

SIGNATURES OF THE YANG-MILLS DECONFINEMENT TRANSITION FROM THE GLUON TWO-POINT CORRELATOR

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- Introduction
- Center Symmetry
- Landau gauge results
- Center-symmetric effective action
- Results
- Conclusion and Outlook

INTRODUCTION

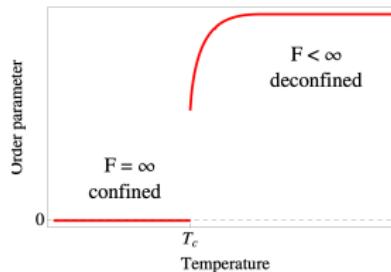
- Insight in the low-energy regime of QCD, especially a signal of the confinement/deconfinement transition.
- Most results are coming from non-perturbative, numerical (lattice) or semi-analytical (FRG and SD) methods.
- From this, we know that:
 - ▶ At some very high temperature T_c , hadrons become free quarks and gluons → quark-gluon plasma.
 - ▶ This transition is related to the breaking of the center symmetry of a gauge group when using pure Yang-Mills theory (infinitely heavy quarks):

$$\mathcal{L} = \frac{1}{4}(F_{\mu\nu}^a)^2$$

CENTER SYMMETRY

- An order parameter for the confinement/deconfinement transition is the Polyakov loop:

$$\mathcal{P} \propto \langle Pe^{i \int_0^\beta d\tau A_0(\tau, x)} \rangle \propto e^{-\beta F}$$



- In the confined phase, F is infinite $\rightarrow \mathcal{P} = 0$. In the deconfined phase, F is finite $\rightarrow \mathcal{P} \neq 0$.
- Under center symmetry $\mathcal{P} \rightarrow Z_N \mathcal{P}$, with Z_N the center elements of the gauge group. So, breaking of the center symmetry signals deconfinement.
- Confirmed by lattice data: second order transition for $SU(2)$, first order for $SU(3)$.

ENCODING OF THE TRANSITION

Polyakov loop:

$$\mathcal{P} \sim \langle Pe^{i \int_0^\beta d\tau A_0(\tau, x)} \rangle \sim e^{-\beta F}.$$

Because the Polyakov loop is related to A_0 , it is expected that the transition is **encoded** in (the tower of)

$$\langle A_0 \rangle, \langle A_0 A_0 \rangle, \dots, \langle A_0^n \rangle.$$

For the appropriate choice of gauge, can the transition be reflected in the **lowest order correlators**?

LANDAU GAUGE CORRELATOR

In principle:

- $\langle A \rangle$ is found by minimizing the effective action $\Gamma[A]$. It represents the state of the system . $\langle A_0 \rangle \rightarrow$ order parameter.
- The two-point correlator derives from the effective action

$$1 \left/ \frac{\partial^2 \Gamma}{\partial A^2} \right|_{A=\langle A \rangle} = \langle AA \rangle_c,$$

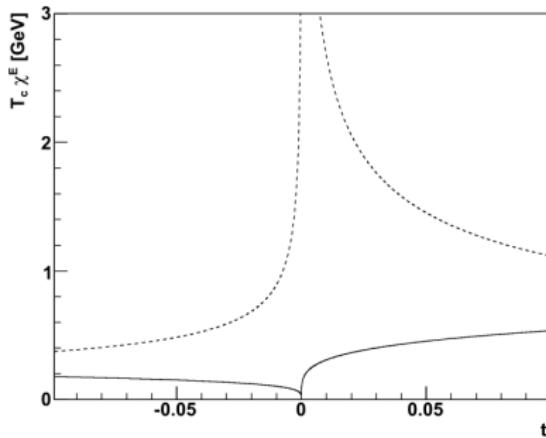
so for $SU(2)$, $\langle A_0 A_0 \rangle$ should diverge at T_c .

In practice:

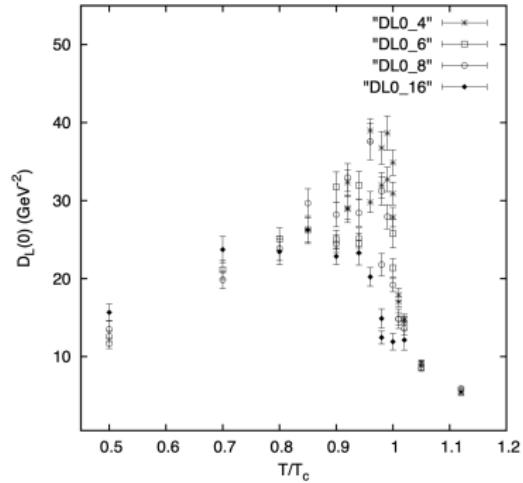
- In the Landau gauge, $\partial_\mu A_\mu = 0$, then $\langle A_0 \rangle = 0$. \rightarrow no order parameter.
- No evidence of divergence of $\langle A_0 A_0 \rangle$ was found on the (gauge-fixed) lattice and in the continuum.

