STATE UNIVERSITY OF NEW YORK

Status of unpolarized TMD extractions from global fits



From Quarks and Gluons to the Internal Dynamics of Hadrons

May 15, 2024

3-*dimensional map* of the internal structure of the nucleon



3-*dimensional map* of the internal structure of the nucleon



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3-*dimensional map* of the internal structure of the nucleon

Non-collinear framework

Nucleon Pol.

Quark Polarization

		U	L	Т
l	J	f_1		
L	-		g_1	
Т	-			h_1



3-*dimensional map* of the internal structure of the nucleon

Non-collinear framework

Nucleon Pol.

Quark Polarization

	U	L	Т
U	f_1		h_1^\perp
L		g_1	h_{1L}^{\perp}
Т	f_{1T}^{\perp}	g_{1T}	$h_1 h_{1T}^{\perp}$



3-*dimensional map* of the internal structure of the nucleon

Non-collinear framework

Quark Polarization

	U	L	Т
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TMD PDFs

 $F(x, \boldsymbol{k}_{\perp}^2, \mu, \zeta)$

3-*dimensional map* of the internal structure of the nucleon

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TMD PDFs

 $F(x, \boldsymbol{k}_{\perp}^2, \mu, \zeta)$

Fraction of longitudinal momentum



TMDs map the distribution of partons inside the nucleon in 3D in momentum space.

They can be extracted through *global fits* There are attempts to calculate them in lattice QCD

Fraction of longitudinal momentum



TMDs map the distribution of partons inside the nucleon in 3D in momentum space.

They can be extracted through *global fits* There are attempts to calculate them in lattice QCD

Are TMDs universal?

Do they depend on x?

Do they depend on the quark flavor?

Transverse momentum

Semi-Inclusive Deep-Inelastic Scattering



$$\begin{split} F_{UU,T}(x,z,|\boldsymbol{q}_{T}|,Q) &= \frac{x}{2\pi} \,\mathcal{H}^{\text{SIDIS}}(Q,\mu) \sum_{a=q,\bar{q}} e_{a}^{2} \int_{0}^{+\infty} d|\boldsymbol{b}_{T}||\boldsymbol{b}_{T}|J_{0}(|\boldsymbol{b}_{T}||\boldsymbol{q}_{T}|) \hat{f}_{1}^{a}(x,b_{T}^{2};\mu,\zeta_{A}) \hat{D}_{1}^{a\to h}(z,b_{T}^{2};\mu,\zeta_{B}) \\ &+ Y_{UU,T}(Q^{2},\mathbf{P}_{hT}^{2}) + \mathcal{O}(M^{2}/Q^{2}) \end{split}$$



Semi-Inclusive Deep-Inelastic Scattering

hadron **A** If $Q^2 \gg M^2$ and $Q^2 \gg q_T^2(P_{hT}^2)$ P_h P_{hT} $\sim zk_{\perp}$ **TMD FF** р k_{\perp} photon quark **TMD PDF** k_{\perp} proton Р $F_{UU,T}(x,z,|\boldsymbol{q}_{T}|,Q) = \frac{x}{2\pi} \mathcal{H}^{\text{SIDIS}}(Q,\mu) \sum_{a=q,\bar{q}} e_{a}^{2} \int_{0}^{+\infty} d|\boldsymbol{b}_{T}||\boldsymbol{b}_{T}|J_{0}(|\boldsymbol{b}_{T}||\boldsymbol{q}_{T}|) \hat{f}_{1}^{a}(x,b_{T}^{2};\mu,\zeta_{A}) \hat{D}_{1}^{a\to h}(z,b_{T}^{2};\mu,\zeta_{B})$ $+Y_{UU,T}(Q^2, \mathbf{P}_{hT}^2) + \mathcal{O}(M^2/Q^2)$

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- The <u>W term</u> dominates in the region where q_T «Q
- The Y term has been excluded in the analysis









TMD in Fourier space

$$\hat{F}(x, b_T^2; \mu, \zeta) = \int \frac{d^2 \boldsymbol{k}_\perp}{(2\pi)^2} e^{i\boldsymbol{b}_T \cdot \boldsymbol{k}_\perp} F(x, k_\perp^2; \mu, \zeta)$$

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$$\hat{f}_1^q(x, b_T^2; \mu, \zeta) = \sum_j C_{q/j}(x, b_*; \mu_{b_*}, \mu_{b_*}^2) \otimes f_1^j(x, \mu_{b_*})$$

