

The **Renormalization** of quasi-DAs: from Mesons to **Baryons**

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OUTLINE

1

Renormalizations

Why Renormalization; Schemes

2

Baryon quasi-DAs on Lattice

Set up; Results; Divergence; Ratio scheme

3

Self renormalization

Self parameterization & Two step fitting

4

Hybrid scheme

Regions; Hybrid renormalization

5

Summary

Quais-DA of Lambda & Proton



Part 1



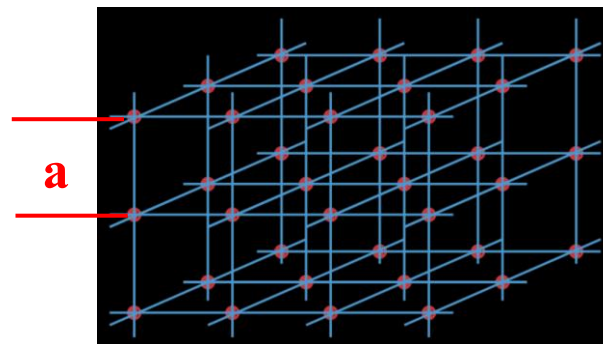
Why Renormalization

Why Renormalization

Part ①

- Scheme Conversion

$$\frac{1}{a} \text{ cut} \rightarrow \text{certain scheme}$$

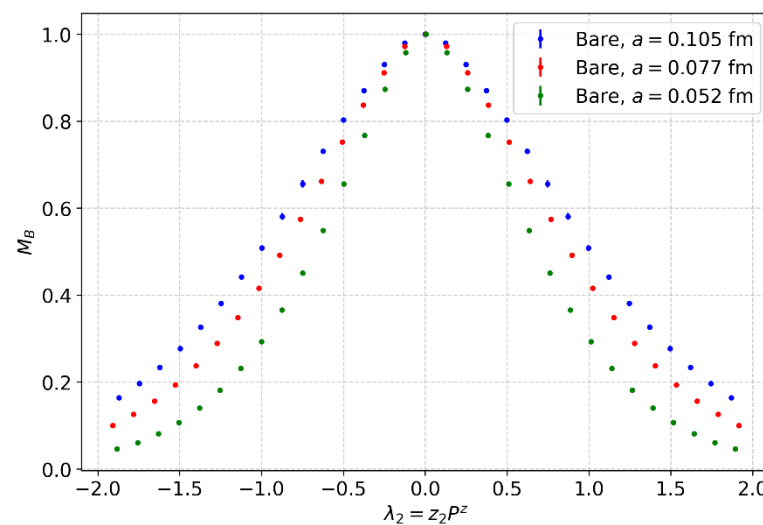
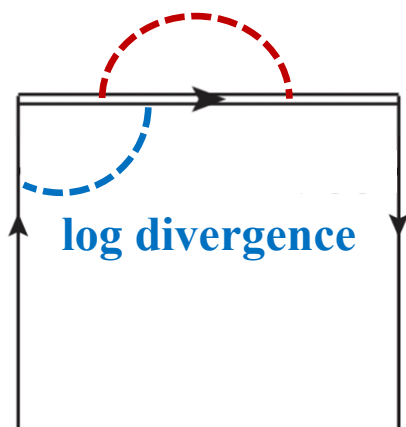


