Fully coherent energy loss in proton-nucleus collisions

#### François Arleo

Subatech, Nantes

Cold Nuclear Matter Effects: from the LHC to the EIC

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### Collinear factorization in pA collisions

A nucleus as an ordinary hadron

$$\frac{\mathrm{d}\sigma_{\mathsf{p}\mathsf{A}}}{\mathrm{d}y\,\mathrm{d}Q} = \sum_{i,j} \int \mathrm{d}x_1 \ f_i^{\mathsf{p}}(x_1,\mu) \int \mathrm{d}x_2 \ \frac{f_j^{\mathsf{A}}(x_2,\mu)}{\mathrm{d}y\,\mathrm{d}Q} + \mathcal{O}\left(\frac{\Lambda_{\mathsf{A}}^n}{Q^n}\right)$$

- Universal (leading twist) nuclear PDF
  - $\triangleright$  could be probed in various processes and collision systems (eA,  $\gamma$ A, pA)
- New scale for power corrections  $(\Lambda_{_{\rm A}} > \Lambda_{_{\rm p}})$ 
  - higher twist processes enhanced in pA collisions (wrt pp)
  - specific processes could spoil the extraction of (universal) nPDF

What to expect for  $f^A$ ? How does it compare to  $f^p$ ?

## nPDF

- Parton distribution functions are modified in nuclei
  - $\blacktriangleright$  Evidence at large x from EMC/NMC measurements in DIS
- Extracted from data global fits
  - nuclear DIS: structure functions  $F_2$
  - ▶ pA collisions: DY, W/Z, jets, hadrons ( $\pi$  at RHIC, D at LHC)
- Leads to  $x_2$  scaling for hadron suppression:  $R_{pA}(x_2, \sqrt{s}) = R_{pA}(x_2)$
- Expected shadowing at small x<sub>2</sub>
  - Leads to hadron suppression at high  $\sqrt{s}$  and large y
  - Amount of shadowing poorly known due to lack of data



Ideally, looking for processes sensitive to PDF only Some requirements (not necessary, but preferable):

- Sufficiently large scale:  $Q \gg Q_s \simeq$  few GeV
  - to ensure factorization to be at work
- ... but not too large to keep some sensitivity
  - +  $f^A/Af^p \simeq 1$  in the Bjorken limit ( $Q^2 \rightarrow \infty$  at fixed x)
- Favor color-neutral probes
  - avoid energy loss effects

Best candidates

Weak bosons, Jets, Drell-Yan

# nNNPDF3.0 (w/ and w/o LHCb D meson data)



- Strong constraints given by forward D-meson data nNNPDF, 2201.12363
- Several other attempts

EPPS21 2112.12462, nCTEQ 2204.09982

Talks by Kusina and Paukkunen, Tue 14

### Nuclear effects beyond nPDF

- Strong light hadron and J/ $\psi$  suppression observed at forward rapidity
- $\bullet$  Smaller suppression in the  $\Upsilon$  channel & much smaller for Drell-Yan



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#### Strong violation of $x_2$ scaling

#### Radiative energy loss in cold nuclear matter

Induced gluon radiation could affect many hard QCD processes!

- Drell-Yan in pA collisions
  - initial-state energy loss
- Hadron production in semi-inclusive DIS  $\gamma^* A \rightarrow h + X$ 
  - final-state energy loss



- Hadron production in pA collisions
  - initial & final state energy loss...and initial-final state interference

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Bow does energy loss depend parametrically on parton energy ?

#### Fully coherent induced gluon spectrum



#### Fully coherent induced gluon spectrum

$$\omega \frac{dI}{d\omega}\Big|_{1\to 1} = \frac{\alpha_s}{\pi} F_c \ln\left(1 + \frac{\hat{q}L}{M_{\perp}^2} \frac{E^2}{\omega^2}\right)$$

FA Peigné Sami 1006.0818 FA Peigné 1212.0434

FA Kolevatov Peigné 1402.1671 Peigné Kolevatov 1405.4241

- First determined in a model, later confirmed in the opacity expansion
- Color factor follows from color algebra:  $F_c = C_R + C_{R'} C_t$ 
  - R(R') = color rep. of the incoming (outgoing) particle

$$\begin{array}{rcl} g \to g & : & F_c = N_c + N_c - N_c = N_c \\ q \to g & : & F_c = C_F + N_c - C_F = N_c \\ q \to q & : & F_c = C_F + C_F - N_c = -1/N_c \quad (<0 \ !) \end{array}$$

• Average energy loss proportional to E !  $\Delta E_{\rm FCEL} \propto \alpha_s \ F_c \ \frac{\sqrt{\hat{q} L}}{M_1} \ E$ 

#### In a nutshell

LPM energy loss (small formation time  $t_f \leq L$ )

$$\Delta E_{_{
m LPM}} \propto lpha_{s} \hat{q} L^{2}$$

- Hadron production in nuclear DIS or Drell-Yan in pA collisions
- Particle suddenly accelerated (e.g. jet in QGP)

