

# Holographic QCD methods for quarkonium production

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This talk is based on:

- **1910.04707** (PRD) (with Ismail Zahed)
- **2106.00722** (PRD) (with Ismail Zahed)
- **2204.08857** (PRD) (with Ismail Zahed)
- **2109.03103** (PRD) (with Yizhuang Liu, Maciej A. Nowak, Ismail Zahed)
- **2401.12162** (PRD) (with Florian Hechenberger, Ismail Zahed)

- The AdS/CFT correspondence can be used to compute correlation functions of local operators [Maldacena (1998); Gubser, Klebanov, Polyakov (1998); Witten (1998)]:

$$Z_{\text{gauge}}(J\mathcal{O}, N_c, \lambda) \equiv Z_{\text{gravity}}(\phi_0, g_5, \alpha'/R^2), \quad \text{where } J \equiv \phi_0.$$

- Correlation functions are evaluated via Witten diagrams in AdS.
- For non-conformal theories with a mass gap (dual to a deformed AdS background), scattering amplitudes can likewise be computed using these Witten diagrams in AdS

$$ds^2 = \frac{R^2}{z^2} (\eta_{\mu\nu} dx^\mu dx^\nu - dz^2), \quad \eta_{\mu\nu} = \text{diag}(1, -1, -1, -1),$$

with  $0 \leq z \leq \infty$ , connects the UV boundary ( $z \rightarrow 0$ ) to the IR ( $z \rightarrow \infty$ ), and mass gap/confinement induced by a background dilaton field  $\phi(z) = \kappa^2 z^2$ .

# Two-point function

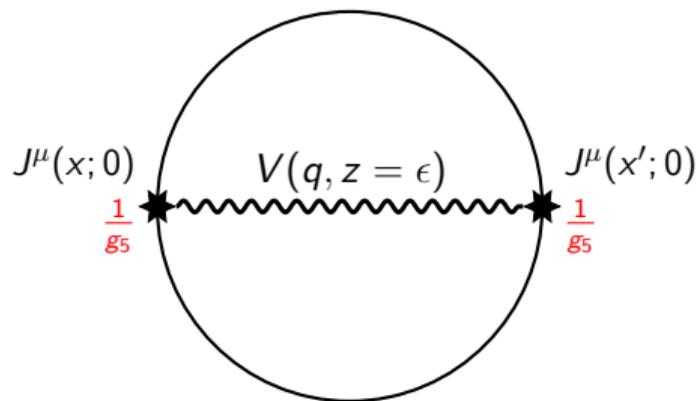


Figure: Witten diagram for the two-point function of the current operator with  $g_5^2 \sim \frac{1}{N_c}$ .

- The bulk-to-boundary propagator for the virtual photon is

$$\mathcal{V}(Q, z) = g_5 \sum_n \frac{F_n \phi_n(z)}{Q^2 + m_n^2} = \Gamma\left(1 + \frac{Q^2}{4\kappa^2}\right) \kappa^2 z^2 \mathcal{U}\left(1 + \frac{Q^2}{4\kappa^2}; 2; \kappa^2 z^2\right),$$

where  $\Gamma$  is the Gamma function and  $\mathcal{U}$  is the Tricomi confluent hypergeometric function.

# Spin-1 (Electromagnetic) Form Factors of Proton

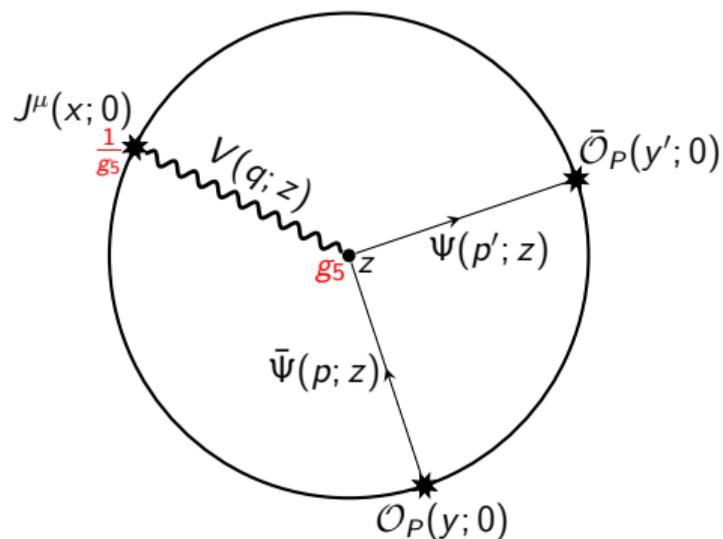


Figure: Witten diagram for EM form factor due to the exchange of vector mesons with  $g_5^2 \sim \frac{1}{N_c}$ .

# Spin-1 (Electromagnetic) Form Factors of Proton

- The scattering amplitude in AdS is

$$S_{Dirac}^{EM}[i, f] = (2\pi)^4 \delta^4(p' - p - q) \frac{1}{g_5} \times g_5 \bar{u}_{s_f}(p') \epsilon_\mu(q) \gamma^\mu u_{s_i}(p) \\ \times \frac{1}{2} \int \frac{dz}{z^{2M}} e^{-\phi} \mathcal{V}(Q, z) (\psi_L^2(z) + \psi_R^2(z)).$$

# Spin-2/0 (Gravitational) Form Factors of Proton

- The gravitational form factors (GFFs) of the proton are defined via the energy-momentum tensor (EMT):

$$\langle p_2 | T^{\mu\nu}(0) | p_1 \rangle = \bar{u}(p_2) \left( A(k) \gamma^{(\mu} p^{\nu)} + B(k) \frac{i p^{(\mu} \sigma^{\nu)\alpha} k_\alpha}{2m_N} + C(k) \frac{k^\mu k^\nu - \eta^{\mu\nu} k^2}{m_N} \right) u(p_1),$$

with  $k = p_2 - p_1$ . Often one writes  $D(k) \equiv 4 C(k)$ .

