

Searches for Lepton Flavor Violation at the EIC

Kaori Fuyuto

LANL → KEK

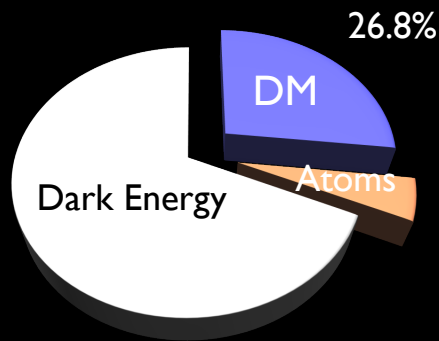


V. Cirigliano, KF, C. Lee, E. Mereghetti, B. Yan, JHEP03(2021)256
S. Banerjee, V. Cirigliano, et al, Snowmass White Paper, 2203.14919
F. Delzanno, KF, S. Gonzalez-Solis, E. Mereghetti, arXiv 2411.13497 (Accepted by JHEP)

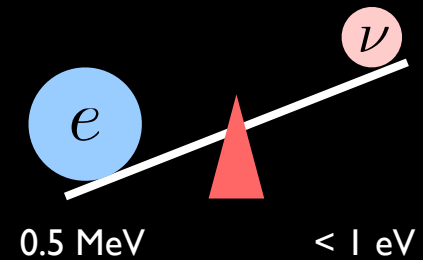
July 22nd, 2025
At Stony Brook U

We still don't know much about our Universe.

What is Dark Matter?



What is the origin of tiny neutrino mass ?



The origin of the present Universe

Why is there more matter than antimatter?

$$\frac{n_b - n_{\bar{b}}}{n_\gamma} = 6.1 \times 10^{-10}$$

✦ Need Physics Beyond the Standard Model

Charged Lepton Flavor Violation

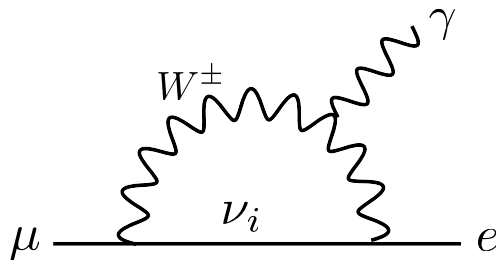
Searches for CLFV are strong tools to probe BSM physics.

*Beyond the minimal extension of the SM

Charged Lepton Flavor Violation

Searches for CLFV are strong tools to probe BSM physics.

Ex) SM + neutrino mass (vSM)



Petcov '77, Marciano-Sanda '77

$$\mathcal{L} = \mathcal{L}_{\text{SM}} + \mathcal{L}_{\nu\text{-mass}}$$

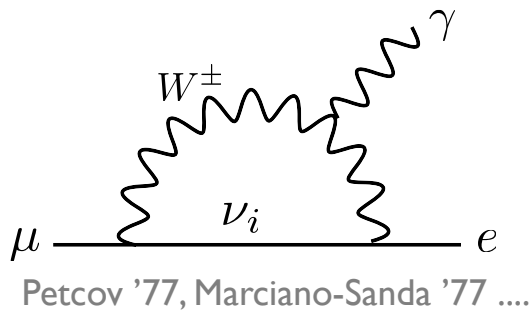
Dirac or Majorana

$$\text{Br}(\mu \rightarrow e\gamma) = \frac{\alpha_{\text{em}}}{32\pi} \left| \sum_{i=2,3} U_{\mu i}^* U_{ei} \frac{\Delta m_{1i}^2}{m_W^2} \right|^2 < 10^{-54} \quad \text{Extremely small!}$$

Charged Lepton Flavor Violation

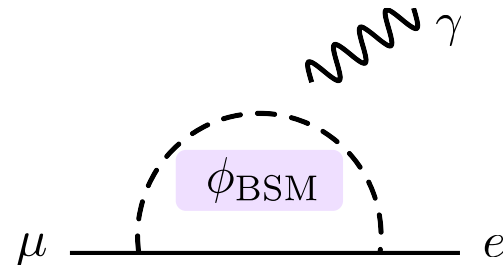
Searches for CLFV are strong tools to probe BSM physics.

Ex) SM + neutrino mass (vSM)



$$\text{Br}(\mu \rightarrow e\gamma) < 10^{-54}$$

$\ll$



$$\text{Br}(\mu \rightarrow e\gamma)_{\text{BSM}}$$

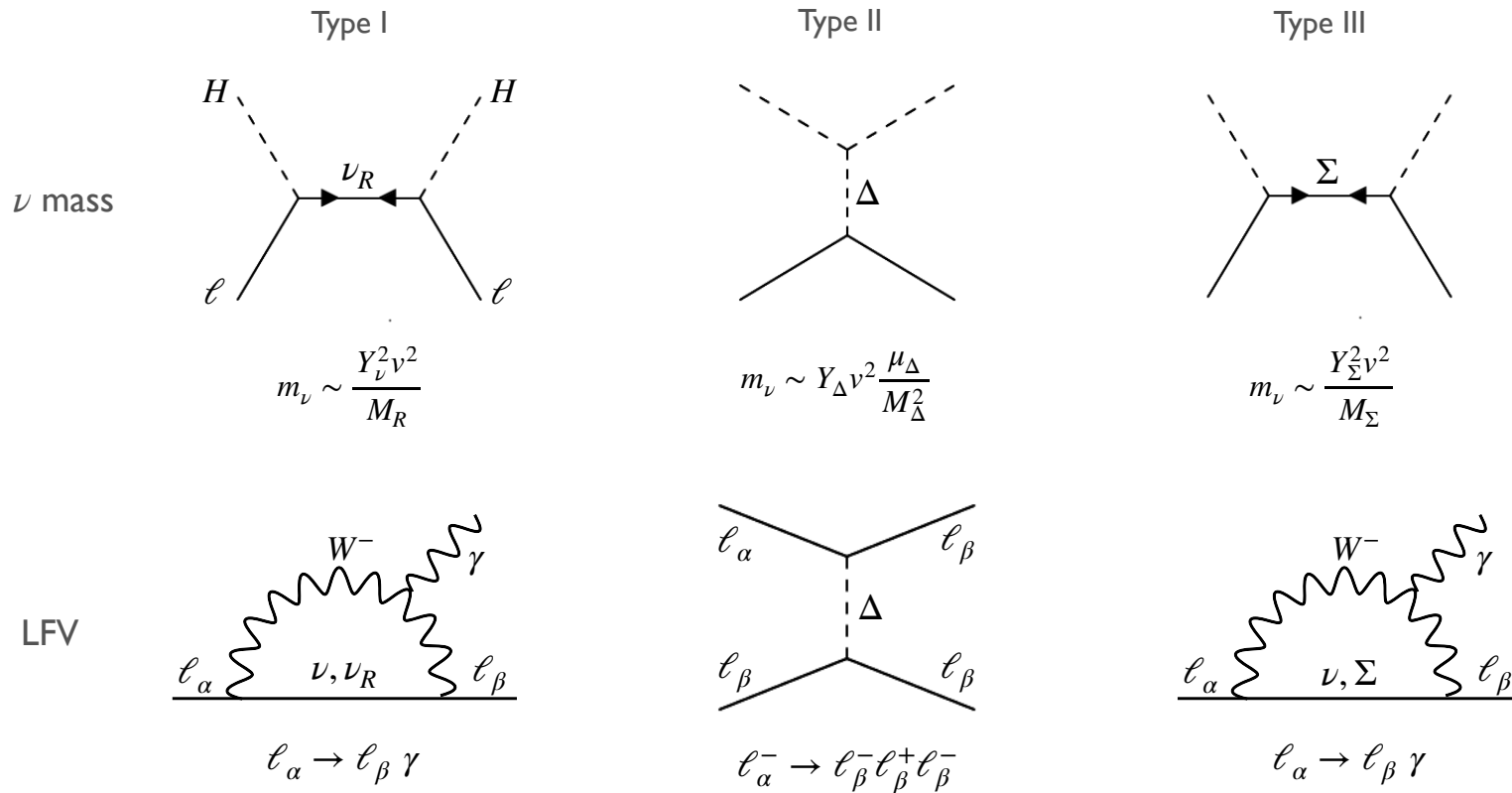
✓ The Observations of CLFV would point to new physics beyond vSM.

*Underlying mechanism of the neutrino mass.

Charged Lepton Flavor Violation

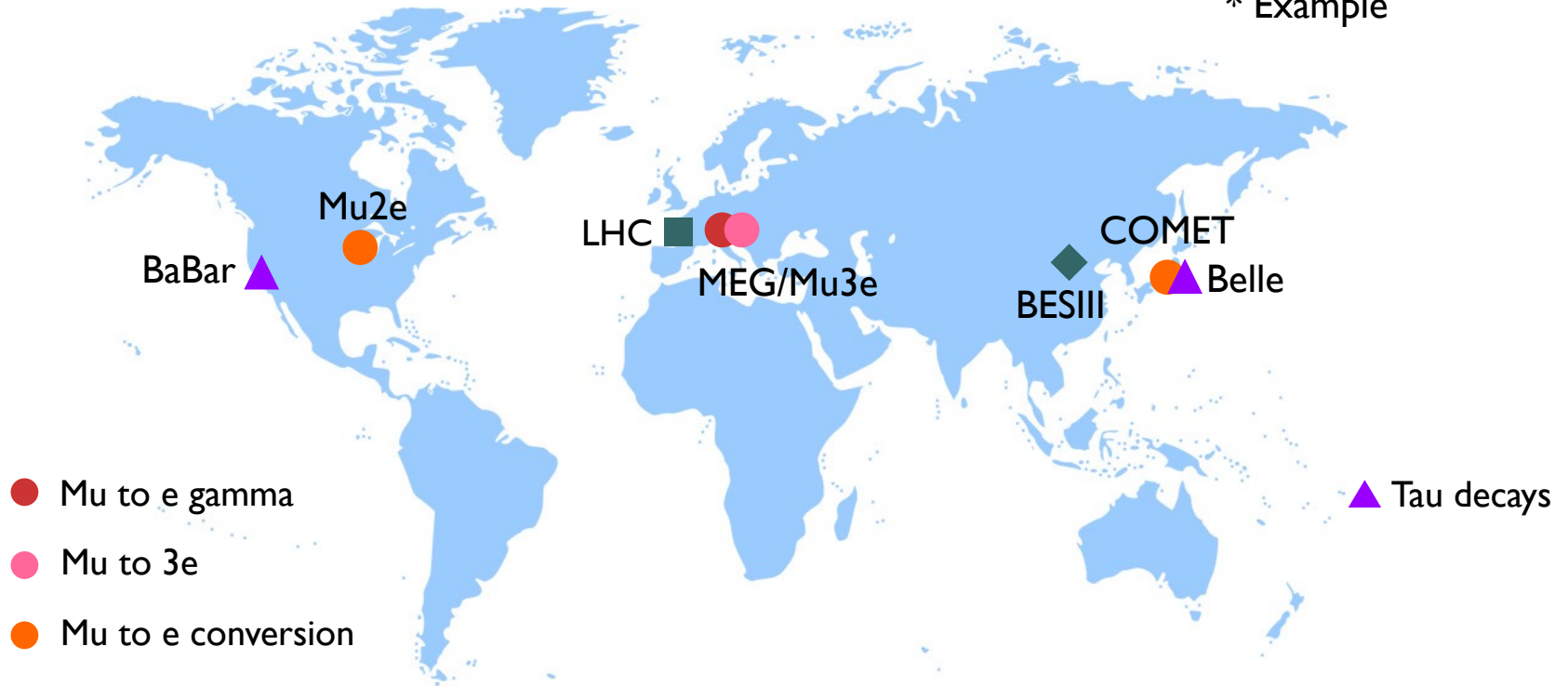
Models that explain neutrino mass usually introduce CLFV at tree or loop level.