Renorm

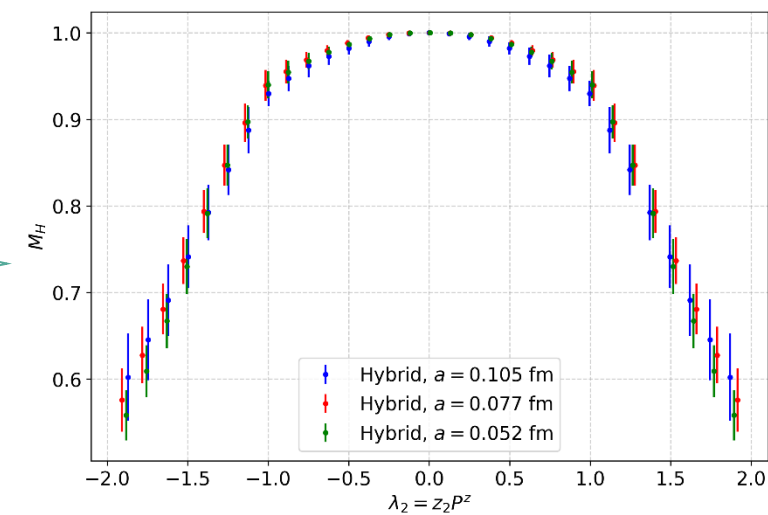


- Divergence Subtraction

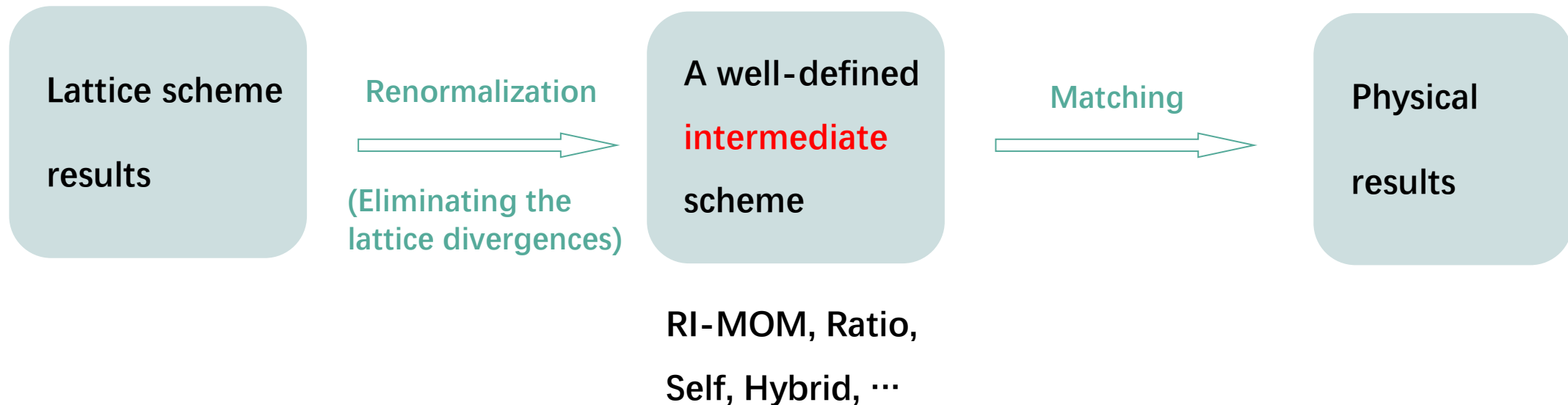
linear divergence



Renorm



Renormalization: **Bridging** the lattice scheme with the physics



- 1) Renormalization requires the ability to **eliminate** the **divergence** caused by the lattice effect.
- 2) The converted scheme is well-defined for effective **matching** or running ...

- **RI/MOM RI/SMOM Scheme**

Martinelli et.al, NPB (1994);

Alexandrou et.al, NPB (2017); Stewart, Zhao, PRD (2018)

He et.al, PRD (2022); Bi et.al, PRD (2023)

- **Ratio Scheme**

Radyushkin et.al, PRD (2017)

- **Self Renormalization and Hybrid Scheme**

Ji et.al, NPB (2021); Huo et.al, NPB (2021)

Pros

Strict subtraction
Well established
Widely applied
...

Easy to use
Well subtraction
at **short-distance**

Cons

Not suitable for **non-local**
(e.g. Gauge dependence,
Signal problem ...)

Extra IR behavior
at **large-distance**

— The best match for the LaMET formalism

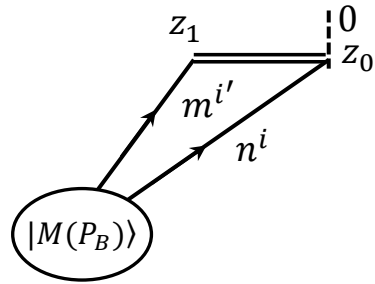


Part 2



• Baryon quasi-DAs on Lattice

➤ Definition of **Meson** Quasi-DAs:

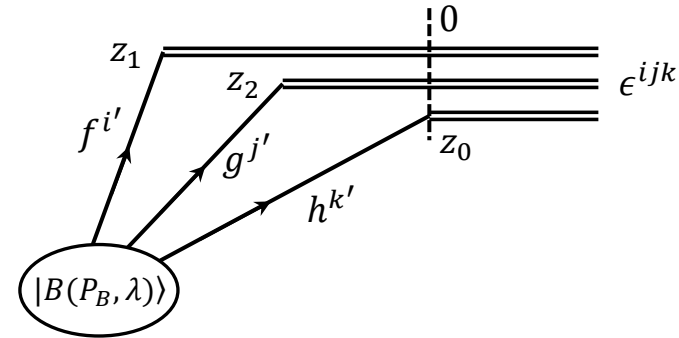


1D



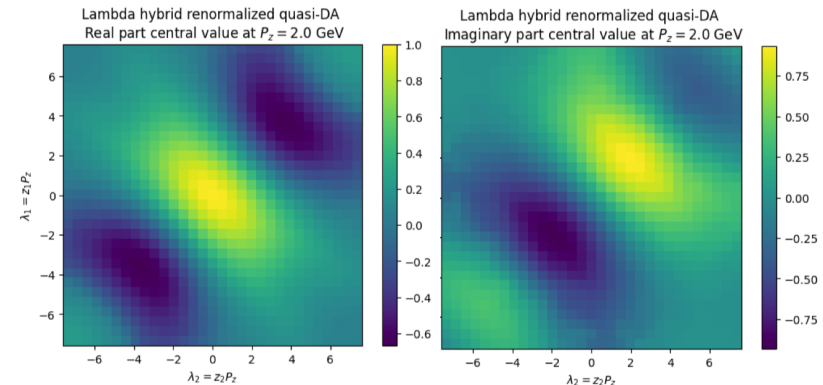
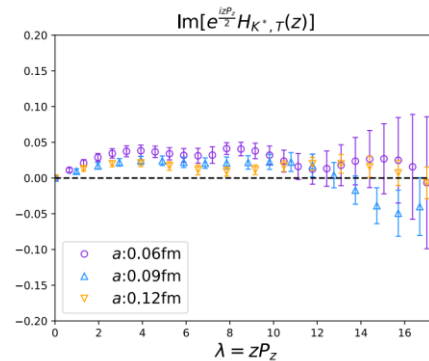
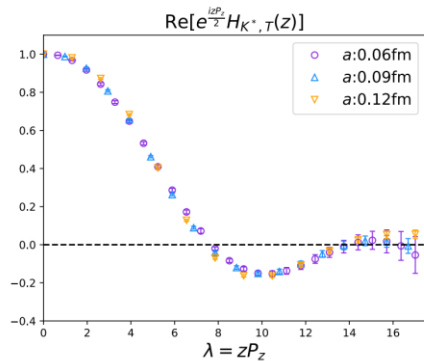
2D