$SU(2)$ LANDAU GAUGE CORRELATORS

Model of the susceptibility

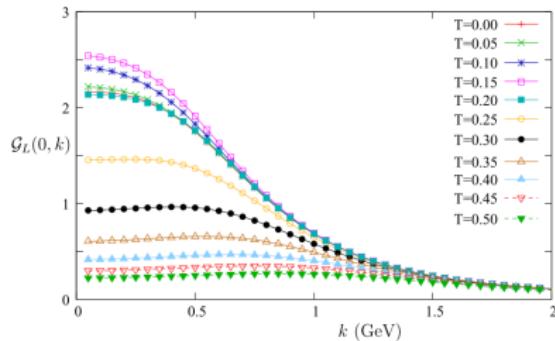


Electric susceptibility (zero momentum
longitudinal propagator)

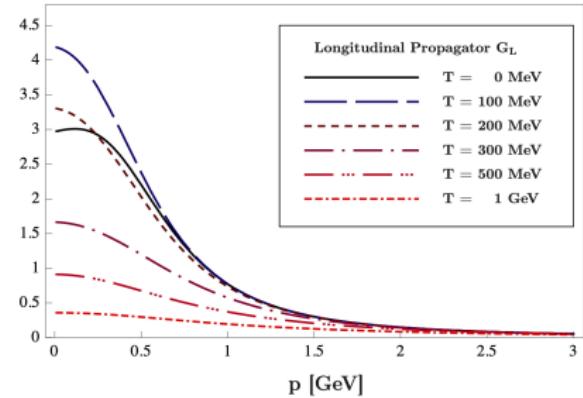


*T. Mendes and A. Cucchieri, PoS LATTICE2014, 183 (2015).

$SU(2)$ LANDAU GAUGE CORRELATORS



Longitudinal gluon propagator



*U. Reinosa, J. Serreau, M. Tissier, N. Wschebor, Phys. Rev. D, (2015)

*L. Fister and J. M. Pawłowski, [arXiv:1112.5440 [hep-ph]](2012).

BACKGROUND FIELD GAUGES

- A priori there is no reason to believe that the Landau gauge will provide the right environment to keep track of the centersymmetry breaking.
- In the *Landau gauge* the effective action is not explicitly center-symmetric $\Gamma[A] \neq \Gamma[A^U]$.
- This should not alter the physical results in principle, but it can lead to problems when approximations (loop calculations) are involved.
- To regain gauge invariance, one solution is to work with the **Background Field Gauges**.

BACKGROUND FIELD GAUGES

- Introducing a background field \bar{A} , the effective action is gauge-invariant $\Gamma_{\bar{A}}[A] = \Gamma_{\bar{A}^U}[A^U]$. Landau-deWitt gauge:

$$\bar{D}_\mu(A_\mu - \bar{A}_\mu) = 0, \text{ with } \bar{D}_\mu \equiv \partial_\mu - [\bar{A}_\mu, \cdot].$$

- In the **Background field effective action** one looks at $\tilde{\Gamma}[\bar{A}]$ by taking $\bar{A} = \langle A \rangle$.
 - The two-point function $\langle AA \rangle_c$ is not directly accessible from $\tilde{\Gamma}[\bar{A}]$,
 - Relies on the strict background independence of $\tilde{\Gamma}[\bar{A}]$, which is not easy to maintain in the presence of truncations.
- We propose the **Center-symmetric Landau gauge**, which fixes $\bar{A} = \bar{A}_c$. Then $\Gamma_{\bar{A}_c}[A] = \Gamma_{\bar{A}_c}[A^{U^c}]$.

