The Polarized EMC Effect

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Polarized Ion Sources and Beams at EIC

10-13 March 2025

Center for Frontiers in Nuclear Science, Stony Brook University



- There is a gap—perhaps a big gap—between traditional picture of a nucleus and a QCD picture
 - this gap manifests in the valence region EMC effect
- Where to start? ⁴He can be consider the lightest "real" nucleus [⁴He_{BE} = 7.1 MeV/A] and EMC effect is fully manifest [$^{208}Pb_{BE} = 7.9$, $^{3}He_{BE} = 2.6$, $^{3}H_{BE} = 2.8 \text{ MeV/A}$]
- ⁴He is a key constituent of nuclei α clustering
 - "standard candle" for QCD and nuclei
- Many foundational QCD questions to address
 - Are the quarks and gluons confined to nucleon-like objects? Does this depend on, e.g., the momentum filter x?
 - What are the quark and gluon mass radii for ⁴He and how does this contrast with the nucleon?
 - What are the pressure and shear forces in ${}^{4}\text{He}?$
- EIC can help bridge this gap



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A More Realistic Impression of ⁴He — Spatial Tomography



2/22

QCD and Imaging of Light Nuclei

- Nuclei provide a QCD laboratory with characteristics not available from protons alone
- Program build around imaging of light nuclei would have tremedous impact and reveal many novel aspects of QCD
 - How is gluon dynamics modified by the nuclear medium?
 - $J \ge 1$ targets \Rightarrow new PDFs, form factors, TMDs, GPDs, etc.
 - Exotic gluonic components from gluon transversity PDFs
 - Color transparency, hidden color, NN correlations, fast quarks
 - Isospin & baryon density effects, e.g., partial restoration of chiral symmetry and possible changes in confinement length scales between quarks and gluons
- Key question: *How does the nucleon-nucleon interaction arise from QCD?*
- EIC's unique capabilities for polarized proton structure should apply equally to nuclei (ideally p, D, ³H, ³He, ⁶Li, ⁷Li, ⁷Be, ..., Ne)



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"No story of modern physics is more intriguing than the history of the theory of nuclear forces." Ruprecht Machleidt, Weinberg's proposal of 1990: A very personal view

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Lithium-6

Matt Stracelor 2013





EMC Effect

- Understanding origin of the EMC effect is critical for a QCD based description of nuclei
 - 40+ years after discovery a broad consensus on explanation is lacking
 - Valence quarks in nucleus carry less momentum than in a nucleon
- Important question: In what processes, and at what energy scales, do quarks and gluons become the effective degrees of freedom in nuclei?
- Modern explanations of EMC effect are based around medium modification of the bound nucleons
 - Is modification caused by *mean-fields* which modify *all* nucleons all of the time or by *SRCs* which modify *some* nucleons some of the time?





long- and short-range NN interactions 4/22

Nucleons in Nuclei

- Nuclei are extremely dense:
 - proton rms radius is $r_p \simeq 0.85$ fm, corresponds hard sphere $r_p \simeq 1.10$ fm
 - ideal packing gives $\rho\simeq 0.13\,{\rm fm}^{-3};$ nuclear matter density is $\rho\simeq 0.16\,{\rm fm}^{-3}$
 - 20% of nucleon volume inside other nucleons nucleon centers $\sim\!2\,\text{fm}$ apart
- For realistic charge distribution 25% of proton charge at distances r > 1 fm (scalar radius?)
- Natural to expect that nucleon properties are modified by nuclear medium
 - in contrast to traditional nuclear physics
- Understanding validity of two viewpoints remains key challenge for nuclear physics a new paradigm or deep insights into color confinement in QCD



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Understanding the EMC Effect

- The puzzle posed by the EMC effect will only be solved by conducting new experiments that expose new aspects
- Measurements should help distinguish between explanations of EMC effect: e.g. whether all nucleons are modified by the medium or only those in SRCs
- Important examples are measurements of the *EMC effect* in polarized structure functions & the flavor dependence of *EMC* effect
- A JLab experiment has been approved to measure the spin structure of ⁷Li
- Flavor dependence has been accessed via JLab DIS experiments on ⁴⁰Ca & ⁴⁸Ca *but parity violating DIS stands to play the pivotal role*



Theory approaches to EMC Effect

• To address origins of EMC effect must determine e.g. nuclear PDFs, TMDs, GPDs:

$$q_{A}(x_{A}) = \frac{P^{+}}{A} \int \frac{d\xi^{-}}{4\pi} e^{ix_{A} P \cdot \xi/A} \left\langle A, P \left| \overline{\psi}_{q}(0) \gamma^{+} \psi_{q}(\xi^{-}, \boldsymbol{\xi}_{T}) \right| A, P \right\rangle \Big|_{\xi^{+}=0}$$

