



Extracting α_s using spin Structure function at ECCE

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Inclusive Deep inelastic scattering



Resolution 4-momentum transfer square $Q^2 = -q^{\mu}q_{\mu} = |\vec{q}|^2 - \omega^2$

Number involved nucleons $x_b = Q^2/2m_p\omega$

Structure functions:

□ $F_{1,2}$: Unpolarized Structure functions □ $g_{1,2}$: Polarized Structure functions

Approaches to extract α_s from spin structure functions

 $\Box Q^2$ -evolution of $g_1(x, Q^2)$: Complex task

- Involves DGLAP global fit,
- non-perturbative inputs: quark and gluon distributions,
- possibly higher-twists for low Q²/ large-x data.

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 $\Box Q^2$ -evolution of moment $\int g_1(x, Q^2) dx$: Simpler

- No *x*-dependence,
- Non-perturbative inputs: more-or-less well measured axial charges a_0 , a_3 and a_8 +) possibly higher-twists for low- Q^2 data).
- Issues: unmeasurable low-*x* contribution, a_0 is Q^2 dependent and may have contribution from gluon ΔG pdf (but not the case in \overline{MS})

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 $\Box Q^2$ -evolution of iso-vector moment $\int g_1^{p-n}(x,Q^2)dx$: Simplest

- Axial charge $a_3 = g_A$ precisely measured ($g_A = 1.2762 \pm 0.0005$)
- DGLAP-evolution known to higher order than single nucleon case
- No gluon contribution.
- Issue: But low-x issue and demands measurement on polarized p and n.

Bjorken Sum Rule

$$\Gamma_1^{p-n} \equiv \int g_1^{p-n} dx = \frac{1}{6} g_A \left[1 - \frac{\alpha_s}{\pi} - 3.58(\frac{\alpha_s}{\pi})^2 - 20.21(\frac{\alpha_s}{\pi})^3 - 175.7(\frac{\alpha_s}{\pi})^4 - \sim 893(\frac{\alpha_s}{\pi})^5 \right]$$

Nucleon's First Nucleon Axial charge pQCD radiative correction spin structure $[\Gamma_1^{p-n}(Q^2 \to \infty)]$ function

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Two possibilities to extract $\alpha_s(M_z)$

 \Box Do an absolute measurement of Γ_1^{p-n} and solve BJSR for $\alpha_s(Q^2)$

- One α_s per Γ_1^{p-n} experimental data point
- Poor systematic uncertainty
- Typical $\frac{\Delta \alpha_s}{\alpha_s} \sim 10\%$



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 $\Box \text{ Measurement of } Q^2 \text{ dependence of } \Gamma_1^{p-n}(Q^2)$

- Need Γ_1^{p-n} at several Q^2 points, Only one (or few) value of α_s
- Good accuracy
- 1990's CERN/SLAC data yielded: $\alpha_s(M_z) = 0.12 \pm 0.009$ *Nucl.Phys. B496 337 (1997)*

Possible future extractions of α_s from $\Gamma_1^{p-n}(Q^2)$



Nucl. Phys. A 1026, 122447



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Simulated data

- DIS events generated using DJANGOH for both *e-p* and *e-³He*
- For e-p: 5x41 GeV, 10x100 GeV, 18x275 GeV
- For e-3He: 5x41 GeV, 10x100 GeV, 18x166 GeV
- Integrated luminosity of 10 fb⁻¹

Note: Neutron structure function is extracted from e-3He using double tagging measurement that significantly minimize the nuclear correction.

