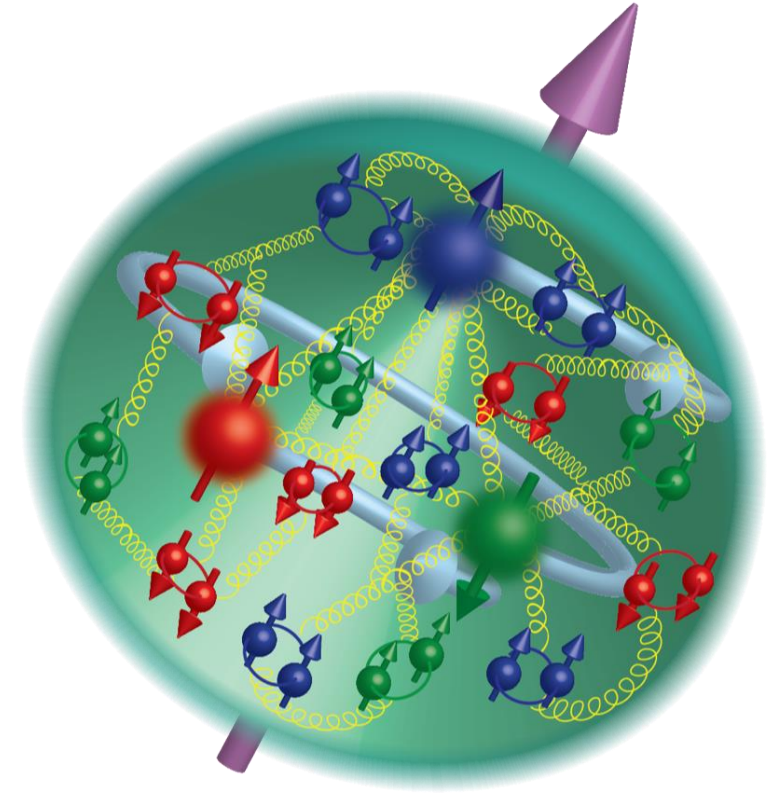
Quarks and gluons carry spin (helicity) and OAM
Naturally there should be spin-orbit coupling

Numbers of quarks and gluons indefinite

Gluon spin and OAM need to be carefully defined
→ Jaffe Manohar scheme

$$\frac{1}{2} = \frac{1}{2} \underset{\text{spin}}{\Delta\Sigma} + \underset{\text{spin}}{\Delta G} + \underset{\text{orbit}}{L_q} + \underset{\text{orbit}}{L_g}$$

Why interesting? Observable consequence?



Quark spin-orbit correlation

Polarized quark GTMD

$$\begin{aligned}\tilde{f}_q(x, \xi, k_\perp, \Delta_\perp) &= \int \frac{d^3 z}{2(2\pi)^3} e^{ixP^+ z^- - ik_\perp \cdot z_\perp} \langle p' s' | \bar{q}(-z/2) W_\pm \gamma^+ \gamma_5 q(z/2) | ps \rangle \\ &= \frac{-i}{2M} \bar{u}(p' s') \left[\frac{\epsilon_{ij} k_\perp^i \Delta_\perp^j}{M^2} G_{1,1}^q + \frac{\sigma^{i+} \gamma_5}{P^+} (k_\perp^i G_{1,2}^q + \Delta_\perp^i G_{1,3}^q) + \sigma^{+-} \gamma_5 G_{1,4}^q \right] u(ps)\end{aligned}$$

Meissner, Metz, Schlegel (2008)

Quark spin-orbit correlation [Lorce, Pasquini \(2011\)](#)

$$C_q = \int_{-1}^1 dx \int d^2 k_\perp \frac{k_\perp^2}{M^2} G_{1,1}^q(x, k_\perp, 0) \sim \langle S^z L^z \rangle$$

x-distribution

$$C_q(x) = \int d^2 k_\perp \frac{k_\perp^2}{M^2} G_{1,1}^q(x, k_\perp, 0)$$

$C_q > 0$ if helicity and OAM are aligned

$C_q < 0$ if they are anti-aligned

Gluon spin-orbit correlation

Polarized gluon GTMD

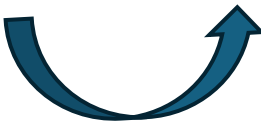
$$\begin{aligned} x\tilde{f}_g(x, \xi, k_\perp, \Delta_\perp) &= i \int \frac{d^3 z}{(2\pi)^3 P^+} e^{ixP^+ z^- - ik_\perp \cdot z_\perp} \langle p' | \tilde{F}^{+\mu}(-z/2) \widetilde{W}_\pm F_\mu^+(z/2) | p \rangle \\ &= \frac{-i}{2M} \bar{u}(p') \left[\frac{\epsilon_{ij} k^i \Delta^j}{M^2} G_{1,1}^g + \frac{\sigma^{i+} \gamma_5}{P^+} (k^i G_{1,2}^g + \Delta^i G_{1,3}^g) + \sigma^{+-} \gamma_5 G_{1,4}^g \right] u(p) \end{aligned}$$

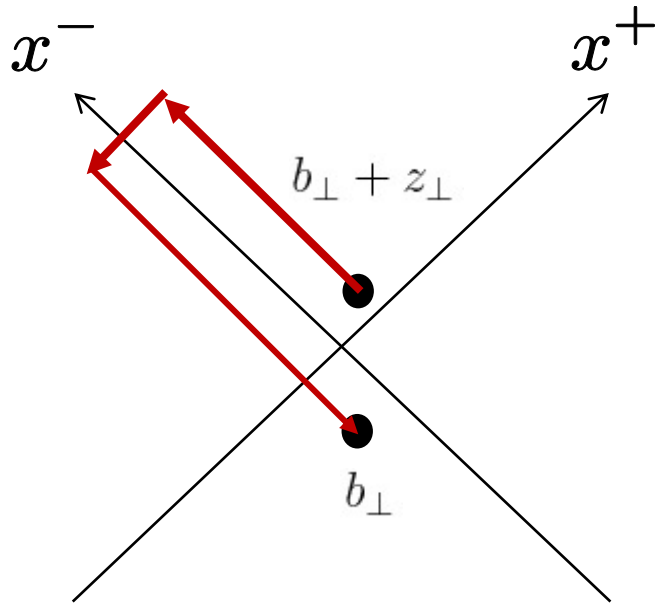
$$xC_g(x) = \int d^2 k_\perp \frac{k_\perp^2}{M^2} G_{1,1}^g(x, k_\perp, 0)$$