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 b_* -prescription

$$\hat{f}_1^q(x, b_T^2; \mu, \zeta) = \sum_j C_{q/j}(x, b_*; \mu_{b_*}, \mu_{b_*}^2) \otimes f_1^j(x, \mu_{b_*})$$
:A

Perturbative TMD at the initial scale

TMD in Fourier space

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:A

Perturbative TMD at the initial scale

$$\times \exp\left\{K(b_*;\mu_{b_*})\ln\frac{\sqrt{\zeta}}{\mu_{b_*}} + \int_{\mu_{b_*}}^{\mu}\frac{d\mu'}{\mu'}\left[\gamma_F - \gamma_K\ln\frac{\sqrt{\zeta}}{\mu'}\right]\right\} : \mathsf{B}$$

Evolution to final scale (of the process)

TMD in Fourier space

$$\hat{F}(x, b_T^2; \mu, \zeta) = \int \frac{d^2 \mathbf{k}_\perp}{(2\pi)^2} e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} F(x, k_\perp^2; \mu, \zeta)$$

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Evolution to final scale (of the process)

$$\times f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : \mathcal{C}$$

Non-perturbative part of the TMD

TMD in Fourier space

$$\hat{F}(x, b_T^2; \mu, \zeta) = \int \frac{d^2 \mathbf{k}_\perp}{(2\pi)^2} e^{i\mathbf{b}_T \cdot \mathbf{k}_\perp} F(x, k_\perp^2; \mu, \zeta)$$

$$\hat{f}_1^q(x, b_T^2; \mu, \zeta) = \sum_j \underbrace{C_{q/j}(x, b_*; \mu_{b_*}, \mu_{b_*}^2)}_{j} \otimes f_1^j(x, \mu_{b_*})$$
:A

Perturbative TMD at the initial scale

Perturbative
$$\times \exp\left\{K(b_*;\mu_{b_*})\ln\frac{\sqrt{\zeta}}{\mu_{b_*}} + \int_{\mu_{b_*}}^{\mu}\frac{d\mu'}{\mu'}\left[\gamma_F - \gamma_K \ln\frac{\sqrt{\zeta}}{\mu'}\right]\right\} : B$$

Evolution to final scale (of the process)

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Non-perturbative part of the TMD

TMD in Fourier space

$$\begin{split} \hat{F}(x, b_T^2; \mu, \zeta) &= \int \frac{d^2 \mathbf{k}_\perp}{(2\pi)^2} e^{i \mathbf{b}_T \cdot \mathbf{k}_\perp} F(x, k_\perp^2; \mu, \zeta) & \text{Collinear extractions} \\ \hat{f}_1^q(x, b_T^2; \mu, \zeta) &= \sum_j \underbrace{C_{q/j}(x, b_*; \mu_{b_*}, \mu_{b_*}^2)}_{j} \otimes \underbrace{f_1^j(x, \mu_{b_*})}_{j} \otimes \underbrace{f_1^j(x, \mu_{b_*})}_{j} &: \mathbf{A} \\ & \text{Perturbative TMD at the initial scale} \\ & \text{Perturbative } \times \exp\left\{\underbrace{K(b_*; \mu_{b_*})}_{k_{b_*}} \ln \frac{\sqrt{\zeta}}{\mu_{b_*}} + \int_{\mu_{b_*}}^{\mu} \frac{d\mu'}{\mu'} \left[\gamma_F - \gamma_K \ln \frac{\sqrt{\zeta}}{\mu'}\right]\right\} &: \mathbf{B} \\ & \text{Evolution to final scale (of the process)} \\ & \times f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} &: \mathbf{C} \end{split}$$