- The bulk metric fluctuations decompose into spin-2 (transvers-traceless part  $h$ ) and spin-0 (traceful part  $f$ ) [Kanitscheider (2008)]:

$$h_{\mu\nu}(k, z) \supset \left[ \epsilon_{\mu\nu}^{TT} h(k, z) \right] + \left[ \frac{1}{3} \eta_{\mu\nu} f(k, z) \right].$$

- For non-degenerate  $2^{++}$  and  $0^{++}$  glueball spectra, the holographic coupling includes both transverse-traceless (spin-2) and scalar (spin-0) fluctuations, respectively.

# Spin-2/0 (Gravitational) Form Factors of Proton

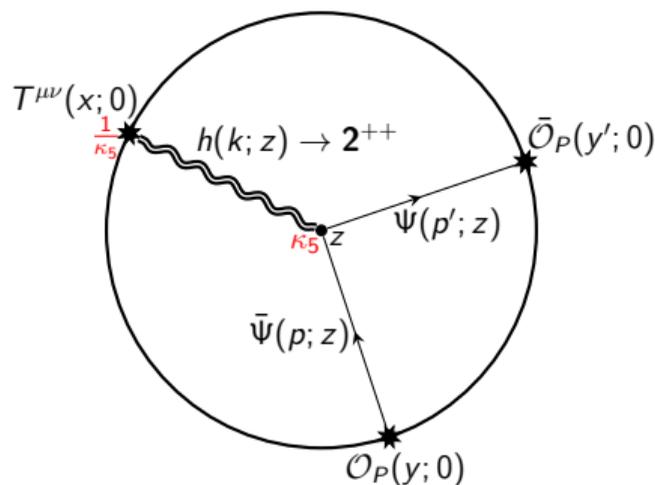
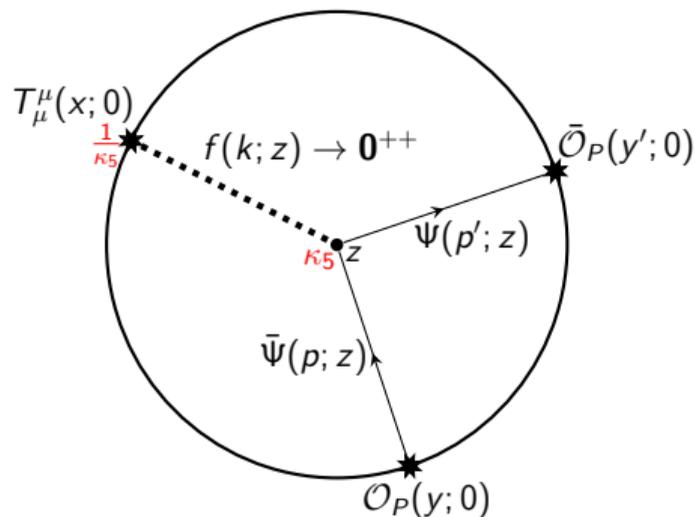


Figure: Witten diagram for spin-2 gravitational form factor due to the exchange  $2^{++}$  glueballs with  $\kappa_5^2 \sim \frac{1}{N_c^2}$ .

$$A(K, \kappa_T) = \frac{1}{2} \int dz \sqrt{g} e^{-\phi} z (\psi_R^2(z) + \psi_L^2(z)) \sum_{n=0}^{\infty} \frac{\sqrt{2} \kappa_5 F_n \psi_n(z)}{K^2 + m_n^2}.$$

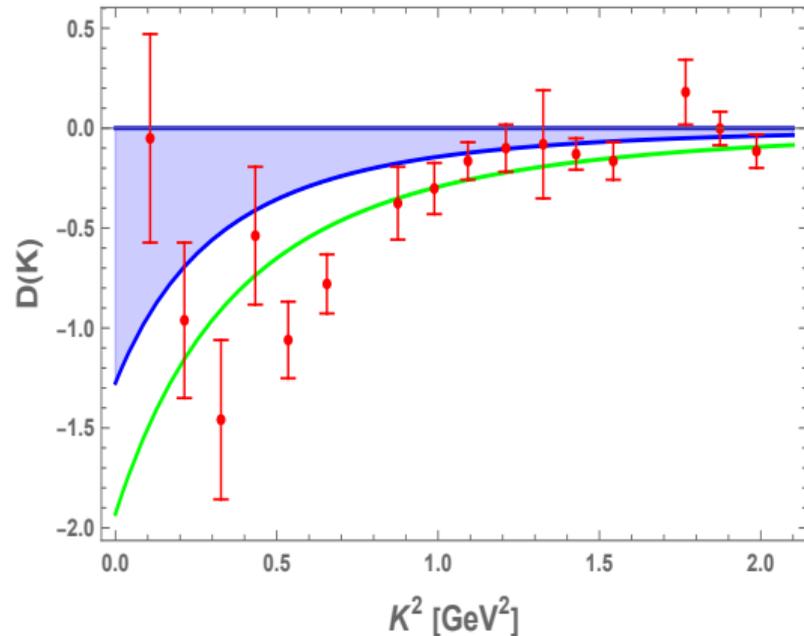
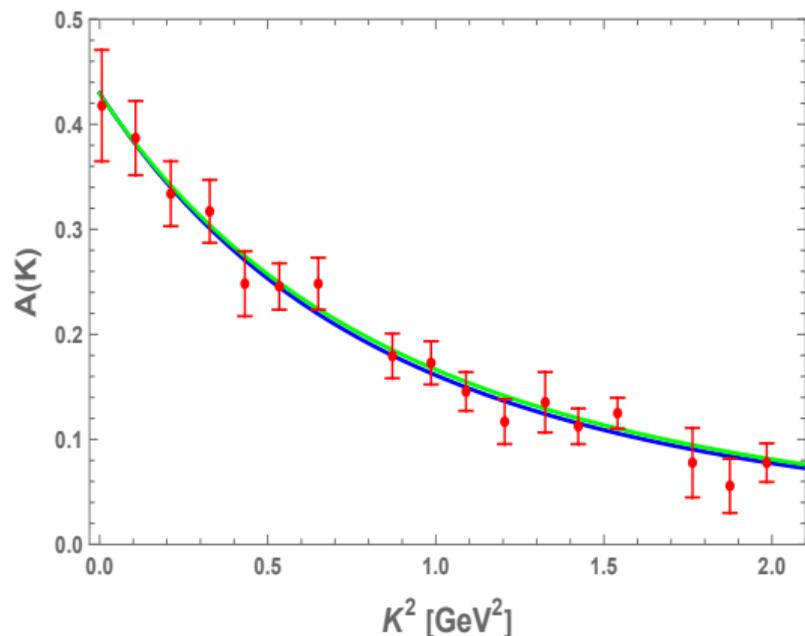
# Holographic Gravitational Form Factors