e.g., A.Abada, et al, JHEP 12 (2007) 061



CLFV searches

* Example



$$\text{BR}(\mu \rightarrow e\gamma) < 3.1 \times 10^{-13}$$

MEG II Collaboration, 2310.12614

$$\text{BR}(\tau \rightarrow e\gamma) < 3.3 \times 10^{-8}$$

BaBar, PRL104 (2010) 021802

$$\text{BR}(\mu^- \text{ Ti} \rightarrow e^- \text{ Ti}) < 6.1 \times 10^{-13}$$

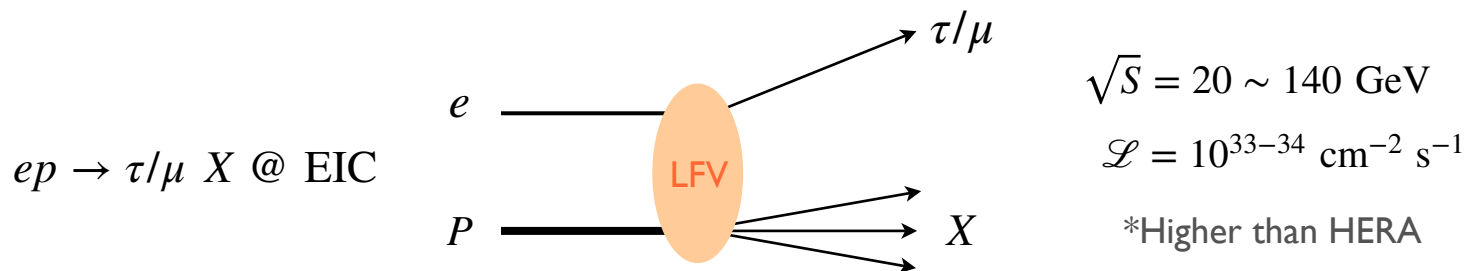
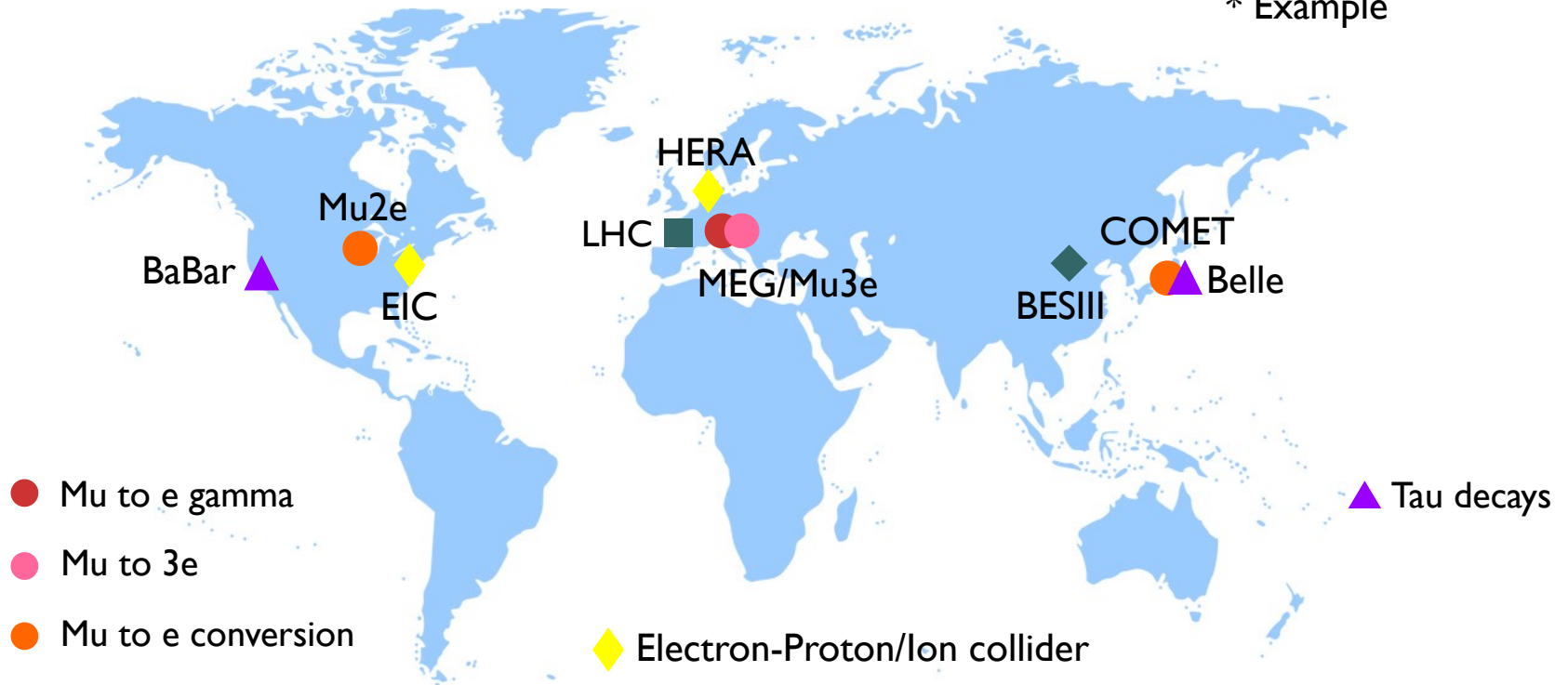
P.Wintz, Conf. Proc. C 980420, 534 (1998).

$$\text{BR}(\tau \rightarrow e\pi^+\pi^-) < 2.3 \times 10^{-8}$$

Belle, PLB719 (2013) 346-353

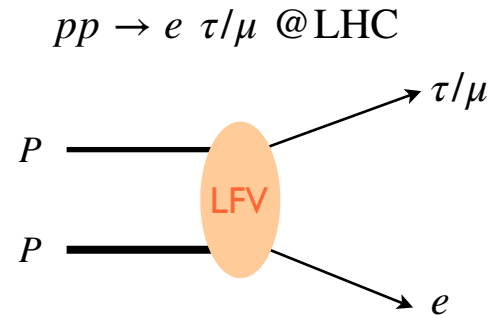
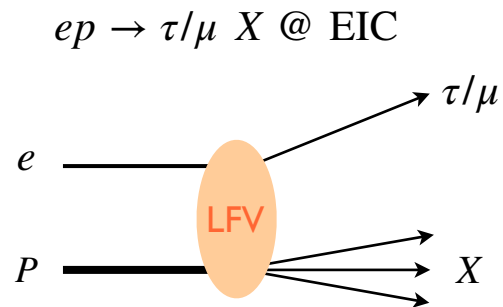
CLFV searches

* Example

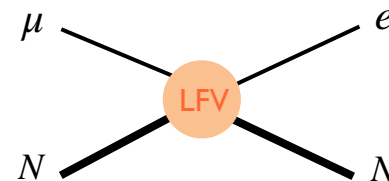
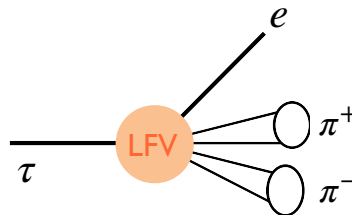
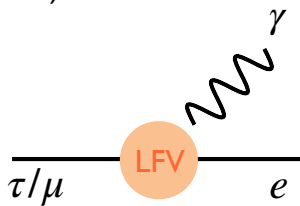


✓ Model-Independent Analysis of CLFV process at low- and high-energy

EIC vs LHC vs Low-Energy CLFV searches



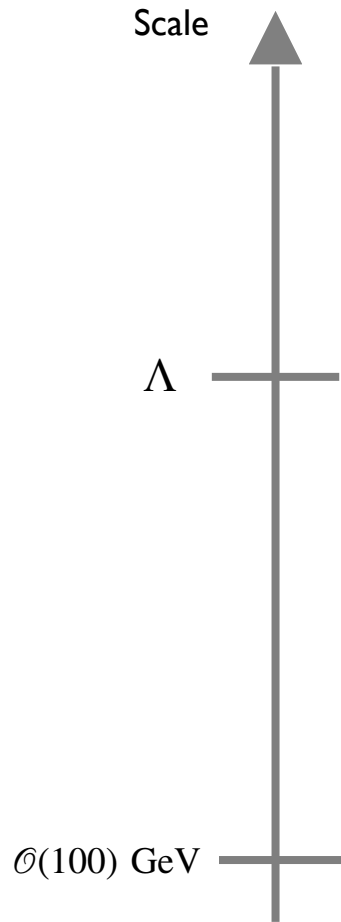
Ex) τ, μ , and meson decays



Model-Independent Analysis

10

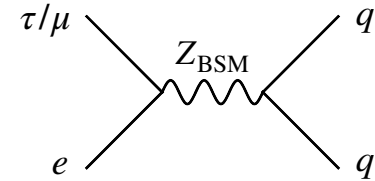
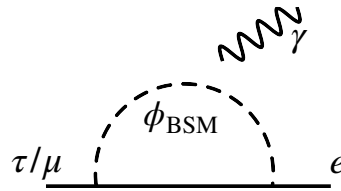
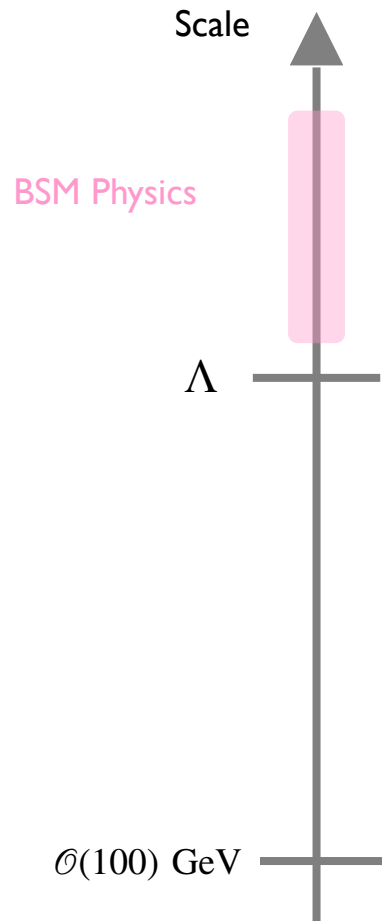
SMEFT : Standard Model Effective Field Theory



Model-Independent Analysis

II

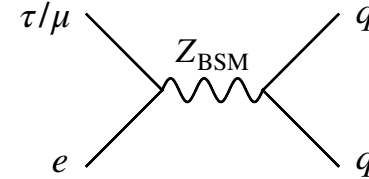
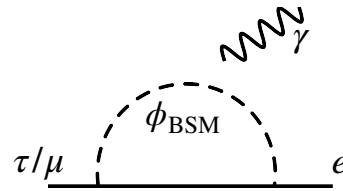
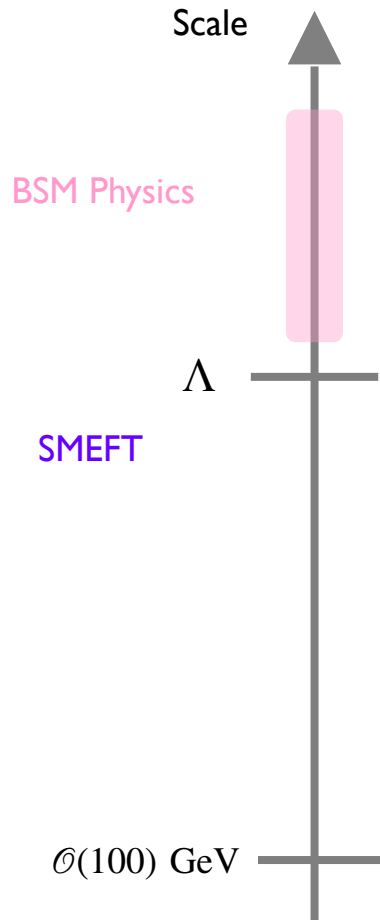
SMEFT : Standard Model Effective Field Theory



Model-Independent Analysis

12

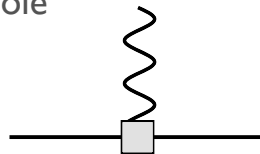
SMEFT : Standard Model Effective Field Theory



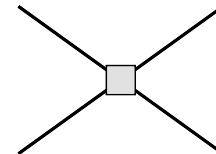
$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \sum_{d=5} \frac{1}{\Lambda^{d-4}} C^{(d)} O^{(d)}$$

Higher Dimensional Operators

Dipole



4 Fermion



All possible interactions based on gauge and Lorentz invariance

✓ EFT can apply to concrete models

4 different types of LFV operators (dim 6)

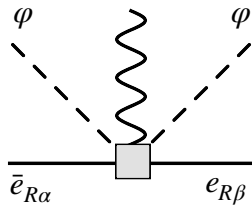
$$\mathcal{L}_{\text{LFV}} = \mathcal{L}_{\psi^2 \varphi^2 D} + \mathcal{L}_{\psi^2 X \varphi} + \mathcal{L}_{\psi^2 \varphi^3} + \mathcal{L}_{\psi^4}$$

X : Gauge boson

ψ : Fermion

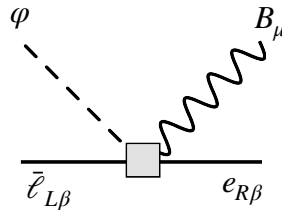
φ : Higgs

* $\mathcal{O}(100)$ independent operators



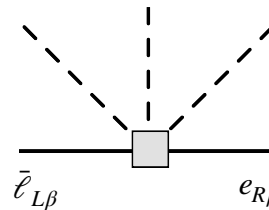
EX) $\varphi^\dagger i \overleftrightarrow{D}_\mu \varphi \bar{e}_{R\alpha} \gamma_\mu e_{R\beta}$

3



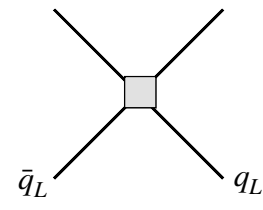
$$\bar{\ell}_{L\alpha} \sigma^{\mu\nu} B_{\mu\nu} \varphi e_{R\beta}$$

2



$$\varphi^\dagger \varphi \bar{\ell}_{L\alpha} \varphi e_{R\beta}$$

1



$$\bar{\ell}_L \gamma^\mu \ell_L \bar{q}_L \gamma_\mu q_L$$

10

x 9 for generic quark flavor

CLFV operators

14

4 different types of LFV operators (dim 6)

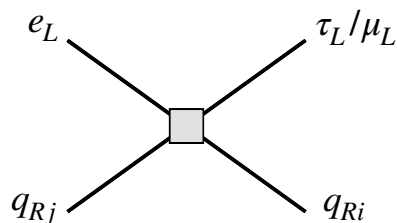
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$$\supset -\frac{4G_F}{\sqrt{2}} \sum_{\substack{\ell = \tau, \mu \\ q = u, d}} [C_{Lq}]_{\ell eij} \bar{\ell}_L \gamma^\mu e_L \bar{q}_{Ri} \gamma_\mu q_{Rj}$$



*Assume a generic quark flavor structure

$$\text{Ex) } [C_{Ld}]_{\tau e} = \begin{pmatrix} [C_{Ld}]_{dd} & [C_{Ld}]_{ds} & [C_{Ld}]_{db} \\ [C_{Ld}]_{sd} & [C_{Ld}]_{ss} & [C_{Ld}]_{sb} \\ [C_{Ld}]_{bd} & [C_{Ld}]_{bs} & [C_{Ld}]_{bb} \end{pmatrix}$$

*Focus on tau-electron case.