➤ Definition of **Baryon** Quasi-DAs:



$$\hat{O}(\vec{x}, t; z) = m_{\alpha}^{i'}(\vec{x} + zn_z, t) \tilde{\Gamma}_{\alpha\beta} W^{i'i}(\vec{x} + zn_z, \vec{x}, t) n_{\beta}^i(\vec{x}, t)$$

$$\begin{aligned} \hat{O}_{\gamma}(\vec{x}, t; z_1, z_2) &= \epsilon^{ijk} W^{ii'}(\infty, \vec{x} + z_1 n_z) f_{\alpha}^{i'}(\vec{x} + z_1 n_z, t) \\ &\times \tilde{\Gamma}_{\alpha\beta} W^{jj'}(\infty, \vec{x} + z_2 n_z, t) g_{\beta}^{j'}(\vec{x} + z_2 n_z) \\ &\times W^{kk'}(\infty, \vec{x}) h_{\gamma}^{k'}(\vec{x}, t) \end{aligned}$$



- Based on CLQCD Ensembles

CLQCD, PRD 109, 054507

- Three different lattice spacing

| Ensemble | Volume | Lattice spacing | π mass | measurement | P^z |
|----------|-------------------|-----------------|------------|-------------|----------------------------|
| C24P29 | $24^3 \times 72$ | 0.105 fm | 293 MeV | 864*4 | 1.96, 2.45, 2.94, 3.43 GeV |
| F32P30 | $32^3 \times 96$ | 0.077 fm | 303 MeV | 777*4 | 1.99, 2.50, 2.99, 3.49 GeV |
| H48P32 | $48^3 \times 144$ | 0.052 fm | 317 MeV | 550*6 | 1.99, 2.48, 2.98, 3.48 GeV |

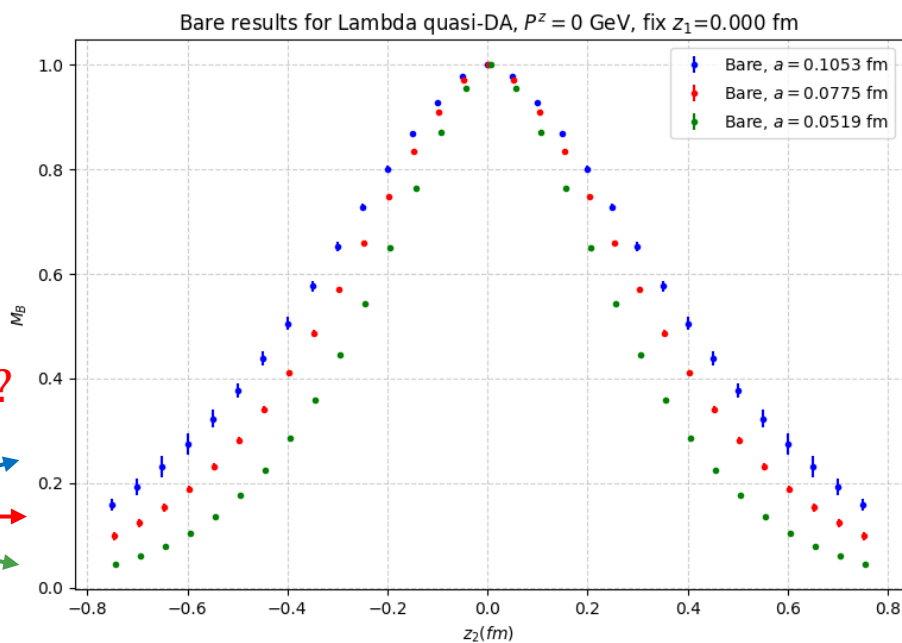
Normed **Bare** quasi-DA $\tilde{\psi}$: $\tilde{\psi}(z_1, z_2, P^Z) = \psi(z_1, z_2, P^Z) / \psi(z_1 = z_2 = 0, P^Z)$

- z-dependence of normed bare $\tilde{\psi}(z_1, z_2, P^Z)$ on different lattice spacing

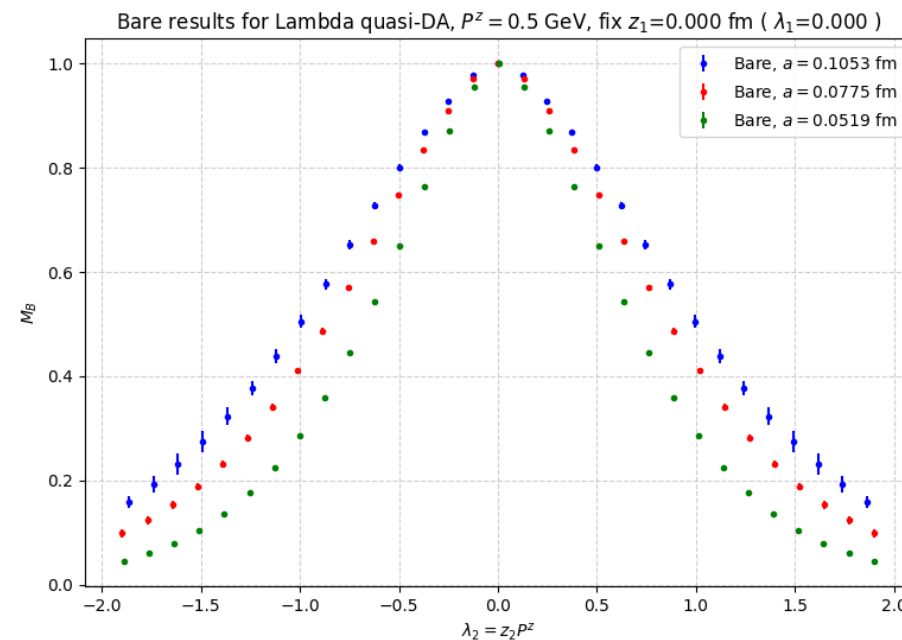
Lambda axial term

fix $z_1=0$ fm

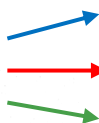
$P^Z=0$ GeV



$P^Z=0.5$ GeV



Divergence?