**Figure:** Witten diagram for scalar gravitational form factor due to the exchange  $0^{++}$  glueballs.

$$D(K, \kappa_T, \kappa_S) = -\frac{4 m_N^2}{3 K^2} \left[ A(K, \kappa_T) - A_S(K, \kappa_S) \right],$$

# Comparison with Lattice Data



Recent lattice QCD results [Pefkou:2021] (red points) compared to holographic fits (blue curves) with  $\kappa_T = 0.388$  GeV,  $\kappa_S = 0.217$  GeV. The green line is a tripole fit to the same lattice data.

# Photoproduction of Heavy Mesons Near Threshold

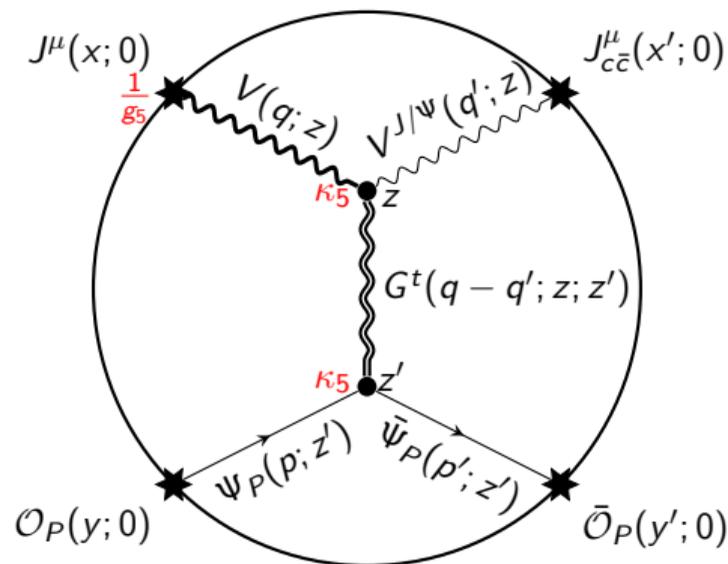


Figure: Witten diagrams for the holographic photo/electroproduction of  $J/\psi$  with  $g_5^2 \sim \frac{1}{N_c}$  and  $\kappa_5^2 \sim \frac{1}{N_c^2}$ .

# Photoproduction of heavy mesons near threshold

- the differential cross section for photoproduction of heavy vector mesons ( $J/\psi$  or  $\Upsilon$ ), near threshold, is given by

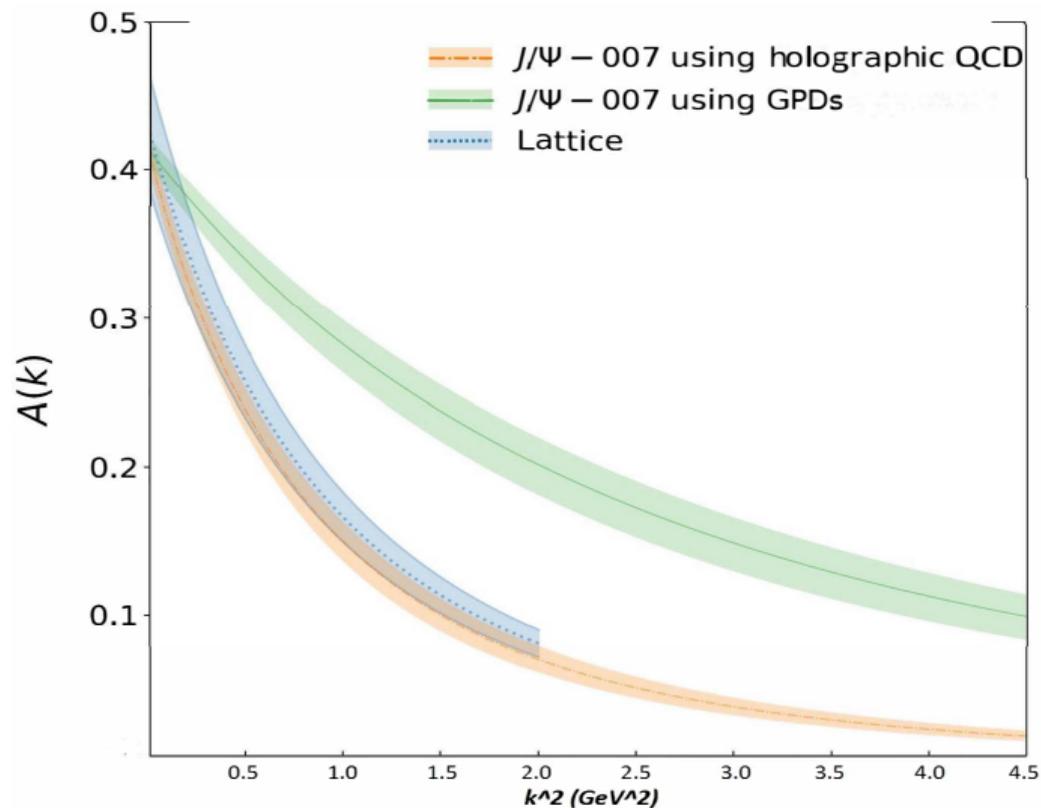
$$\begin{aligned} \frac{d\sigma}{dt} &= \mathcal{N}^2 \times [A(t) + \eta^2 D(t)]^2 \\ &\times \frac{1}{A^2(0)} \times \frac{1}{32\pi(s - m_N^2)^2} \times F(s, t, M_V, m_N) \times \left(1 - \frac{t}{4m_N^2}\right) \end{aligned}$$