Fully coherent energy loss (large formation time  $t_f \gg L$ )

$$\Delta E_{\rm FCEL} \propto \alpha_s \ F_c \ \frac{\sqrt{\hat{q} \ L}}{M_{\perp}} \ E \quad (\gg \Delta E_{\rm LPM})$$

Important for hadron production in pA collisions at all energies

- Needs color in both initial & final state
  - ${\scriptstyle \blacktriangleright}$  no effect on W/Z nor Drell-Yan, no effect in DIS
- Power suppressed: negligible when  $M_{\perp} \gg \sqrt{\hat{q} \, L} \sim Q_s$

#### Goal

- Explore phenomenological consequences of coherent energy loss
- Approach as simple as possible with the least number of assumptions
- Observable: quarkonium suppression in pA (and AA) collisions

Physical picture and assumptions



- Color neutralization happens on long time scales:  $t_{\text{octet}} \gg t_{\text{hard}}$
- In-medium rescatterings do not resolve the octet Q ar Q pair
- Hadronization happens outside of the nucleus:  $t_\psi\gtrsim L$

### Model

#### Energy shift

$$\frac{1}{A}\frac{\mathrm{d}\sigma_{\mathrm{pA}}^{\psi}}{\mathrm{d}E}\left(E,\sqrt{s}\right) = \int_{0}^{\varepsilon_{\mathrm{max}}} \mathrm{d}\varepsilon \,\mathcal{P}(\varepsilon,E) \,\frac{\mathrm{d}\sigma_{\mathrm{pp}}^{\psi}}{\mathrm{d}E}\left(E+\varepsilon,\sqrt{s}\right)$$

- pp cross section fitted experimental data
- $\mathcal{P}(\epsilon):$  quenching weight related to the induced gluon spectrum

$$\mathsf{P}(\epsilon) \simeq rac{dI(\epsilon)}{d\omega} \exp\left\{-\int_{\epsilon}^{\infty} d\omega rac{dI}{d\omega}
ight\}$$

- Length L given by Glauber model
- Transport coefficient

$$\hat{q}(x) = \frac{4\pi^2 \alpha_s C_R}{N_c^2 - 1} \rho x G(x) = \hat{q}_0 \left(\frac{10^{-2}}{x}\right)^{0.3} \quad ; \ \hat{q}_0 = 0.05 - 0.09 \text{ GeV}^2/\text{fm}$$

▶  $\hat{q}_0$  range in agreement with LPM energy loss and nuclear broadening studies, corresponds to  $Q_s \simeq 1.3$ –1.8 GeV at LHC at mid-rapidity

## LHC predictions

FA Peigné 1204.4609 1212.0434 FA Kolevatov Peigné Rustamova 2003.06337



• Moderate effects ( $\sim$  20%) around mid-rapidity, smaller at y < 0

- Large effects above  $y \gtrsim 2-3$
- Very good agreement with LHC data as well with all other  $J/\psi$  data at lower energy (NA3, E866, PHENIX...)

François Arleo (Subatech)

## From quarkonium to light hadron production

Which differences from quarkonium to single light hadron production?



- Final state made of two partons at leading order
  - Partons produced with opposite and large transverse momenta  $K_1 \simeq K_2 \gg \sqrt{\hat{q}L}$  and energy fractions  $\xi$  and  $1 \xi$
- Use medium-induced gluon spectrum associated to  $2 \rightarrow 2$  scattering
  - Final state in different color representations R with probability  $\rho_{\rm R}(\xi)$
- Hadronization:  $z \neq 1$

### From quarkonium to light hadron production

Which differences from quarkonium to single light hadron production?

$$p = \underbrace{\sum_{a} \underbrace{\sum_{b} \underbrace{k_{2}}{}}_{A \times \underbrace{b}} \underbrace{\sum_{b} \underbrace{k_{2}}{}}_{C} \underbrace{k_{1} \times \underbrace{k_{2}}{}}_{K_{1} \times \underbrace{k_{2}}{}} \underbrace{k_{2} \times -K_{1}}_{L}$$

Assuming a smooth variation of  $\rho$  and  $R_{pA}^{h}$  with  $\xi$ 

$$\begin{split} R_{\rm pA}^{\rm h}(y,p_{\perp}) &\simeq \sum_{\rm R} \rho_{\rm R}(\xi) \, R_{\rm pA}^{\rm R}(y,p_{\perp}) \\ R_{\rm pA}^{\rm R}(y,p_{\perp}) &= \int_{0}^{\delta_{\rm max}} {\rm d}\delta \, \hat{\mathcal{P}}_{\rm R}\left(x,\frac{\sqrt{\hat{q}L}\langle z\rangle}{M_{\xi}}\right) \, \frac{{\rm d}\sigma_{pp}^{\rm h}(y+\delta,p_{\perp})}{{\rm d}y \, {\rm d}p_{\perp}} \Big/ \frac{{\rm d}\sigma_{pp}^{\rm h}(y,p_{\perp})}{{\rm d}y \, {\rm d}p_{\perp}} \end{split}$$