OUR SETUP

- We work in the Landau-deWitt gauge with a background field \bar{A}_μ :

$$\bar{D}_\mu(A_\mu - \bar{A}_\mu) = 0, \text{ with } \bar{D}_\mu \equiv \partial_\mu - [\bar{A}_\mu, \cdot].$$

- We take \bar{A} and $\langle A \rangle$ in the temporal direction, $\propto \delta_{\mu 0}$, and along the diagonal color directions (σ^3 for $SU(2)$, (λ^3, λ^8) for $SU(3)$), so that $\Gamma[A, \bar{A}] \propto V(A, \bar{A})$. We write $\langle A \rangle = \delta_{\mu 0} \frac{T}{g} r \frac{\sigma^3}{2}$ and $\bar{A} = \delta_{\mu 0} \frac{T}{g} \bar{r} \frac{\sigma^3}{2}$.
- E.g. in $SU(2)$, under center transform $r \rightarrow 2\pi - r$ so centersymmetric value is $r_c = \pi$. For $SU(3)$, $r_c = (3/4\pi, 0)$.
- We fix $\bar{r} = r_c$: Center-symmetric Landau gauge. Center-symmetric phase when $r = r_c$, → order parameter.

CURCI-FERRARI MODEL

We have computed $\langle A \rangle$ and $\langle A(0, p)A(0, -p) \rangle$ up to first loop order in the finite temperature Curci-Ferrari model:

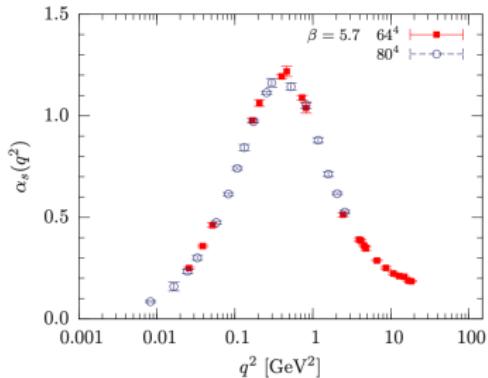
$$S = S_{YM} + S_{gf} + \int_{x,\tau} \frac{m^2}{2} (A_\mu^a - \bar{A}_\mu^a)^2$$

Several motivations:

- Perturbative gauge-fixed Yang-Mills **breaks down** at low energies (Landau pole, Gribov copies...), there is no analytical model for this region.
- A gluon mass term seems to dominate the (unknown) gauge-fixed action in the IR; decoupling behaviour on the lattice. CF could be an **effective model**.
- The CF model is renormalizable, avoids the Landau pole and lifts the degeneracy between Gribov copies.

CURCI-FERRARI MODEL

We use perturbation theory in the non-perturbative region. Actually, the coupling might not be that large in the deep IR.



*I.L. Bogolubsky, E.M. Ilgenfritz, M. Muller-Preussker, and A. Sternbeck. Phys.Lett., B676:69–73, 2009.

RESULTS - T_c (MeV)

	Lattice	FRG-BG ¹	CF-BG, 1-lp ²	CF-BG, 2-lp ³	CF-CS, 1-lp ⁴
SU(2)	295	230	238	284	265
SU(3)	270	275	185	254	267

BG: Background effective action

CS: Centersymmetric Landau gauge

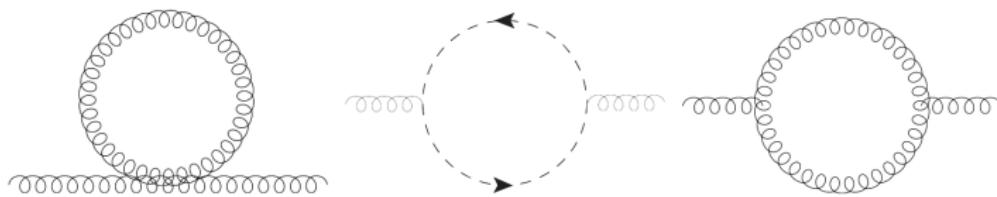
¹L. Fister and J. M. Pawłowski, Phys.Rev. D88 (2013) 045010

²U. Reinosa, J. Serreau, M. Tissier and N. Wschebor, Phys.Lett. B742 (2015) 61-68.

³U. Reinosa, J. Serreau, M. Tissier and N. Wschebor, Phys.Rev. D93 (2016) 105002.