Non-perturbative part of the TMD

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$$F_{UU,T}(x,z,|\boldsymbol{q}_{T}|,Q) \sim \int_{0}^{+\infty} d|\boldsymbol{b}_{T}||\boldsymbol{b}_{T}|J_{0}(|\boldsymbol{b}_{T}||\boldsymbol{q}_{T}|)\hat{f}_{1}^{a}(x,b_{T}^{2};\mu,\zeta_{A})\hat{D}_{1}^{a\to h}(z,b_{T}^{2};\mu,\zeta_{B})$$



$$F_{UU,T}(x,z,|\boldsymbol{q}_{T}|,Q) \sim \int_{0}^{+\infty} d|\boldsymbol{b}_{T}||\boldsymbol{b}_{T}|J_{0}(|\boldsymbol{b}_{T}||\boldsymbol{q}_{T}|)\hat{f}_{1}^{a}(x,b_{T}^{2};\mu,\zeta_{A})\hat{D}_{1}^{a\to h}(z,b_{T}^{2};\mu,\zeta_{B})$$

$$F_{UU}^{1}(x_{A},x_{B},|\boldsymbol{q}_{T}|,Q) \sim \int_{0}^{+\infty} d|\boldsymbol{b}_{T}||\boldsymbol{b}_{T}|J_{0}(|\boldsymbol{b}_{T}||\boldsymbol{q}_{T}|)\hat{f}_{1}^{a}(x_{A},b_{T}^{2};\mu,\zeta_{A})\hat{f}_{1}^{\bar{a}}(x_{B},b_{T}^{2};\mu,\zeta_{B})$$





GLOBAL FITs

Available TMD fitting frameworks

https://github.com/MapCollaboration/NangaParbat



\equiv README.md

Nanga Parbat is a fitting framework aimed at the determination of the non-perturbative component of TMD distributions.

https://teorica.fis.ucm.es/artemide/

arTeMiDe

Download

You can obtain NangaParbat directly from the github repository:

https://github.com/MapCollaboration/NangaParbat

For the last development branch you can clone the master code:

git clone git@github.com:MapCollaboration/NangaParbat.git







News

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Archive of older links/news.

Recent version/release can be found in repository

Download





If you have found mistakes, or have suggestions/questions, please, contact us.

Some extra materials can be found on <u>Alexey's web-page</u>

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Articles, presentations & supplementary materials

Extra pictures for the paper arXiv:1902.08474

About us & Contacts

Link to the text in Inspire.

Archive of older links/news

Seminar of A.Vladimirov in Pavia 2018 on TMD evolution

Available Global Fits

	Accuracy	SIDIS	DY	N of points	χ²/N _{data}
Pavia 2017 Bacchetta, Delcarro, et al., JHEP 06 (2017)	NLL			8059	1.55
SV 2019 Scimemi, Vladimirov, JHEP 06 (2020)	N ³ LL ⁻		~	1039	1.06
MAPTMD22 Bacchetta, Bertone, et al., JHEP 10 (2022)	N³LL⁻	~	~	2031	1.06

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MAPTMD22 global fit

- Global analysis of Drell-Yan and SIDIS data sets: 2031 data points
- Perturbative accuracy: N³LL⁻
- Number of fitted parameters: 21
- Extremely good description: $\chi^2 / N_{data} = 1.06$

Differences in recent global fits

MAPTMD22 vs SV19
Differences in recent global fits

MAPTMD22 vs SV19

Criteria of data selection
 2031 vs 1039 included data

Differences in recent global fits

MAPTMD22 vs SV19

- Criteria of data selection
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- Implementation of TMD evolution

CSS framework vs zeta-prescription

Differences in recent global fits

MAPTMD22 vs SV19

- Criteria of data selection
 2031 vs 1039 included data
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CSS framework vs zeta-prescription

Nonperturbative parameterization
Gaussian + wGaussian vs complicated function





Drell-Yan data

484

DY fixed-target + collider



Drell-Yan data

484

SIDIS data 1547

DY fixed-target + collider HERMES + COMPASS



Drell-Yan data

SIDIS data 1547

HERMES + COMPASS

Total number of data 2031

Perturbative TMD at the initial scale

$$\hat{f}_1^q(x, b_T^2; \mu, \zeta) = \sum_j C_{q/j}(x, b_*; \mu_{b_*}, \mu_{b_*}^2) \otimes f_1^j(x, \mu_{b_*})$$
:A

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Collinear distributions

:A

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Collinear distributions

Input for PDFs: MMHT2014



Harland-Lang, Martin, Motylinski, Thorne, EPJ C 75 (2015)

:A

Perturbative TMD at the initial scale

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Collinear distributions

Input for PDFs: MMHT2014



Input for FFs: DSS14 - DSS17



De Florian, Sassot, Hepele, Hernandez-Pinto, Stratmann, PRD 91 (2015) De Florian, Hepele, Hernandez-Pinto, Sassot, Stratmann, PRD 95 (2017)