Decay mode		Upper limit (90 % C . L.)	
$\tau \rightarrow e\pi^+\pi^-$	uu/dd/ss	2.3×10^{-8}	Belle PLB719(2013)346
$\tau \rightarrow e\pi^0$		8×10^{-8}	Belle PLB648(2007)341
$\tau \rightarrow e\eta$		9.2×10^{-8}	Belle PLB648(2007)341
$\tau \rightarrow e\eta'$		1.6×10^{-7}	Belle PLB648(2007)341
$\tau \rightarrow eK_S$	ds/sd	2.6×10^{-8}	Belle PLB692(2010)4
$\tau \rightarrow e\pi^+K^-$		3.7×10^{-8}	Belle PLB719(2013)346
$\tau \rightarrow e\pi^-K^+$		3.1×10^{-8}	Belle PLB719(2013)346
$B^0 \rightarrow e^\pm\tau^\mp$	db/bd	1.6×10^{-5}	Belle PRD104(2021)9
$B^+ \rightarrow \pi^+e^+\tau^-$		7.4×10^{-5}	BaBar PRD86(2012)012004
$B^+ \rightarrow \pi^+e^-\tau^+$		2.0×10^{-5}	BaBar PRD86(2012)012004
$B^+ \rightarrow K^+e^+\tau^-$	sb/bs	1.53×10^{-5}	Belle PRL130(2023)26 261802
$B^+ \rightarrow K^+e^-\tau^+$		1.5×10^{-5}	Belle PRL130(2023)26 261802

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- Certain combinations of CLFV operators can be bounded.

$$\text{Ex) } \text{BR}(\tau \rightarrow e\pi^+\pi^-) \simeq 0.5 \times \left| [C_{Lu}]_{uu} - [C_{Ld}]_{dd} \right|^2$$

A. Celis, V. Cirigliano, E. Passemar, PRD89(2014)095014

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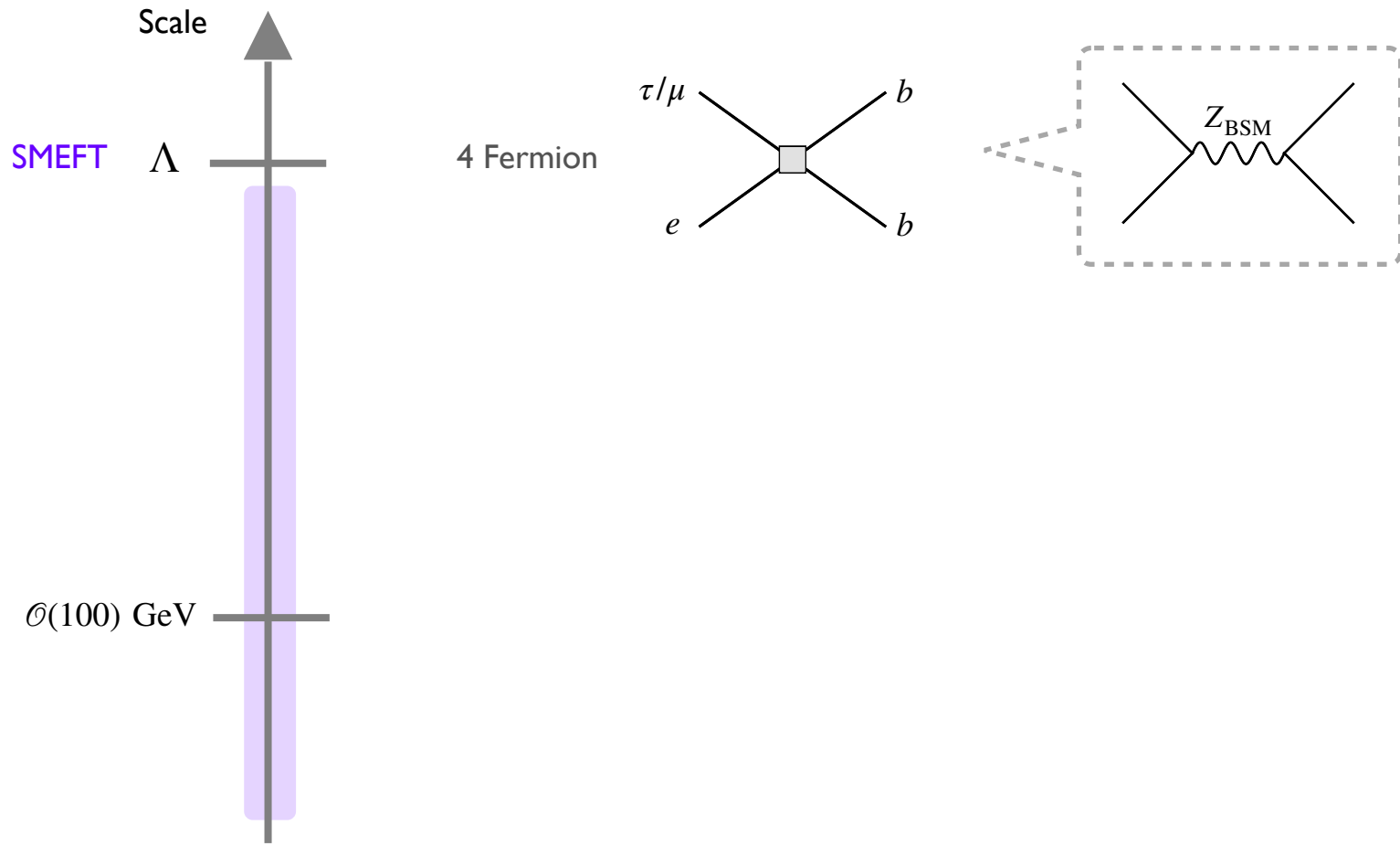
- Quark-flavor conserving processes are generated by light quarks operators

$$[C_{Lu}]_{\tau e} = \begin{pmatrix} [C_{Lu}]_{uu} & [C_{Lu}]_{uc} & [C_{Lu}]_{ut} \\ [C_{Lu}]_{cu} & [C_{Lu}]_{cc} & [C_{Lu}]_{ct} \\ [C_{Lu}]_{tu} & [C_{Lu}]_{tc} & [C_{Lu}]_{tt} \end{pmatrix} \quad [C_{Ld}]_{\tau e} = \begin{pmatrix} [C_{Ld}]_{dd} & [C_{Ld}]_{ds} & [C_{Ld}]_{db} \\ [C_{Ld}]_{sd} & [C_{Ld}]_{ss} & [C_{Ld}]_{sb} \\ [C_{Ld}]_{bd} & [C_{Ld}]_{bs} & [C_{Ld}]_{bb} \end{pmatrix} \quad \text{How?}$$

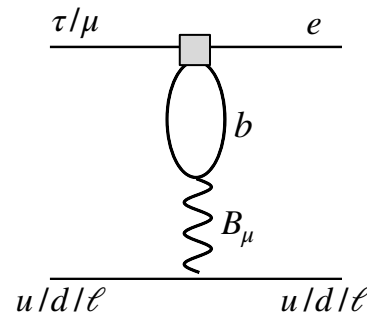
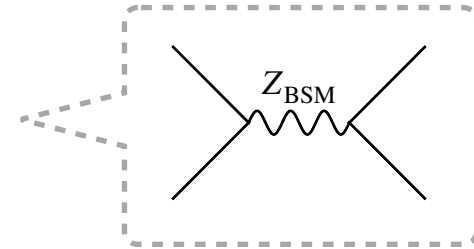
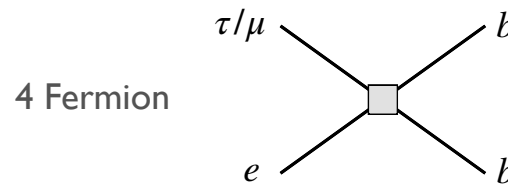
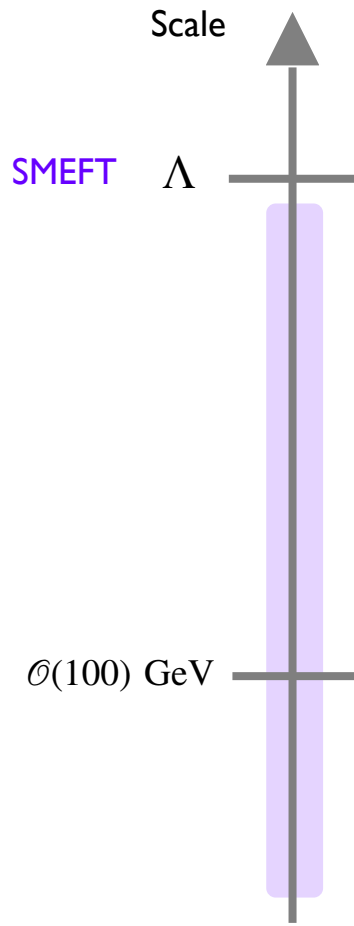
Model-Independent Analysis

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Scale Running Effects based on the Renormalization Group Equations

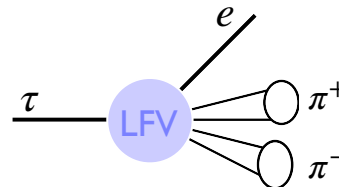


Scale Running Effects based on the Renormalization Group Equations



$$\text{Ex)} \quad \mu \frac{d}{d\mu} [C_{Lu}]_{dd} = \frac{4}{3} N_c \frac{g_1^2}{(4\pi)^2} y_d^2 [C_{Ld}]_{bb}$$

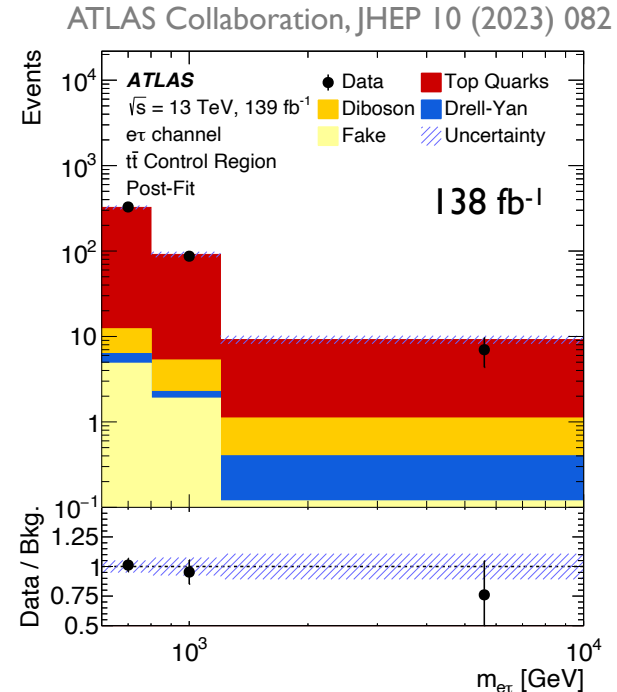
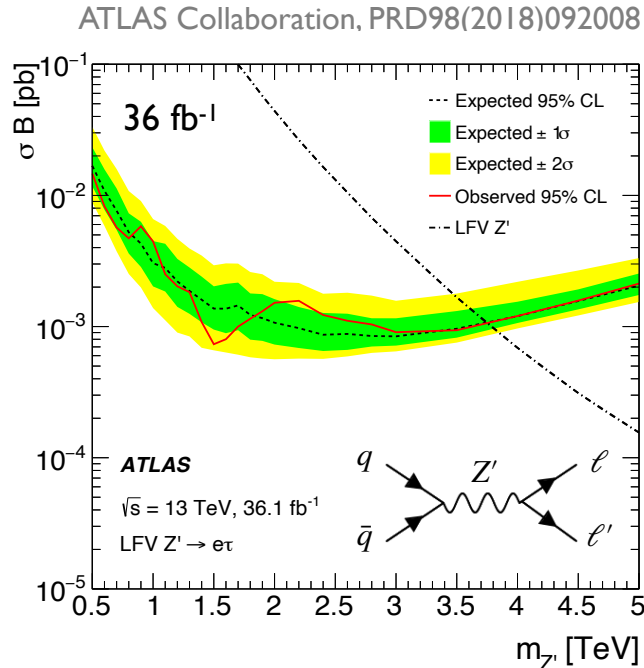
For example,



$$[C_{Lu}]_{uu}, [C_{Ld}]_{dd} \sim \frac{g^2}{(4\pi)^2} [C_{Ld}]_{bb}$$

Loop effect $\sim \mathcal{O}(10^{-3})$

LHC search



- Bound on CLFV top decay by ATLAS with 79.8 fb⁻¹: $\text{BR}(t \rightarrow q\ell\ell') < 1.86 \times 10^{-5}$ (95 % CL.)