$C_g(x)$ is odd. The first moment vanishes $\int dx C_g(x) = 0$

OAM and spin-orbit correlation

$$L_q(x) = \int dk_{\perp} db_{\perp} b_{\perp} \times k_{\perp} f_q(x, k_{\perp}, b_{\perp}) \quad C_q(x) = \int dk_{\perp} db_{\perp} b_{\perp} \times k_{\perp} \tilde{f}_q(x, k_{\perp}, b_{\perp})$$


 γ_5 rotation



$$k_{\perp} \rightarrow \partial^{\mu} \rightarrow D^{\mu} - i \frac{1}{D^{+}} F^{+\mu}$$

Staple-shaped Wilson line

→ Gauge invariant canonical OAM [YH \(2011\)](#)

Twist structure of OAM PDF

YH, Yoshida (2012)

$$\begin{aligned}
 L_{can}^q(x) = & x \int_x^{\epsilon(x)} \frac{dx'}{x'} (H_q(x') + E_q(x')) - x \int_x^{\epsilon(x)} \frac{dx'}{x'^2} \tilde{H}_q(x') && \text{Wandzura-Wilczek part} \\
 & - x \int_x^{\epsilon(x)} dx_1 \int_{-1}^1 dx_2 \Phi_F(x_1, x_2) \mathcal{P} \frac{3x_1 - x_2}{x_1^2 (x_1 - x_2)^2} \\
 & - x \int_x^{\epsilon(x)} dx_1 \int_{-1}^1 dx_2 \tilde{\Phi}_F(x_1, x_2) \mathcal{P} \frac{1}{x_1^2 (x_1 - x_2)^2} .
 \end{aligned}$$

genuine twist-3

$$\Phi_F \sim \langle P' | \bar{\psi} \gamma^+ F^{+i} \psi | P \rangle$$

$$M_F \sim \langle P' | F^{+\mu} F^{+i} F_{\mu}^+ | P \rangle$$

$$\begin{aligned}
 L_{can}^g(x) = & \frac{x}{2} \int_x^{\epsilon(x)} \frac{dx'}{x'^2} (H_g(x') + E_g(x')) - x \int_x^{\epsilon(x)} \frac{dx'}{x'^2} \Delta G(x') \\
 & + 2x \int_x^{\epsilon(x)} \frac{dx'}{x'^3} \int dX \Phi_F(X, x') + 2x \int_x^{\epsilon(x)} dx_1 \int_{-1}^1 dx_2 \tilde{M}_F(x_1, x_2) \mathcal{P} \frac{1}{x_1^3 (x_1 - x_2)^2} \\
 & + 2x \int_x^{\epsilon(x)} dx_1 \int_{-1}^1 dx_2 M_F(x_1, x_2) \mathcal{P} \frac{2x_1 - x_2}{x_1^3 (x_1 - x_2)^2}
 \end{aligned}$$

Twist structure of spin-orbit correlation

YH, Schoenleber (2024)

See also,
Rajan, Engelhardt, Liuti (2017)
for the quark part

$$C_q(x) = x \int_x^{\epsilon(x)} \frac{dx'}{x'^2} x' \Delta q(x') - x \int_x^{\epsilon(x)} \frac{dx'}{x'^2} q(x') \\ - x \int_x^{\epsilon(x)} dx_1 \int_{-1}^1 dx_2 \frac{\Psi_{qF}(x_1, x_2)}{x_1 - x_2} P \frac{3x_1 - x_2}{x_1^2(x_1 - x_2)} \\ - x \int_x^{\epsilon(x)} dx_1 \int_{-1}^1 dx_2 \tilde{\Psi}_{qF}(x_1, x_2) P \frac{1}{x_1^2(x_1 - x_2)},$$

$$C_g(x) = x \int_x^{\epsilon(x)} \frac{dx'}{x'^2} x' \Delta G(x') - 2x \int_x^{\epsilon(x)} \frac{dx'}{x'^2} G(x') - 4x \sum_q \int_x^{\epsilon(x)} \frac{dx'}{x'^3} \int dX \tilde{\Psi}_{qF}(X, x') \\ + 4x \int_x^{\epsilon(x)} dx_1 \int dx_2 P \frac{\tilde{N}_F(x_1, x_2)}{x_1^3(x_1 - x_2)} + 4x \int_x^{\epsilon(x)} dx_1 \int dx_2 \frac{N_F(x_1, x_2)}{x_1^3(x_1 - x_2)} P \frac{2x_1 - x_2}{x_1 - x_2}$$

Ψ_F, N_F partly related to ETQS and three-gluon distributions relevant to transverse SSA

2 spin sum rules, 1 momentum sum rule?

Spin

$$\frac{1}{2} = \frac{1}{2} \sum_q (A_{q+\bar{q}} + B_{q+\bar{q}}) + \frac{1}{2}(A_g + B_g) \quad \text{Ji (1996)}$$

$$= \frac{1}{2} \Delta \Sigma + \Delta G + L_q + L_g \quad \text{Jaffe, Manohar (1990)}$$

Momentum

$$1 = \sum_q A_{q+\bar{q}} + A_g \quad \text{Feynman (1969)}$$

2 spin sum rules, 2 momentum sum rules!