# Color dependence

FA Peigné 2003.01987 FA Cougoulic Peigné 2003.06337



- Rapidity dependence reminiscent of quarkonium suppression
- Significant suppression, especially in the 27 color state
- Color-averaged suppression similar to that of an octet
- Effects weaken at large  $p_{\perp}$

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## Predictions at LHC



- Significant effects
  - More pronounced at larger y (measurable e.g. by LHCb)
  - Persists up to  $p_{\perp} \simeq 10 \,\, {
    m GeV}$
- All scattering processes can be computed (here most important ones)
- Similar in magnitude to saturation/nPDF effects

#### Comparison to data



• Precise baseline in agreement with ALICE  $\pi^{\pm}/K^{\pm}$  & CMS h<sup>±</sup> data

- brings constraints on other physical effects
- disagreement with  $p/\bar{p}$  data

## From light hadrons to D-mesons

Above model can be simply extended to the case of D mesons

FA Jackson Peigné 2107.05871



- Accounts for half of the suppression
  - Room for nPDF effects
- Small relative uncertainty

### Prompt photons

- Prompt photons in pp/pA collisions have long been seen as a good probes of gluon (n)PDF
- However... photons also affected by FCEL due to recoiling jet

FA, Bourgeais, Jackson, in progress

At leading order

- Compton scattering  $(qg \rightarrow \gamma q)$  color triplet final state
- Annihilation  $(q\bar{q} 
  ightarrow \gamma g)$  , color octet final state



- FCEL affects D-meson production in pA collisions...but nPDF are extracted using D-meson data ignoring FCEL
- How to extract nPDF reliably given FCEL?
  - $\scriptstyle \bullet$  Use color neutral probes in pA collisions: DY, W/Z
  - Use large- $Q^2$  measurements: jets, top quarks
  - Use DIS data
  - ... or include FCEL in nPDF global fits

## Reweighting nPDF, w/ and w/o FCEL

Given a new data set, PDF can be conveniently reweighted X Ignore FCEL :

 $\mathcal{P}(\mathbf{f}_{\mathsf{A}} | \, \mathsf{pQCD} \cap \mathsf{world} \,\, \mathsf{data})$ 

- Statistically good' fits can be obtained, including LHCb data
- Strong constraints... but unreliable result
- Include FCEL

 $\mathcal{P}(\mathbf{f}'_{\mathsf{A}}|\operatorname{pQCD}\cap\operatorname{\mathsf{FCEL}}\cap\operatorname{world}\operatorname{\mathsf{data}})$ 

- Part of the nuclear dependence cannot be attributed to nPDF
- Different physical processes with different scaling properties
- Resulting nPDF extracted from data will not be the same:  $f'_A \neq f_A$

## Reweighting nPDF, w/ and w/o FCEL

Given a new data set, PDF can be conveniently reweighted



 $\mathcal{P}(f'_{\mathsf{A}} | \mathsf{FCEL} \cap \mathsf{LHCb} \mathsf{ data})$ 

 $\mathcal{P}(f_{\mathsf{A}} \mid \mathsf{no} \mathsf{FCEL} \cap \mathsf{LHCb} \mathsf{ data})$ 

 $f'_{A} \neq f_{A}$ 

## A first attempt

Reweighting nPDF w/ and w/o FCEL using J/ $\!\psi$  LHC data



FA Avez Hard Probes 2023 & work in preparation

- Large difference between nPDF alone and nPDF + FCEL
- Similar results for all nPDF sets
- Lesser gluon shadowing when FCEL is included

#### Atmospheric neutrinos

- HE cosmic rays (protons) collide on  $\langle A \rangle \simeq 14.5 \Rightarrow$  Air shower
  - Similar to pA collisions at LHC
- D-meson decay responsible for the atmospheric  $\nu$  flux

Lipari 1993, Gondolo Ingelman Thunman 1996



## Depletion of atmospheric neutrinos from FCEL

$$R_{\nu} \equiv \Phi_{\nu}^{\mathsf{FCEL}} / \Phi_{\nu}^{\mathsf{FCEL}} \simeq \int_{0}^{1} \mathrm{d}z \, z^{\gamma} P(z)$$
 FA Jackson Peigné 2112.10791  

$$x_{2} \qquad x_{2} \qquad x_{2} \qquad x_{2} \qquad x_{2} \qquad x_{2} \qquad x_{3} \qquad x_$$

- Sizeable depletion of atmospheric  $\nu$ 's
- Similar in size to nPDF / saturation effects
- Reduced background for searches of astrophysical  $\nu$ 's

François Arleo (Subatech)

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#### • FCEL predicted from first principles with small uncertainty

- ► Δ∝E
- Leads to QCD factorization breaking

#### • Affects significantly hadron production in pA collisions

- Applied to light hadrons, D-mesons, quarkonia, prompt photons
- ... and atmospheric neutrino flux through D-meson decays

#### • Crucial impact on nPDF extractions

- D-meson/quarkonium data cannot be used in global fits, unless FCEL is properly taken into account
- Electron-ion collider best place to probe nPDF