with the normalization factor  $\mathcal{N}$  defined as

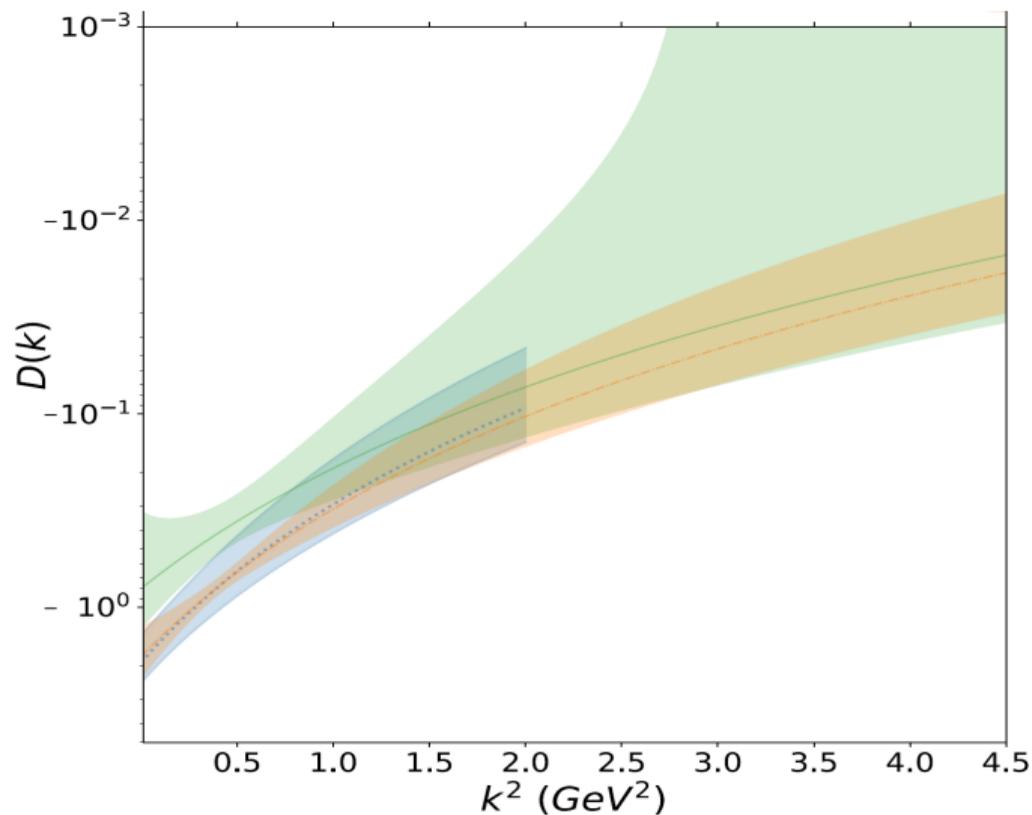
$$\mathcal{N}^2 \equiv e^2 \times \left(\frac{f_V}{M_V}\right)^2 \times \mathbb{V}_{h\gamma^* J/\psi}^2 \times (2\kappa_5^2)^2 \times A^2(0) = 7.768^2 \text{ nb/GeV}^6$$

- note that  $F(s, t) \sim s^4 \sim 1/\eta^4$  with the amplitude  $\mathcal{A} \sim s^2 \times A(t) + s^0 \times D(t)$  as expected from  $2^{++}$  and  $0^{++}$  glueball t-channel exchanges

# Extraction of the $2^{++}$ glueball contribution



# Extraction of the $0^{++}$ glueball contribution



# Electroproduction of heavy mesons near threshold

- the differential cross section for electroproduction of heavy vector mesons ( $J/\psi$  or  $\Upsilon$ ), near threshold, is given by

$$\frac{d\sigma(s, t, Q, M_{J/\psi}, \epsilon_T, \epsilon'_T)}{dt} \propto \mathcal{I}^2(Q, M_{J/\psi}) \times \left(\frac{s}{\kappa^2}\right)^2 \times [A(t) + \eta^2 D(t)]^2$$
$$\frac{d\sigma(s, t, Q, M_{J/\psi}, \epsilon_L, \epsilon'_L)}{dt} \propto \frac{1}{9} \times \frac{Q^2}{M_{J/\psi}^2} \times \mathcal{I}^2(Q, M_{J/\psi}) \times \left(\frac{s}{\kappa^2}\right)^2 \times [A(t) + \eta^2 D(t)]^2$$

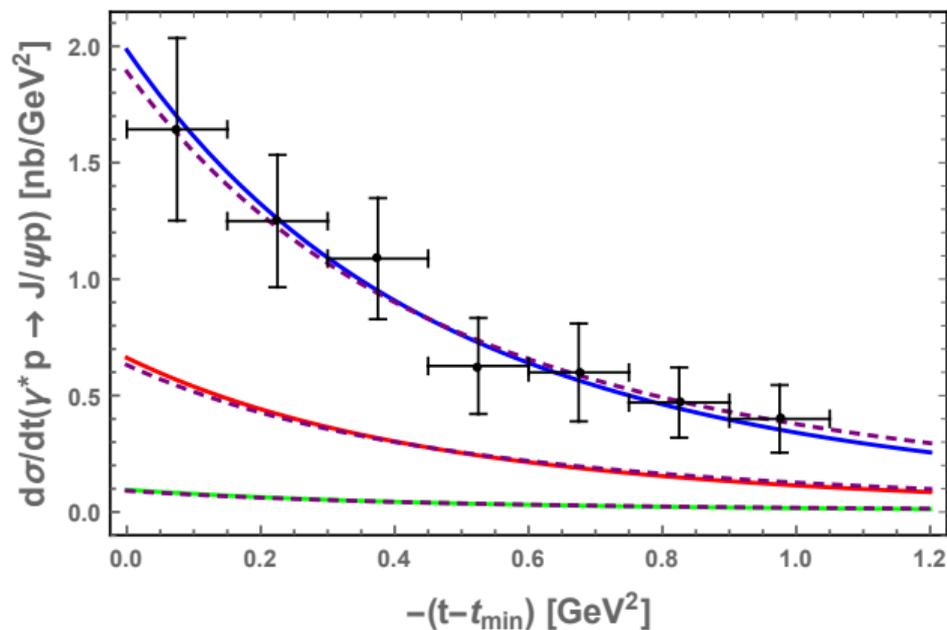
# Electroproduction of heavy mesons near threshold