Spin

$$\frac{1}{2} = \frac{1}{2} \sum_q (A_{q+\bar{q}} + B_{q+\bar{q}}) + \frac{1}{2}(A_g + B_g) \quad \text{Ji (1996)}$$

$$= \frac{1}{2} \Delta\Sigma + \Delta G + L_q + L_g \quad \text{Jaffe, Manohar (1990)}$$

Momentum

$$1 = \sum_q A_{q+\bar{q}} + A_g \quad \text{Feynman (1969)}$$

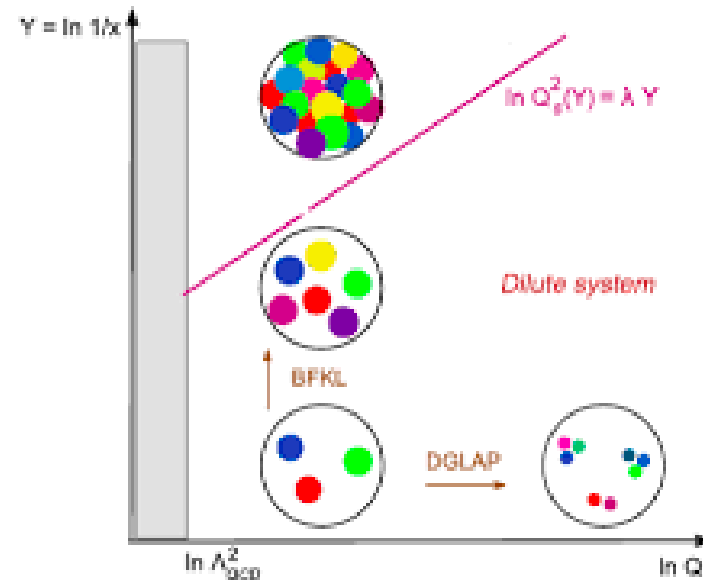
$$= \Delta\Sigma^{(3)} + \frac{1}{2} \Delta G^{(3)} - 3C_q^{(2)} - \frac{3}{2} C_g^{(2)} \quad \text{YH, Schoenleber (2024)}$$

$$- \frac{3}{2} \int dx dx' \left[\sum_q \left(\frac{2x}{x-x'} \tilde{\Psi}_{Fq}(x, x') + \Psi_{Fq}(x, x') \right) - \frac{\tilde{N}_F(x, x')}{x-x'} \right] + \sum_q \frac{m_q}{M} H_{1q}^{(2)}$$

Spin-orbit correlation at small-x

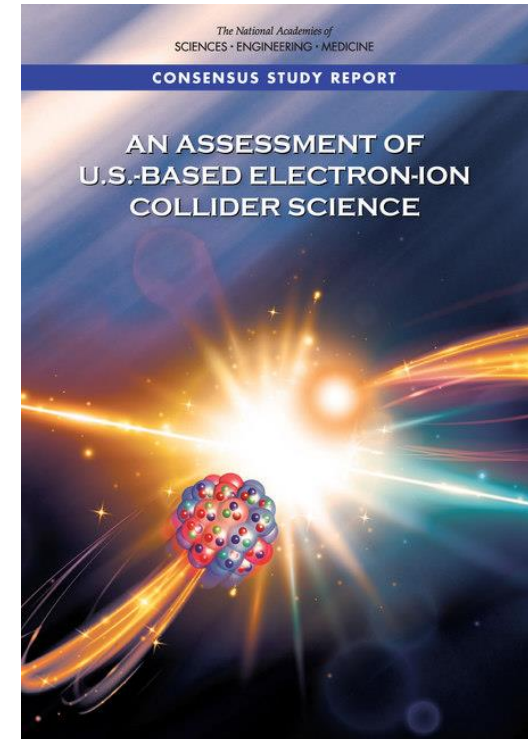
Gluon saturation at small-x:
one of the core topics of EIC

Naively, anything related to helicity, OAM are
subleading at small-x



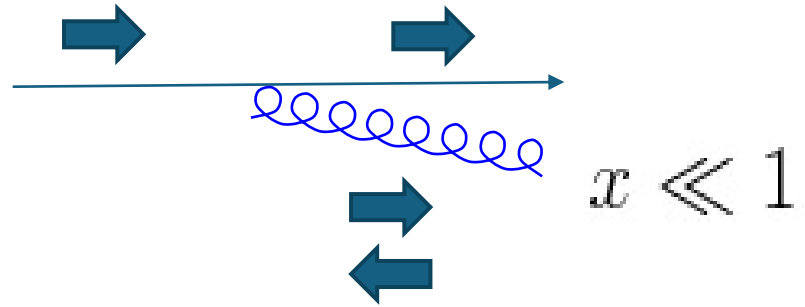
Finding 1: An EIC can uniquely address three profound questions about nucleons—protons—and how they are assembled to form the nuclei of atoms:

- How does the mass of the nucleon arise?
- How does the spin of the nucleon arise?
- What are the emergent properties of dense systems of gluons?



Intuitive argument

Imagine a very energetic quark emits a soft gluon

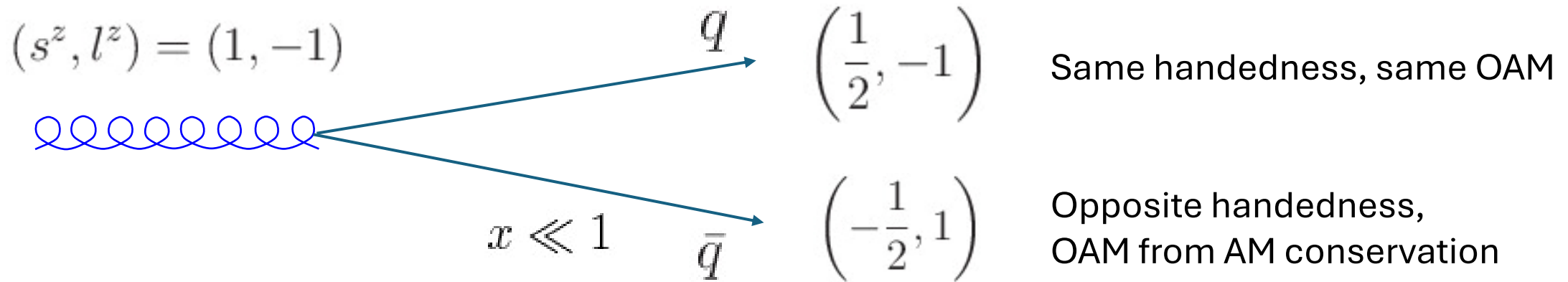


Quark spin and momentum unchanged.