Bare quasi-DA results: Linear divergence

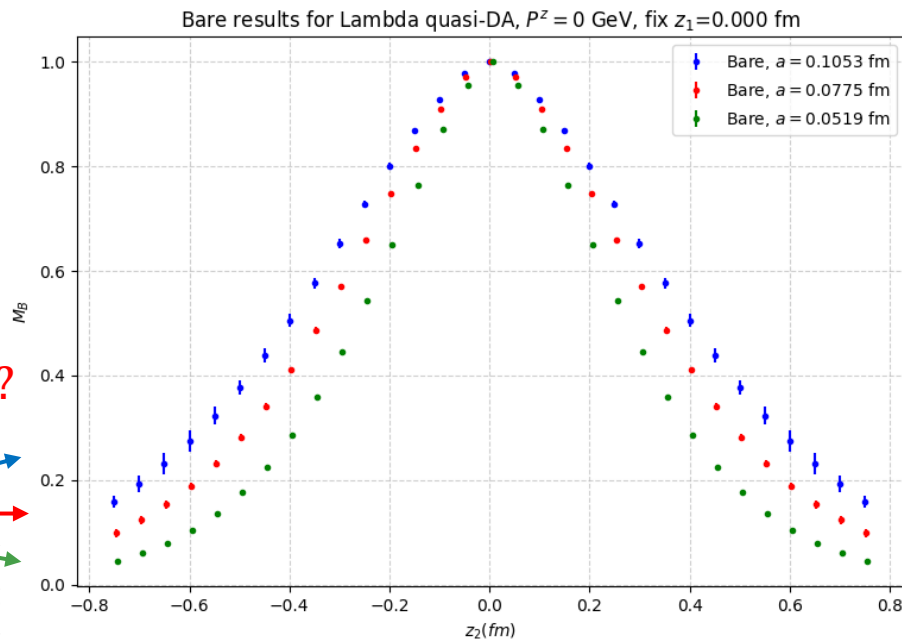
Part ②

Linear divergence form for Normed Bare quasi-DA $\tilde{\psi}$:

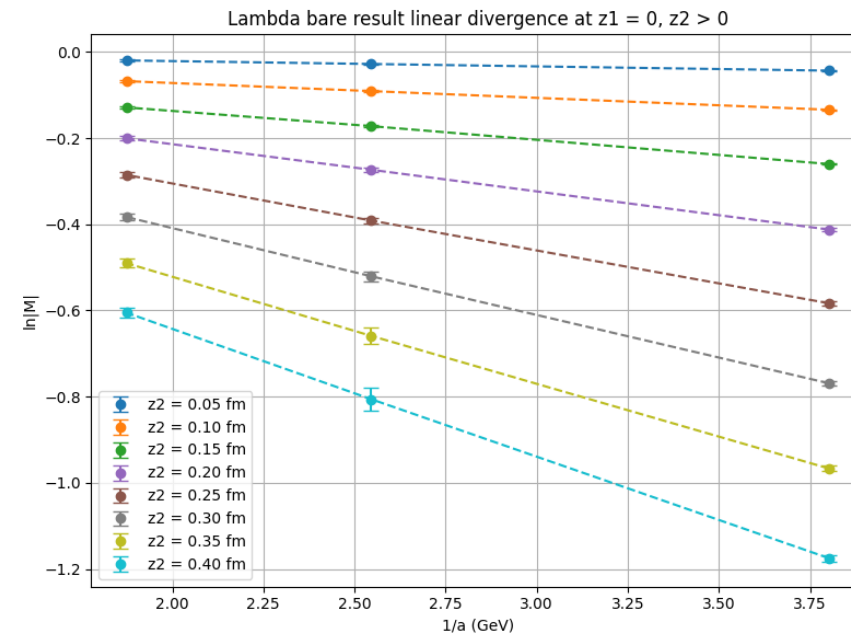
$$\tilde{z} = \begin{cases} |z_1 - z_2| & z_1 z_2 < 0 \\ \max(z_1, z_2) & z_1 z_2 \geq 0 \end{cases}$$

$$M(z_1, z_2; P_z; a) = \exp \left[\left(\frac{k}{a \ln(a \Lambda_{\text{QCD}})} + m_0 \right) \tilde{z} \right] \times m(z_1, z_2; P_z; a)$$