- where we defined the transition form factor that controls the  $Q$  dependence as

$$\mathcal{I}(Q, M_{J/\psi}) = \frac{\mathcal{I}(0, M_{J/\psi})}{\frac{1}{6} \times \left( \frac{Q^2}{4\kappa_{J/\psi}^2} + 3 \right) \left( \frac{Q^2}{4\kappa_{J/\psi}^2} + 2 \right) \left( \frac{Q^2}{4\kappa_{J/\psi}^2} + 1 \right)},$$

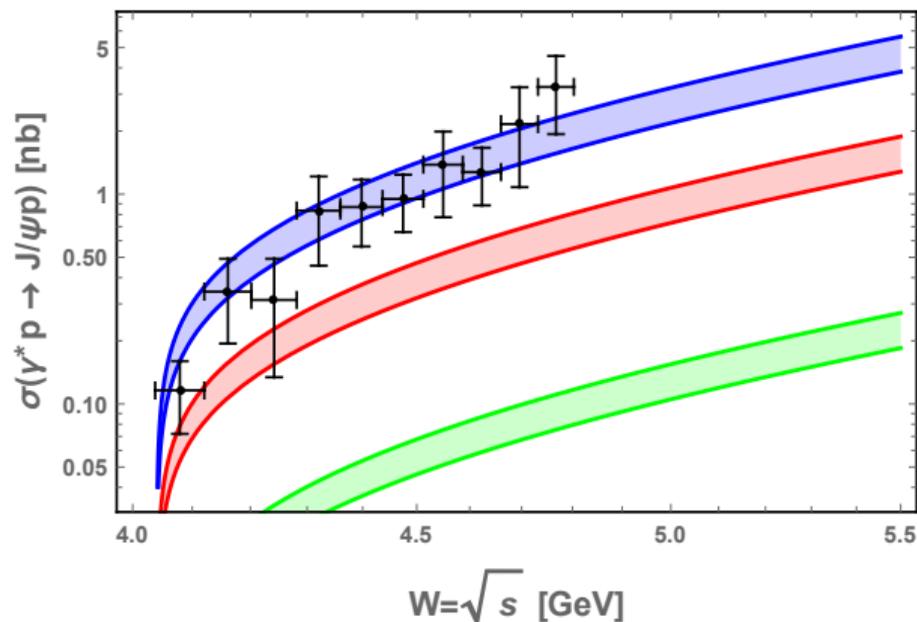
$$\text{with } \mathcal{I}(0, M_{J/\psi}) = \frac{g_5 f_{J/\psi}}{4M_{J/\psi}}$$

# Electroproduction of heavy mesons near threshold



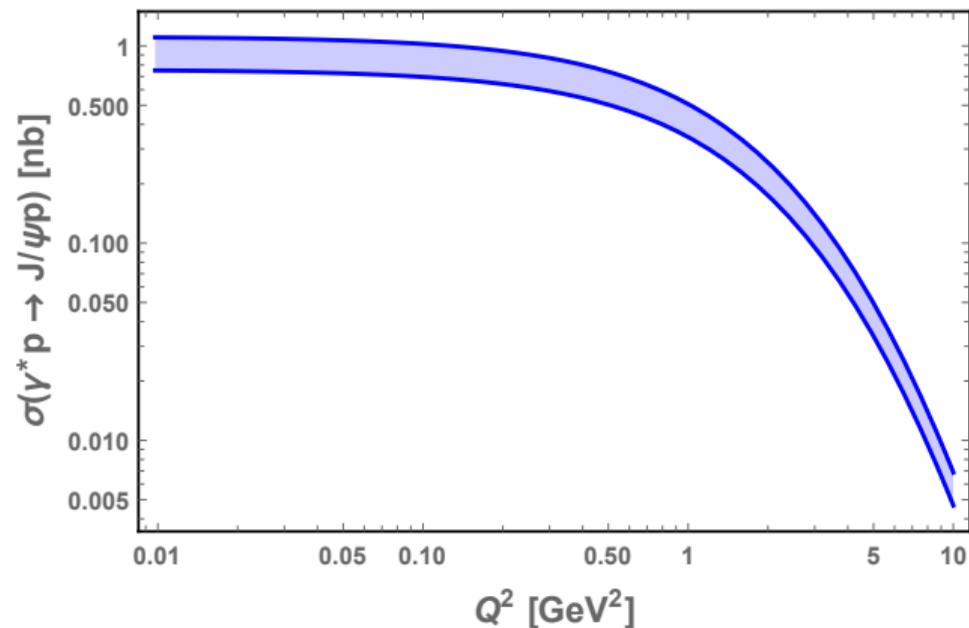
**Figure:** The variation of the total differential cross section with  $t$  and  $Q^2$  for  $s = 21 \text{ GeV}^2$ . The blue curve is for  $Q = 0$ . The red curve is for  $Q = 1.2 \text{ GeV}$ . The green curve is for  $Q = 2.2 \text{ GeV}$ . The data is from GlueX collaboration at JLab in 2019.

# Electroproduction of heavy mesons near threshold



**Figure:** The variation of the total cross section (near threshold) with  $Q^2$  and  $\sqrt{s}$ . The blue band is for  $Q^2 = 0$  (the data is from GlueX in 2019), the red band is for  $Q^2 = 1.2^2 \text{ GeV}^2$ , the green band is for  $Q^2 = 2.2^2 \text{ GeV}^2$ .