• Common to approximate using convolution formalism

$$q_A(x_A) = \sum_{\alpha} \int_0^A dy_A \int_0^1 dz \ \delta(x_A - y_A z) \ f_A^{\alpha}(y_A) \ q_{\alpha}(z)$$

- $\alpha = (bound)$ protons, neutrons, pions, deltas ...
- $q_{\alpha}(z)$ PDFs of quarks q in bound hadron α
- $f_{\alpha}(y_A)$ PDFs of hadron in nucleus



Nucleon Momentum Distributions in Nuclei



- Modern GFMC or VMC nucleon momentum distributions have significant high momentum tails
 - indicates presence of SRCs: ${\sim}20\%$ for ^{12}C
- Light-cone momentum distribution:

$$f(y_A) = \int rac{\mathrm{d}^3 ec{p}}{(2\pi)^3} \; \delta\left(y_A - rac{p^+}{P^+}
ight) \;
ho(p)$$

• Naive SRCs introduce effect of opposite sign to EMC effect



The Nambu–Jona-Lasinio Model



- NJL model is a Lagrangian based covariant QFT, exhibits dynamical chiral symmetry breaking, & aspects of confinement
 - Calculations proceed via bound state equations: gap, Bethe-Salpeter, Faddeev, ...
- Quark confinement is implemented via proper-time regularization
 - Quark propagator: $[p m + i\varepsilon]^{-1} \rightarrow Z(p^2)[p M + i\varepsilon]^{-1}$
 - Wave function renormalization vanishes at quark mass-shell: $Z(p^2 = M^2) = 0$
- Finite density calculations are possible at mean-field level with interactions in $\sigma, \omega, \rho, \ldots$ channels
 - Effective *NN* potential is derived via hadronization methods and calculations are done self-consistently
 - Model exhibits correct saturation of nuclear matter is symmetry energy





Nucleon Electromagnetic Form Factors

• Nucleon = quark+diquark

• Form factors given by Feynman diagrams:



- Calculation satisfies electromagnetic gauge invariance; includes
 - dressed quark–photon vertex with ρ and ω contributions
 - contributions from a pion cloud

[ICC, W. Bentz and A. W. Thomas, Phys. Rev. C 90, 045202 (2014)]





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Nucleon guark distributions

• Nucleon = quark+diquark



• Covariant, correct support; satisfies sum rules, Soffer bound & positivity

 $\langle q(x) - \bar{q}(x) \rangle = N_q, \quad \langle x u(x) + x d(x) + \ldots \rangle = 1, \quad |\Delta q(x)|, \quad |\Delta_T q(x)| \leq q(x)$

[ICC, W. Bentz and A. W. Thomas, Phys. Lett. B 621, 246 (2005)]





Spin and Gluon EMC Effects

- To solve puzzle of EMC effect need new observables, e.g., gluon and spin EMC effects
 - Can help distinguish between different explanations of the EMC effect
 - Mean-field and SRC make different predictions for spin EMC effect
- The gluon EMC effect can be defined as

 $R_g(x) = \frac{g_A(x)}{Z g_p(x) + N g_n(x)}$

- Analogous definition for gluon spin EMC effect, with, $Z \rightarrow P_p$ and $N \rightarrow P_n$
- Results obtained in mean-field model that describes the EMC effect and predicts spin EMC effect
 - Gluons are generated purely perturbatively
 - Provides a baseline for comparison and understanding of future measurements

[X. G. Wang, W. Bentz, ICC, and A. W. Thomas, J. Phys. G 49, (2022)]



12/22

Polarized EMC Effect – Update

- Proposal "The EMC Effect in Spin Structure Functions" for ⁷Li completed jeopardy process in PAC 48 (2020)
 - scientific rating went from B⁺ to A⁻ almost unheard of thanks to a lot of work from Will Brooks and Sebastian Kuhn
- Spin/Polarized EMC effect experiments are just measurements of the spin structure function(s) of a nucleus exactly analogous to nucleon DIS
- Polarized EMC effect provides insight into QCD effects in nuclei

 $\Delta R(x) = \frac{g_{1A}(x)}{g_{1A}^{\text{naive}}(x)} = \frac{g_{1A}(x)}{P_{p} \, g_{1p}(x) + P_{n} \, g_{1n}(x)}$

- $P_p \& P_n$ effective polarizations of protons/neutrons in nucleus
- JLab will hopefully soon run polarized ⁷Li DIS experiment
 - Ideal target should have spin dominated by protons and small A
 - Candidate nuclei include ³H ($J = \frac{1}{2}$), ⁷Li ($J = \frac{3}{2}$), ¹¹B ($J = \frac{3}{2}$), ...







