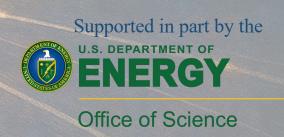
# From Ultra-Peripheral to Fixed Target: Low energy photonuclear physics at the EIC

Daniel Brandenburg



Fixed-target eXperiment at EIC (FIXE) initiative September 29, 2025





#### **Outline**

- Physics motivation
  - Hadron Mass generation & QCD Trace Anomaly
- Near-threshold photonuclear production
- EIC collider mode
- Fixed Target mode
  - Photon source and approximate rates
  - Sub-threshold production
  - Lighter VM production
- Utilization of the ePIC instrumentation
- Summary & Conclusions

#### Trace Anomaly and the Origin of mass

>99% of all visible mass in the Universe is built from nucleons Higgs mass accounts for a small portion, with majority resulting from the strong force self-interactions

$$T^{\mu}_{\mu} = \frac{\beta(g)}{2g}G^2 + (1 + \gamma_m)\overline{\psi}m\psi,$$

Gluons 
$$(T_g^{\mu\nu})$$
 Quarks  $(T_q^{\mu\nu})$ 

$$\langle P|T^{\mu\nu}_{q,G}(0)|P\rangle = 2P^{\mu}P^{\nu}A_{q,G}(0) + 2M^2g^{\mu\nu}\bar{C}_{q,G}(0),$$

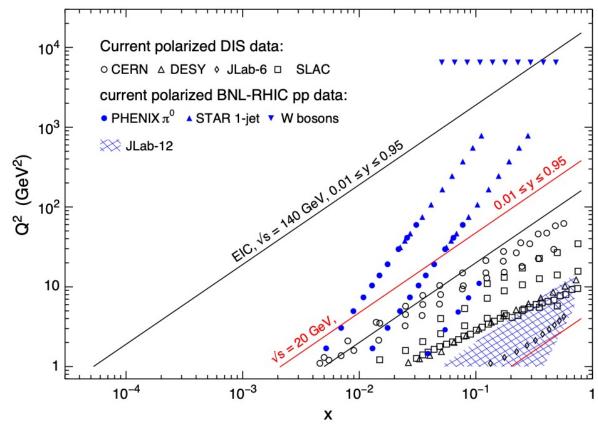
#### Threshold Vector Meson Production

- Probe gluonic structure of the nucleon near confinement scale (especially for heavier VMs like J/psi and Upsilon)
- Access gluonic gravitational form factors (GFFs)
- Connect to the proton mass decomposition and QCD trace anomaly
- Unique sensitivity in Upsilon, J/ψ, φ, and ρ near threshold, each in contribute different constraints

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Wang et al. (2023) — arXiv:2308.04644, 2304.07964
Hatta et al. (2025) — arXiv:2501.12343
Du, Xie et al. (2020) — arXiv:2009.08345
Winney (2023) — arXiv:2305.01449
Kim et al. (2024) — arXiv:2411.12187
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## Why Threshold Production?

- Valence (large-x) gluons & transition from non-perturbative to perturbative QCD
- Near threshold, the produced VM is slow in the target rest frame
- Amplitude is dominated by multigluon exchange with x ≥ 0.1, where constraints are poor.
- mapping energy and tdependence here tests models (Regge/Pomeron, GPDs, VMD/dispersion, NRQCD).



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#### Threshold Vector Meson Production

• In  $\gamma p \rightarrow VM p$ :

$$W_{thr} = m_p + m_{VM}$$

Vector Meson	Invariant Mass [GeV]	Threshold [GeV]
ρ	0.775	1.713
$\phi$	1.019	1.957
$J/\psi$	3.097	4.035
Υ	9.46	10.398

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# First Measurement of Near-Threshold $J/\psi$ Exclusive Photoproduction off the Proton

