

Polarized PDFs (hPDF): JAM perspective

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In collaboration with
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**Extracting the Strong Coupling at the EIC
and other Future Colliders**

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$$\mathcal{L}_{\text{QCD}} = \sum_q \bar{\psi}_q (i \not{D} - m_q) \psi_q - \frac{1}{2} \text{Tr}[G_{\mu\nu} G^{\mu\nu}]$$

Recent developments

Sign of gluon polarization

Zhou, NS, Melnitchouk '22

Karpie, Whitehill, Melnitchouk, Monahan, Orginos, Qiu, Richards, NS, Zafeiropoulos '23

de Florian, Forte, Vogelsang '24

Hunt-Smith, Cocuzza, Melnitchouk, NS, Thomas, White '24

Small x formalism for hPDF global analysis

Adamiak, Kovchegov, Melnitchouk, Pitonyak, NS, Sievert '21

Adamiak, Baldonado, Kovchegov, Melnitchouk, Pitonyak, NS, Sievert, Tarasov, Tawabutr '23

Adamiak, Baldonado, Kovchegov, Li, Melnitchouk, Pitonyak, NS, Sievert, Tarasov, Tawabutr '25

Approximate NNLO hPDF global analysis

Borsa, Stratmann, Vogelsang, de Florian, Sassot, '25

Cruz-Martinez, Hasenack, Hekhorn, Magni, Nocera, Rabemananjara, Rojo, Sharma,

van Seeventer '25

Bertone, Chiefs, Nocera '24

Transverse Momentum

Bacchetta, Bongallino, Cerutti, Radici, Rossi, '24

hPDFs at high x

Cocuzza, Hunt-Smith, Melnitchouk, NS, Thomas,



Hadron structure

Polarized sea, sea asymmetries, gluon polarization, spin puzzle, Higher twist effects, high x asymptotics, intrinsic transverse momentum

Formalisms

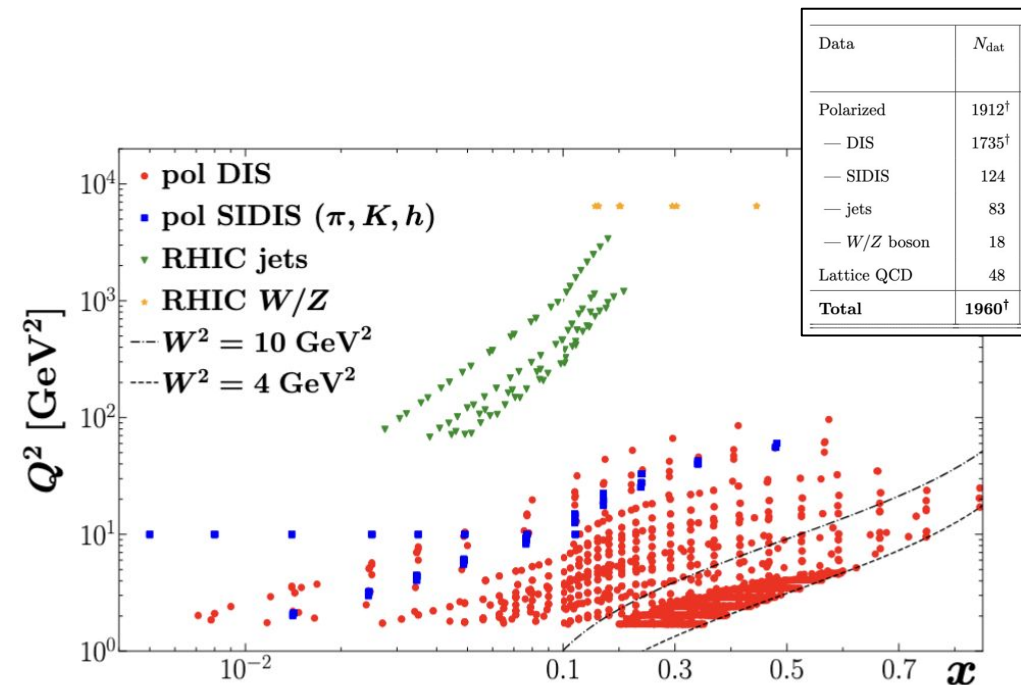
Higher order corrections, dipole formalism, target mass corrections

Methodologies

MC based approach, ANN, simultaneous extraction framework (key observables SIDIS), LQCD input

hPDFs at High x

Cocuzza, Hunt-Smith, Melnitchouk, NS, Thomas, (in prep.)



Motivations

Reconstruction of hPDFs at very high x
Simultaneous reconstruction of PDFs, hPDFs and FFs
Constrain subleading power corrections in DIS
Inclusion of LQCD

Ingredients

Perturbative accuracy: NLO, ZMVFS (NLL for PDFs and α_s)
TMC corrections using AOT
Nuclear smearing for 3He
NP modeling for power corrections

Methodology

Traditional parametrization
Mellin space techniques: no K factor approach
Ensemble approach for UQ
Multi-step strategy

Theory setup for pol. inclusive DIS

Leading Twist Collinear factorization

$$g_1^{\text{LT}}(x, Q^2) = \frac{1}{2} \sum_q e_q^2 [\Delta C_q^{\text{DIS}} \otimes \Delta q^+ + 2\Delta C_g^{\text{DIS}} \otimes \Delta g]$$

Target Mass corrections

Moffat, Roger, Melnitchouk NS, Steffens '19

$$g_1^{\text{TMC}}(x, Q^2) = \frac{1}{\rho^2} g_1^{\text{LT}}(x_N, Q^2) + \frac{2(\rho - 1)}{\rho^2} \int_{x_N}^1 \frac{dz}{z} g_1^{\text{LT}}(z, Q^2)$$

$$g_2^{\text{TMC}}(x, Q^2) = -\frac{1}{\rho^2} g_1^{\text{LT}}(x_N, Q^2) + \frac{2}{(1 + \rho)\rho^2} \int_{x_N}^1 \frac{dz}{z} g_1^{\text{LT}}(z, Q^2) \quad \rho^2 = 1 + \frac{4M^2 x^2}{Q^2}$$

Higher Twist effects

$$g_1^{\text{HT}} = \frac{c_1^{\text{HT}}}{Q^2} \quad g_2^{\text{HT}} = c_2^{\text{HT}}$$

Final expressions

$$g_i = g_i^{\text{TMC}} + g_i^{\text{HT}} \quad i = 1, 2$$

For Light nuclei DIS

$$g_i^A(x, Q^2) = \sum_N \left[\Delta f_{ij}^{N/A} \otimes g_j^N \right] (x, Q^2)$$