Harland-Lang, Martin, Motylinski, Thorne, EPJ C 75 (2015)

:A

 $f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\}$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\}$$

$$f_{1NP}(x, b_T^2) \propto \text{F.T. of} \left(e^{-\frac{k_\perp^2}{g_{1A}}} + \lambda_B k_\perp^2 e^{-\frac{k_\perp^2}{g_{1B}}} + \lambda_C e^{-\frac{k_\perp^2}{g_{1C}}} \right)$$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\}$$

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 $g_1(x) = N_1 \frac{(1-x)^{\alpha} x^{\sigma}}{(1-\hat{x})^{\alpha} \hat{x}^{\sigma}}$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\}$$

$$f_{1\mathrm{NP}}(x, b_T^2) \propto \mathrm{F.T.} \text{ of } \left(e^{-\frac{k_\perp^2}{g_{1A}}} + \lambda_B k_\perp^2 e^{-\frac{k_\perp^2}{g_{1B}}} + \lambda_C e^{-\frac{k_\perp^2}{g_{1C}}} \right)$$

$$g_1(x) = N_1 \frac{(1-x)^{\alpha} x^{\sigma}}{(1-\hat{x})^{\alpha} \hat{x}^{\sigma}}$$

$$D_{1\mathrm{NP}}(x, b_T^2) \propto \mathrm{F.T.} \text{ of } \left(e^{-\frac{P_\perp^2}{g_{3A}}} + \lambda_{FB} k_\perp^2 e^{-\frac{P_\perp^2}{g_{3B}}} \right)$$

$$g_3(z) = N_3 \frac{(z^{\beta} + \delta)(1-z)^{\alpha}}{(\hat{z}^{\beta} + \delta)(1-\hat{z})^{\alpha}}$$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\}$$

$$f_{1NP}(x, b_T^2) \propto F.T. \text{ of } \left(e^{-\frac{k_\perp^2}{g_{1A}}} + \lambda_B k_\perp^2 e^{-\frac{k_\perp^2}{g_{1B}}} + \lambda_C e^{-\frac{k_\perp^2}{g_{1C}}} \right)$$

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$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\}$$

Bacchetta, Gamberg, Goldstein, et al., PLB 659 (2008) Bacchetta, Conti, Radici, PRD 78 (2008) Pasquini, Cazzaniga, Boffi, PRD 78 (2008) Matevosyan, Bentz, Cloet, Thomas, PRD 85 (2012) Burkardt, Pasquini, EPJA (2016) Grewal, Kang, Qiu, Signori, PRD 101 (2020)

$$\begin{split} f_{1\mathrm{NP}}(x,b_T^2) &\propto \mathrm{F.T. of} \left(e^{-\frac{k_\perp^2}{g_{1A}}} + \lambda_B k_\perp^2 e^{-\frac{k_\perp^2}{g_{1B}}} + \lambda_C e^{-\frac{k_\perp^2}{g_{1C}}} \right) \\ g_1(x) &= N_1 \frac{(1-x)^\alpha x^\sigma}{(1-\hat{x})^\alpha \hat{x}^\sigma} \\ D_{1\mathrm{NP}}(x,b_T^2) &\propto \mathrm{F.T. of} \left(e^{-\frac{P_\perp^2}{g_{3A}}} + \lambda_{FB} k_\perp^2 e^{-\frac{P_\perp^2}{g_{3B}}} \right) \\ g_3(z) &= N_3 \frac{(z^\beta + \delta)(1-z)^\gamma}{(\hat{z}^\beta + \delta)(1-\hat{z})^\gamma} \\ g_K(b_T^2) &= -g_2^2 \frac{b_T^2}{4} \end{split}$$

11 parameters for TMD PDF + 1 for NP evolution + 9 for TMD FF = 21 free parameters

High-Energy Drell-Yan beyond NLL

 $Q\sim 100~{\rm GeV}$



Bacchetta, Bertone, Bissolotti, et al., JHEP 07 (2020)

SIDIS observables beyond NLL

High-Energy Drell-Yan beyond NLL

 $Q\sim 100~{\rm GeV}$



Bacchetta, Bertone, Bissolotti, et al., JHEP 07 (2020)

SIDIS observables beyond NLL

 $Q \sim 2 \,\,\mathrm{GeV}$

High-Energy Drell-Yan beyond NLL

 $Q\sim 100~{\rm GeV}$



Bacchetta, Bertone, Bissolotti, et al., JHEP 07 (2020)



Bacchetta, Bertone, Bissolotti, et al., JHEP 07 (2020)





Bacchetta, Bertone, Bissolotti, et al., JHEP 07 (2020)

The description considerably worsens at higher orders!!