# Electroproduction of heavy mesons near threshold



**Figure:** The variation of the total cross section with  $Q^2$  (near threshold),  $s=W^2 = 4.4^2 \text{ GeV}^2$ .

# Electroproduction of heavy mesons near threshold

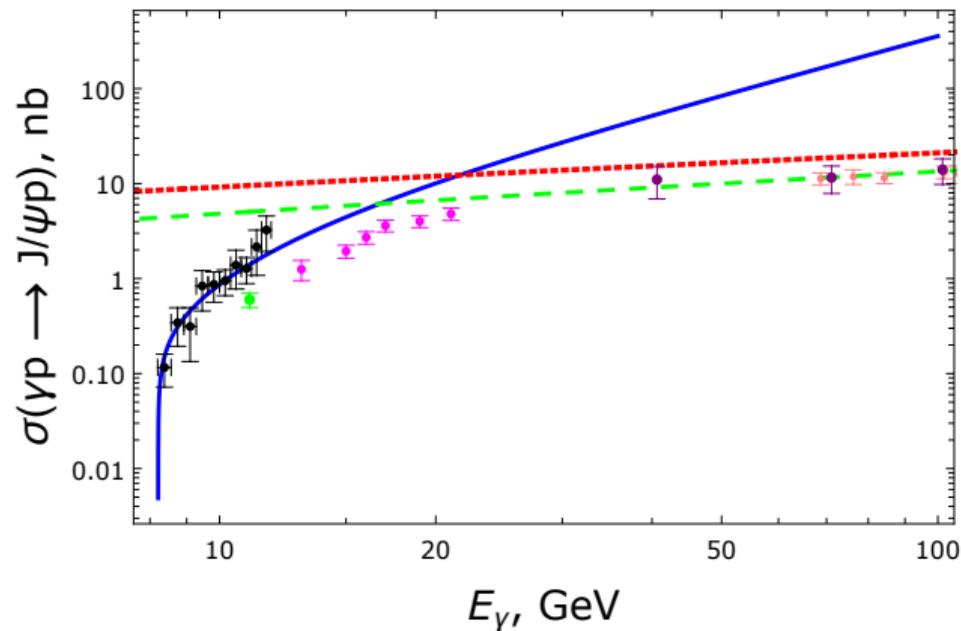


Figure: The total cross section for  $J/\psi$  photoproduction.

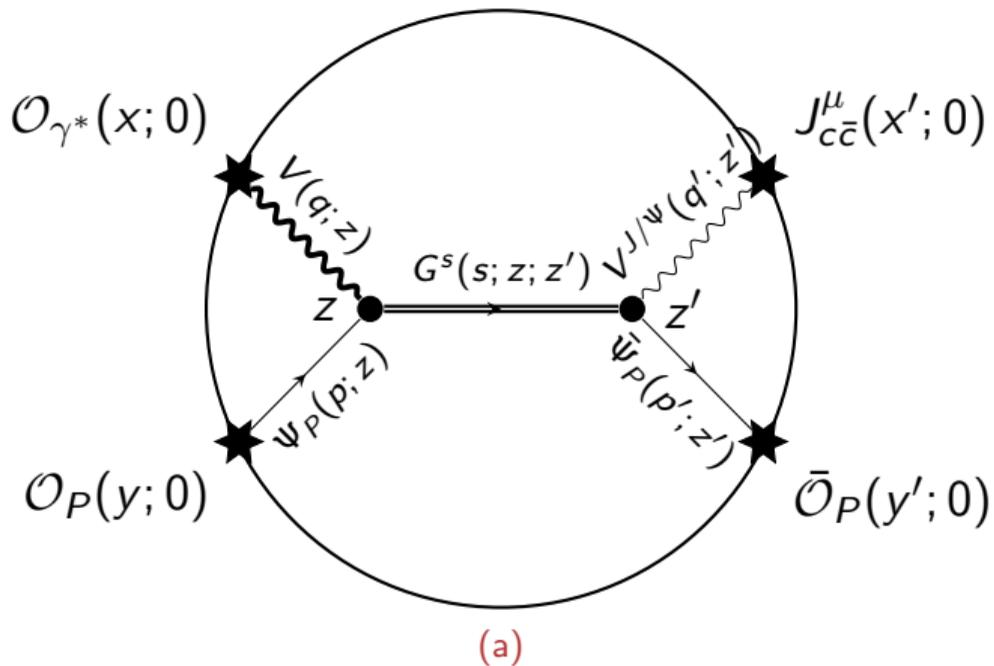
# Pentaquarks at JLab and EIC

- The masses and total widths of the three charm pentaquark states reported by LHCb [Aaij:2019] are

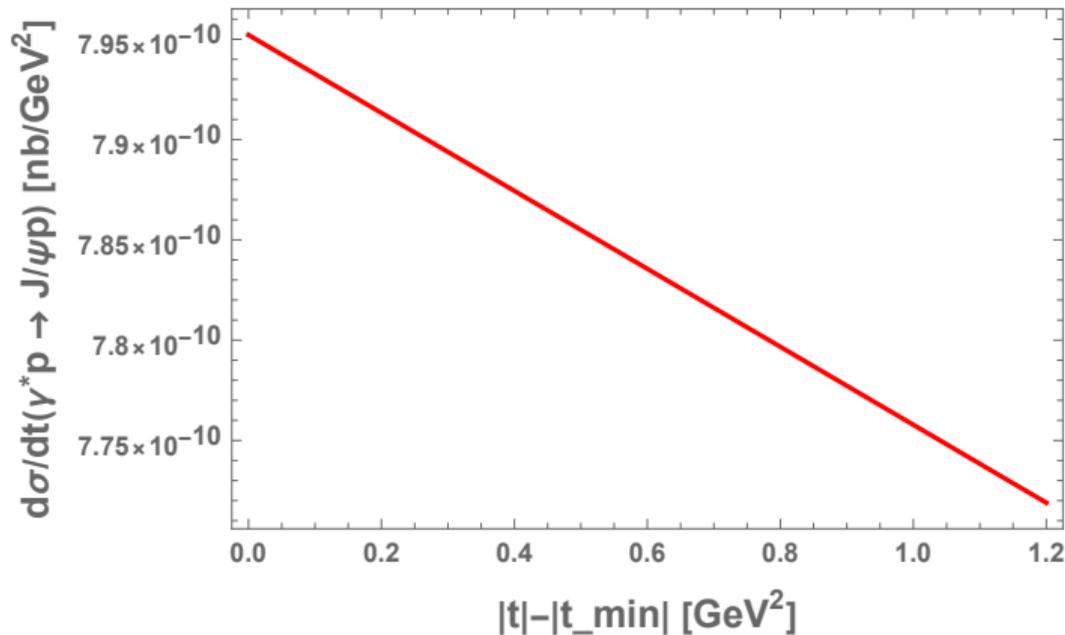
$$m_{P_c} = 4311.9 \pm 0.7 \text{ MeV} \quad \Gamma_{P_c} = 9.8 \pm 2.7 \text{ MeV}$$

$$m_{P_c} = 4440.3 \pm 1.3 \text{ MeV} \quad \Gamma_{P_c} = 20.6 \pm 4.9 \text{ MeV}$$