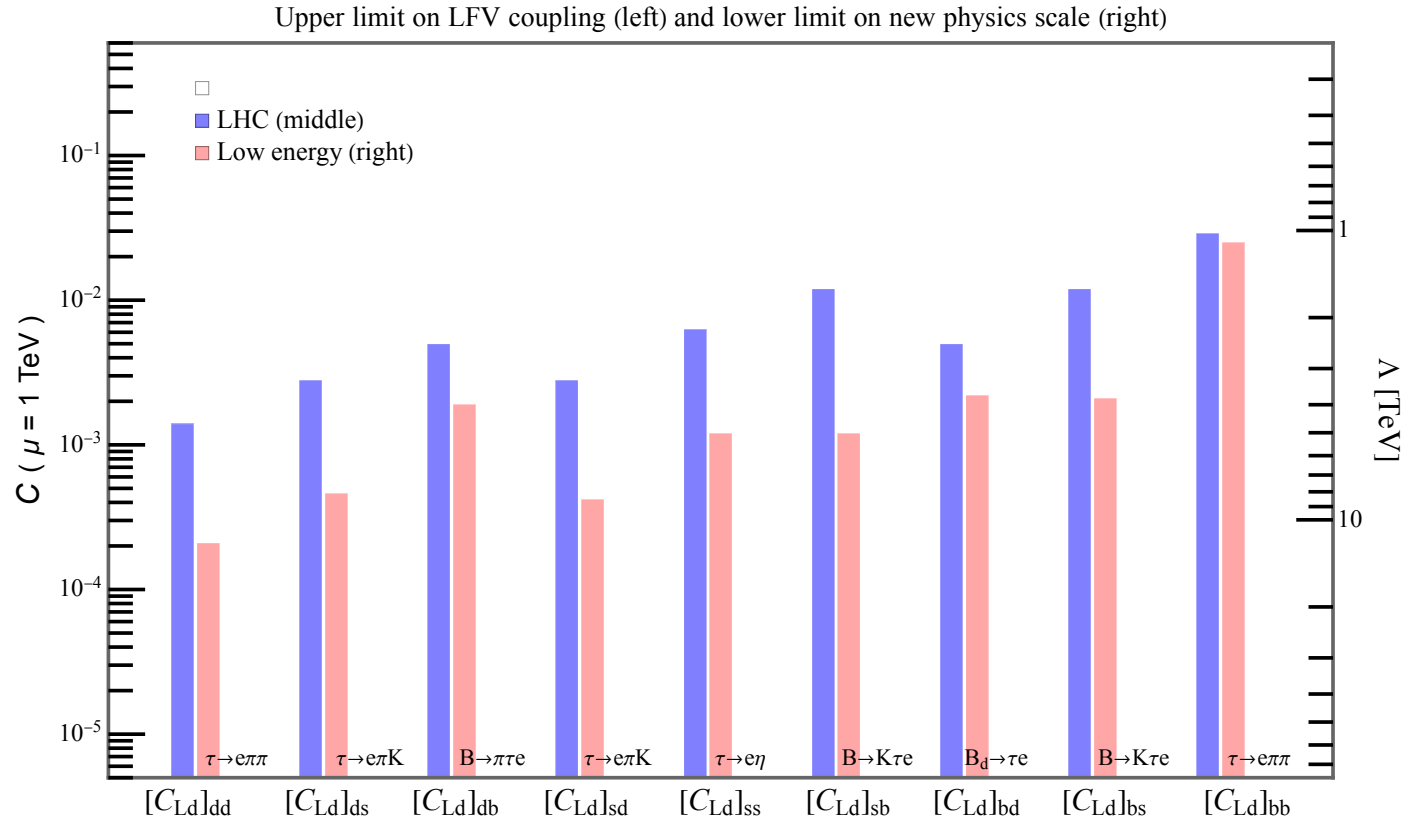
ATLAS collaboration, ATLAS-CONF-2018-044

- ATLAS published $pp \rightarrow \ell\ell'$ bounds in high-mass final states using 36 fb⁻¹

'22 ATLAS and '23 CMS results with 138 and 139 fb⁻¹ ATLAS JHEP 10 (2023) 082
CMS JHEP 05 (2023) 227

Existing bounds

* Single Operator Analysis

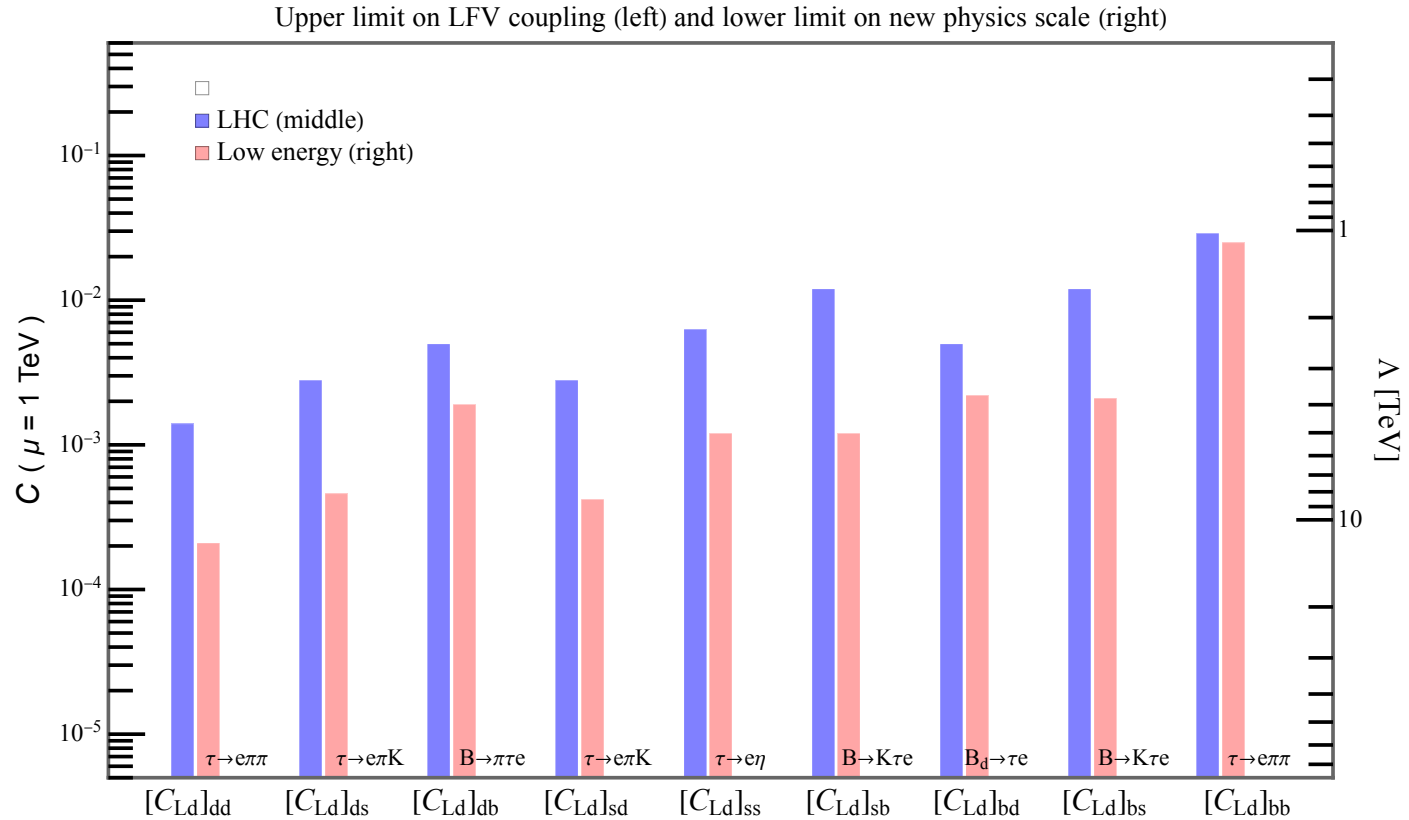


$$[C_{Ld}]_{ij} \bar{\tau}_L \gamma^\mu e_L \bar{d}_{Ri} \gamma_\mu d_{Rj}$$

- Operators with d-type quarks sector well constrained by low-energy
- PDF and loop suppression in $[C_{Ld}]_{bb}$

Existing bounds

* Single Operator Analysis

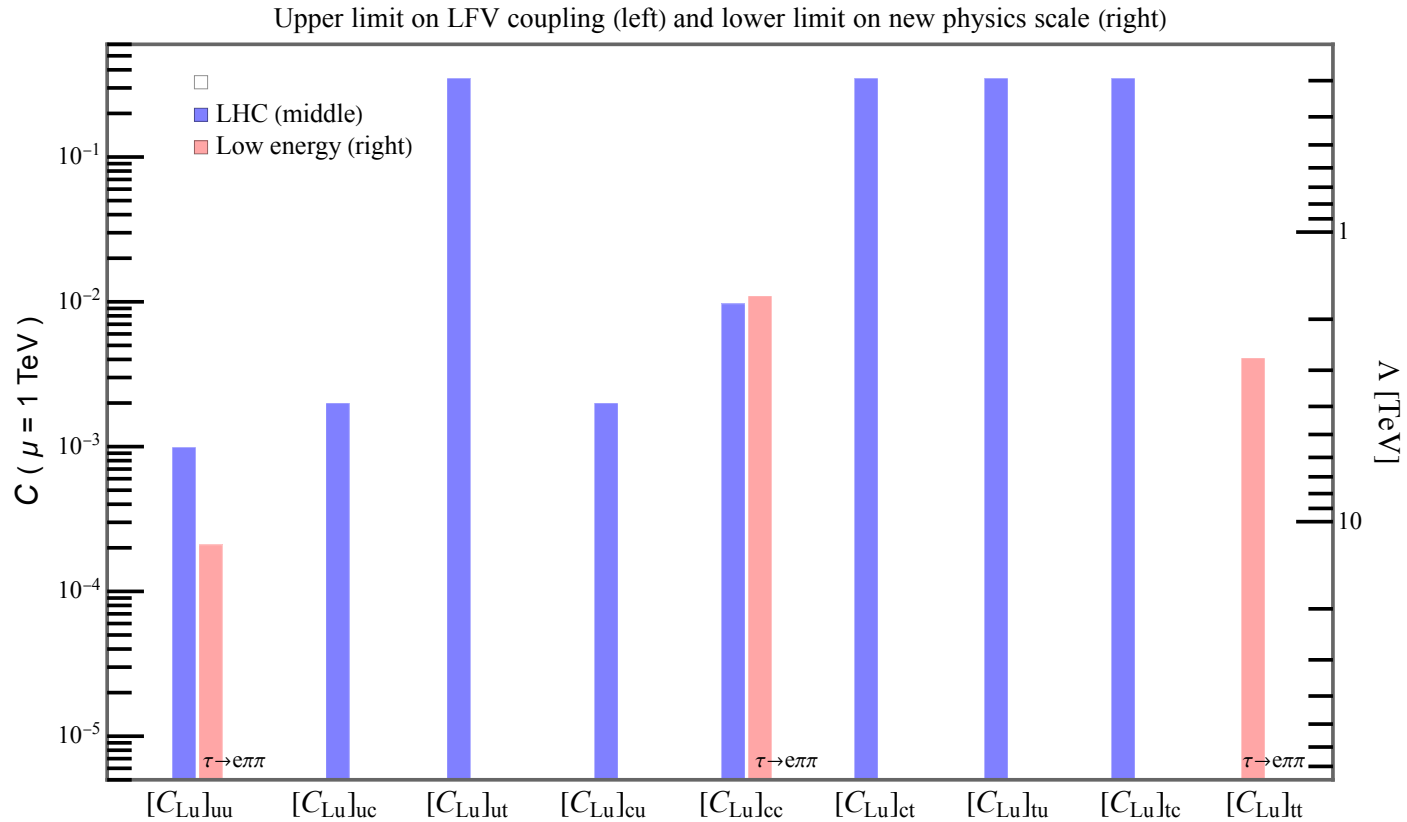


$$\frac{4G_F}{\sqrt{2}}C = \frac{1}{\Lambda^2}$$

Lower bound on the scale : $\Lambda \gtrsim 10$ TeV from $[C_{Ld}]_{dd}$

Existing bounds

* Single Operator Analysis



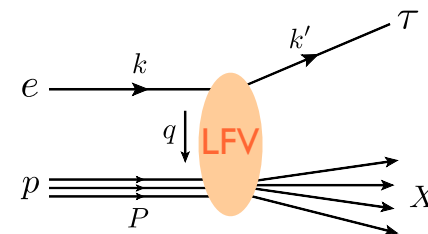
$$[C_{Lu}]_{ij} \bar{\tau}_L \gamma^\mu e_L \bar{u}_{Ri} \gamma_\mu u_{Rj}$$

- Less constrained by low energy than d-type operators
- Strong bound on $[C_{Lu}]_{tt}$ from $\tau \rightarrow e\pi^+\pi^-$

EIC Analysis

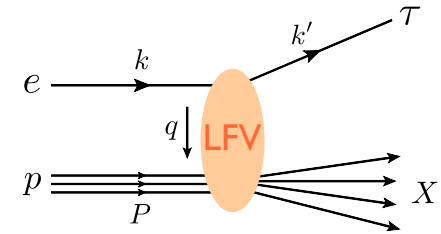
- Cross sections : $\mathcal{O}(1 - 10)$ pb at $\sqrt{S} = 141$ GeV
e.g., 19 pb for $[C_{Lu}]_{uu}$ and 0.8 pb for $[C_{Ld}]_{bb}$

- Major backgrounds
 - 1) Neutral Current $ep \rightarrow ej$
 - 2) Charged Current $ep \rightarrow \nu_e j$



EIC Analysis

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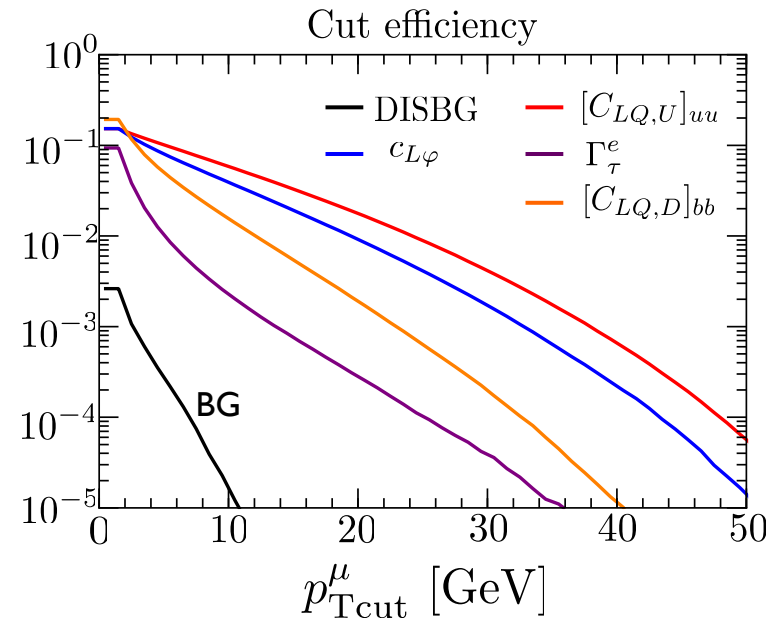
- Promising ID channel

$$\text{BR}(\tau \rightarrow e \bar{\nu}_e \nu_\tau) = 17.82 \%$$

$$\checkmark \text{BR}(\tau \rightarrow \mu \bar{\nu}_\mu \nu_\tau) = 17.39 \%$$

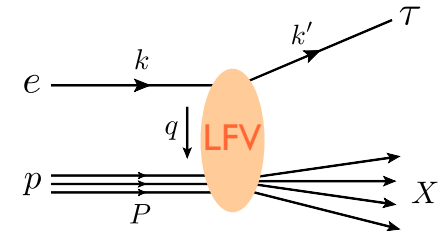
$$\text{BR}(\tau \rightarrow X_h \nu_\tau) = 64.8 \%$$

$$p_T^\mu > 10 \text{ GeV}, \quad E_T > 15 \text{ GeV}, \quad p_T^{j_1} > 20 \text{ GeV}$$