COMPASS multiplicities (one of many bins)



COMPASS multiplicities (one of many bins)



Discrepancy of an almost constant factor

Normalization issue confirmed also in other analyses from different collaborations



Normalization issue confirmed also in other analyses from different collaborations

22 $\mathbf{20}$ Q^2 10^{-3} Q^2 = 20.00 GeV= 22.10 G0.24 < z < 0.30 $\mathbf{\bar{x}_{Bj}} = 0.291$ $x_{Bj} = 0.157$ 0.30 < z < 0.4014 10^{0} 0.40 < z < 0.50 10° 0.65 < z < 0.70▼ 10^{-2} 10^{-3} ${
m Q}^2=8.30\,\,{
m GeV}^2$ $Q^2=9.30~{
m GeV}^2$ $Q^2 = 9.80 ~{
m GeV}^2$ $Q^2 = 11.00 \,\, \mathrm{GeV}$ $\mathbf{x}_{\mathrm{Bi}} = 0.077$ $\mathbf{x}_{\mathrm{Bj}} = 0.045$ $x_{Bj} = 0.149$ $\mathrm{x_{Bj}}=0.254$ 2 4 6 10 18 10^0 $q_T (GeV)$ 10 10^{-2} 10^{-3} $Q^2=3.50\,\,{\rm GeV}^2$ $Q^2=4.10~{
m GeV}^2$ $\rm Q^2=5.30~GeV^2$ $\mathbf{2.5}$ $\mathbf{x}_{\mathrm{Bi}} = 0.017$ $\mathbf{x}_{\mathrm{Bi}} = 0.026$ $x_{Bi} = 0.043$ $\mathbf{x}_{\mathrm{Bi}} = 0.075$ $\mathbf{x}_{Bi} = 0.133$ 4 6 2 $q_T ~(GeV)$ 10⁴ 2.0 10^{-1} 10^{-2} 10^{-3} $Q^2=1.80\,\,{\rm GeV}^2$ $Q^2=2.10~{
m GeV}^2$ $Q^2=2.30\,\,{
m GeV}^2$ $\mathrm{Q}^2=2.30~\mathrm{GeV}^2$ $Q^2=2.30\,\,{
m GeV}^2$ $\mathrm{Q}^2=2.50~\mathrm{GeV}^2$ 1.5 $\mathbf{x_{Bj}} = 0.007$ $x_{Bj} = 0.011$ $x_{Bj} = 0.016$ $\mathbf{x}_{\mathbf{Bj}} = 0.025$ $\mathbf{x_{Bj}} = 0.042$ $\mathbf{x}_{\mathrm{Bj}} = 0.068$ 10^{0} $F_{UU} = \sum \mathcal{H}_q \text{ F.T.} \left\{ \tilde{D}_{h/q}(z, z \boldsymbol{b}_{\perp}; Q) \ \tilde{f}_{q/P}(x, \boldsymbol{b}_{\perp}; Q) \right\}$ 10^{-1} 1.0 10^{-2} 10^{-3} 0.5large q_T corrections + power suppressed terms q_T (GeV) q_T (GeV) q_T (GeV) q_T (GeV) q_T (GeV) N_i $\rightarrow^{x_{Bj}}$

Sun, Isaacson, Yuan, Yuan, IJNP A (2014) Gonzalez-Hernandez, PoS DIS2019 (2019)

Normalization issue confirmed also in other analyses from different collaborations



Sun, Isaacson, Yuan, Yuan, IJNP A (2014)

Gonzalez-Hernandez, PoS DIS2019 (2019)

Vladimirov, JHEP 12 (2023)

The situation is worse for Q ~ 2–4 GeV which are typical for Semi-Inclusive Deep-Inelastic Scattering (SIDIS). In this case the **problem with normalization is of order of factor 2-3**

Normalization issue confirmed also in other analyses from different collaborations



Sun, Isaacson, Yuan, Yuan, IJNP A (2014)

Gonzalez-Hernandez, PoS DIS2019 (2019)

Vladimirov, JHEP 12 (2023)