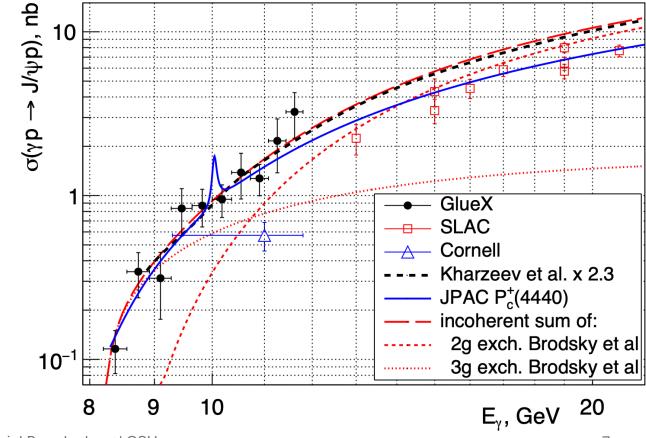
A. Ali<sup>10</sup>, M. Amaryan<sup>22</sup>, E. G. Anassontzis<sup>2</sup>, A. Austregesilo<sup>3</sup>, M. Baalouch<sup>22</sup>, F. Barbosa<sup>14</sup>, J. Barlow<sup>7</sup>, A. Barnes<sup>3</sup>, E. Barriga<sup>7</sup> et al. (GlueX Collaboration)

Phys. Rev. Lett. **123**, 072001 – **(2019)** 

Phys. Rev. C **108**, 025201 – (**2023**)

Phys. Rev. Lett. **134**, 201903 – **(2025)** 

- First measurements at Jlab in the last few years
- Similar trend but tension with some lattice calculations and models



## First Measurement of Near-Threshold $J/\psi$ Exclusive Photoproduction off the Proton

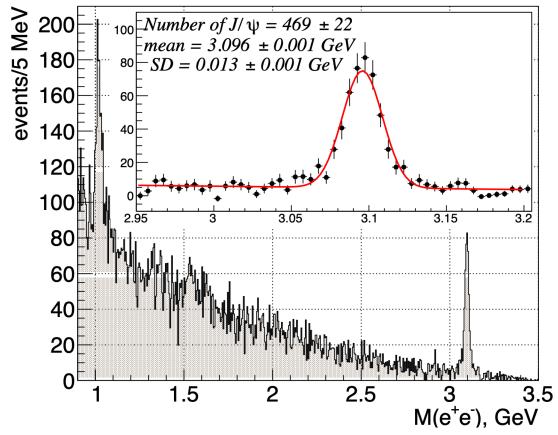
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#### Hadron Mass: Priority for the EIC

Discussed extensively in Yellow Report, White papers, etc.

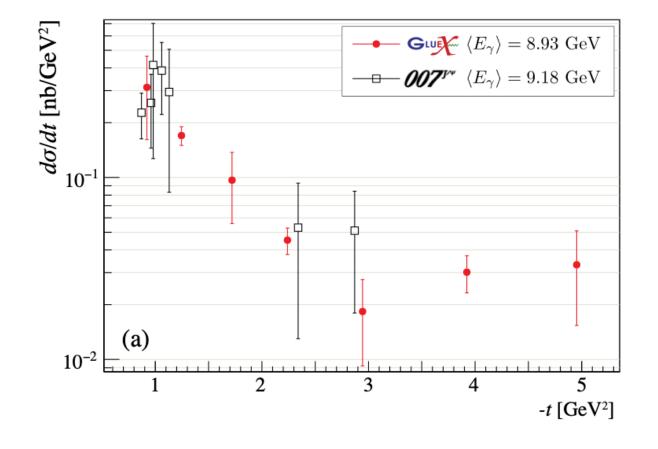
- Exclusive processes: DVCS & vector mesons (YR, Vol II)
- Heavy quarkonia: sensitivity to gluon GPDs & GFFs
- Influence forward detector concepts: Roman pots, recoil tagging (YR, Vol III)

$$W^{2} \approx y \, s_{ep} + m_{p}^{2} - Q^{2}$$
$$y = \frac{E_{\gamma}}{E_{e}}$$