EIC Analysis

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e.g., 19 pb for $[C_{Lu}]_{uu}$ and 0.8 pb for $[C_{Ld}]_{bb}$



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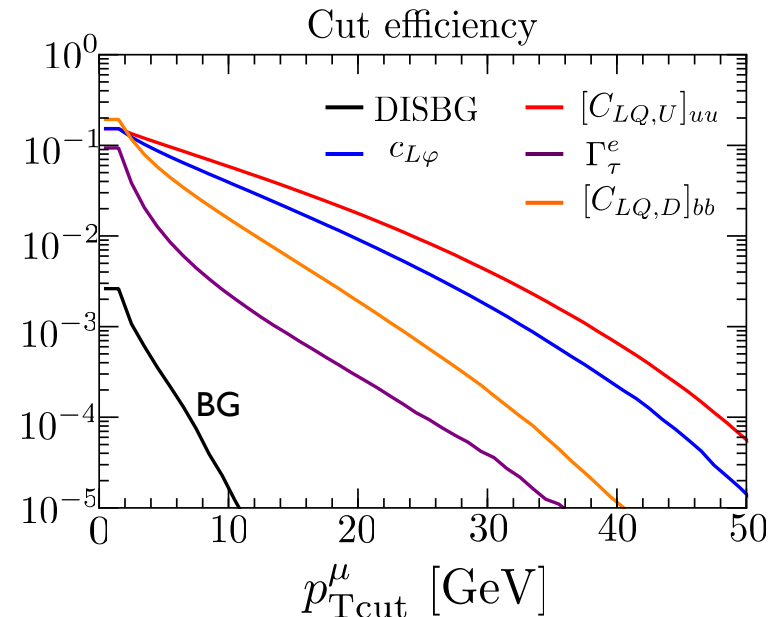
$$\checkmark \text{BR}(\tau \rightarrow \mu \bar{\nu}_\mu \nu_\tau) = 17.39 \%$$

$$\text{BR}(\tau \rightarrow X_h \nu_\tau) = 64.8 \%$$

E.g., Cut Efficiencies

6 - 10% for valence quarks

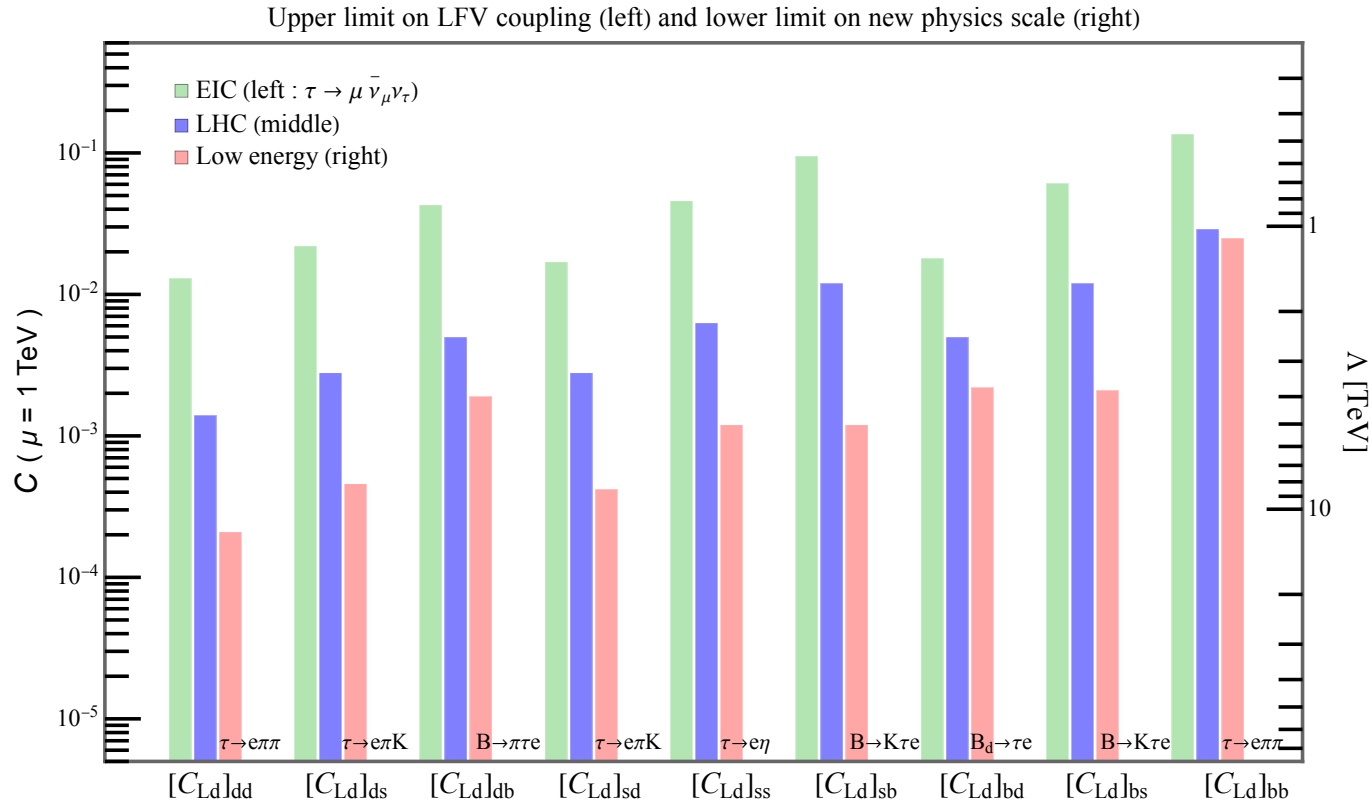
A few % for heavy quarks



EIC vs Current limits

*Single Operator Analysis

$\sqrt{s} = 141 \text{ GeV}$, $\mathcal{L} = 100 \text{ fb}^{-1}$ @EIC



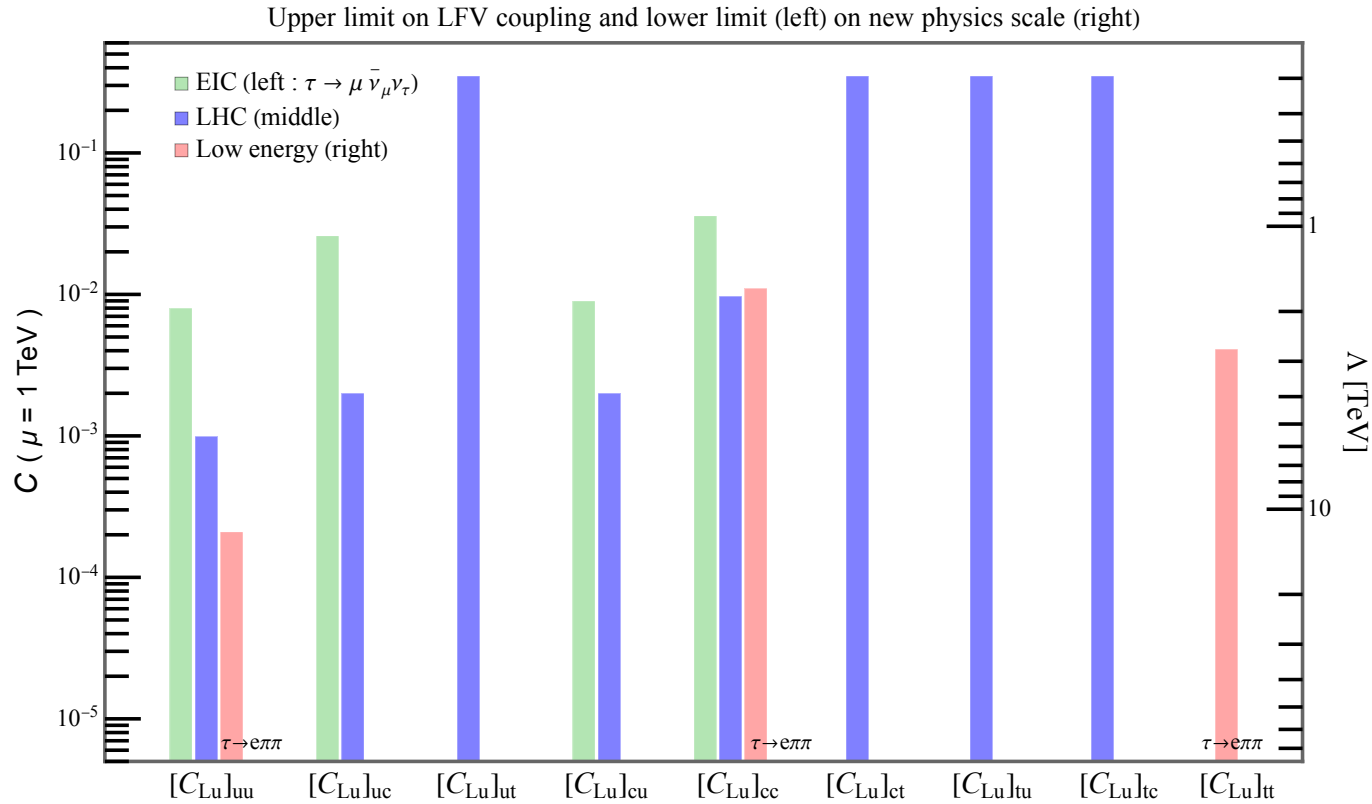
$$[C_{Ld}]_{ij} \bar{\tau}_L \gamma^\mu e_L \bar{d}_{Ri} \gamma_\mu d_{Rj}$$

- Overall, stronger limits from low-energy and LHC
- Possibility that the EIC can compete is in $[C_{Ld}]_{bb}$

EIC vs Current limits

*Single Operator Analysis

$\sqrt{s} = 141 \text{ GeV}$, $\mathcal{L} = 100 \text{ fb}^{-1}$ @EIC

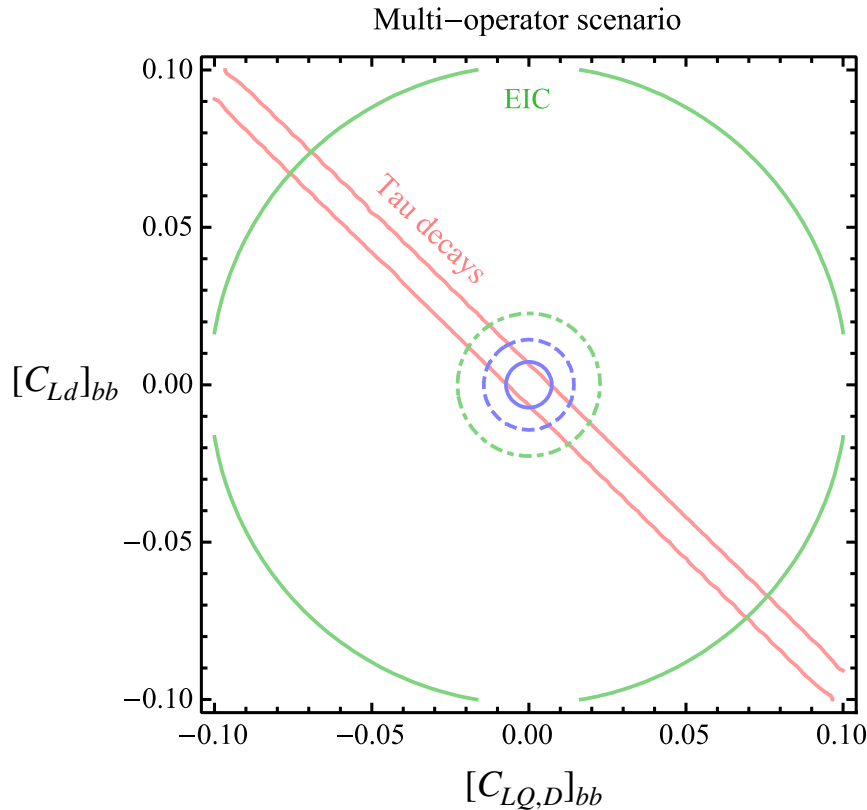


$$[C_{Lu}]_{ij} \bar{\tau}_L \gamma^\mu e_L \bar{u}_{Ri} \gamma_\mu u_{Rj}$$

- Overall, stronger limits from low-energy and LHC
- Possibility that the EIC can compete is in $[C_{Ld}]_{bb}$ and $[C_{Lu}]_{cc}$

Multi-operator scenario

S. Banerjee, V. Cirigliano, et al,
Snowmass White Paper, 2203.14919



*Case with 8 nonzero CLFV operators

Z couplings + down-type 4F operators

$$\mathcal{L}_{\text{LFV}} \supset -\frac{g_2}{c_W} \left(c_{L\varphi}^{(1)} + c_{L\varphi}^{(3)} \right) \bar{\tau}_L \gamma^\mu Z_\mu e_L$$

$$-\frac{4G_F}{\sqrt{2}} \sum_{a=d,s,b} [C_{Ld}]_{aa} \bar{\tau}_L \gamma^\mu e_L \bar{d}_{Ra} \gamma_\mu d_{Ra}$$

$$-\frac{4G_F}{\sqrt{2}} \sum_{a=d,s,b} [C_{LQ,D}]_{aa} \bar{\tau}_L \gamma^\mu e_L \bar{d}_{La} \gamma_\mu d_{La}$$

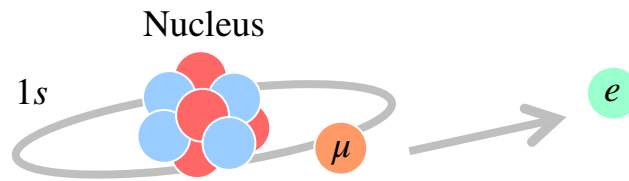
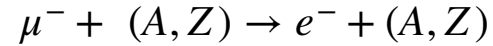
- Collider probes are necessary to close the free direction.