$P^Z = 0 \text{ GeV}$



Log scale



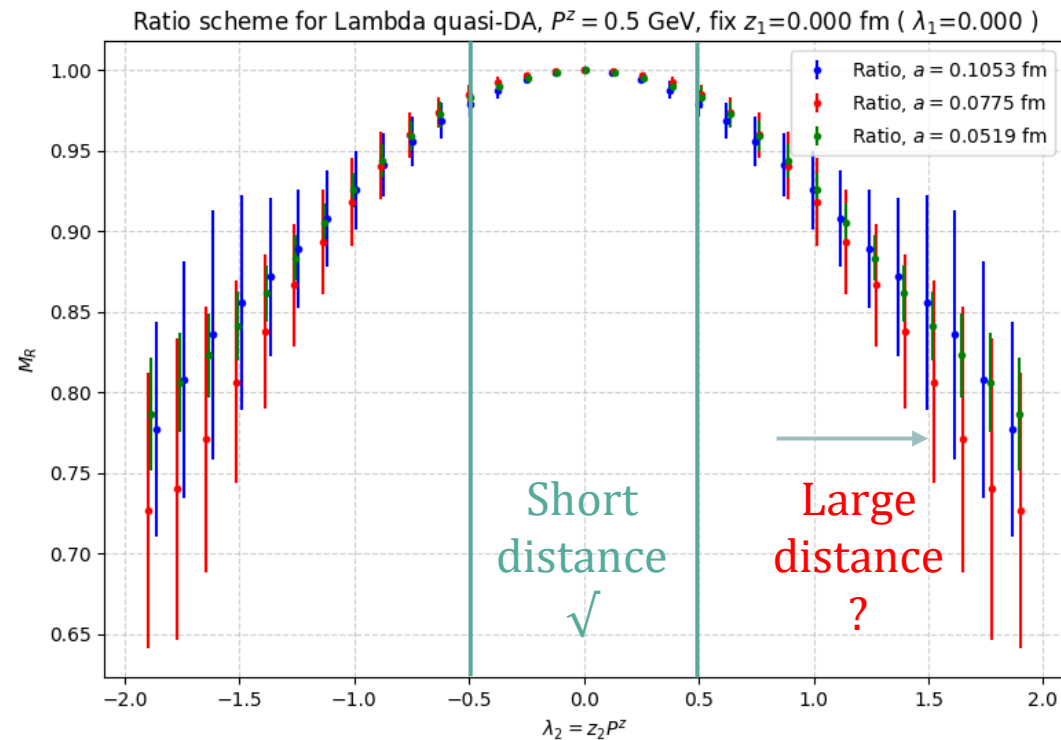
Ratio scheme quasi-DA $\tilde{\psi}^{\text{ratio}}$: $\tilde{\psi}^{\text{ratio}}(z_1, z_2, P^Z) = \tilde{\psi}(z_1, z_2, P^Z) / \tilde{\psi}(z_1, z_2, P^Z = 0)$

- z-dependence of ratio scheme quasi-DA $\tilde{\psi}^{\text{ratio}}(z_1, z_2, P^Z)$ on different lattice spacing

Lambda axial term

$P^Z = 0.5$ GeV

fix $z_1 = 0$ fm



$$\tilde{\psi}(z_1, z_2, P^Z) / \tilde{\psi}(z_1, z_2, P^Z = 0)$$

gives extra IR divergence at large distance



Part 3



Self renormalization

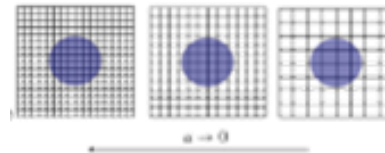
Self renorm: Use Lattice data to renorm them-selves

- 1) Lattice perturbation theory gives a parametrization of the Lattice divergences

$$\hat{M}_{\overline{\text{MS}}}(z_1, z_2, 0, P^z, \mu) = \frac{\hat{M}(z_1, z_2, 0, P^z, a)}{Z_R(z_1, z_2, a, \mu)} \quad \tilde{z} = \begin{cases} |z_1 - z_2| & z_1 z_2 < 0 \\ \max(z_1, z_2) & z_1 z_2 \geq 0 \end{cases}$$

$$Z_R(z_1, z_2, a, \mu) = \exp \left[\left(\frac{k}{a \ln[a\Lambda_{\text{QCD}}]} - \underline{m_0} \right) \tilde{z} + \frac{\gamma_0}{b_0} \ln \left[\frac{\ln[1/(a\Lambda_{\text{QCD}})]}{\ln[\mu/\Lambda_{\text{MS}}]} \right] + \ln \left[1 + \frac{d}{\ln(a\Lambda_{\text{QCD}})} \right] + f(z_1, z_2) a^2 \right]$$

- 2) With data from different lattice ensembles we can fit the terms with lattice spacing dependence



- 3) Remaining terms can be obtained by matching with the perturbative result at the short-distance region.

| Ensemble | Volume | Lattice spacing |
|----------|-------------------|-----------------|
| C24P29 | $24^3 \times 72$ | 0.1052 fm |
| F32P30 | $32^3 \times 96$ | 0.0775 fm |
| H48P32 | $48^3 \times 144$ | 0.0520 fm |

Self-renormalization: 2-step fitting

$$\tilde{z} = \begin{cases} |z_1 - z_2| & z_1 z_2 < 0 \\ \max(z_1, z_2) & z_1 z_2 \geq 0 \end{cases}$$

Part ③

- **Step-1** (fitting the lattice spacing dependence)

$$\ln M(z_1, z_2; P_z = 0; a) = \frac{k}{a \ln(a\Lambda_{\text{QCD}})} \tilde{z} + g(z_1, z_2) + \underline{f(z_1, z_2)} a^2 + \frac{\gamma_0}{b_0} \ln \frac{\ln(1/a\Lambda_{\text{QCD}})}{\ln(\mu/\Lambda_{\overline{\text{MS}}})} + \ln \left[1 + \frac{d}{\ln(a\Lambda_{\text{QCD}})} \right]$$

b_0

fix $\Lambda_{\text{QCD}} = 0.2 \text{ GeV}$

- **Step-2** (matching the perturbative result)

$$g(z_1, z_2) - \ln Z_{\overline{\text{MS}}}(z_1, z_2; \mu, \Lambda_{\overline{\text{MS}}}) = \underline{m_0} \tilde{z} + \underline{b_0}$$

$$Z_{\overline{\text{MS}}}(z_1, z_2; \mu, \Lambda_{\overline{\text{MS}}}) = 1 + \frac{\alpha_s C_F}{2\pi} \left[\frac{7}{8} \ln \frac{z_1^2 \mu^2 e^{2\gamma_E}}{4} + \frac{7}{8} \ln \frac{z_2^2 \mu^2 e^{2\gamma_E}}{4} + \frac{3}{4} \ln \frac{(z_1 - z_2)^2 \mu^2 e^{2\gamma_E}}{4} + 4 \right]$$