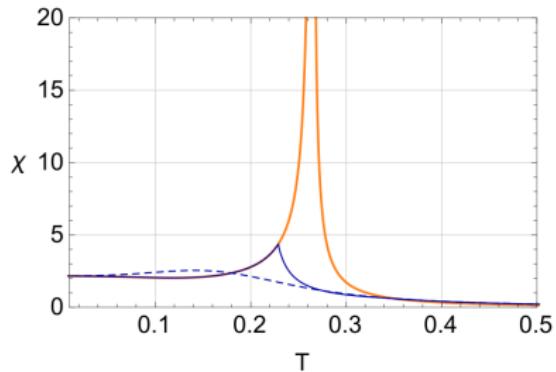
⁴DvE, U. Reinosa, J. Serreau and M. Tissier, SciPost Phys. **12**, 087 (2022)

FEYNMAN DIAGRAMS GLUON PROPAGATOR



- Calculated in finite temperature trough Matsubara techniques:
 $\int d^d Q \rightarrow T \sum_q \int d^{d-1} q$
- We calculated the spatial integral with Feynman techniques, the Matsubara sums numerically.

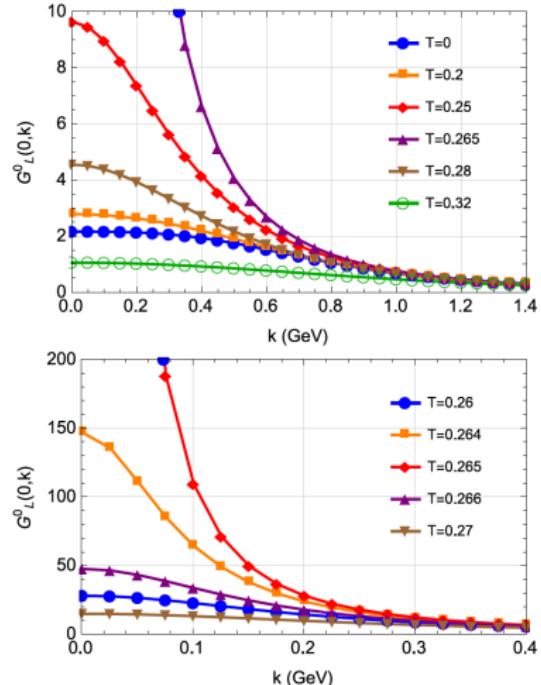
RESULTS: SU(2) GLUON PROPAGATOR



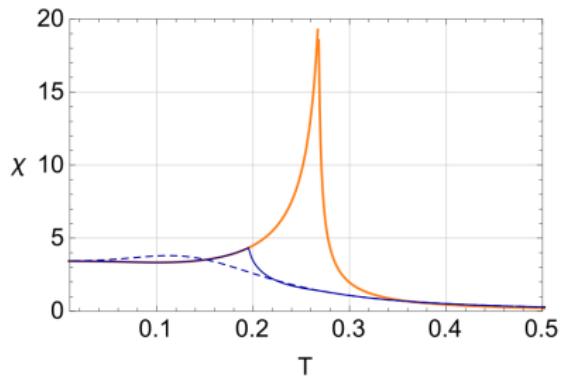
Landau gauge
Background field effective action
Centersymmetric Landau gauge

$$m=0.68 \text{ GeV}, \mu=1 \text{ GeV}, g=7.5$$

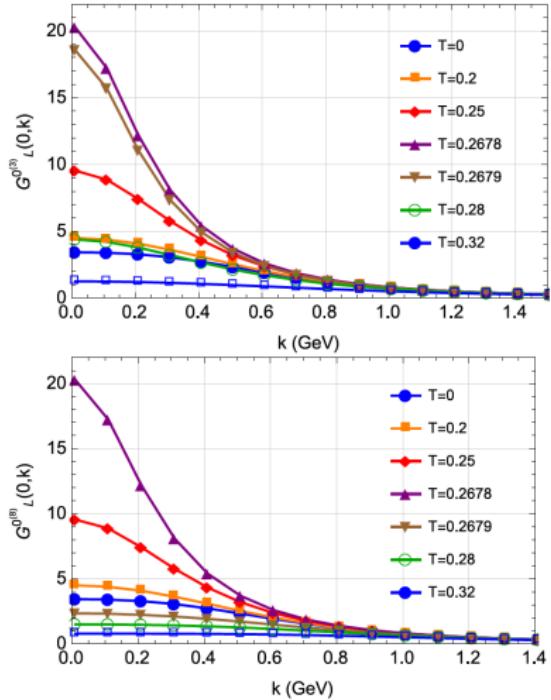
*DvE, U. Reinosa, J. Serreau and M. Tissier, SciPost Phys. 12, 087 (2022).