The situation is worse for Q ~ 2–4 GeV which are typical for Semi-Inclusive Deep-Inelastic Scattering (SIDIS). In this case the **problem with normalization** *is of order of factor 2-3*

In contrast with SV19 analysis

situation at low energy scale

 $\langle Q \rangle = 3 \text{ GeV}$







MAP22 work solution

MAP22 work solution

SIDIS multiplicity

$$M(x, z, P_{hT}, Q) = \frac{d\sigma}{dx dQ dz dP_{hT}} \left/ \frac{d\sigma}{dx dQ} \right|$$

MAP22 work solution

SIDIS multiplicity

$$M(x, z, P_{hT}, Q) = \frac{d\sigma}{dx dQ dz dP_{hT}} \left/ \frac{d\sigma}{dx dQ} \frac{d\sigma}{dx dQ dz} \right|$$

Collinear SIDIS cross section
MAP22 work solution

SIDIS multiplicity

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Collinear SIDIS cross section

Good agreement theory/data



Khalek, Bertone, Nocera, et al., PRD 104 (2021)

MAP22 work solution

SIDIS multiplicity

$$M(x, z, P_{hT}, Q) = \frac{d\sigma}{dx dQ dz dP_{hT}} \left/ \frac{d\sigma}{dx dQ} \frac{d\sigma}{dx dQ} \frac{d\sigma}{dx dQ dz} \right|_{hT}$$

Collinear SIDIS cross section

Normalization of prediction such that

$$\int d\mathbf{P_{hT}}W(x, z, Q, \mathbf{P_{hT}}) = \frac{d\sigma}{dxdQdz}$$

Piacenza, PhD thesis (2020)

Good agreement theory/data



Khalek, Bertone, Nocera, et al., PRD 104 (2021)

MAP22 work solution

SIDIS multiplicity

SIDIS multiplicity
$$M(x, z, P_{hT}, Q) = \frac{d\sigma}{dx dQ dz dP_{hT}} / \frac{d\sigma}{dx dQ}$$

Collinear SIDIS cross section $\frac{d\sigma}{dx dQ dz}$

Normalization of prediction such that

Good agreement theory/data



Khalek, Bertone, Nocera, et al., PRD 104 (2021)

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Normalization of prediction such that

Good agreement theory/data



Calculable before the fit

Khalek, Bertone, Nocera, et al., PRD 104 (2021)

MAP22: Results for SIDIS data



MAP22: Results for DY data



MAP22: Extracted TMD PDFs



Non-trivial dependence on the variable x

Need more data to better constrain small-x region (EIC?)

MAPTMD22 vs SV19



Contrary to collinear extractions, we are still far from getting a good compatibility between two different TMD extractions

Comparison between recent global fits

Collins-Soper kernel:

Comparison between recent global fits

Collins-Soper kernel: kernel of the rapidity evolution equation

$$\frac{\partial \ln \hat{f}_1(x, b_T; \mu, \zeta)}{\partial \ln \sqrt{\zeta}} = K(b_T, \mu)$$

Collins-Soper kernel: kernel of the rapidity evolution equation

$$\frac{\partial \ln \hat{f}_1(x, b_T; \mu, \zeta)}{\partial \ln \sqrt{\zeta}} = K(b_T, \mu)$$

perturbatively calculable

Collins-Soper kernel: kernel of the rapidity evolution equation



Bollweg, Gao, Mukherjee, et al., PLB 852 (2024)

Collins-Soper kernel: kernel of the rapidity evolution equation



Bollweg, Gao, Mukherjee, et al., PLB 852 (2024)

Compatibility between the most recent extractions

What happened in the

last 2 years?

VOI, XNO. 2990,	NEW-YORK, SUNDAY, APRIL 21, 1961					PRICE TWO CENTS.
	Accuracy	SIDIS	DY	N of points	χ²/N _{data}	Flavor Dependence
MAPTMD22 Bacchetta, Bertone, et al., JHEP 10 (2022)	N³LL ⁻	~	~	2031	1.06	×

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PDF bias Bury, Hautmann, JHEP 10 (2022)	N ³ LL	×	~	507	1.12,	

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ART23 Moss, Scimemi, JHEP 05 (2024)	N ⁴ LL [–]	×		627	0.97	~

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HSO approach Aslan, Boglione, preprint: 2401.14266	NLO	×	Low Q only	130, 52	1.04, 1.68	×

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MAPTMD24?	N ³ LL	~	~	2031	?	~

and he that and more how management to like standing steen which the Music of the Rowsen R. solt. But the self-ining threes phints the average available, and Deschour done Processia. want chast, and conceptions, as house income and low provide and we to and the s

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MAPTMD24?