- **Re- matrix element & renorm factor**

$$M_R(z_1, z_2; P_z = 0) = \exp \left(g(z_1, z_2) - m_0 \tilde{z} - b_0 \right)$$

$$Z_R(z_1, z_2; a) = \exp \left[\frac{k}{a \ln(a\Lambda_{\text{QCD}})} \tilde{z} + m_0 \tilde{z} + f(z_1, z_2) a^2 + \frac{\gamma_0}{b_0} \ln \frac{\ln(1/a\Lambda_{\text{QCD}})}{\ln(\mu/\Lambda_{\overline{\text{MS}}})} + \ln \left[1 + \frac{d}{\ln(a\Lambda_{\text{QCD}})} \right] \right]$$

Self fitting step-1

Part ③

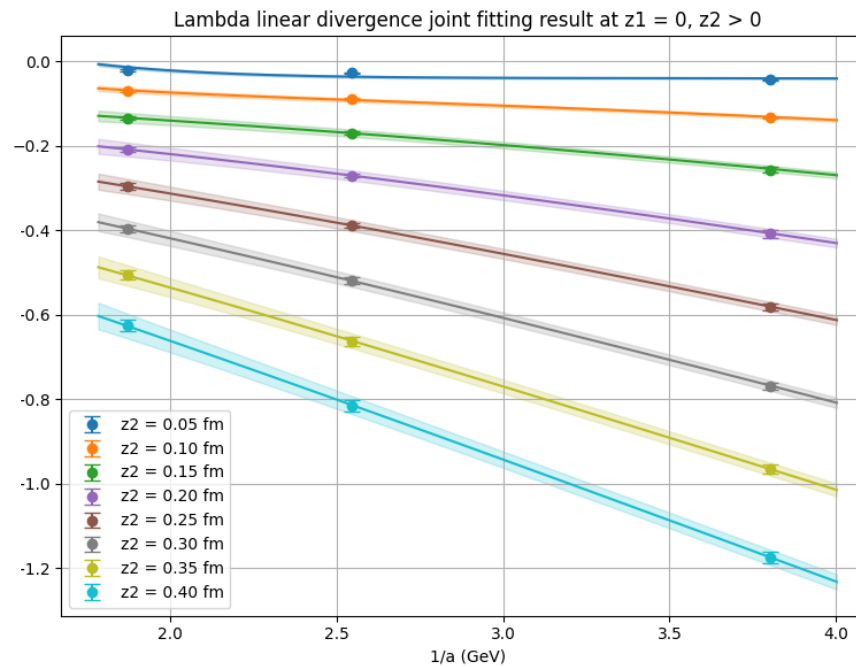
$$\ln M(z_1, z_2; P_z = 0; a) = \frac{k}{a \ln(a\Lambda_{\text{QCD}})} \tilde{z} + g(z_1, z_2) + \underline{f(z_1, z_2)} a^2 + \frac{\gamma_0}{b_0} \ln \frac{\ln(1/a\Lambda_{\text{QCD}})}{\ln(\mu/\Lambda_{\overline{\text{MS}}})} + \ln \left[1 + \frac{d}{\ln(a\Lambda_{\text{QCD}})} \right]$$

Step-1 fit result

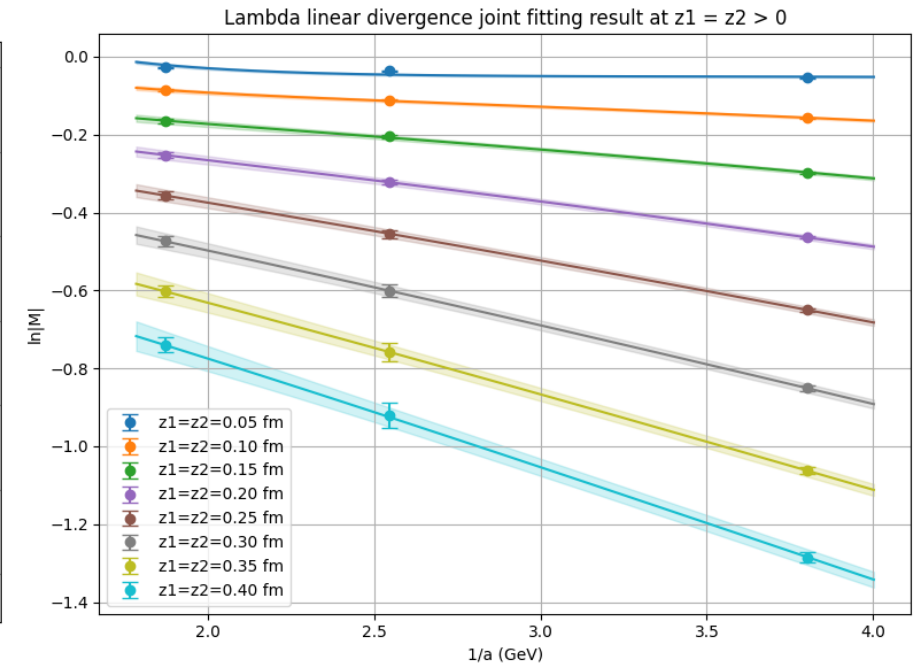
Interpolate to $a_0 = 0.05$ fm

$z_1, z_2: 0.15 \sim 0.70$ fm

Fit result $z_1=0$ (log scale)



Fit result $z_1=z_2$ (log scale)