Parametrization & optimization

$$T(x, \mu^2) = \frac{N}{\mathcal{M}} x^\alpha (1-x)^\beta (1 + \gamma\sqrt{x} + \eta x)$$



Input hpdfs + power corrections

$$\Delta u = \underline{\Delta u_v} + \Delta \bar{u}$$

$$\Delta d = \underline{\Delta d_v} + \Delta \bar{d}$$

$$\Delta \bar{u} = \underline{\Delta S} + \underline{\delta \bar{u}}$$

$$\Delta \bar{d} = \underline{\Delta S} + \underline{\delta \bar{d}}$$

$$\Delta s = \Delta \bar{s} = \Delta S$$

$$\underline{c_i^{\text{HT}}}$$



Input pdfs +
power
corrections



Input pion
FFs



Input Kaon
FFs

Ensemble approach: data resampling + optimization

Multi-step strategy: sequential prior-posterior loops with increasing data aggregation

Quality metrics: reduced chi2 + Z scores based on chi2 distribution

Sensitivity on DIS cuts

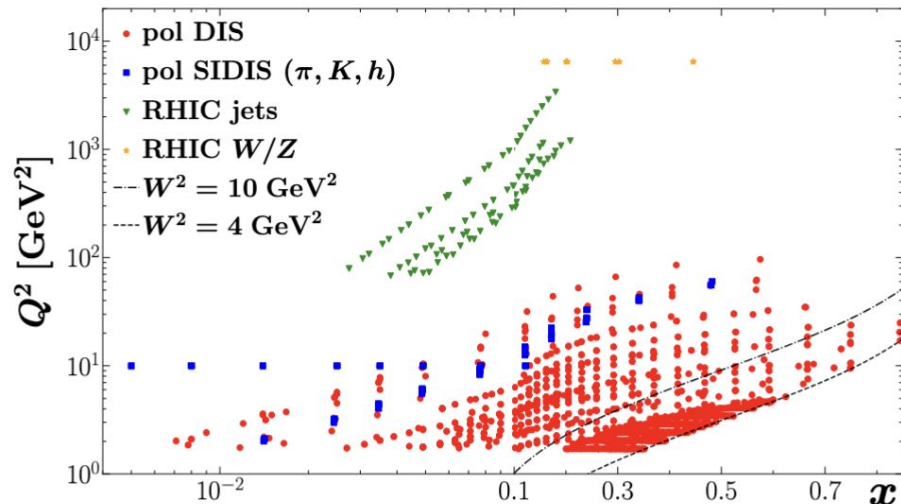


W_{\min}^2 (GeV 2)	3.5	4	5	6	10
N_{dat}	2002	1735	1287	1008	689
χ_{red}^2	1.10	1.02	0.97	1.00	1.00
Z-score	+3.11	+0.60	-0.84	+0.02	+0.00

Multiple analysis with different W_{\min} cuts

$W_{\min} \sim 4$ is the minimum we can push to obtain stable results

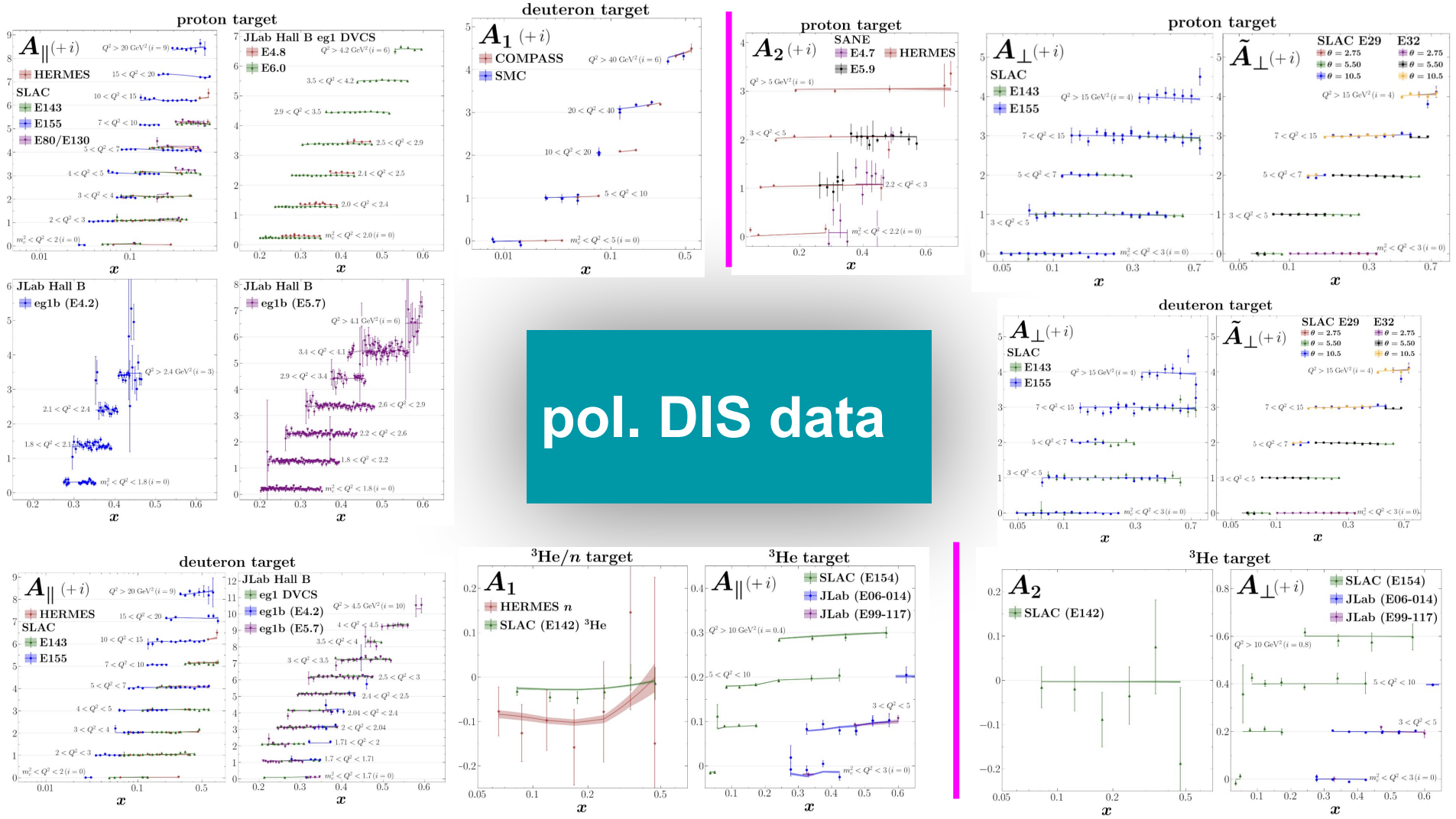
Increase in χ^2 for $W_{\min} \sim 3.5$ due to proximity to resonance region