Study of flavor-dependent behavior of TMDs through a global fit



• The internal dynamics of quark in 3D can be studied in terms of transverse-momentum-dependent distributions (TMDs)

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- There are only few extractions of TMDs through global fits of Drell-Yan and SIDIS data (MAPTMD22: state-of-the-art)

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- SIDIS theoretical calculations are affected by a normalization issue (MAPTMD22: first attempt of work solution)

- The internal dynamics of quark in 3D can be studied in terms of transverse-momentum-dependent distributions (TMDs)
- There are only few extractions of TMDs through global fits of Drell-Yan and SIDIS data (MAPTMD22: state-of-the-art)
- SIDIS theoretical calculations are affected by a normalization issue (MAPTMD22: first attempt of work solution)

 Nowadays, studies on the extraction of flavor-dependent TMDs are in progress (up to now, only on Drell-Yan data)



Resummation of large logs

Resummation of large logs

$$S_{\text{pert}}(\mu_b, \mu) = 1 + \sum_{k=0}^{\infty} R_{N^k \text{LL}}$$

$$S_{\text{pert}}(\mu_b,\mu) = 1 + \sum_{k=0}^{\infty} \sum_{n=1+\lfloor k/2 \rfloor}^{\infty} \left(\frac{\alpha_S(\mu)}{4\pi}\right)^n \sum_{k=1}^{2n} L^{2n-k} R^{(n,2n-k)}$$

Resummation of large logs

$$S_{\text{pert}}(\mu_b, \mu) = 1 + \sum_{k=0}^{\infty} R_{N^k \text{LL}}$$
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Resummation of large logs

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Accuracy	H and C	K and γ_F	γκ	PDF and a_s evol.
LL	0	-	1	-
NLL	0	1	2	LO
NLL'	1	1	2	NLO
NNLL	1	2	3	NLO
NNLL'	2	2	3	NNLO
N ³ LL	2	3	4	NNLO
N ³ LL'	3	3	4	N ³ LO

Bacchetta, Bertone, Bissolotti, et al., JHEP 07 (2020)

TMD handbook, Boussarie, et al., 2023

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : C$$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : C$$
$$\mu_b = \frac{2e^{-\gamma_E}}{|\boldsymbol{b}_T|}$$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : C$$

$$\mu_b = \frac{2e^{-\gamma_E}}{|\boldsymbol{b}_T|} \quad \xrightarrow{b_T \gg 1} \quad 0$$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : C$$

 $\mu_b = \frac{2e^{-\gamma_E}}{|\mathbf{b}_T|} \xrightarrow{b_T \gg 1} 0 \quad \alpha_S(\mu_b) \to +\infty$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : C$$

$$\mu_b > \mu \qquad \infty \qquad \xleftarrow{b_T \ll 1} \qquad \mu_b = \frac{2e^{-\gamma_E}}{|\boldsymbol{b}_T|} \qquad \xleftarrow{b_T \gg 1} \qquad 0 \qquad \alpha_S(\mu_b) \to +\infty$$
$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : C$$

$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : C$$

 b_* -prescription



$$f_{NP}(x, b_T^2) \exp\left\{g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}}\right\} : C$$

$$b_{\min} = \frac{2e^{-\gamma E}}{\mu}$$

Collins, Soper, Sterman, Nucl. Phys. B250 (1985) Collins, Gamberg, et al., PRD (2016) Bacchetta, Echevarria, Mulders, et al., JHEP 11 (2015)



b

Bacchetta, Echevarria, Mulders, et al., JHEP 11 (2015)

$$\begin{split} f_{NP}(x, b_T^2) \exp \left\{ g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}} \right\} &: \mathcal{C} \\ \hline \mu_b > \mu & & \longleftarrow & \mu_b = \frac{2e^{-\gamma_E}}{|b_T|} & \xrightarrow{b_T \gg 1} & 0 & \alpha_S(\mu_b) \to +\infty \\ \mathbf{b}_* \text{-prescription} & & & \mathbf{b}_{\mathrm{max}} = 2e^{-\gamma_E} & & & \mathbf{b}_{\mathrm{max}} \\ b_{\mathrm{min}} = \frac{2e^{-\gamma_E}}{\mu} & & & \mathbf{b}_{\mathrm{min}} & & \mathbf{b}_{\mathrm{min}} \\ & & & \mathbf{collins}, \text{Soper, Sterman, Nucl. Phys. B250 (1985)} \\ & & & & \mathbf{collins}, \text{Gamberg, et al., PRD (2016)} \\ \end{split}$$