Parameters:

| | | | |
|-----------|-------------|---|----------|
| k | 0.7792 (27) | [| 0 ± 10] |
| g_-14_-10 | 1.020 (39) | [| 0 ± 10] |
| f_-14_-10 | 1.2 (5.9) | [| 0 ± 10] |
| g_-14_-9 | 1.066 (40) | [| 0 ± 10] |
| f_-14_-9 | 0.4 (5.5) | [| 0 ± 10] |
| g_-14_-8 | 1.100 (43) | [| 0 ± 10] |
| f_-14_-8 | -0.3 (5.4) | [| 0 ± 10] |
| g_-14_-7 | 1.124 (46) | [| 0 ± 10] |
| f_-14_-7 | -0.7 (5.4) | [| 0 ± 10] |
| g_-14_-6 | 1.135 (47) | [| 0 ± 10] |
| f_-14_-6 | -0.6 (5.4) | [| 0 ± 10] |
| g_-14_-5 | 1.133 (46) | [| 0 ± 10] |
| f_-14_-5 | -0.2 (5.3) | [| 0 ± 10] |
| g_-14_-4 | 1.119 (44) | [| 0 ± 10] |
| f_-14_-4 | 0.5 (5.1) | [| 0 ± 10] |
| g_-14_4 | 1.209 (43) | [| 0 ± 10] |
| f_-14_4 | 6.0 (6.2) | [| 0 ± 10] |
| g_-14_5 | 1.253 (49) | [| 0 ± 10] |

Self fitting step-2

Part ③

$$g(z_1, z_2) - \ln Z_{\overline{\text{MS}}}(z_1, z_2; \mu, \Lambda_{\overline{\text{MS}}}) = \underline{m_0} \tilde{z} + \underline{b_0}$$

$$Z_{\overline{\text{MS}}}(z_1, z_2; \mu, \Lambda_{\overline{\text{MS}}}) = 1 + \frac{\alpha_s C_F}{2\pi} \left[\frac{7}{8} \ln \frac{z_1^2 u^2 e^{2\gamma_E}}{4} + \frac{7}{8} \ln \frac{z_2^2 u^2 e^{2\gamma_E}}{4} + \frac{3}{4} \ln \frac{(z_1 - z_2)^2 \mu^2 e^{2\gamma_E}}{4} + 4 \right]$$

Step-2 fit result

z_1, z_2 : 0.05 ~ 0.20 fm

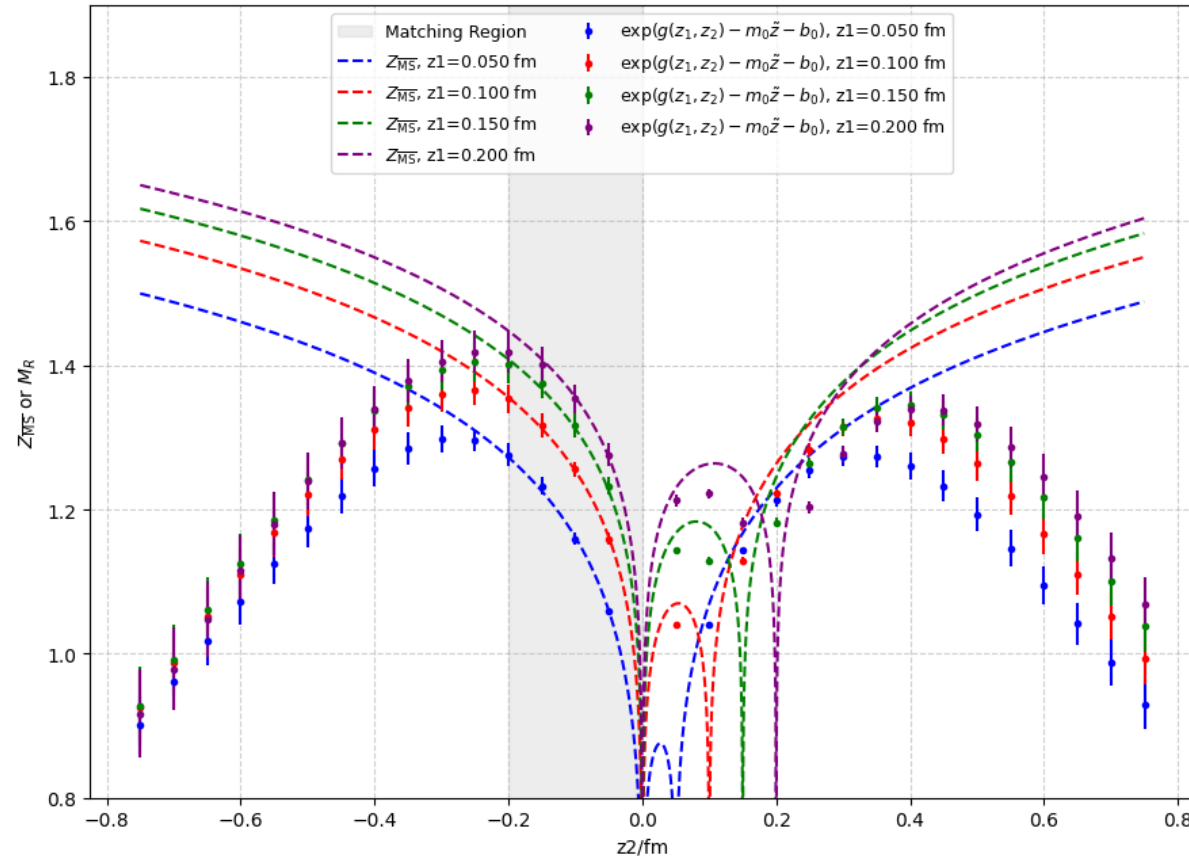
Parameters:

| | | | | |
|-------|-------------|---|------------|---|
| m_0 | 1.254 (11) | [| 0 ± 10 |] |
| b_0 | 0.1093 (23) | [| 0 ± 10 |] |