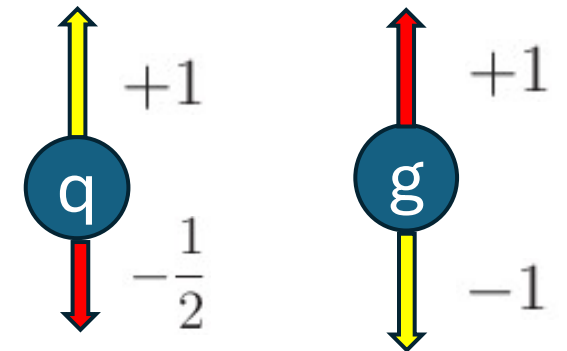
From angular momentum conservation, the total angular momentum of the emitted gluon must be zero

$$(s^z, l^z) = (\pm 1, \mp 1)$$

Imagine the emitted soft gluon further splits into a $q\bar{q}$ pair



Helicity and OAM are always in opposite directions
Remarkably, only $L^z = \pm 1$ states appear in this argument



Spin-orbit **anti**-correlation

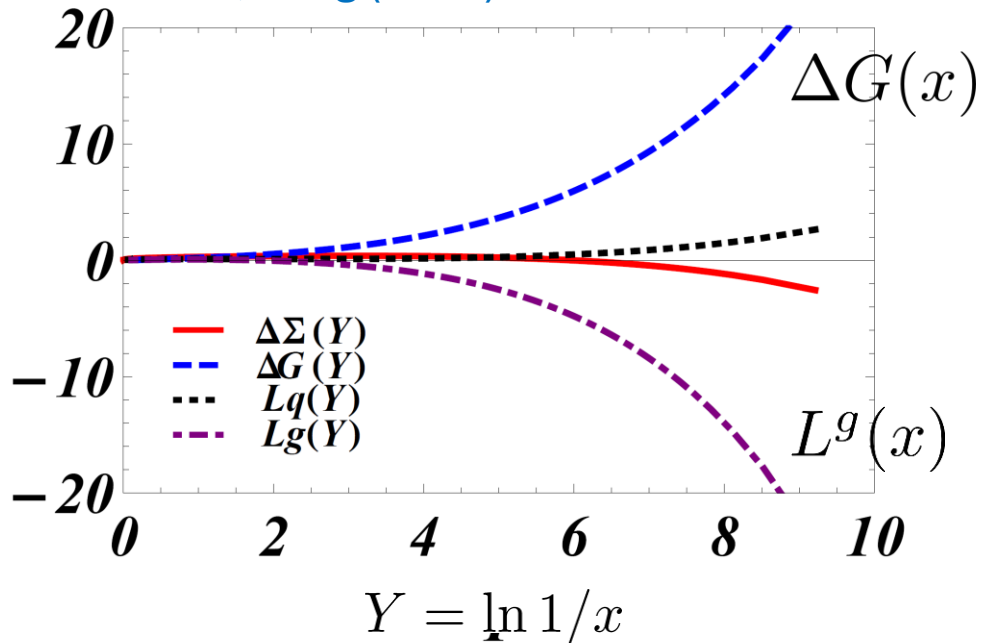
Helicity-OAM cancellation at small-x

If $\Delta q(x) \sim \Delta G(x) \sim \frac{1}{x^c}$, then

Boussarie, YH, Yuan (2019)

$$\Delta q(x) \approx -\frac{1}{1+c} L_q(x), \quad \Delta G(x) \approx -\frac{2}{1+c} L_g(x),$$

YH, Yang (2018)



More detailed analysis [Kovchegov, Manley \(2023\)](#)
[Manley \(2024\)](#)

In fact, nucleon polarization is not crucial.
The correlation exists even in
unpolarized/spinless hadrons

Gluon spin-orbit coupling at small-x

$$\frac{i}{x} \int \frac{d^3 z}{(2\pi)^3 P^+} e^{ixP^+ z^- - ik_\perp \cdot z_\perp} \langle p' | 2\text{Tr}[W_+ \tilde{F}^{+\mu}(-z/2) W_\pm F_\mu^+(z/2)] | p \rangle = -i \frac{\epsilon_{ij} k_\perp^i \Delta_\perp^j}{M^2} C_g^{[+\pm]}(x, \xi, k_\perp, \Delta_\perp),$$