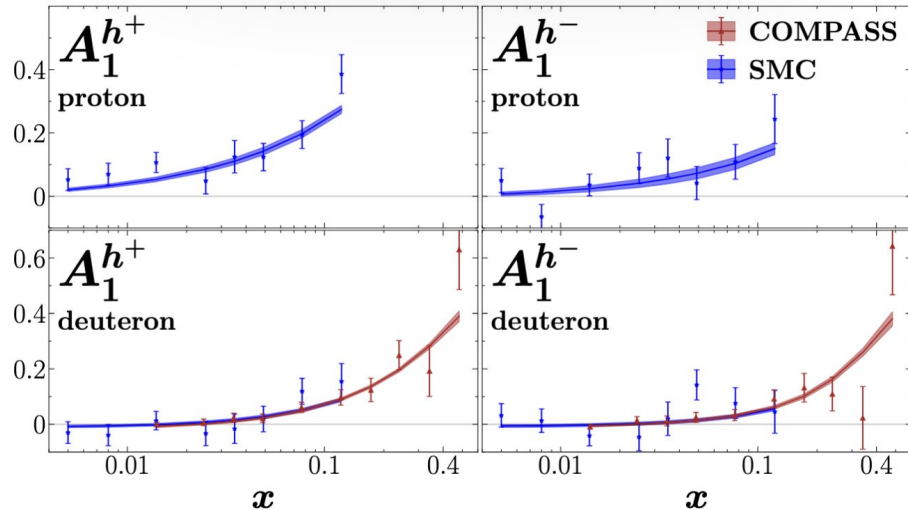
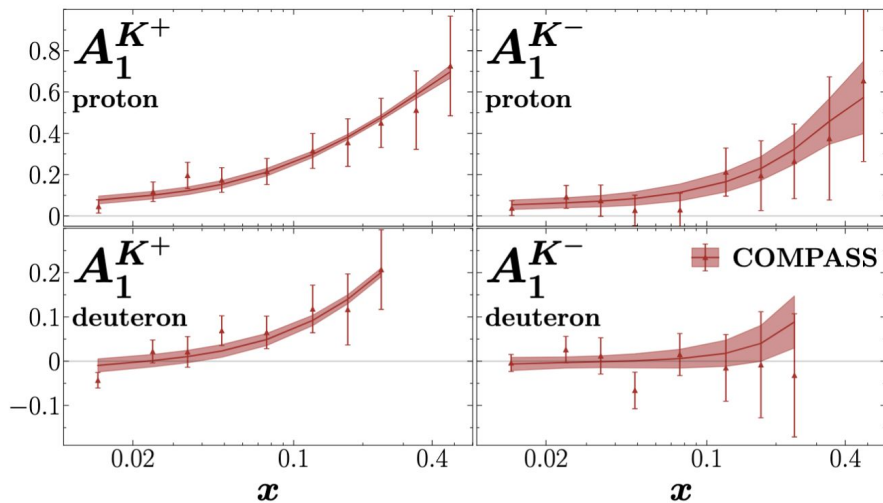
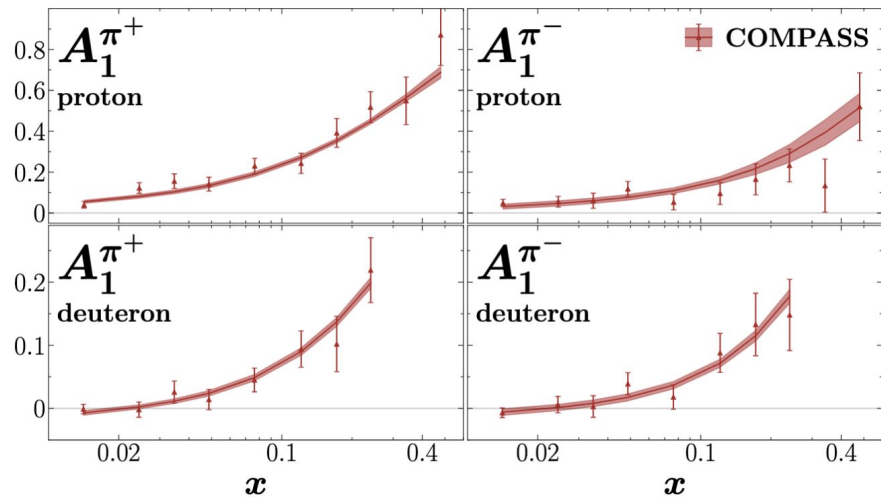
Role of HT effects

Data	N_{dat}	χ_{red}^2 (Z-score)		
		$W^2 > 10 \text{ GeV}^2$ (no HT)	$W^2 > 4 \text{ GeV}^2$ (no HT)	$W^2 > 4 \text{ GeV}^2$ (with HT)
Polarized	1912 [†]	0.86 (−3.13)	1.01 (+0.63)	0.98 (−0.61)
— DIS	1735 [†]	1.01 (+0.28)	1.04 (+1.15)	1.03 (+0.79)
— SIDIS	124	0.82 (−1.51)	0.85 (−1.21)	0.77 (−1.95)
— jets	83	0.83 (−1.08)	0.84 (−1.06)	0.83 (−1.08)
— W/Z boson	18	0.78 (−0.61)	0.82 (−0.47)	0.67 (−1.03)
Lattice QCD	48	0.61 (−2.15)	0.57 (−2.44)	0.58 (−2.40)
Total	1960[†]	0.93 (−1.54)	1.01 (+0.20)	0.99 (−0.35)