*Need a numerical approach to deal with multi-coupling scenario

What about $e \rightarrow \mu$ case?

What about $e \rightarrow \mu$ case?

Muon to Electron conversion in Muonic Atom



*Mono-energetic electron

$$E_e = m_\mu - B_\mu - E_{\text{rec}}$$

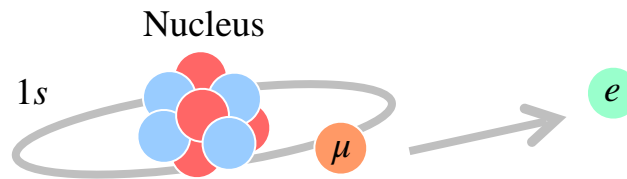
$$B_\mu \simeq \frac{(Z\alpha_{\text{em}})^2}{2} m_\mu$$

What about $e \rightarrow \mu$ case?

Muon to Electron conversion in Muonic Atom

$$\mu^- + (A, Z) \rightarrow e^- + (A, Z)$$

*Mono-energetic electron



$$E_e = m_\mu - B_\mu - E_{\text{rec}}$$

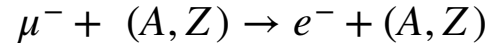
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Y. Kuno (for the COMET Collaboration), Prog.Theor. Phys. 2013, 022C01 (2013).

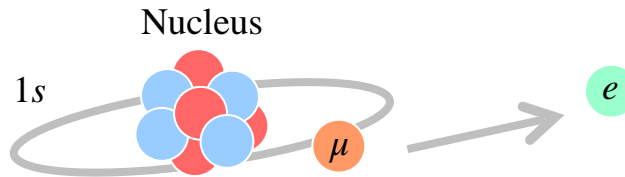
Process	$\Gamma_{\text{capt}} \times 10^{-15} \text{ [MeV]}$	Upper Bound
$\mu^- + \text{Ti} \rightarrow e^- + \text{Ti}$	1.705	6.1×10^{-13}
$\mu^- + \text{Au} \rightarrow e^- + \text{Au}$	8.603	7×10^{-13}

What about $e \rightarrow \mu$ case?

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$\mu^- + \text{Al} \rightarrow e^- + \text{Al}$	0.463	Expected $\mathcal{O}(10^{-17})$

Aluminum Target
@ Mu2E and COMET

* Improvement by 4 orders of magnitude

Low-Energy Muon and Meson Decay

34

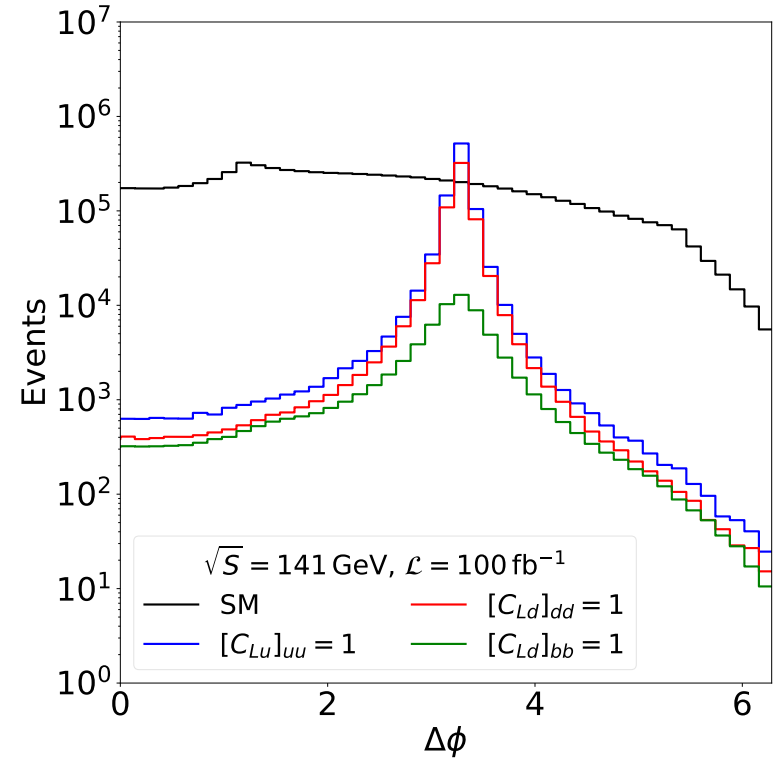
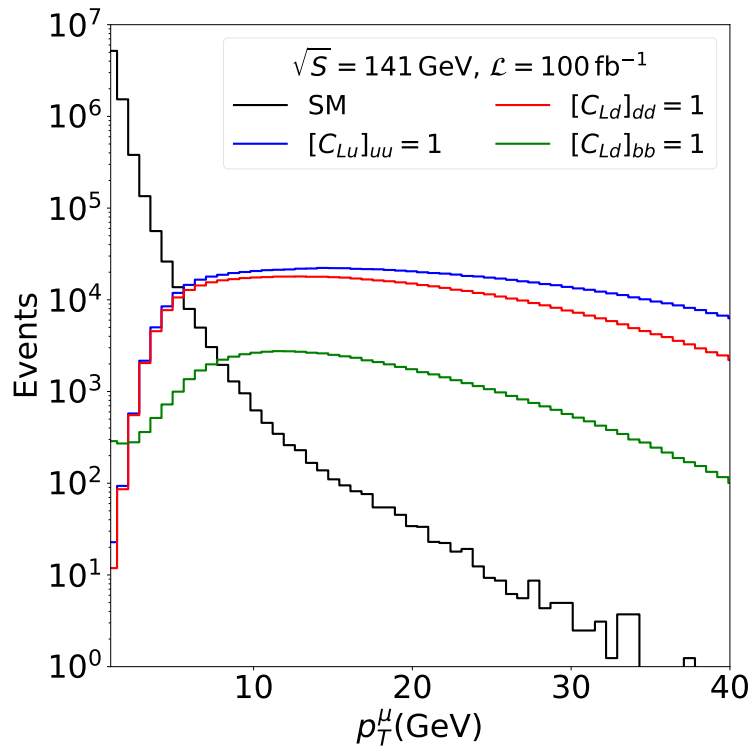
Decay mode		Upper limit (90 % C . L.)	
$\mu + \text{Ti} \rightarrow e + \text{Ti}$		6.1×10^{-13}	SINDRUM – II CPC 980420 (1998) 534
$\mu \rightarrow 3e$		1.0×10^{-12}	SINDRUM NPB 299 (1988)1
$K_L^0 \rightarrow e^\pm \mu^\mp$		4.7×10^{-12}	E871 PRL81(1998)5734
$K_L^0 \rightarrow \pi^0 e^\pm \mu^\mp$	ds/ds	7.6×10^{-11}	KTeV PRL100(2008)131803
$K^+ \rightarrow \pi^+ e^+ \mu^-$		6.6×10^{-11}	NA62 PRL127(2021)131802
$K^+ \rightarrow \pi^+ e^- \mu^+$		1.3×10^{-11}	E865 PRD72(2005)012005
$D^0 \rightarrow e^\pm \mu^\mp$		1.3×10^{-8}	LHCb PLB754(2016)167
$D^0 \rightarrow \pi^0 e^\pm \mu^\mp$		8.0×10^{-7}	BaBar PRD101(2020)112003
$D^+ \rightarrow \pi^+ e^+ \mu^-$	uc/cu	2.1×10^{-7}	LHCb JHEP06(2021)044
$D^+ \rightarrow \pi^+ e^- \mu^+$		2.2×10^{-7}	LHCb JHEP06(2021)044
$D_s^+ \rightarrow K^+ e^+ \mu^-$		7.9×10^{-7}	LHCb JHEP06(2021)044
$D_s^+ \rightarrow K^+ e^- \mu^+$		5.6×10^{-7}	LHCb JHEP06(2021)044
$B^0 \rightarrow e^\pm \mu^\mp$		1.0×10^{-9}	LHCb JHEP03(2018)078
$B^+ \rightarrow \pi^+ e^\pm \mu^\mp$	db/bd	1.7×10^{-7}	BaBar PRL99(2007)051801
$B^+ \rightarrow K^+ e^- \mu^+$		6.4×10^{-9}	LHCb PRL123(2019)241802
$B^+ \rightarrow K^+ e^+ \mu^-$	sb/bs	7.0×10^{-9}	LHCb PRL123(2019)241802
$B_s \rightarrow e^\pm \mu^\mp$		5.4×10^{-9}	LHCb JHEP03(2018)078

Low-Energy Muon and Meson Decay

35

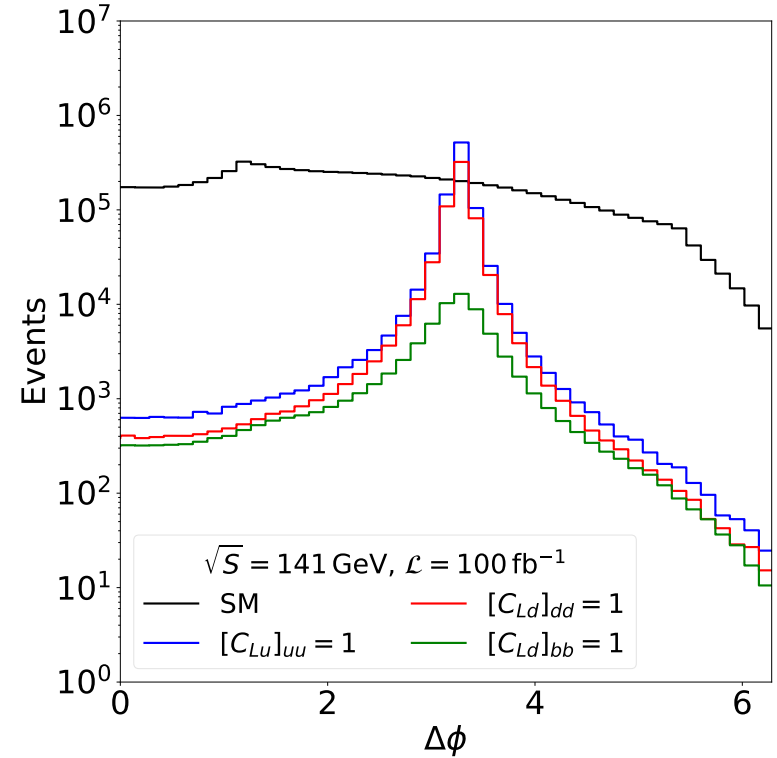
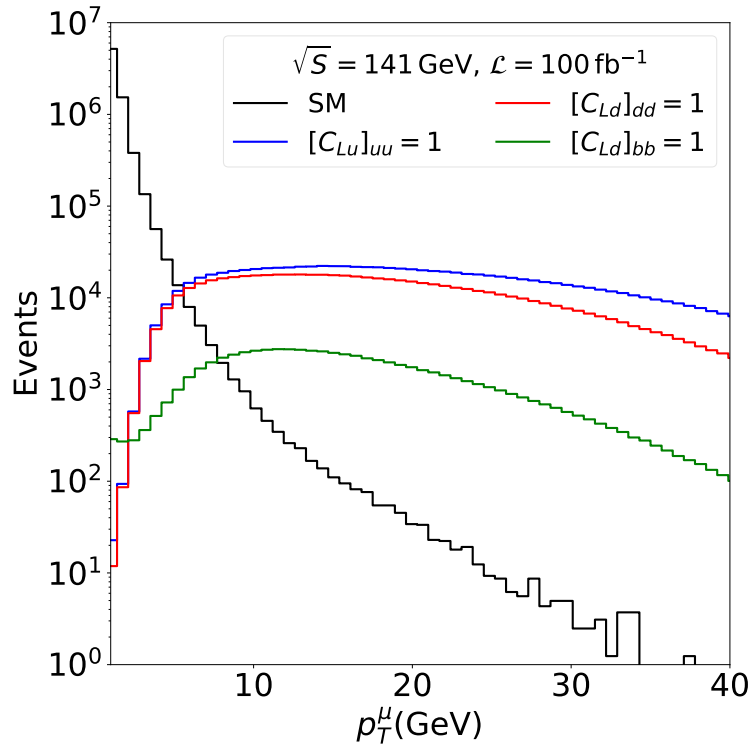
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Require one muon with high p_T^μ and one jet, not to have electrons with high p_T^e



SM BGs (black line) : Small p_T^μ and little azimuthal correlation between muon and jet

Require one muon with high p_T^μ and one jet, not to have electrons with high p_T^e