There are two inequivalent configurations of Wilson lines

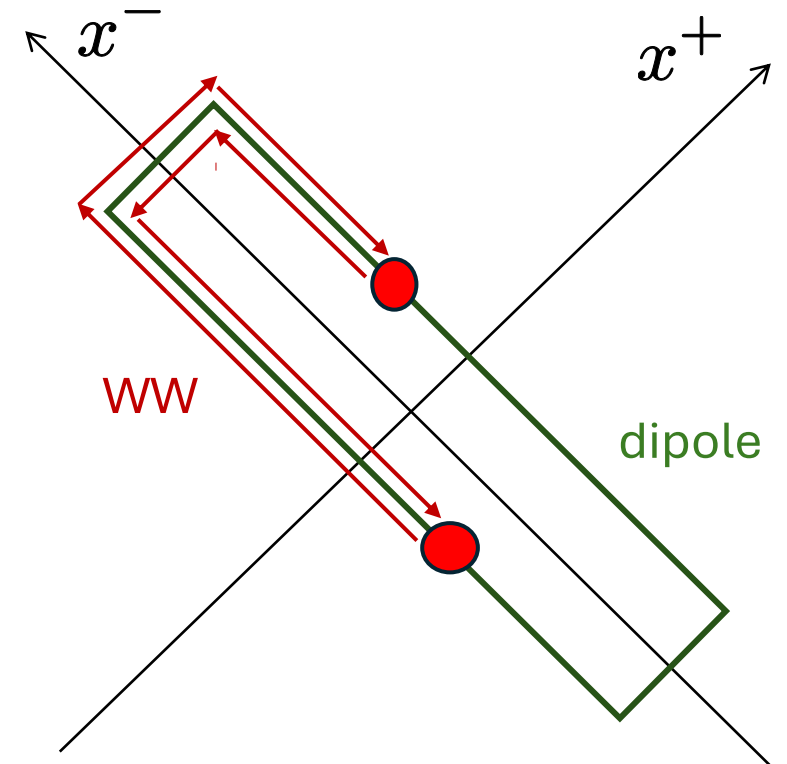
Weiszacker-Williams type

Dipole type

Bomhof, Mulders, Pijlman (2006)

Dominguez, Marquet, Xiao, Yuan (2011)

Approximate $e^{ixP^+ z^-} \approx 1$ (eikonal approximation)



Dipole gluon

$$\frac{x C_g^{\text{dip}}(x, k_\perp)}{M^2} = -\frac{2N_c}{\alpha_s} \int \frac{d^2 w_\perp d^2 z_\perp}{(2\pi)^4} e^{-ik_\perp \cdot (z_\perp - w_\perp)} \frac{\langle p | \frac{1}{N_c} \text{Tr} U(w_\perp) U^\dagger(z_\perp) - 1 | p \rangle}{\langle p | p \rangle}$$

cf. Boer, van Daal, Mulders, Petreska (2018)

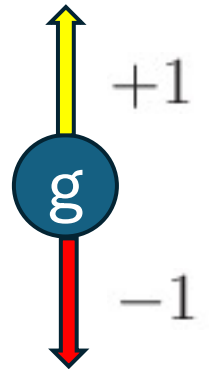
WW gluon

$$k_\perp^2 \frac{C_g^{\text{WW}}(x, k_\perp)}{M^2} = -f_g^{\text{WW}}(x, k_\perp) - \frac{C_F}{\pi \alpha_s x} \int \frac{d^2 b_\perp d^2 r_\perp}{(2\pi)^3} e^{-ik_\perp \cdot r_\perp} \partial_i^r D(r_\perp) \partial_i^r \left(\frac{1 - e^{\frac{N_c}{C_F} D(r_\perp)}}{D(r_\perp)} \right)$$

$$C_g^{\text{dip}}(x) = C_g^{\text{WW}}(x) = -G(x)$$

$$-1 \times 1 = -1$$

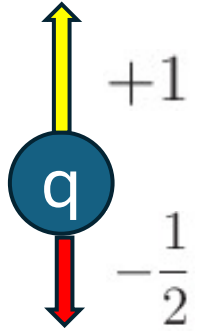
times the number of gluons



In fact, $C_g^{\text{WW}}(x) = C_g^{\text{dip}}(x)$ exactly for all values of x .

Quark spin-orbit coupling at small-x

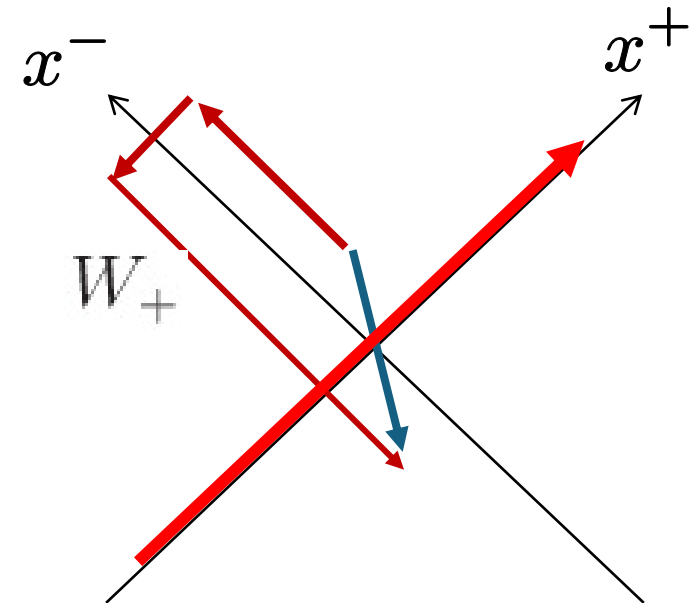
$$\int \frac{d^3 z}{2(2\pi)^3} e^{ixP^+ z^- - ik_\perp \cdot z_\perp} \langle p' | \bar{\psi}(-z/2) \gamma^+ \gamma_5 W_\pm \psi(z/2) | p \rangle = -i \frac{\epsilon_{ij} k_\perp^i \Delta_\perp^j}{M^2} C_q(x, \xi, k_\perp, \Delta_\perp)$$



$$\frac{C_q(x, k_\perp)}{M^2} = \frac{N_c S_\perp}{8\pi^4 x k_\perp^2} \int d^2 k_{g\perp} (k_\perp - k_{g\perp}) \cdot k_\perp \frac{\ln \frac{k_\perp^2}{(k_\perp - k_{g\perp})^2}}{k_\perp^2 - (k_\perp - k_{g\perp})^2} \frac{\langle p | \left(\frac{1}{N_c} \text{Tr} U U^\dagger - 1 \right) (k_{g\perp}) | p \rangle}{\langle p | p \rangle}$$

$$C_q(x) = -\frac{1}{2} q(x)$$

$$-\frac{1}{2} \times 1 = -\frac{1}{2} \text{ times the number of quarks}$$



Quantum entanglement of spin and OAM

Implement perfect spin-orbit **anti**-correlation

‘Bell states’

$$|\Psi^+\rangle = \frac{1}{\sqrt{2}}(|+\rangle_s |-\rangle_l + |-\rangle_s |+\rangle_l), \quad |\Psi^-\rangle = \frac{1}{\sqrt{2}i}(|+\rangle_s |-\rangle_l - |-\rangle_s |+\rangle_l)$$