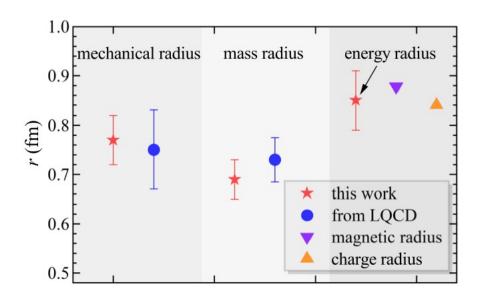
With EIC beam configurations (5,10,18)x(41, 100, 275) and  $y_{min} \approx 0.01$ , EIC collider mode can reach  $W_{min} \approx 3$ -5 GeV

#### **Experimental Access**

- Total cross section slope  $\rightarrow$  gluon pressure & shear (GFF  $C_g(t)$ )
- Threshold normalization  $\rightarrow$  gluon energy-momentum fraction (GFF  $A_g(0)$ )
- |t|-dependence of differential cross section → transverse distributions and subnucleonic dynamics

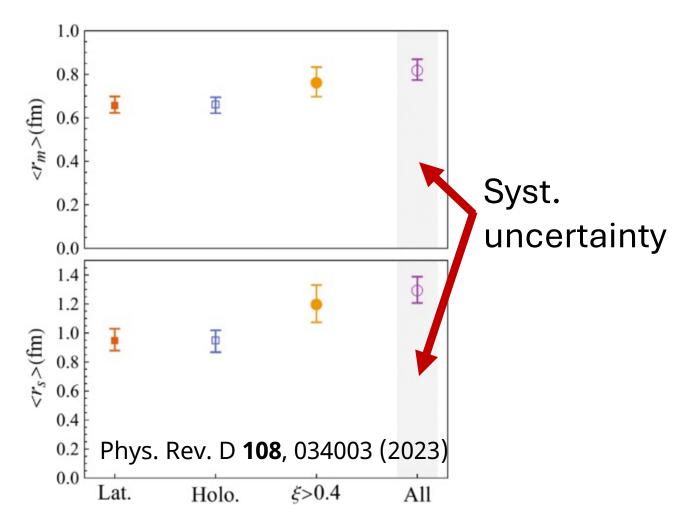


## Proton (gluon) Radius



Chin.Phys.C 48 (2024) 5, 054102

Recent data is narrowing in, but still lacks precision



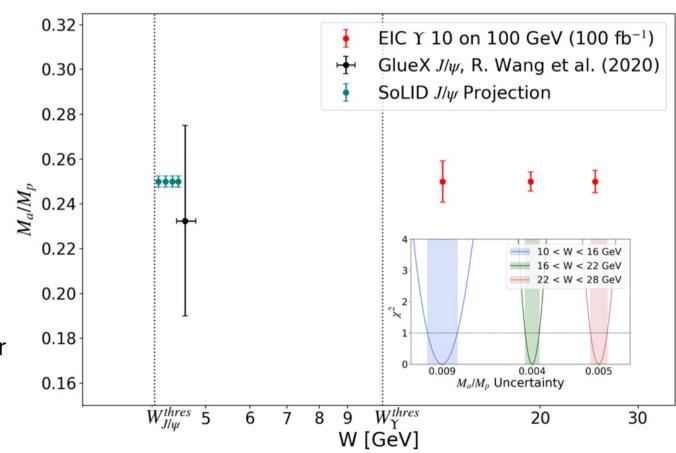
#### EIC Collider mode (ep + eA)

$$W^{2} \approx y \, s_{ep} + m_{p}^{2} - Q^{2}$$
$$y = \frac{E_{\gamma}}{E_{e}}$$

J/psi may be challenging in collider mode and will likely require Q^2 >> 0, will also depend on energy resolution for y measurements

EIC collider mode will be most ideal for threshold Upsilon photoproduction (Q^2~0)