Mean-Field Calculations of Polarized EMC Effect

- Several relativistic mean-field calculations of polarized EMC effect
 - all calculations find polarized EMC same size or larger than EMC effect
- Large polarized EMC effect results because in-medium quarks are more relativistic (*M*^{*} < *M*)
 - quark lower components are enhanced
 - in-medium we find that quark spin is converted to orbital angular momentum









EMC and Polarized EMC in Spin-1 Nuclei

- Of the 250+ stable nuclei only 3 have spin-one: 2 H (deuteron), 6 Li, and 14 N
- Spin of each of these nuclei us primarily carried by the two valence nucleon
- For ⁶Li nucleons in $p_{3/2}$ and ¹⁴N nucleons in $p_{1/2}$, so each nucleon has one unit of orbitial angular momentum

$$\Psi_{^{6}\mathrm{Li}}^{11} = \sqrt{\frac{3}{10}} \left[p^{3/2} n^{-1/2} \right] + \sqrt{\frac{3}{10}} \left[p^{-1/2} n^{3/2} \right] - \sqrt{\frac{2}{5}} \left[p^{1/2} n^{1/2} \right]$$

• For deuteron nucleon is *s*-wave, so very interesting to compare polarized EMC effect for these nuclei





Polarized EMC Effect for Spin-One Nuclei

• Recall that the polarized EMC effect is defined by

 $\Delta R(x) = \frac{g_{1A}(x)}{g_{1A}^{\text{naive}}(x)} = \frac{g_{1A}(x)}{P_{p} \, g_{1p}(x) + P_{n} \, g_{1n}(x)}$

• P_p & P_n effective polarizations of protons/neutrons

 $14 \mathrm{N}$

0.2

0.4

0.0

-0.1

-0.3

-0.4 L

B -0.2

- ⁶Li polarized EMC effect similar to that of ⁷Li
 - effective polarizations: $P_p = P_n = 1/3$
- In ¹⁴N valence nucleons in p_{1/2} state
 - causes zero-crossing in numerator of $\Delta R(x)$



The Deuteron

- The deuteron is the simplest nucleus naively consisting of a proton + neutron with 2.2 MeV binding
 - however deuteron is greater than sum of its parts, having many properties not found in either of its primary constituents
 - deuteron is also finally tuned making it an interesting target to isolate QCD effects
- Unique properties of deuteron:
 - a quadrupole moment and gluon transversity PDF
 - many TMDs and GPDs associated with tensor polarization
- Additional spin-independent leading-twist PDF called $b_1^q(x)$

$$b_1(x) = e_q^2 \left[b_1^q(x) + b_1^{\bar{q}}(x)
ight], \quad \int_0^1 dx \left[b_1^q(x) - b_1^{\bar{q}}(x)
ight] = 0$$

- Need tensor polarized target to measure $b_1(x) (\text{HERMES})$
 - impossible to explain HERMES data with only bound nucleon degrees of freedom need exotic QCD states, 6q bags, etc.
 - Hall C has an approved proposal (J.-P. Chen, et al.)





TMDs of Spin-1 Targets

- A spin-1 target can have tensor polarization $[\lambda=0]$
 - 3 additional *T*-even and 7 additional *T*-odd quark TMDs compared to nucleon
- Analogous situation for gluon TMDs
 - to fully expose role of quarks and gluons in nuclei need polarized nuclear targets (transverse and longitudinal) with all spin projections, e.g., for J = 1: ²H, ⁶Li
- Spin 4-vector of a spin-one particle moving in z-direction, with spin quantization axis S = (S_T, S_L), reads: S^μ(p) = (p_z/m_h S_L, S_T, p₀/m_h S_L)



- longitudinal polarization $\implies \boldsymbol{S}_T = 0, S_L = 1$; transverse $\implies |\boldsymbol{S}_T| = 1, S_L = 0$
- Associated quark correlation function:

$$\left\langle \gamma^{+} \right\rangle_{\boldsymbol{S}}^{(\lambda)}(x,\boldsymbol{k}_{T}) \equiv f(x,\boldsymbol{k}_{T}^{2}) - \frac{3\lambda^{2} - 2}{2} \left[\left(S_{L}^{2} - \frac{1}{3} \right) \theta_{LL}(x,\boldsymbol{k}_{T}^{2}) + \frac{(\boldsymbol{k}_{T} \cdot \boldsymbol{S}_{T})^{2} - \frac{1}{3}\boldsymbol{k}_{T}^{2}}{m_{h}^{2}} \theta_{TT}(x,\boldsymbol{k}_{T}^{2}) + S_{L} \frac{\boldsymbol{k}_{T} \cdot \boldsymbol{S}_{T}}{m_{h}} \theta_{LT}(x,\boldsymbol{k}_{T}^{2}) \right]_{18/22}$$