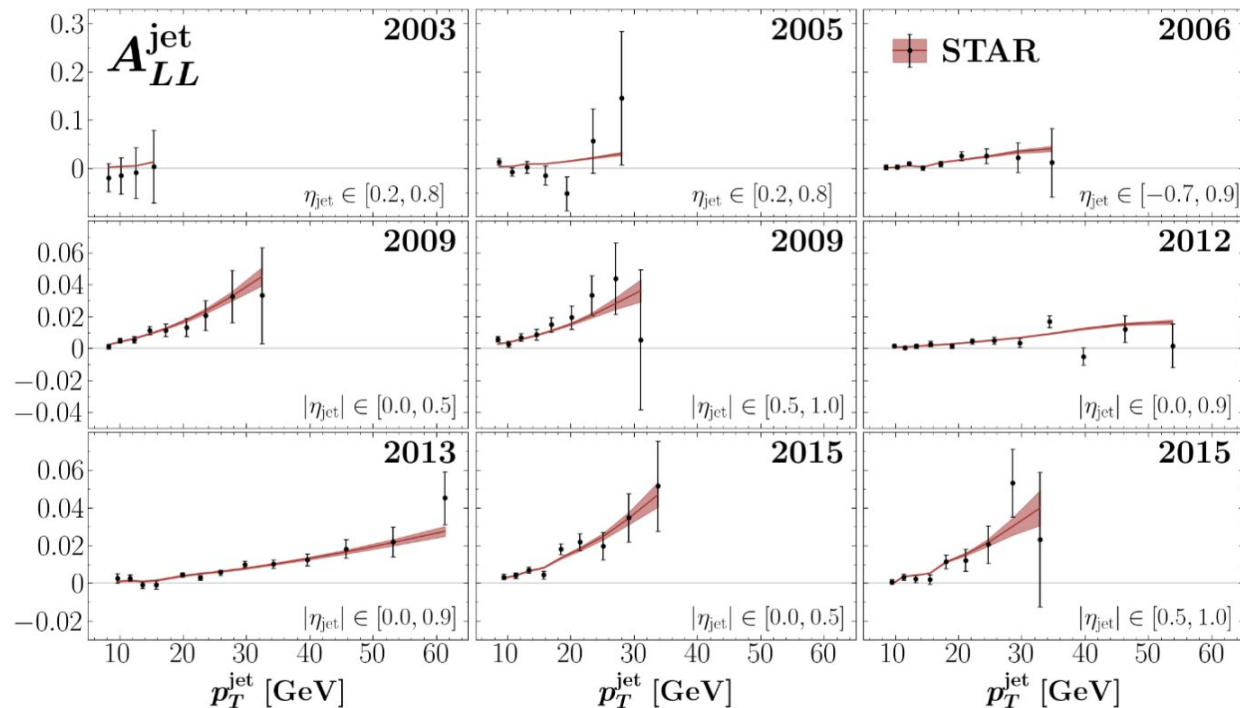
Stable global agreement when lowering $W_{2\text{min}}$
Inclusion of HT effects gives marginal difference with the quality metrics
At present, we see no strong evidence to support sizable HT effects



pol. SIDIS data

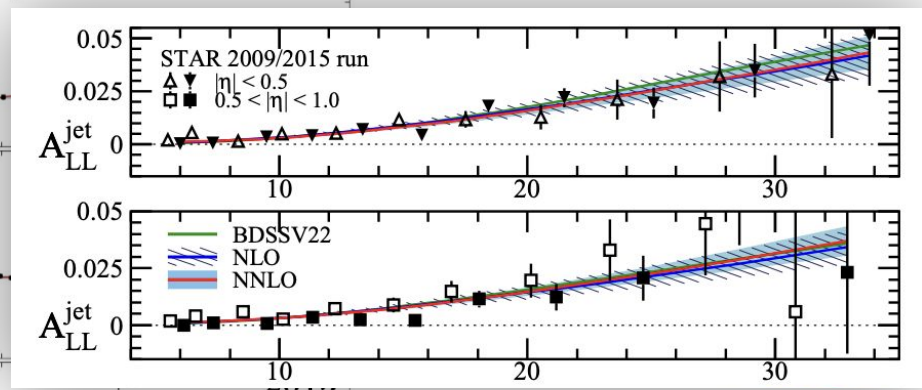
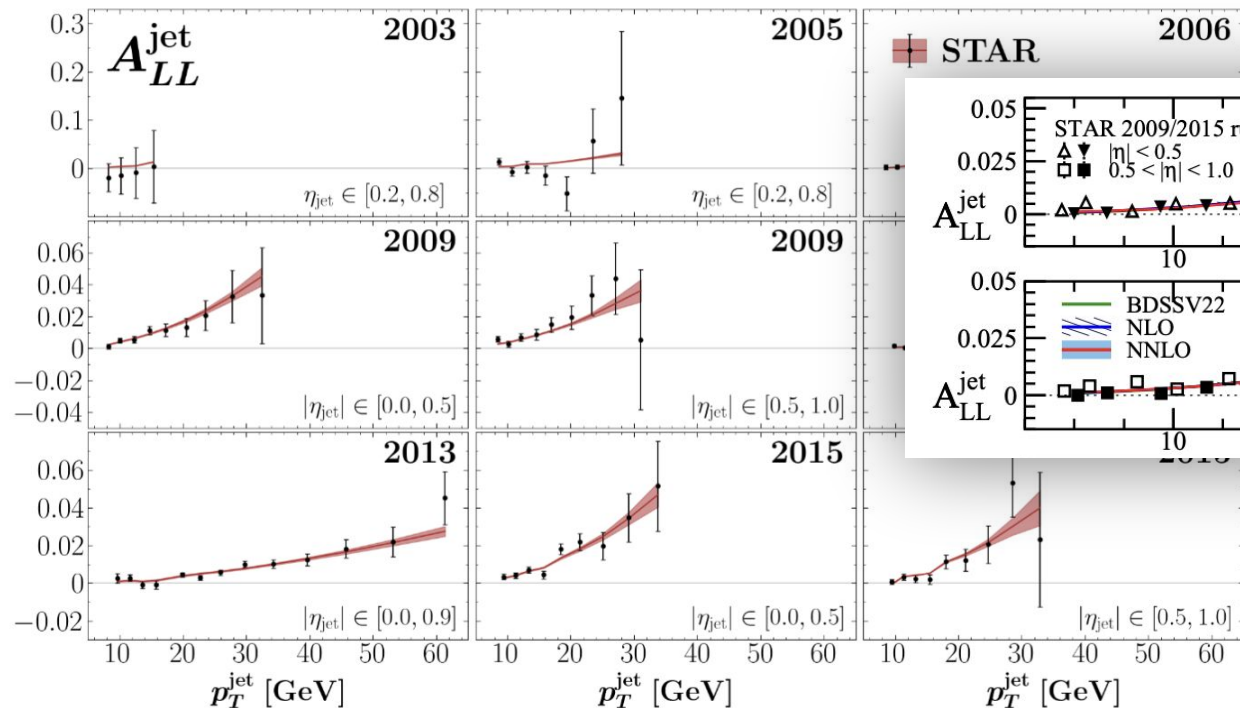