9/29/25

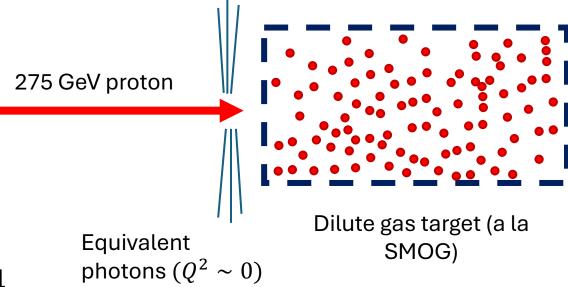


EIC Yellow Report https://arxiv.org/pdf/2103.05419<sub>12</sub>

## Low-Energy Photonuclear w/Fixed Target

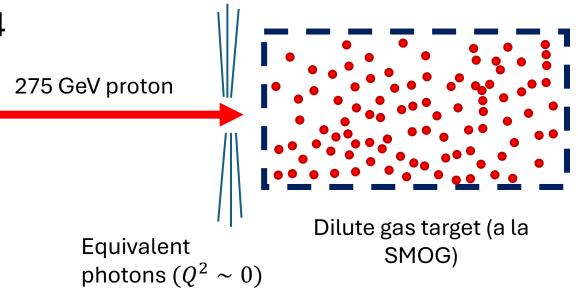
- 275 GeV proton beam
- EPA photon  $k \le \frac{\gamma \hbar c}{R}$
- $W_{max} \approx \sqrt{m_p^2 + 2m_p k_{max}} \approx 12.5$  GeV

Able to access threshold J/psi (4.04 GeV) and Upsilon (10.4 GeV)



## Low-Energy Photonuclear w/Fixed Target

Able to access threshold J/psi (4.04 GeV) and Upsilon (10.4 GeV)



Production Rate:

$$L_{\gamma p} = I_p \, n_t \, N_{\gamma}$$

$$R \approx L_{\gamma p} \times \sigma_{\gamma p \to J/\psi p}$$

$$n_t \approx 10^{13} \text{cm-2}, \sigma_{\gamma p \to J/\psi p} \approx 1\text{-3 nb}$$
  $\mathcal{O}(100-1000) J/\psi$  per day

Complement EIC Collider measurements (electroproduction  $Q^2 > 0$ ) and Jlab experiments

#### Sub-Threshold production and A-scan

$$k_{\mathrm{t}h} = \frac{(M_A + M_V)^2 - M_A^2}{2M_A} = M_V + \frac{M_V^2}{2M_A}$$

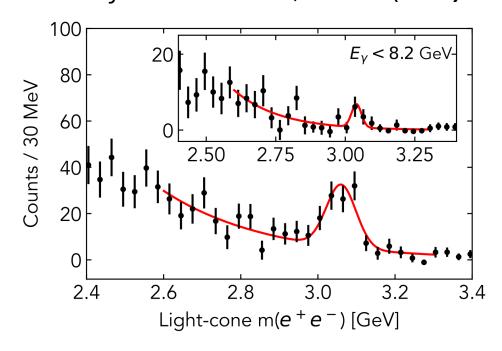
 $\approx 8.21$  GeV on a proton to  $\approx 3.35$  GeV on  $^{20}$ Ne

Production threshold decreases for heavier targets. Access production below the free proton threshold

Nucleus	Plane-wave	Measured	Statistical	Systematic
	cross section	cross section	uncertainty	uncertainty
<sup>2</sup> H	0.24 nb	0.23 nb	$0.07~\mathrm{nb}$	0.04 nb
$^4{ m He}$	0.22 nb	0.33 nb	0.06 nb	$0.05~\mathrm{nb}$
$^{12}\mathrm{C}$	0.24 nb	0.25 nb	$0.05~\mathrm{nb}$	$0.05~\mathrm{nb}$

EIC fixed target would allow scan of heavier A And first measurements of sub-threshold Upsilon