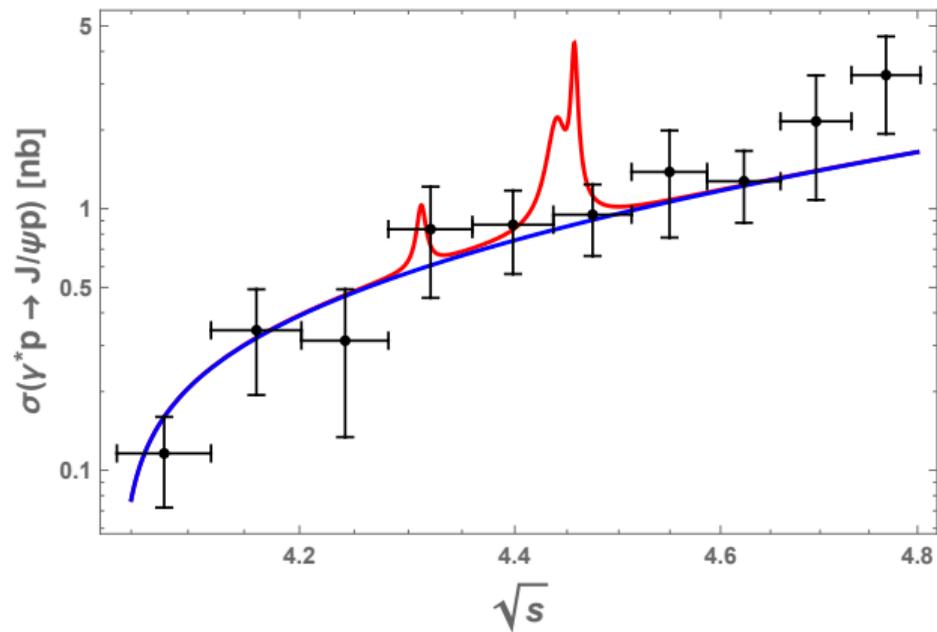
$$m_{P_c} = 4457.3 \pm 0.6 \text{ MeV} \quad \Gamma_{P_c} = 6.4 \pm 2.0 \text{ MeV}$$



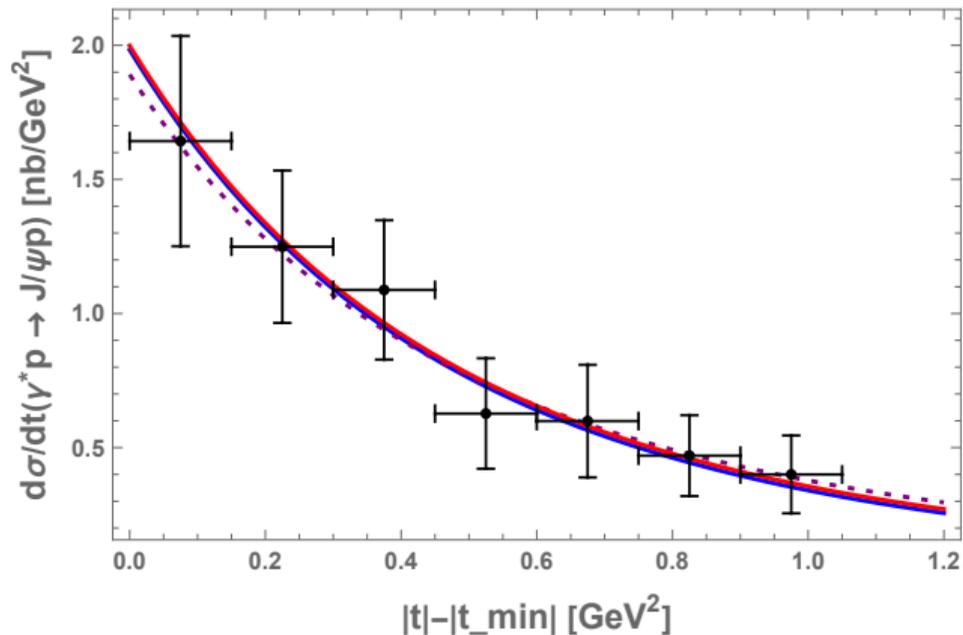
**Figure:** Witten diagram for the holographic **s-channel** electroproduction of  $J/\Psi$ .



**Figure:** s-channel contribution to the photo-production differential cross section for  $V = J/\Psi$ , including all three charm pentaquark contributions.



**Figure:** Total cross section for  $V = J/\psi$  photo-production: the blue-solid curve is our t-channel contribution, the red-solid curve is the sum of t- and s-channel contribution showing the three holographic pentaquarks times  $\mathcal{N}_s = 2.0 \times 10^6$  to make them visible, and the data is from [GlueX:2019].



**Figure:** s-channel plus t-channel contributions to the photo-production differential cross section for  $V = J/\Psi$ , including all three charm pentaquark contributions. The blue-solid and dashed-purple curves are the holographic t-channel contributions. The solid-red line is the total s- and t-channel contribution with the s-channel contribution multiplied by  $\mathcal{N}_s = 2.0 \times 10^7$  to make it slightly visible. The data are from [GlueX:2019] at  $\sqrt{s} = 4.6$  GeV.