leading twist		quark operator		
		γ^+	$\gamma^+\gamma_5$	$\gamma^+\gamma^i\gamma_5$
target polarization	U	$f_1 = \bigcirc$ unpolarized		$h_1^{\perp} = { { } { { } { } { } { } { } { } } { } } { } } { } } } { } } } { } } } } } } { } { } { } { } { } { } { } { } { } { } { } { } { } { } { } } } $
	L		$g_1 = \longrightarrow - \longleftrightarrow$ helicity	$h_{1L}^{\scriptscriptstyle \perp} = \underbrace{ \swarrow }_{\text{worm gear } 1} - \underbrace{ \checkmark }_{\text{worm gear } 1}$
	т	$f_{1T}^{\perp} = \underbrace{\bullet}_{\text{Sivers}}^{\bullet} - \underbrace{\bullet}_{\text{Sivers}}^{\bullet}$	$g_{1T} = \underbrace{\bigstar}_{\text{worm gear 2}} - \underbrace{\bigstar}_{\text{worm gear 2}}$	$h_{1} = \underbrace{\begin{pmatrix} \bullet \\ \bullet \\ transversity \end{pmatrix}}_{transversity}$ $h_{1T}^{\perp} = \underbrace{\begin{pmatrix} \bullet \\ \bullet \\ pretzelosity \end{pmatrix}}_{pretzelosity}$
	HENSOR	$\left. \begin{array}{c} \theta_{LL}(x, \boldsymbol{k}_{T}^{2}) \\ \\ \theta_{TT}(x, \boldsymbol{k}_{T}^{2}) \end{array} \right\} \\ \theta_{LT}(x, \boldsymbol{k}_{T}^{2}) \end{array}$	$egin{aligned} g_{1TT}(x,oldsymbol{k}_T^2)\ g_{1LT}(x,oldsymbol{k}_T^2) \end{aligned}$	$egin{array}{l} h_{1LL}^{\perp}(x,m{k}_{T}^{2}) \ h_{1TT}, \ h_{1TT}^{\perp} \ h_{1LT}, \ h_{1LT}^{\perp}, \ h_{1LT}^{\perp} \end{array}$

Spin-1 Target TMDs – with Nucleon Analogs



helicity







transversity



[Yu Ninomiya, ICC and Wolfgang Bentz, Phys. Rev. C 96, no.4, 045206 (2017)]

Spin-1 Target TMDs – Tensor Polarization

- Calculations assume point-like nucleons but nevertheless show tensor polarized TMDs have some surprising features
- TMDs $\theta_{LL}(\times \mathbf{k}_T^2) \& \theta_{LT}(\times \mathbf{k}_T^2)$ identically vanish at x = 1/2 for all \mathbf{k}_T^2
 - x = 1/2 corresponds to zero relative momentum between (the two) constituents, that is, s-wave contributions
 - therefore $\theta_{LL} \& \theta_{LT}$ primarily receive contributions from $L \ge 1$ components of the wave function – sensitive to orbital angular momentum
- Features hard to determine from a few moments challenge for traditional lattice QCD methods





[Yu Ninomiya, ICC and Wolfgang Bentz, Phys. Rev. C 96, no.4, 045206 (2017)]

Deuteron GPDs

- The deuteron has a rich GPD structure
- The impact parameter PDFs provide a spatial tomography for various *x* slices
 - tensor polarized along z-axis clear donut shape
 - longitudinally polarized along x-axis clear dumbbell shape
- These quantities provide an interesting connection to traditional nuclear physics results for the deuteron
 - nuclear spatial densities have donut and dumbbell shapes
- Does the gluon donut align with the quark donut – does this change with x? Incredible insight into NN interaction possible





J. Carlson, R. Schiavalla, Rev. Mod. Phys. **70** 743 (1998)

J. L. Forrest *et al.* Phys. Rev. **C54** 646 (1996)

Conclusion and Outlook

- Tremendous opportunity for EIC to transform understanding of QCD in nuclei
 - GPDs and TMDs of light nuclei
 - medium effects on gluon structure via J/ψ production
 - Anti-shadowing region and its A dependence
 - $b_1(x)$ and gluon transversity in deuteron and ⁶Li
- Key physics questions: How does the *NN* interaction arise from QCD? How do quark/gluon confinement length scales change in medium?
- Can explore these questions by imaging light nuclei and comparing quarks and gluons for slices in x, k_T^2 , and b_T^2
 - correlations between quarks and gluons in nuclei provide insights into color confinement