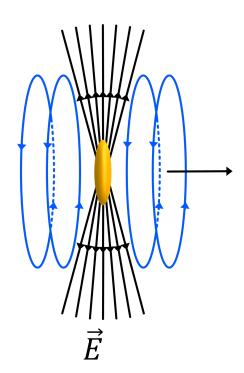
Phys. Rev. Lett. 134, 201903 (2025)



Phys. Rev. D 108, 054018 (2023)

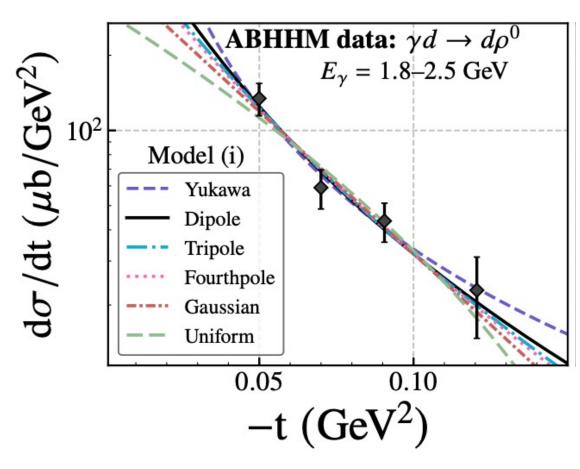
#### $\mathbb{Z}^2$ Enhancement from Nuclear EM Fields

- Nuclear beams benefit from the  $\mathbb{Z}^2$  enhancement of the WW photon flux
- However, in fixed target mode, 100 GeV beam gives  $W_{max} < 3$  GeV  $\rightarrow$  insufficient for J/psi threshold production (even on heavy nuclear targets)
- Lighter mesons can be probed ( $W_{th} \approx 2$  GeV) for  $\rho$ ,  $\phi$
- J/psi (and Upsilon) received most attention because heavy VM are more directly related to the trace anomaly than light
- Other interesting physics in lighter mesons  $(
  ho,\phi)$



## Why study $\rho$ , $\phi$ ?

- Relation to trace anomaly is less direct, but benefit from much higher rate (than J/psi) and Z^2 photon flux enhancement
- Vector mesons (ρ, φ) probing Reggeon vs Pomeron
  - Extraction of SDME to explore nonnatural contributions + look for SCHC violations
  - Strong handle on the nonperturbative end of the spectrum; tests of VMD and dispersion relations close to threshold
- Test models of nuclear shadowing at intermediate x

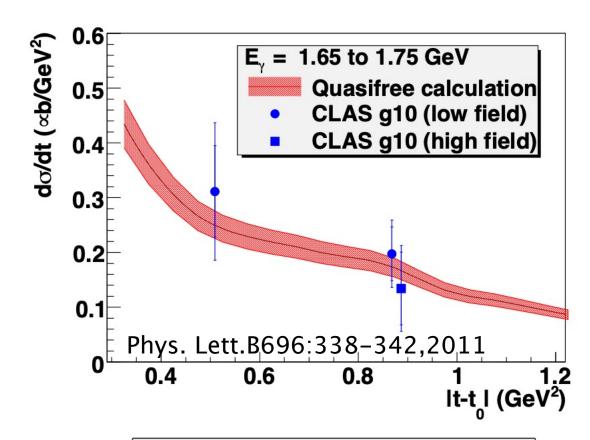


Revisiting the deuteron mass radius

arXiv:2504.10023 [hep-ph]