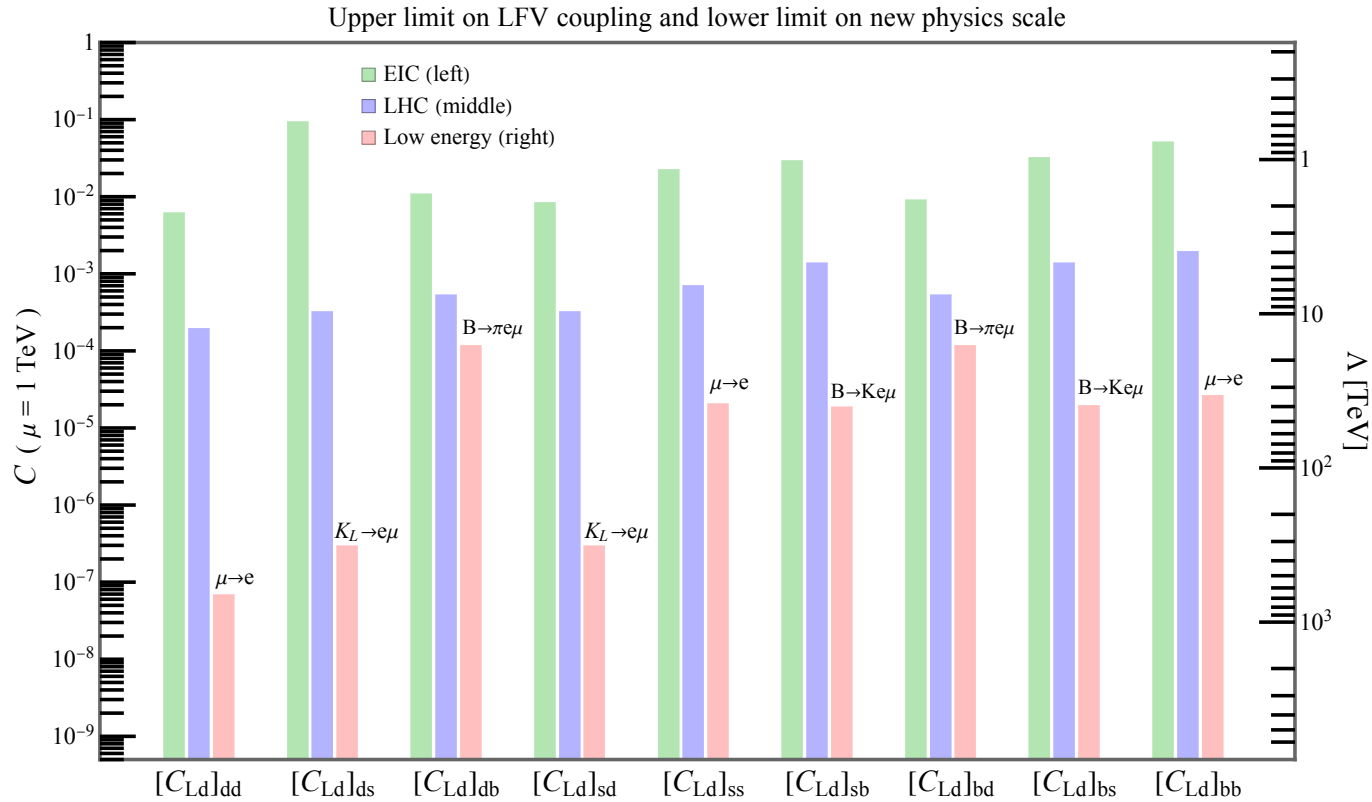
$$p_T^\mu > 20 \text{ GeV}, \quad p_T^e < 5 \text{ GeV}, \quad p_T^j > 20 \text{ GeV}, \quad 2 < \Delta\phi < 4$$

EIC vs Current limits

F. Delzanno, KF, S. Gonzalez-Solis, E. Mereghetti
arXiv 2411.13497 (Accepted by JHEP)

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* Single Operator Analysis

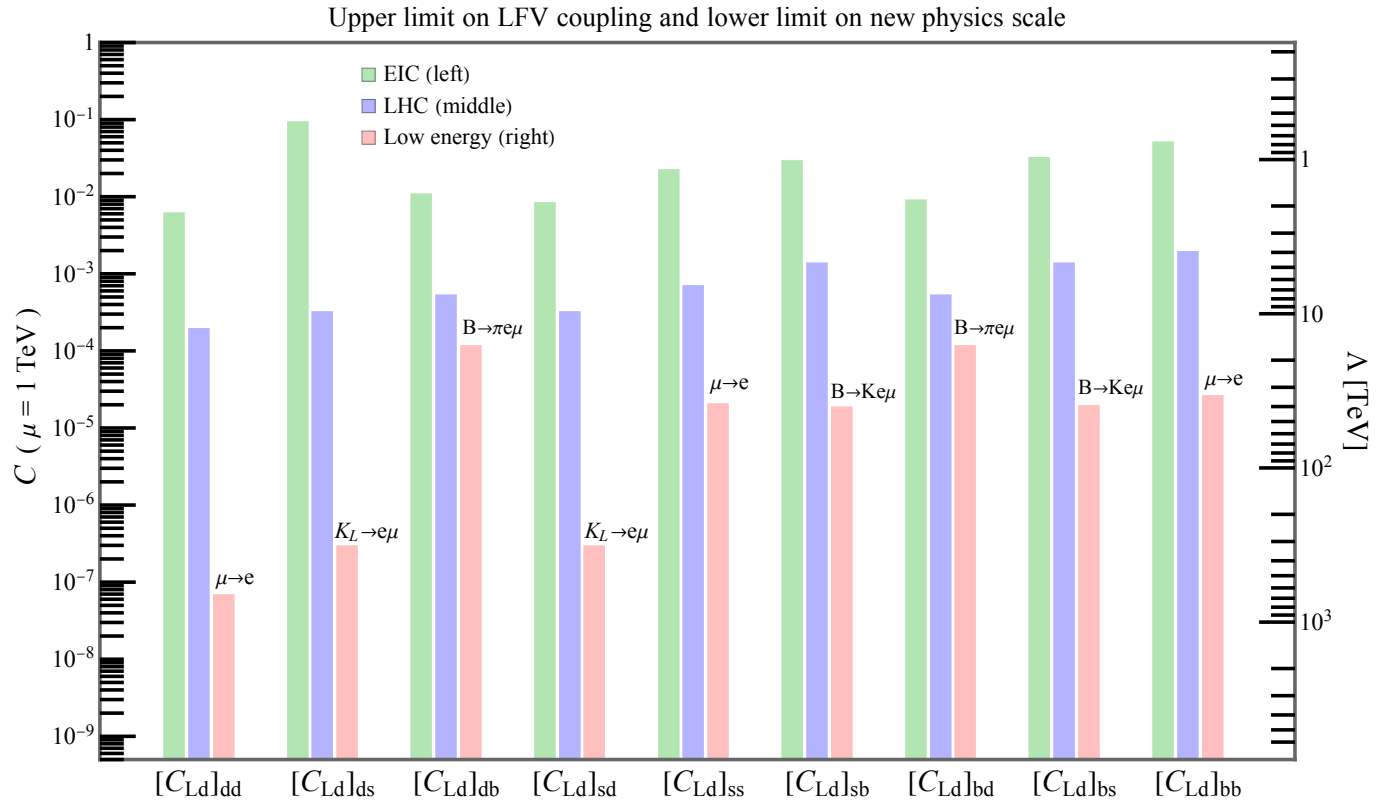


$[C_{Ld}]_{ij} \bar{\tau}_L \gamma^\mu e_L \bar{d}_{Ri} \gamma_\mu d_{Rj}$ • $\mu \rightarrow e$ conversion currently gives strong bound

EIC vs Current limits

F. Delzanno, KF, S. Gonzalez-Solis, E. Mereghetti
arXiv 2411.13497 (Accepted by JHEP)

* Single Operator Analysis



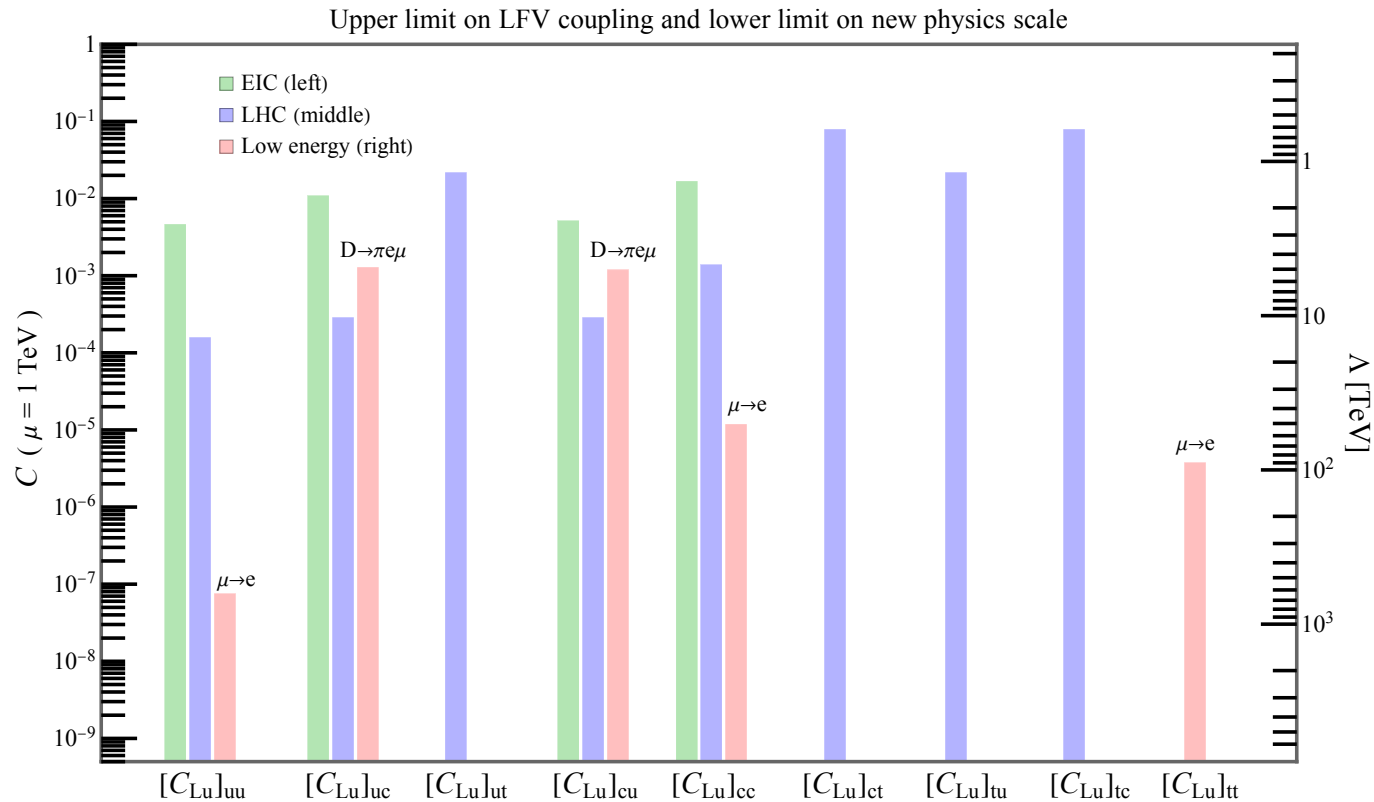
$$\frac{4G_F}{\sqrt{2}}C = \frac{1}{\Lambda^2}$$

Lower bound on the scale : $\Lambda \gtrsim \mathcal{O}(100)$ TeV from $[C_{Ld}]_{dd}$

EIC vs Current limits

F. Delzanno, KF, S. Gonzalez-Solis, E. Mereghetti
arXiv 2411.13497 (Accepted by JHEP)

* Single Operator Analysis



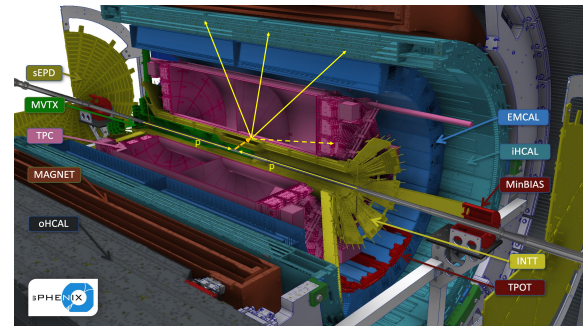
$[C_{Lu}]_{ij} \bar{\tau}_L \gamma^\mu e_L \bar{u}_{Ri} \gamma_\mu u_{Rj}$ • A factor of 10 weaker bound on $[C_{Lu}]_{cu}$ at the EIC

Study on Experimental Side

Thanks to Ming Liu!

Study tau tagging using sPHENIX p+p data and utilizing AI/ML algorithms that can optimize the tagging efficiency.

- Displaced tau-decay vertex
- Narrow jet from tau hadronic decay
- Single muon/hadron from tau decays

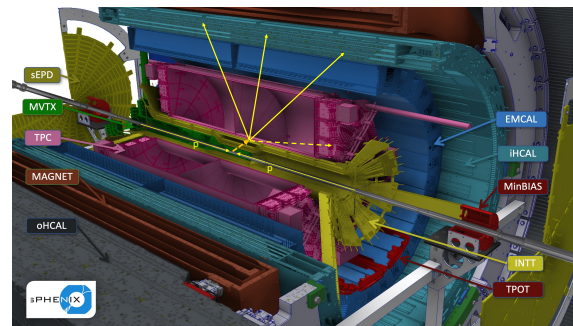


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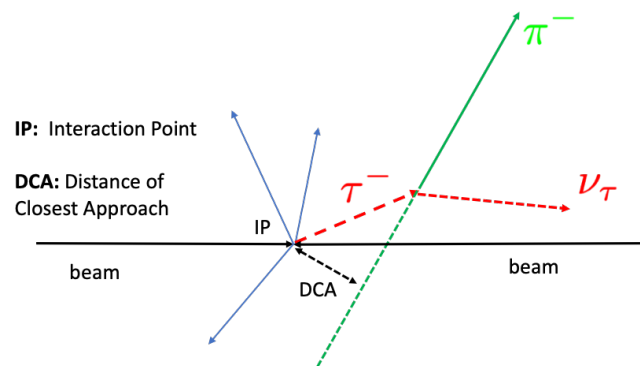


sPHENIX : proton+proton collisions data from Run2024

- Identify the displaced τ decay vertex from the collision point (called primary vertex)
- Simulation of $p+p \rightarrow \text{Upsilon} \rightarrow \text{Tau} + \text{Tau}^-$