# Holographic search for exotics at JLab and EIC

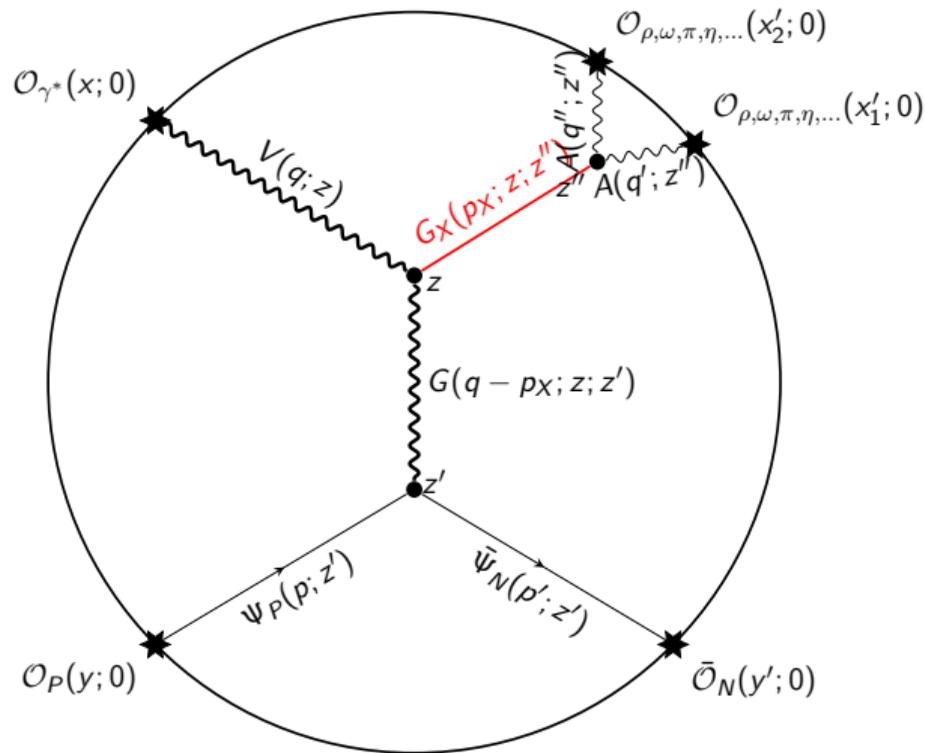


Figure: A general t-channel Witten diagram for exotics at JLab and EIC.

$$V_\mu = V_\mu^a T^a + \frac{\tilde{V}_\mu}{\sqrt{4}} \mathbf{1}_{4 \times 4}$$

$$V_\mu = \begin{pmatrix} \frac{\rho_\mu^0}{\sqrt{2}} + \frac{\omega_\mu}{\sqrt{6}} + \frac{\phi_\mu}{\sqrt{12}} & \rho_\mu^+ & K_\mu^{*+} & D_\mu^{*0} \\ \rho_\mu^- & -\frac{\rho_\mu^0}{\sqrt{2}} + \frac{\omega_\mu}{\sqrt{6}} + \frac{\phi_\mu}{\sqrt{12}} & K_\mu^{*0} & D_\mu^{*+} \\ K_\mu^{*-} & \bar{K}_\mu^{*0} & -\frac{2\omega_\mu}{\sqrt{6}} + \frac{\phi_\mu}{\sqrt{12}} & D_\mu^{*-} \\ \bar{D}_\mu^{*0} & D_\mu^{*-} & \bar{D}_\mu^{*0} & -\frac{3\phi_\mu}{\sqrt{12}} \end{pmatrix} + \frac{\tilde{V}_\mu}{\sqrt{4}} \mathbf{1}_{4 \times 4}$$

and

$$B = B^a T^a + \frac{\tilde{B}}{\sqrt{4}} \mathbf{1}_{4 \times 4}$$

$$B = \begin{pmatrix} \frac{\Sigma^0}{\sqrt{2}} + \frac{\Lambda}{\sqrt{6}} + \frac{B^0}{\sqrt{12}} & \Sigma^+ & p & \Sigma_c^{++} \\ \Sigma^- & -\frac{\Sigma^0}{\sqrt{2}} + \frac{\Lambda}{\sqrt{6}} + \frac{B^0}{\sqrt{12}} & n & \Sigma_c^+ \\ \Xi^- & \Xi^0 & -\frac{2\Lambda}{\sqrt{6}} + \frac{B^0}{\sqrt{12}} & \Xi_c^0 \\ \Xi_c^- & \Xi_c^0 & \Omega_c & -\frac{3B^0}{\sqrt{12}} \end{pmatrix} + \frac{\tilde{B}}{\sqrt{4}} \mathbf{1}_{4 \times 4}$$

$$\mathcal{L}_{\text{int}} = \frac{1}{2} g_5 f^{abc} \left( \partial_\mu V_\nu^a - \partial_\nu V_\mu^a \right) V^{\mu b} V^{\nu c} - \frac{1}{4} g_5^2 f^{abc} f^{ade} V_\mu^b V_\nu^c V^{\mu d} V^{\nu e}.$$

$$\mathcal{L}_{\text{CS}} \subset g_5 \epsilon^{MNPQR} \text{Tr} \left[ A_M \tilde{F}_{NP}(V) \tilde{F}_{QR}(V) + A_M \tilde{F}_{NP}(A) \tilde{F}_{QR}(A) \right].$$

$$\mathcal{L}_{\text{int}} = -i \frac{g}{2} f^{abc} V_\mu^a \bar{B}^c \gamma^\mu B^b$$

- **AdS/CFT Approach:** Witten diagrams in a confining AdS background capture QCD correlators and scattering amplitudes.
- **Form Factors:** Proton gravitational form factors arise from metric fluctuations (corresponding to  $2^{++}$ , and  $0^{++}$  glueballs), matching lattice and experimental data.
- **Near-Threshold Production:** Holographic amplitudes for photoproduction of heavy quarkonia ( $J/\psi$ ,  $\Upsilon$ )
- **Electroproduction &  $Q^2$ :** Virtual photon effects enter through a transition form factor with multi-pole behavior, describing electroproduction cross sections at various kinematics.
- **Pentaquarks (s-channel):** Possible s-channel resonances (pentaquarks) appear near threshold in addition to dominant t-channel glueball exchanges, offering predictions for JLab and EIC.

# Thank You!