pol pp Jets



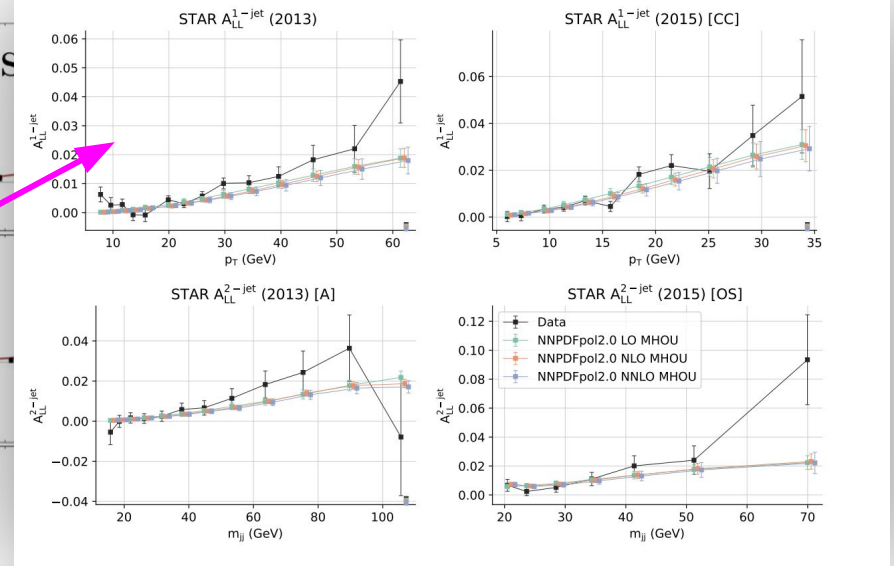
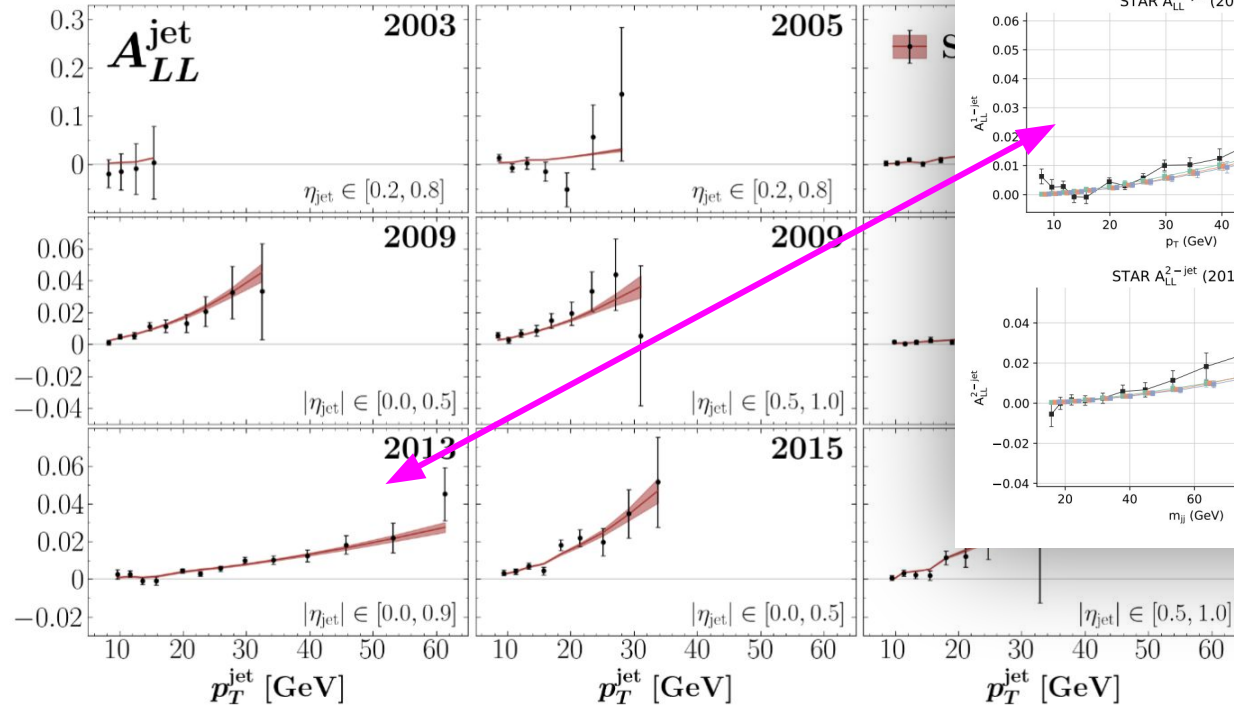
χ_{red}^2 (Z-score)		
Data	N_{dat}	$W^2 > 4 \text{ GeV}^2$ (no HT)
Polarized	1912 [†]	1.01 (+0.63)
— DIS	1735 [†]	1.04 (+1.15)
— SIDIS	124	0.85 (−1.21)
— jets	83	0.84 (−1.06)
— W/Z boson	18	0.82 (−0.47)
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Total	1960[†]	1.01 (+0.20)