Maximally entangled state realized on each soft gluon!

$$\langle S^z \rangle = \langle L^z \rangle = 0 \quad \text{but} \quad \langle S^z L^z \rangle = -1$$

True nature of the system encoded in correlations

QED example

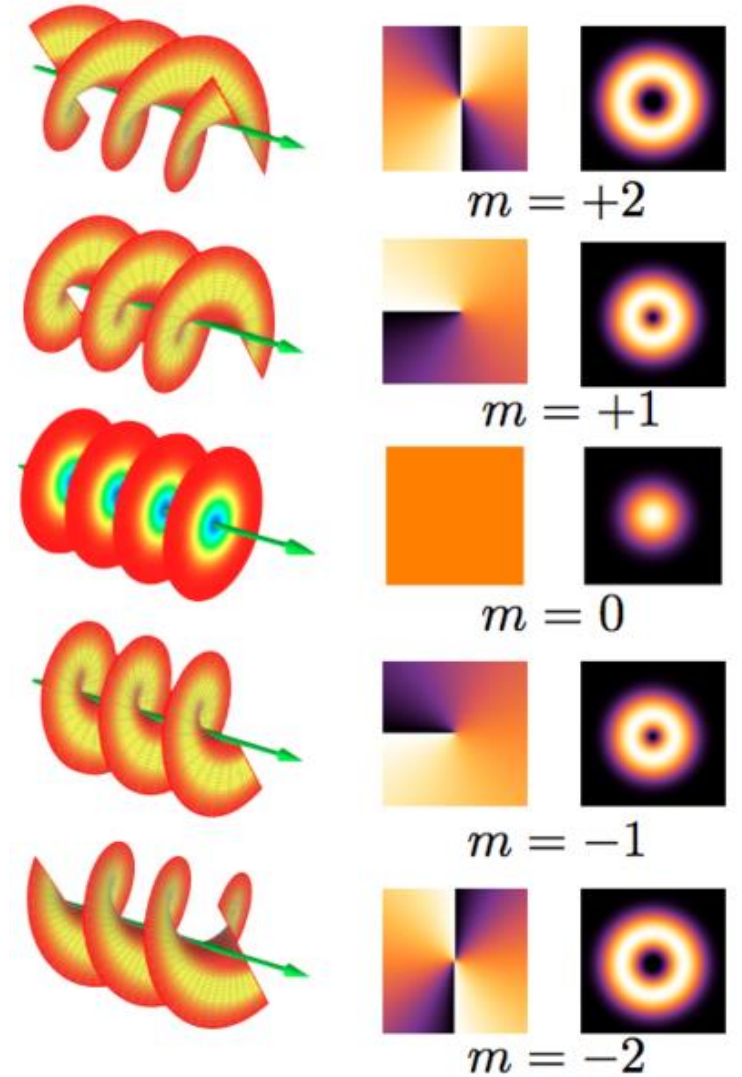
Photon OAM

$$|\pm\rangle_l \sim e^{\pm i\phi}$$

e.g., Laguerre-Gaussian beam

$$|\Psi^+\rangle \sim (1, i)e^{-i\phi} + (1, -i)e^{i\phi} \sim (\cos \phi, \sin \phi)$$

$$|\Psi^-\rangle \sim -i \left((1, i)e^{-i\phi} - (1, -i)e^{i\phi} \right) \sim (-\sin \phi, \cos \phi)$$



Quantum entanglement of the spin and orbital angular momentum of photons using metamaterials

TOMER STAV , ARKADY FAERMAN , ELHANAN MAGUID , DIKLA OREN , [...], AND MORDECHAI SEGEV

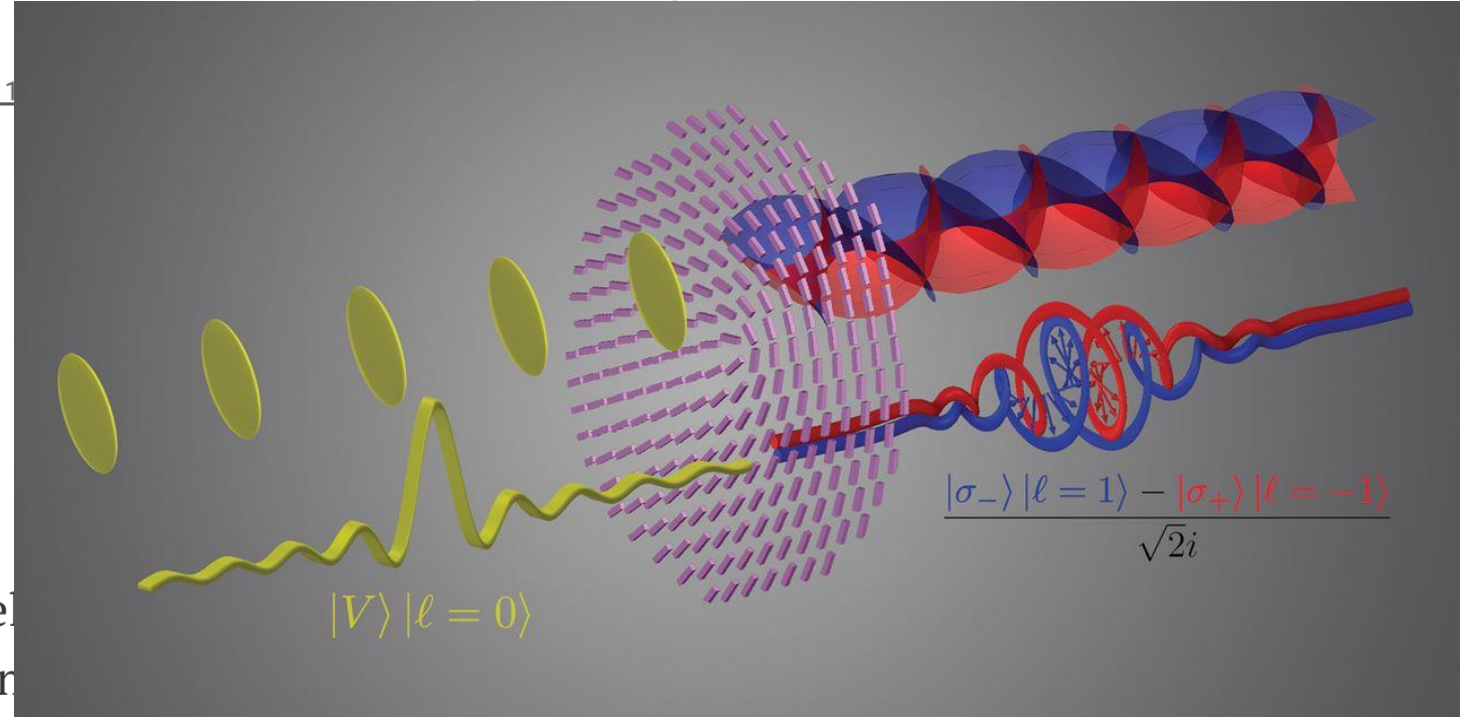
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SCIENCE 14 Sep 2018 Vol 361, Issue 6407 pp. 1101-1104 DOI: 10.1126/science.1257578