$$\hat{f}_1(x, b_T^2; \mu, \zeta) = \left[\frac{\hat{f}_1(x, b_T^2; \mu, \zeta)}{\hat{f}_1(x, b_*(b_T^2); \mu, \zeta)}\right] \hat{f}_1(x, b_*(b_T^2); \mu, \zeta) \equiv f_{\rm NP}(x, b_T^2; \zeta) \hat{f}_1(x, b_*(b_T^2); \mu, \zeta)$$

0.2

0

0

1

 $|b_T| \; [{
m GeV^{-1}}]$

 $b_{st}(b_T^2)$

3

 $|b_T|$

 $\mathbf{2}$

$$\begin{split} f_{NP}(x, b_T^2) \exp \left\{ g_K(b_T^2) \ln \frac{\sqrt{\zeta}}{\sqrt{\zeta_0}} \right\} &: \mathsf{C} \\ \hline \mu_b > \mu & \infty & \longleftarrow & \mu_b = \frac{2e^{-\gamma_E}}{|b_T|} & \stackrel{b_T \gg 1}{\longrightarrow} & 0 & \alpha_S(\mu_b) \to +\infty \\ \mathbf{b}_* \text{-prescription} \\ b_{\max} = 2e^{-\gamma_E} & & \mathbf{perturbative} \\ b_{\min} = \frac{2e^{-\gamma_E}}{\mu} & & \mathbf{perturbative} & \mathbf{Non-perturbative} \\ \mathbf{collins, Soper, Sterman, Nucl. Phys. B250 (1985)} & & \mathbf{colling} \\ \end{split}$$

Collins, Gamberg, et al., PRD (2016) Bacchetta, Echevarria, Mulders, et al., JHEP 11 (2015)

$$\hat{f}_{1}(x, b_{T}^{2}; \mu, \zeta) = \left[\frac{\hat{f}_{1}(x, b_{T}^{2}; \mu, \zeta)}{\hat{f}_{1}(x, b_{*}(b_{T}^{2}); \mu, \zeta)}\right] \hat{f}_{1}(x, b_{*}(b_{T}^{2}); \mu, \zeta) \equiv \underbrace{f_{\mathrm{NP}}(x, b_{T}^{2}; \zeta)}_{\mathrm{NP}(x, b_{T}^{2}; \zeta)} \hat{f}_{1}(x, b_{*}(b_{T}^{2}); \mu, \zeta)$$

0.2 -

 $b_{st}(b_T^2)$

 $|b_T|$

MAPTMD22 — Error analysis

Error propagation 250 Montecarlo replicas



MAPTMD22 — Error analysis

Error propagation 250 Montecarlo replicas **Correlation matrix** Hints of the appropriateness of the chosen functional form



Comparison with SV19

Scimemi, Vladimirov, arXiv:1912.06532

Drell-Yan

SIDIS



Comparison with SV19

Scimemi, Vladimirov, arXiv:1912.06532

Drell-Yan

SIDIS



Results obtained within the arTeMiDe framework

include (m/Q)include (M/Q)include (q_T/Q) in kinematics include (q_T/Q) in x_S, z_S

Scimemi, Vladimirov, arXiv:1912.06532

Results obtained within the arTeMiDe framework

#	
# #	 PARAMETERS OF TMDX-SIDIS
<i>"</i> *10	
*p1 T	initialize TMDX-SIDIS module
*A	Main definitions
*p1 L0	Order of coefficient function
*р2 Т	Use transverse momentum corrections in kinematics
*р3 Т	Use target mass corrections in kinematics
*p4 T	Use product mass corrections in kinematics
*р5 Т	Use transverse momentum corrections in x1 and z1

Scimemi, Vladimirov, arXiv:1912.06532

Results obtained within the arTeMiDe framework



Scimemi, Vladimirov, arXiv:1912.06532

Results obtained within the arTeMiDe framework



Results obtained within the arTeMiDe framework



This is NOT a constant factor