pol pp Jets

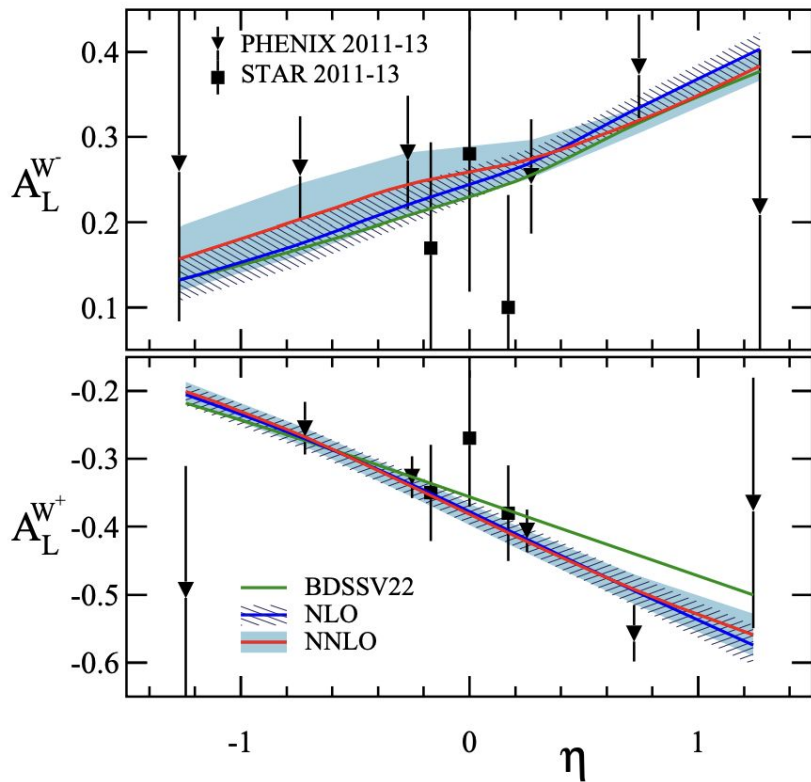
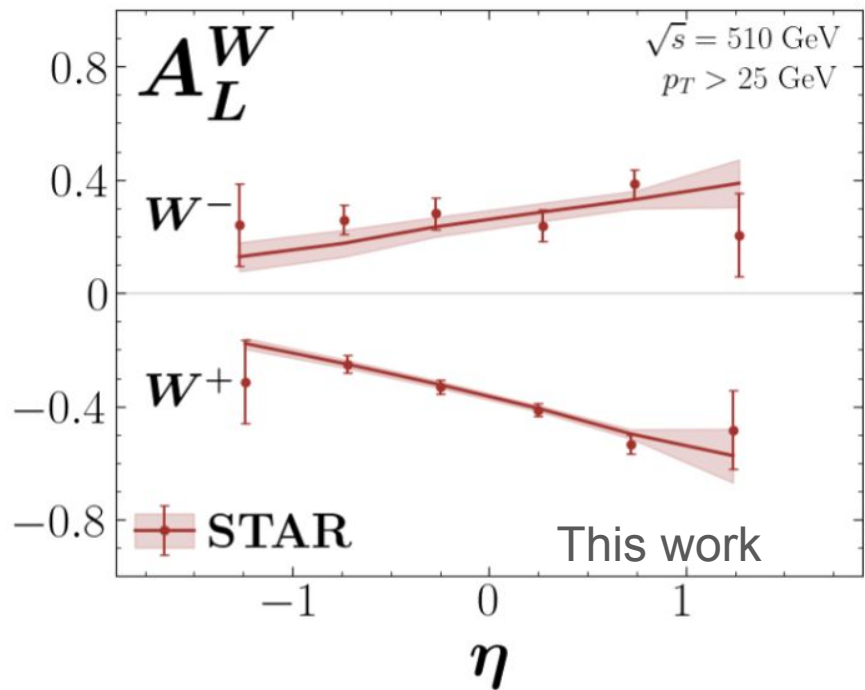


Borsa, Stratmann, Vogelsang, de Florian, Sassot, '25

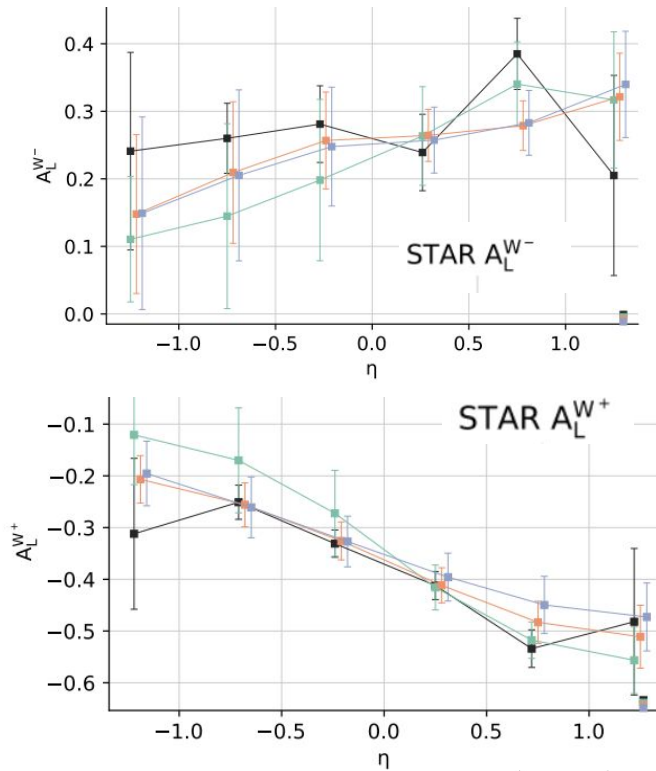
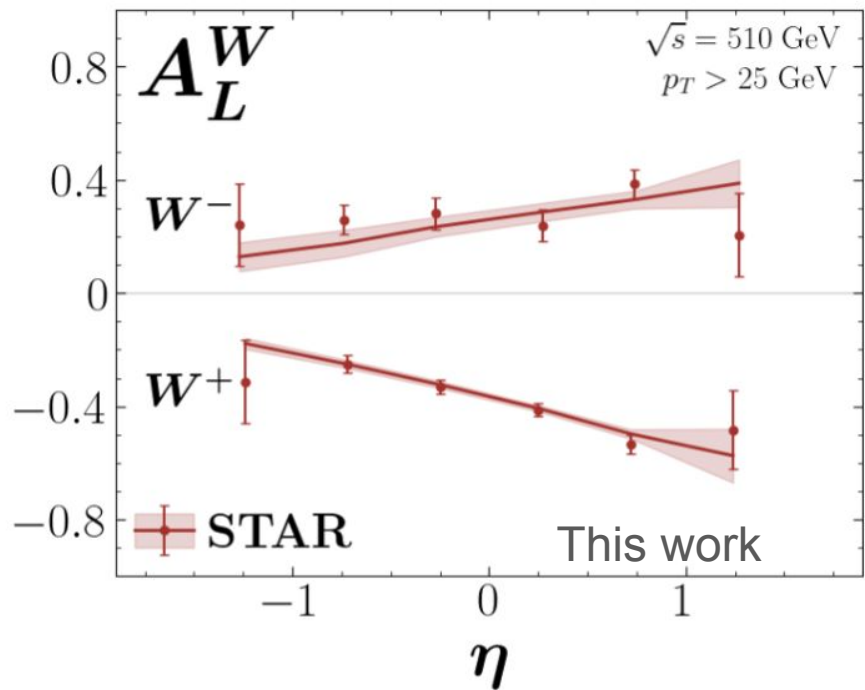
pol pp Jets