Abstract

Metamaterials constructed from deep subwavelength phenomena ranging from negative refractive index to general relativity, and superresolution imaging. More recently, metamaterials have been suggested as a new platform for quantum optics. We present the use of a dielectric metasurface to generate entanglement between the spin and orbital angular momentum of photons. We demonstrate the genera-



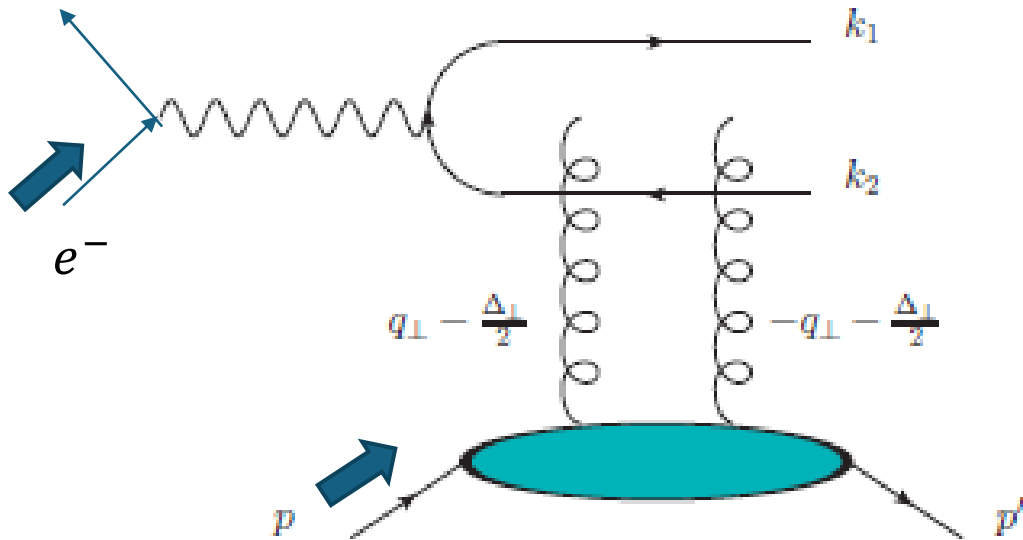
In QCD, spin-orbit entanglement is a default property of soft gluons!

Connection to experiment

Quark and gluon GTMDs $G_{1,1}$ appeared in certain exclusive reactions,
e.g., [Bhattacharya, Metz, Zhou \(2017\)](#)
but no quantitative estimate made.

Longitudinal double spin asymmetry in diffractive dijets
→ previously proposed as a signal of **gluon OAM**

[Bhattacharya, Boussarie, YH, \(2022\)](#)