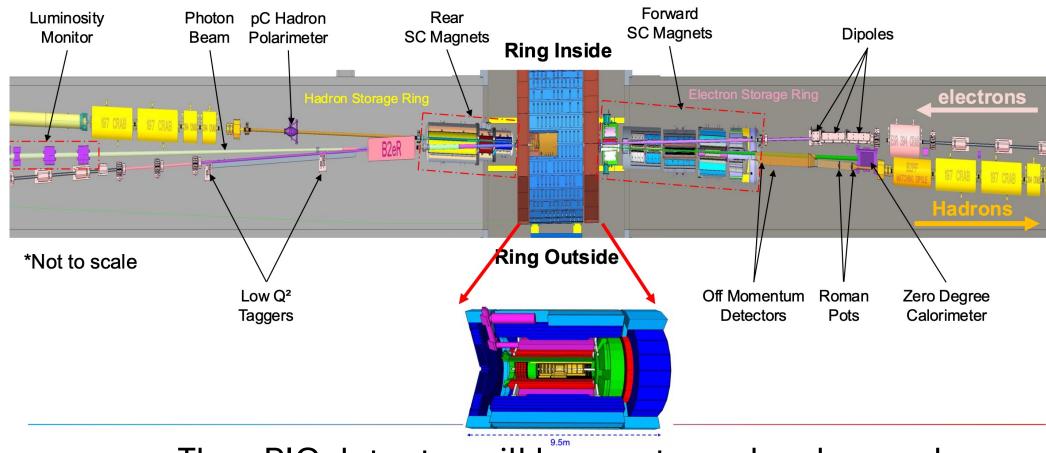
## Why study $\rho$ , $\phi$ ?

- $\phi$  probes the strangeness content and is largely unconstrained at large-x
- OZI → enhanced sensitivity to gluonic mechanisms and to strangeness content of the nucleon (possible s \(\bar{s}\) admixtures)
- Coherent φ on light nuclei → Adependence and small-|t|
  slopes for gluon distributions at
  large x.



$E_{\gamma}$	$ t-t_0 $	$a \tau$	stat.	
(GeV)	$(\mathrm{GeV}^2)$	$(\mu {\rm b}/{\rm GeV}^2)$	uncer.	uncer.
1.65-1.75	0.509	0.31	0.084	0.094
1.65-1.75	0.887	0.20	0.049	0.037
1.65-1.75	0.924	0.13	0.066	0.041

#### Use of the ePIC Detector



The ePIC detector will have extremely advanced array of far-forward/backward capabilities

#### Use of the ePIC Detector

- Zero-Degree Calorimeters (both beams): neutron/photon detection; define Coulomb-excitation tags and impact-parameter classes for ion beams.
- Far-forward hadron system / Roman Pots: detect intact p and diffracted nucleons; measure t with good resolution.
- Off-momentum electron taggers / low-Q² taggers: select quasi-real photons from the electron beam; define clean γ-induced samples.
- Forward tracking + calorimetry: reconstruct low-p\_T vector mesons, provide exclusivity vetoes, and measure SDMEs.
- Global vetoes & TOF: suppress peripheral hadronic backgrounds; enforce exclusivity.

#### Thoughts on using the electron beam

- Fine photon energy control is needed for certain aspects of the nearthreshold physics, i.e. structures in cross section
- Premiere experiments use thin radiators or inverse Compton scattering -> fine control on quasi mono-energetic photon beams

#### At the EIC:

- Tagged electron -> access virtual photons to lower Q^2 and higher  $E_{\gamma}$  (for displaced target with ePIC acceptance)
- Thin radiator + photon tagger -> create a next generation version of current Jlab experiments
- Inverse Compton Scattering -> utilize expertise and R&D for the polarimeters
- Larger initial investment, but more topics accessible with fine photon energy control + higher rate

#### Summary & Conclusions

Rich physics from vector Meson production near threshold

Access to QCD trace anomaly and the gluonic gravitational Form

**Factors** 

$$T^{\mu}_{\mu} = \frac{\beta(g)}{2g}G^2 + (1+\gamma_m)\overline{\psi}m\psi,$$

- SMOG-like fixed target at EIC would allow complementary nearthreshold measurements of  $\rho, \phi, J/\psi$
- Extend EIC collider kinematic reach to highest x values
- Equivalent Photons from proton beam allow threshold  $J/\psi$ , while lighter mesons benefit from Z^2 enhancement