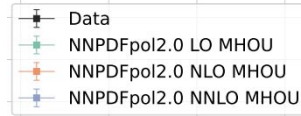
Cruz-Martinez, Hasenack, Hekhorn, Magni,
Nocera, Rabemananjara, Rojo, Sharma,
van Seeventer '25



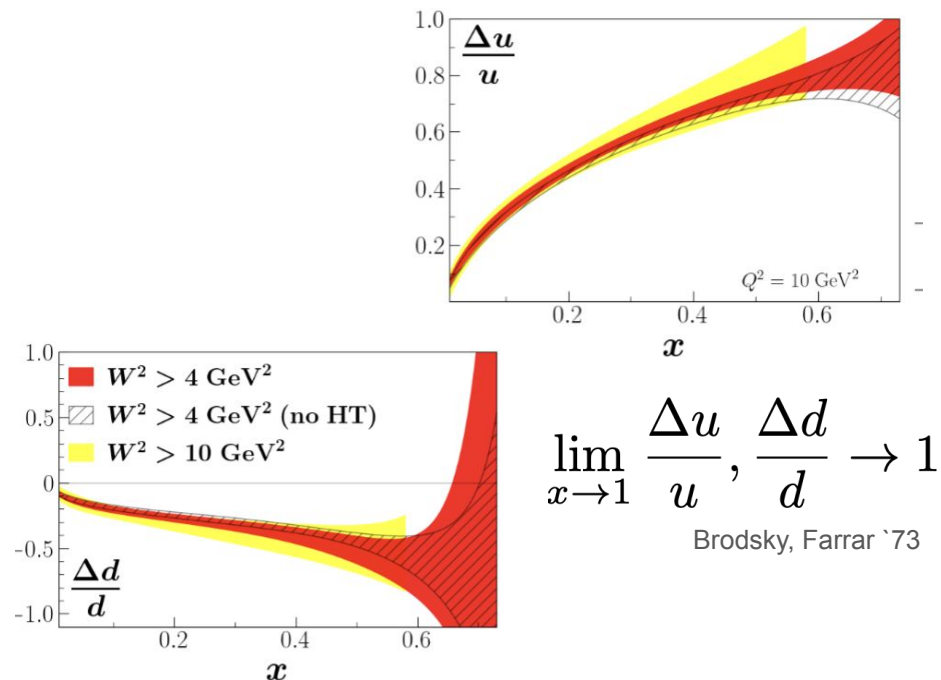
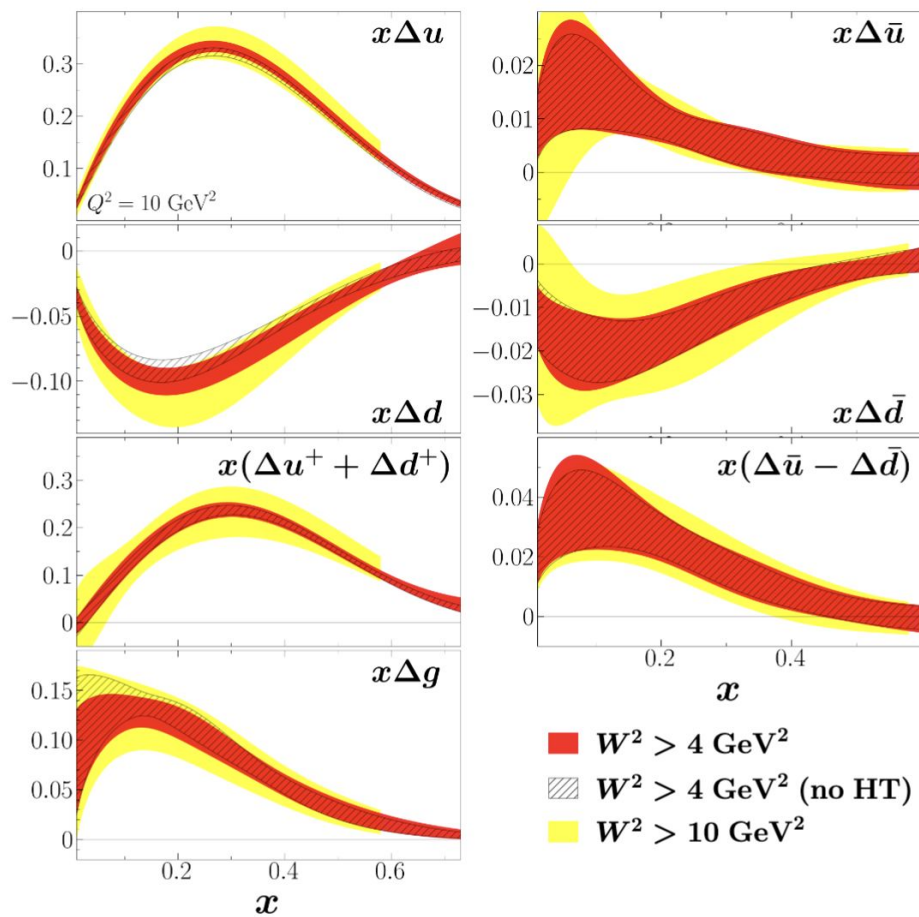
pol W⁺/W⁻



Cruz-Martinez, Hasenack, Hekhorn, Magni,
Nocera, Rabemananjara, Rojo, Sharma,
van Seeventer '25