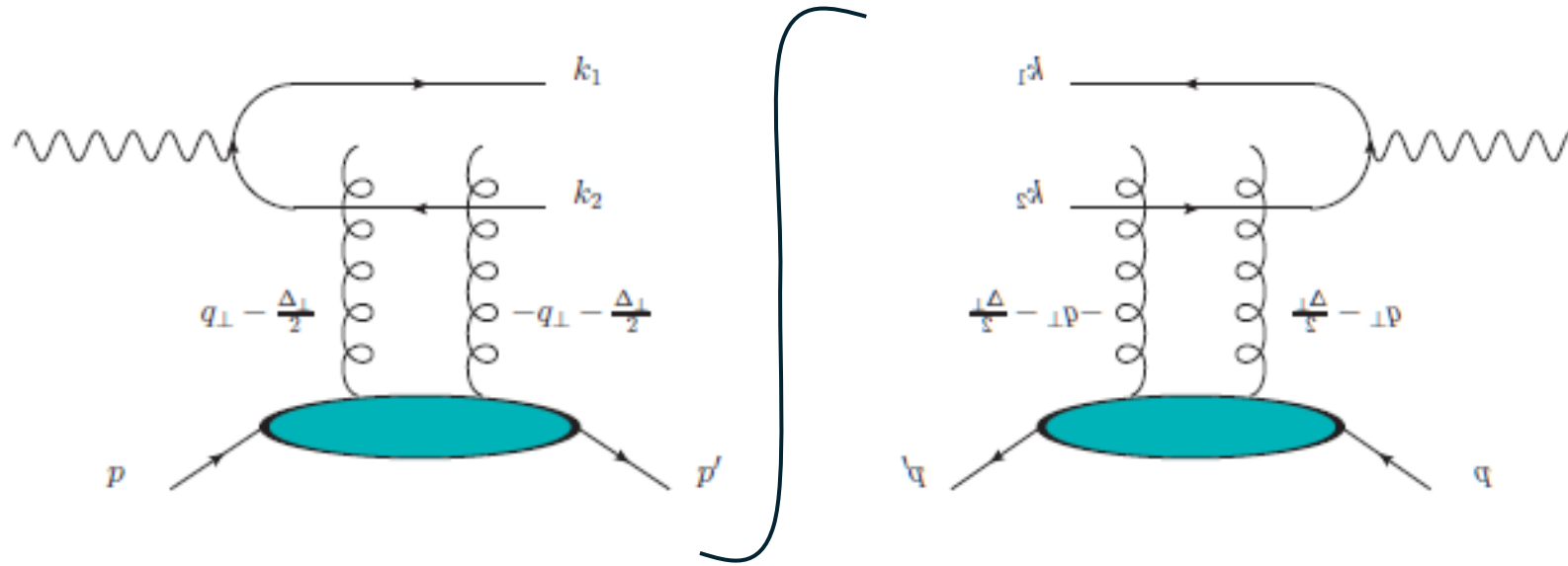


$$L^z \sim b_\perp \times k_\perp$$

conjugate to Δ_\perp
proton recoil momentum

correlated with jet
transverse momentum

Spin, orbit, and spin-orbit



\mathcal{H}_g

\mathcal{H}_g

$\tilde{\mathcal{H}}_g$

\mathcal{L}_g

$\tilde{\mathcal{H}}_g$

\mathcal{C}_g

OAM

helicity

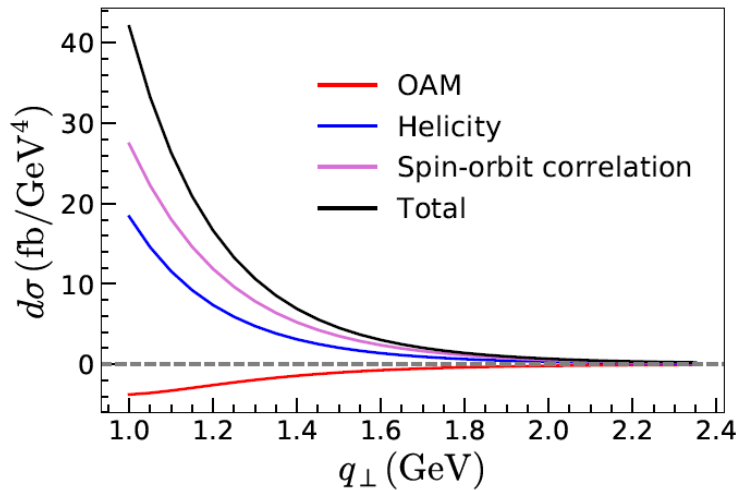
spin-orbit
correlation

computed in our
2022 paper

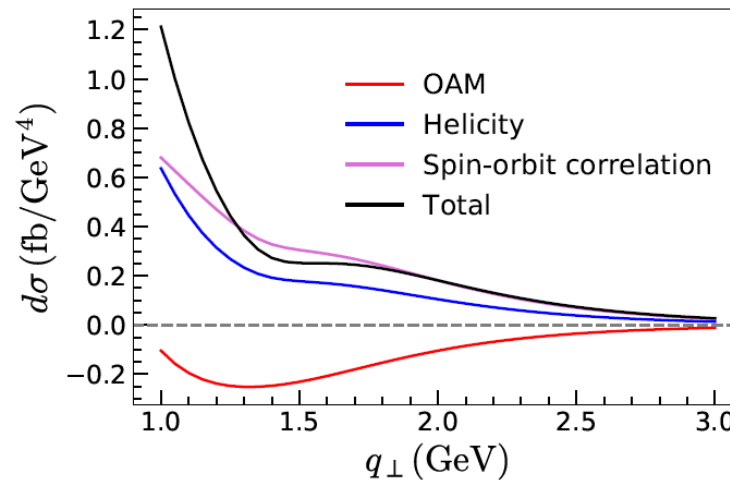
→ previously neglected

Prediction at the EIC (revised)

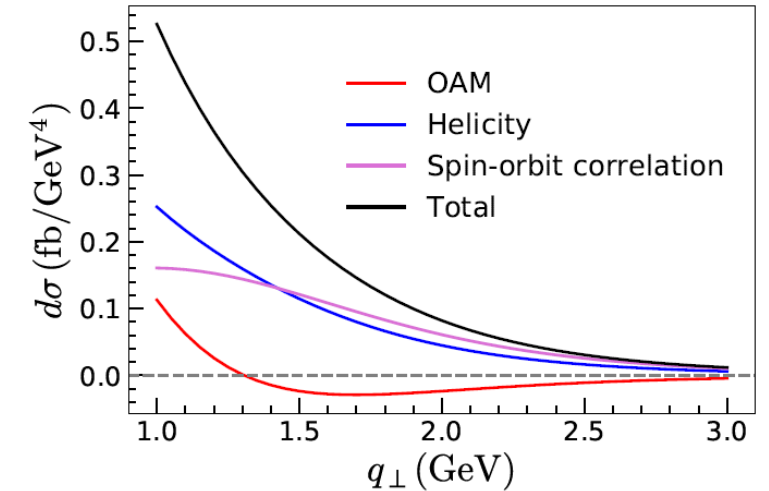
Bhattacharya, Boussarie, YH (2024)



$$Q^2 = 2.7 \text{ GeV}^2$$



$$Q^2 = 4.8 \text{ GeV}^2$$



$$Q^2 = 10 \text{ GeV}^2$$

In practice, jet reconstruction at low-Pt is challenging.

Re-formulate the process as semi-inclusive **diffractive** DIS (SIDDIS) [YH, Xiao, Yuan \(2022\)](#)

Conclusions

Quark and gluon spin-orbit correlations analyzed in QCD

New momentum sum rule, analog of Jaffe-Manohar spin sum rule

Novel emergent property of dense systems of gluons uncovered

Quantum entanglement between spin and OAM

→ Connection to QIS? EIC?

Finding 1: An EIC can uniquely address three profound questions about nucleons—protons—and how they are assembled to form the nuclei of atoms:

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