EIC: e+proton collisions

Simulations of ep collisions from the ePIC detector setup



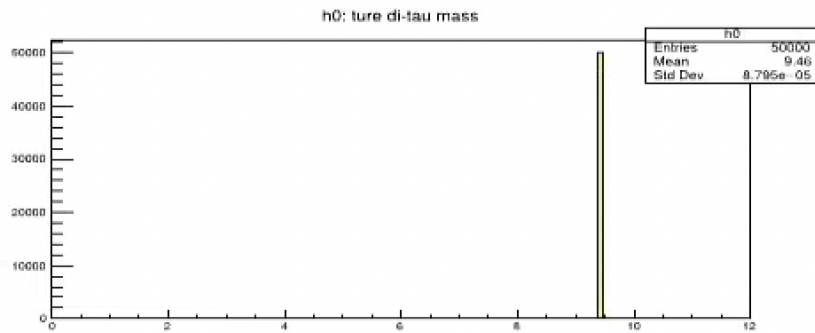
Study on Experimental Side

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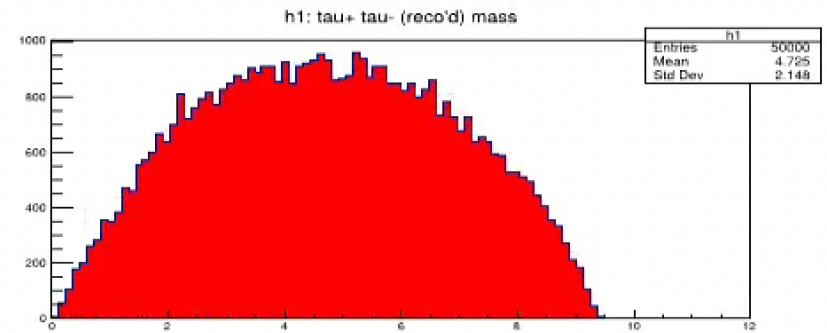
43

Tau Tagging @ sPHENIX: $pp \rightarrow \Upsilon \rightarrow \tau^- \tau^+$

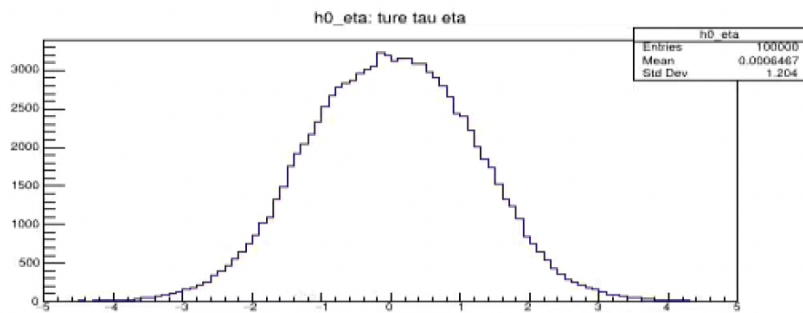
Invariant mass of Tau+ and Tau-:



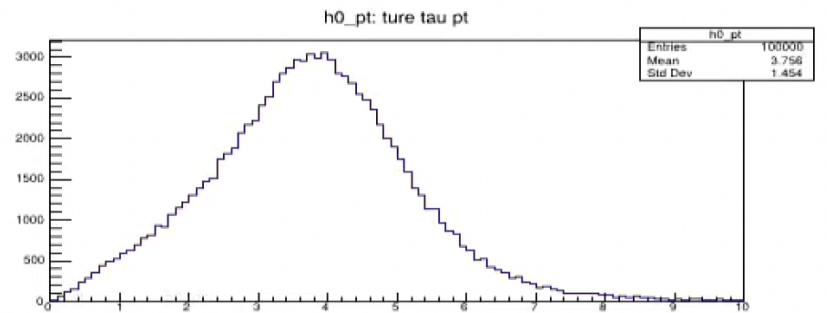
True Tau kinematics (p,e):
Mass = 9.46GeV



Reconstructed Tau kinematics (p,e),
missing neutrinos: <mass> ~5 GeV



η



p_T

Study on Experimental Side

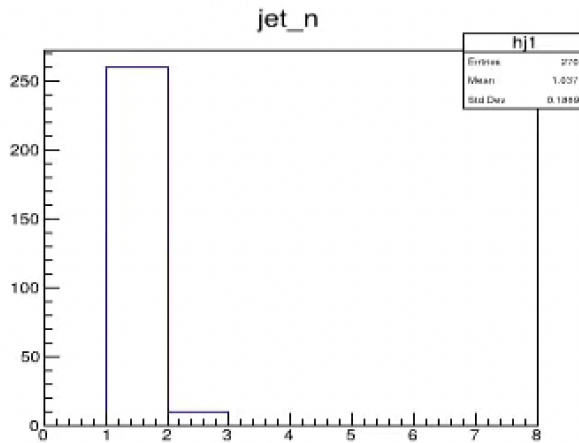
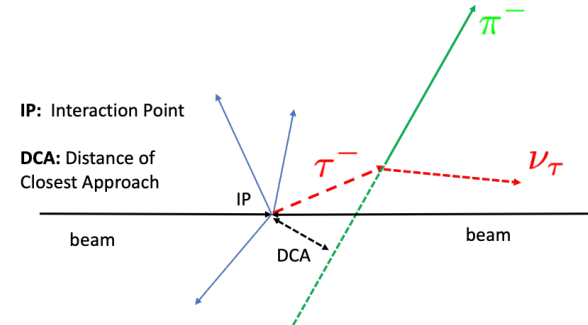
Thanks to Ming Liu!

44

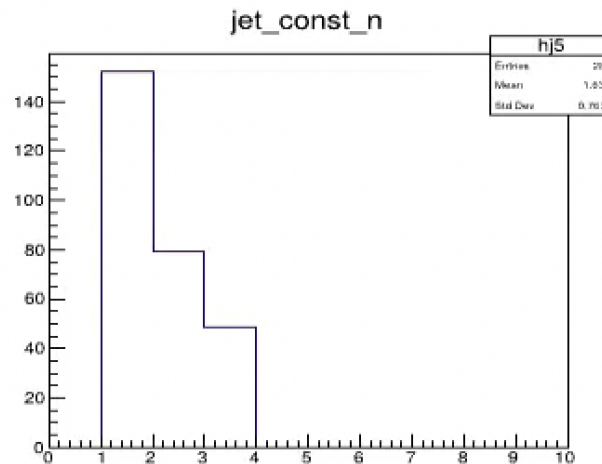
Tau sample selection: $p_T(\text{ture}) > 5\text{GeV}$

Tau-jet reconstruction:

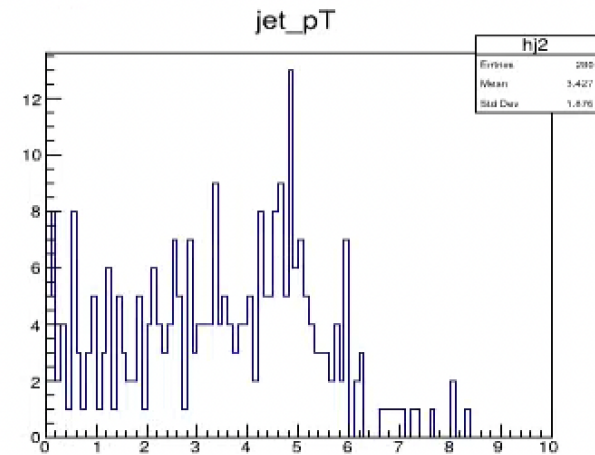
- 1) reconstructed from “visible tracks” with $\text{DCA} > 0$
- 2) anti- k_T , $R = 1.0$



Mostly single jet:
Tracks with $\text{DCA} > 0$



of particles in jets



“Reconstructed” jet p_T ,
w/o neutrinos

Summary

✦ Searches for Lepton Flavor Violations are Powerful Probes of BSM Physics.

- Systematic Analysis based on SMEFT
 - The RGEs allow to constrain CLFV heavy quark operators
- Operators involving b and c in $e - \tau$ case are promising at the EIC
- Collider searches are essential in multi-operator scenarios
- Strong bound in $e - \mu$ case especially from $\mu \rightarrow e$ conversion

Outlook/Discussion

- Multi-Operator Analysis using Machine Learning
- Experimental study of tau lepton tagging
- LNV / Sterile Neutrino Searches at the EIC

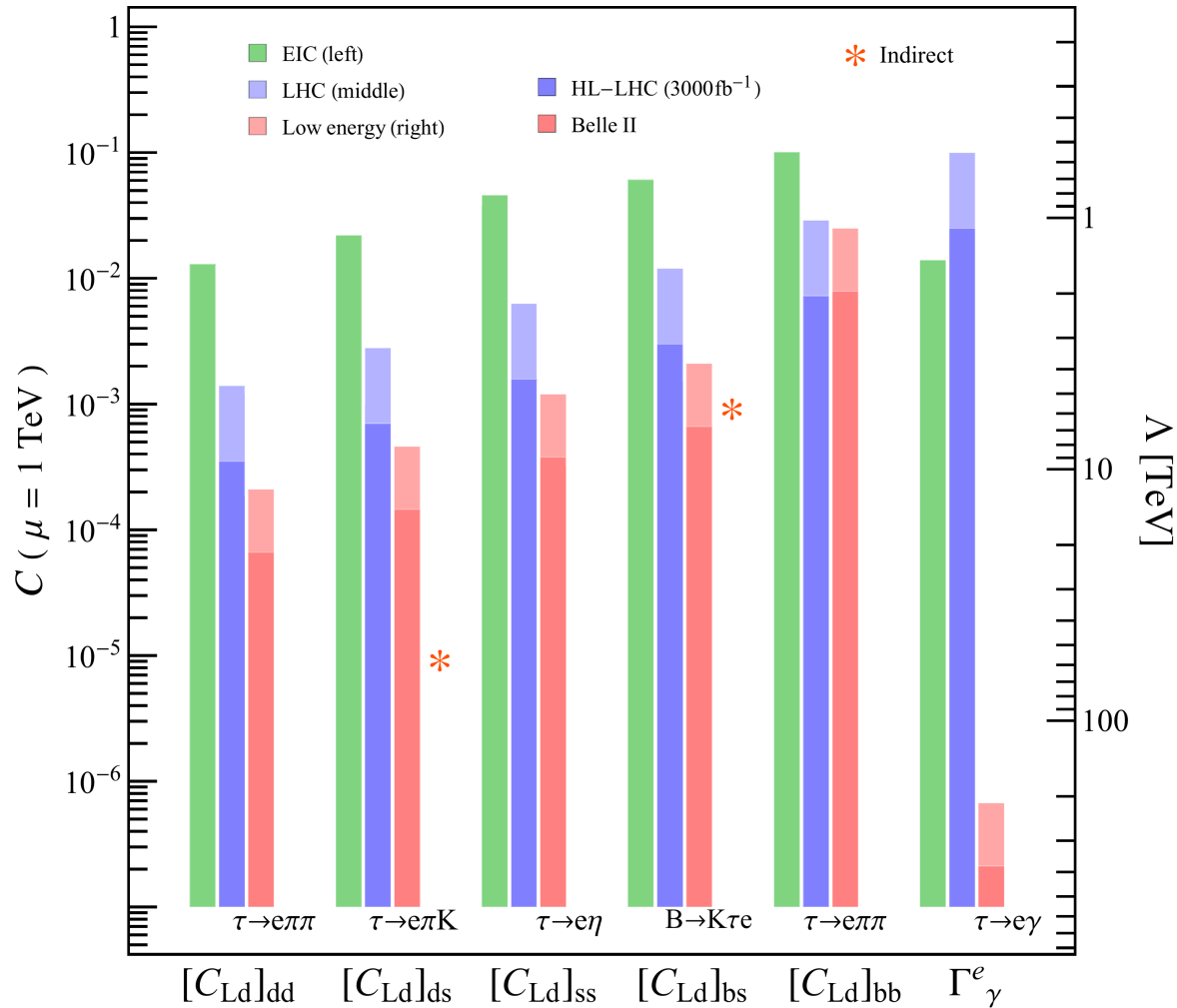
Thank you for your attention!



Backup slides

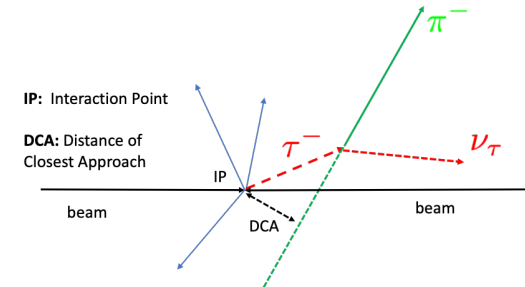
Apple to Apple

S. Banerjee, V. Cirigliano, et al,
Snowmass White Paper, 2203.14919



How to tag Tau leptons experimentally?

- Exploit tau-lepton decay topology
 - Displaced track(s) from tau decay
 - Pencil-like isolated track(s)
- Study tau-lepton tagging algorithms using sPHENIX p+p data,
 - Identify displaced tracks with silicon pixel detectors (MVTX)
 - Tag pencil-like “jets” with EMCal and HCal
- EIC/ePIC detector simulations
 - Tau-lepton tagging with algorithms developed from sPHENIX data



Search for:

- “Tau-> nv_tau + pi+”, BR = 11.5%,
displaced isolated single track
- “Tau -> pi+ 2pi0 + nv_tau”, BR = 9.5%,
pencil-like jets
- “Tau-> muon/e + nv_mu/e + nv_tau”, BR = 17.8%,
displaced isolated single track

Benchmark performance:

$$\Upsilon \rightarrow \tau^+ + \tau^-$$

broad mass of di-hadron and di-lepton of “tau candidates”

Decay mode	Resonance	$\mathcal{B}$ (%)
Leptonic decays		35.2
$\tau^- \rightarrow e^- \bar{\nu}_e \nu_\tau$		17.8
$\tau^- \rightarrow \mu^- \bar{\nu}_\mu \nu_\tau$		17.4
Hadronic decays		64.8
$\tau^- \rightarrow h^- \nu_\tau$		11.5
$\tau^- \rightarrow h^- \pi^0 \nu_\tau$	$\rho(770)$	25.9
$\tau^- \rightarrow h^- \pi^0 \pi^0 \nu_\tau$	$a_1(1260)$	9.5
$\tau^- \rightarrow h^- h^+ h^- \nu_\tau$	$a_1(1260)$	9.8
$\tau^- \rightarrow h^- h^+ h^- \pi^0 \nu_\tau$		4.8
Other		3.3

*Study tau tagging algorithms using sPHENIX p+p data and utilize ML to optimize the tagging efficiency.