Reconstructed hPDFs

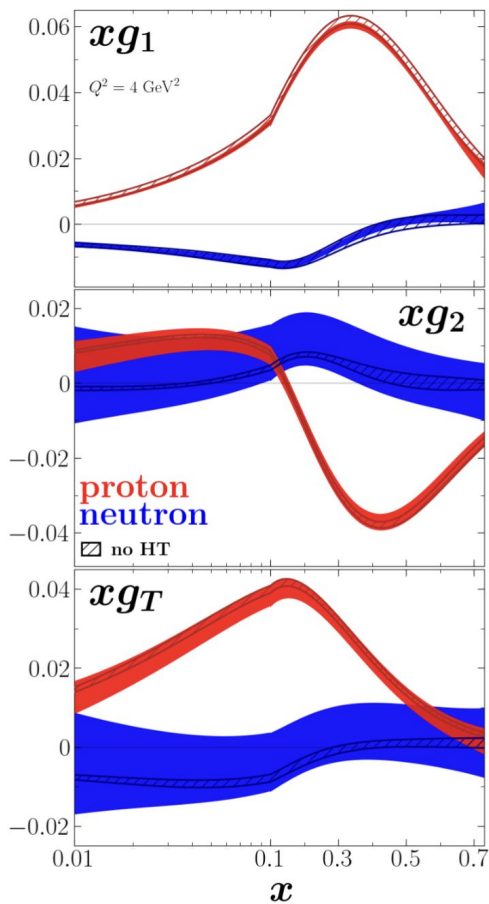


$$\lim_{x \rightarrow 1} \frac{\Delta u}{u}, \frac{\Delta d}{d} \rightarrow 1$$

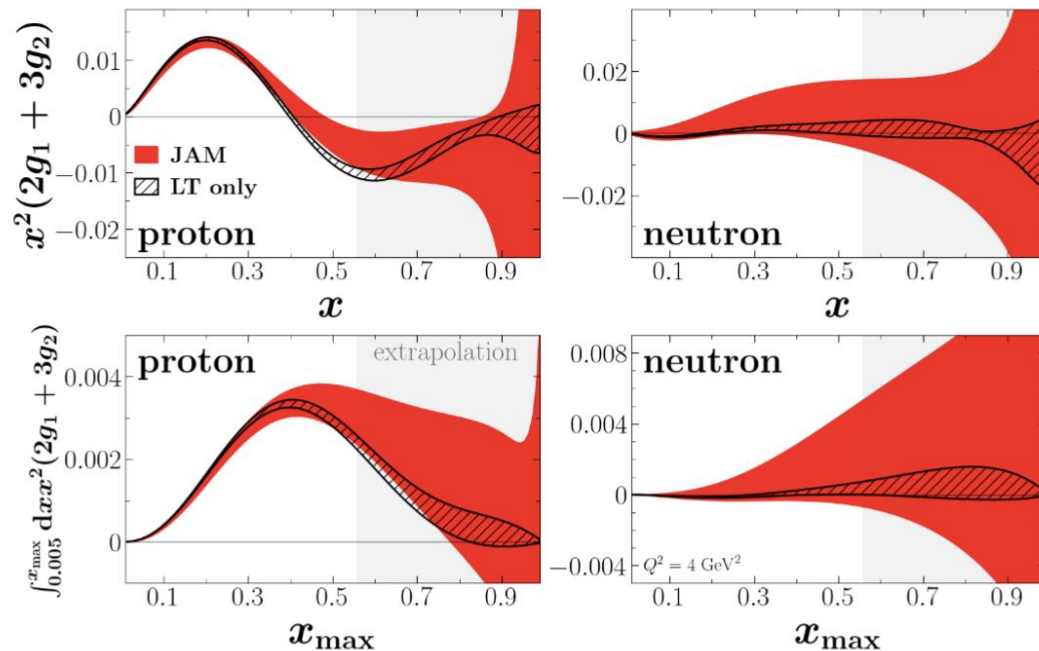
Brodsky, Farrar '73

Inclusion of high x DIS data reduces significantly the uncertainties on hPDF
 Reconstruction of hPDFs is stable regardless of HT effects
 Sign change of $\Delta d/d$ at high x is still not confirmed empirically

d2 matrix element (twist 3)

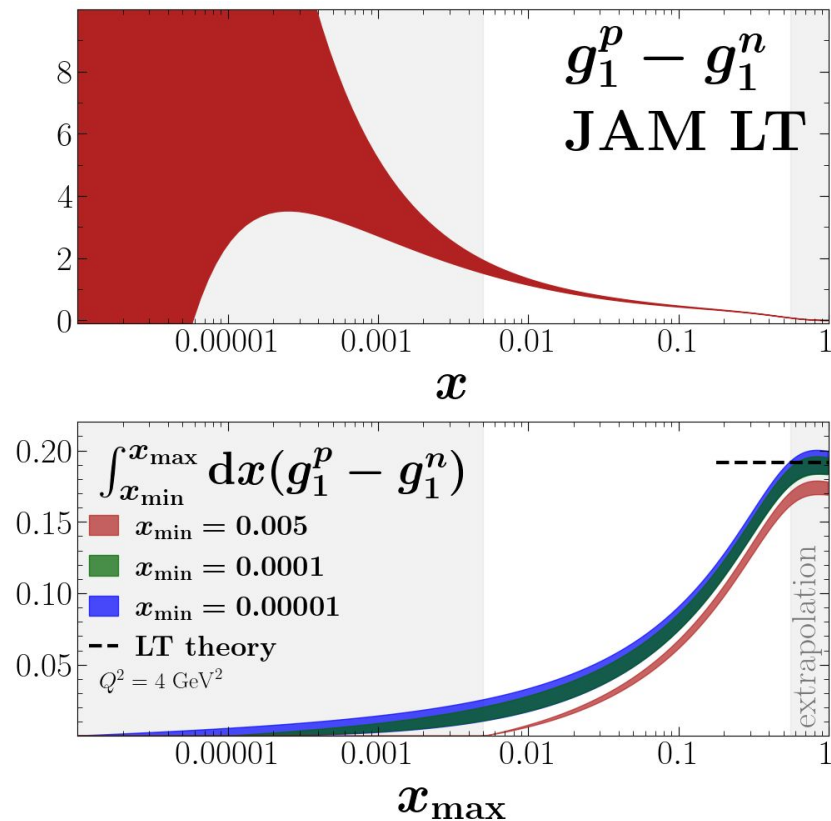


$$d_2(Q^2) = \int_0^1 dx x^2 [2g_1(x, Q^2) + 3g_2(x, Q^2)]$$



The present data constraints HT effects up to $x \sim 0.5$.
Beyond this, large uncertainties are present preventing to elucidate the genuine twist 3 effects.
Neutron d_2 is even less constrained and more data at high x in ^3He is needed

Bjorken sum rule (BJSR)



$$\int_0^1 dx [g_1^p(x, Q^2) - g_1^n(x, Q^2)] = \frac{g_A}{6} \left(1 - \frac{\alpha_S(Q^2)}{\pi} \right)$$

The approximated reconstruction of BJSR is very close to analytic results at NLO.

Current uncertainties are still too large to be competitive to extract α_S .

Constraining g_1^p and g_1^n at smaller values x (eg EIC) has the potential to offer competitive constraints on α_S .

Summary & outlook

Performed a comprehensive QCD global analysis by simultaneously extracting PDFs, hPDFs, FFs, and DIS HT effects.

Based on stability of the results, it is possible to lower the W^2 cuts down to 4 GeV^2 .

HT in protons are found to be negligible within uncertainties. For neutrons, the uncertainties are still too large to make any conclusions.

New estimates of BJSR has been provided. Compatible with analytic results, uncertainties still large to be competitive for constraining α_S , need more data at small x (EIC)