□ Monte Carlo simulation using ECCE configuration

- Including detector effects: acceptance, resolution, efficiencies
- Provide psesudo data for analysis
- Data analysis
 - Asymmetries: $A_{\parallel,\perp}$, A_1
 - Polarized structure function: g_1^p , g_1^n
 - Bjorken Sum Rule: $[\Gamma_1^{p-n}(Q^2)]$
 - Ready for fitting

Uncertainties

□ Statistics

□ Systematics

- Detector effects,
- Beam polarimetries,
- Radiative corrections,
- Missing high- and low-*x* part,
- PDF parameterizations;
- Negligible: neutron information

EIC: generated pseudo-data





Measured fraction of the Bjorken sum $\Gamma_1^{p-n}(Q^2)$



Fit and procedure:

• Main fit function: Bjorken sum approximant at N⁴LO with α_s at 4-loop (i.e β_3), for main result.

$$\int \Gamma_1^{p-n} = \frac{1}{6} g_A \left[1 - \frac{\alpha_s}{\pi} - 3.58 (\frac{\alpha_s}{\pi})^2 - 20.21 (\frac{\alpha_s}{\pi})^3 - 175.7 (\frac{\alpha_s}{\pi})^4 \right]$$

$$\begin{split} \alpha_s^{\overline{\text{MS}}}(Q) &= \frac{4\pi}{\beta_0 \ln(Q^2/\Lambda_s^2)} \left[1 - \frac{\beta_1}{\beta_0^2} \frac{\ln(\ln(Q^2/\Lambda_s^2))}{\ln(Q^2/\Lambda_s^2)} + \frac{\beta_1^2}{\beta_0^4 \ln^2(Q^2/\Lambda_s^2)} (\ln^2(\ln(Q^2/\Lambda_s^2)) - \ln(\ln(Q^2/\Lambda_s^2)) - 1 + \frac{\beta_2\beta_0}{\beta_1^2}) + \frac{\beta_1^3}{\beta_0^6 \ln^3(Q^2/\Lambda_s^2)} (-\ln^3(\ln(Q^2/\Lambda_s^2)) + \frac{5}{2} \ln^2(\ln(Q^2/\Lambda_s^2)) + 2\ln(\ln(Q^2/\Lambda_s^2)) - \frac{1}{2} - 3\frac{\beta_2\beta_0}{\beta_1^2} \ln(\ln(Q^2/\Lambda_s^2)) + \frac{\beta_3\beta_0^2}{2\beta_1^3}) + \frac{\beta_1^4}{\beta_0^8 \ln^4(Q^2/\Lambda_s^2)} (\ln^4(\ln(Q^2/\Lambda_s^2)) - \frac{13}{3} \ln^3(\ln(Q^2/\Lambda_s^2)) - \frac{3}{2} \ln^2(\ln(Q^2/\Lambda_s^2)) + 4\ln(\ln(Q^2/\Lambda_s^2)) + \frac{7}{6} + \frac{7}{6} + \frac{3\beta_2\beta_0}{\beta_1^2} (2\ln^2(\ln(Q^2/\Lambda_s^2)) - \ln(\ln(Q^2/\Lambda_s^2)) - 1) - \frac{\beta_3\beta_0^2}{\beta_1^3} \left(2\ln(\ln(Q^2/\Lambda_s^2)) + \frac{1}{6} \right) \right] \end{split}$$

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a)

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• Secondary fit at N⁵LO and α_s at 5-loop, for pQCD truncation uncertainty.

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- Secondary fit at N⁵LO and α_s at 5-loop, for pQCD truncation uncertainty.
- Systematically vary fit Q² range to minimize total uncertainty: Low Q² points have high α_s sensitivity but larger pQCD truncation error. High Q² points have smaller α_s sensitivity but smaller pQCD error. May not be worth including the lowest and/or highest Q² points. (Not worth using all points for statistics sake since stat. error is negligible.)
- 2-parameter fit:
 - Λ_s is the free parameter of interest. From it, we obtain $\alpha_s(M_z)$.
 - \circ g_A Well-known but left as a free to account for normalization uncertainties.

 $1.31\% = 0.83\%(\text{exp.}) \oplus 0.64\%(\text{truncation}) \oplus 0.78\%(\text{polarimetries})$



Compared to other DIS results and world average (from PDG)



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Conclusion:

- Realistic simulation shows that EIC can yield a competitive measurement.
- Just one method. Other extractions will be available, e.g.:
- Global fits (unpolarized and polarized)
- Inclusive neutral current reactions (EIC+HERA). S. Cerci, *et al.* EPJC, 83(11):1011, 2023: $\Delta \alpha_s(M_Z) = 0.4\%$