RESULTS: SU(3) GLUON PROPAGATOR



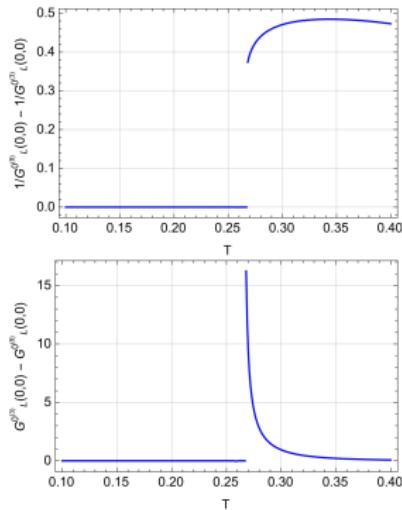
Landau gauge
Background field effective action
Centersymmetric Landau gauge



$$m=0.54 \text{ GeV}, \mu=1 \text{ GeV}, g=4.9$$

*DvE, U. Reinosa, J. Serreau and M. Tissier, SciPost Phys. 12, 087 (2022).

SU(3) PROPAGATOR DIFFERENCE

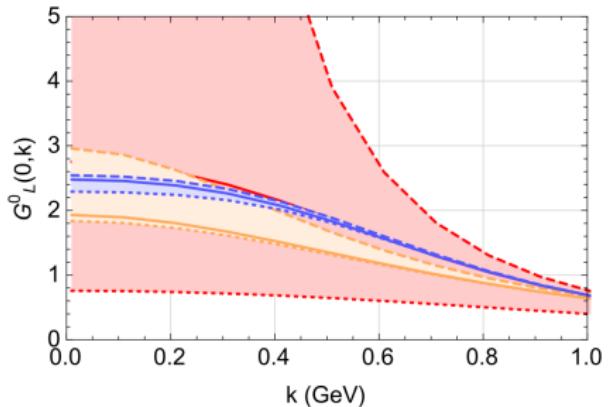


Confirmed by gauge-fixed lattice data, in collaboration with O. Oliveira and P. Silva.

CONCLUSION AND OUTLOOK

- We have performed, for the first time, calculations of the gluon one- and two-point correlator in the centersymmetric Landau gauge.
- We find a good agreement with lattice data for T_c .
- We find that for $SU(2)$, the deconfinement transition is signaled by a **divergence** of the longitudinal gluon propagator for $k \rightarrow 0$.
- For $SU(3)$, the difference between the propagators in the neutral color mode is an order parameter for the transition.
- Ideas for future works: RG improvement, transversal propagator and dynamically generated mass [with D. Dudal and D. Vercauteren].

RESULTS: SU(2) GLUON PROPAGATOR

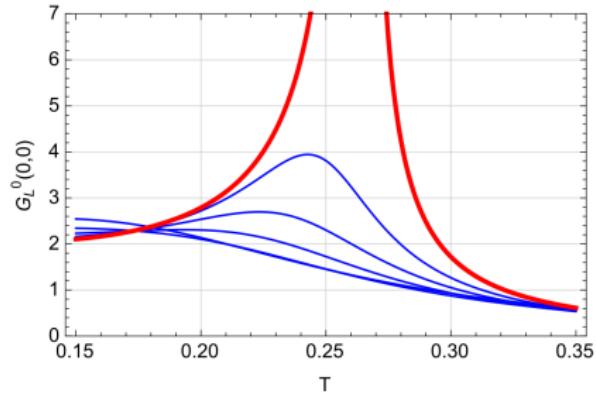


Landau gauge

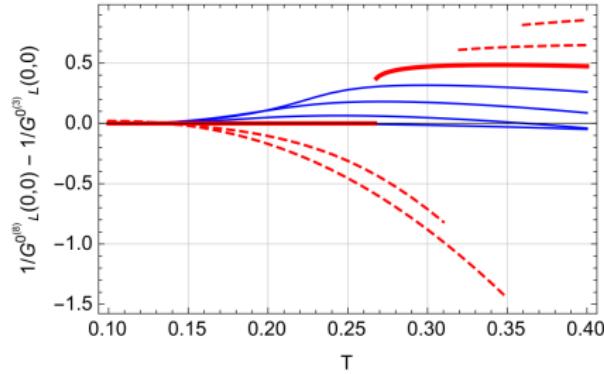
Background Field Effective action

Center-symmetric Landau Gauge

RESULTS: SU(2) GLUON PROPAGATOR



SU(3) PROPAGATOR DIFFERENCE



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