Fit:

| key | y[key] | f(p)[key] |
|--------|-------------|----------------|
| g_-4_1 | 0.6632 (58) | 0.6636 (11) |
| g_-4_2 | 0.7873 (73) | 0.7897 (15) |
| g_-4_3 | 0.8843 (88) | 0.8905 (20) |
| g_-4_4 | 0.960 (10) | 0.9808 (25) * |
| g_-3_1 | 0.5688 (45) | 0.56330 (87) * |
| g_-3_2 | 0.7001 (60) | 0.6943 (11) |
| g_-3_3 | 0.8029 (74) | 0.7981 (15) |
| g_-3_4 | 0.8843 (88) | 0.8905 (20) |
| g_-2_1 | 0.4438 (39) | 0.44705 (95) |
| g_-2_2 | 0.5897 (45) | 0.58605 (87) |
| g_-2_3 | 0.7001 (60) | 0.6943 (11) |
| g_-2_4 | 0.7873 (73) | 0.7897 (15) |
| g_-1_1 | 0.2911 (23) | 0.2924 (13) |

Lambda result of matching the lattice matrix elements to the MSbar scheme perturbative expression



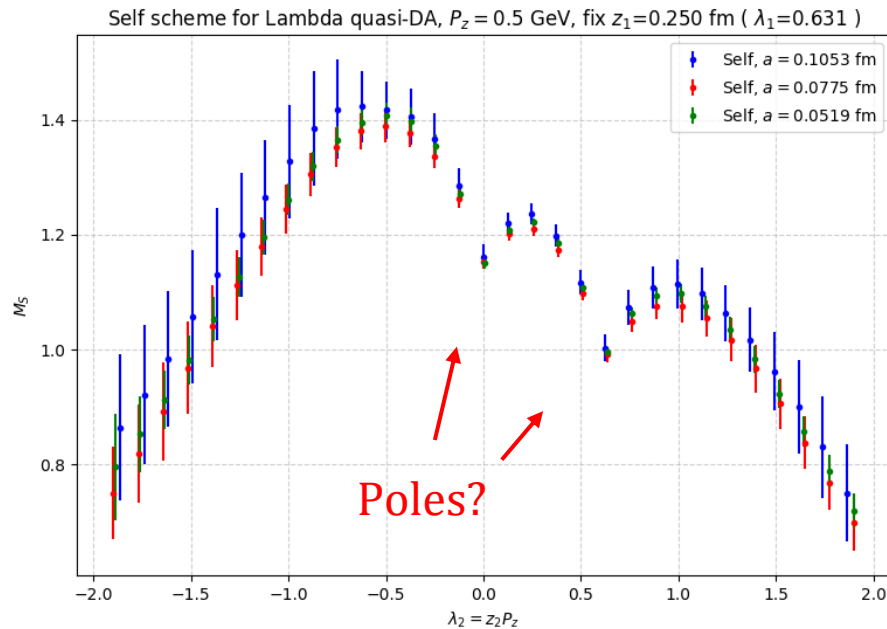
Self scheme quasi-DA $\tilde{\psi}^{\text{self}}$:

$$M_R(z_1, z_2; P_z) = \frac{M(z_1, z_2; P_z; a)}{Z_R(z_1, z_2; a)}$$

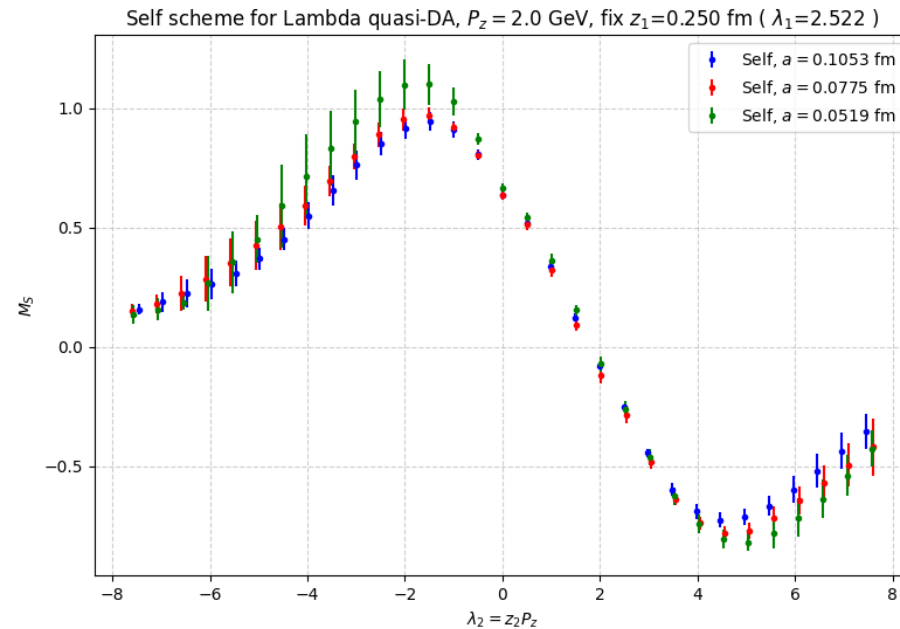
- z-dependence of quasi-DA $\tilde{\psi}^{\text{self}}(z_1, z_2, P^Z)$ on different lattice spacing

Lambda Axial term fix $z_1=0.250$ fm

$P^Z=0.5$ GeV



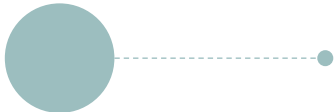
$P^Z=2.0$ GeV



\overline{MS} : $\ln(\mu^2 z_1^2)$, $\ln(\mu^2 z_2^2)$, and $\ln(\mu^2 (z_1 - z_2)^2)$

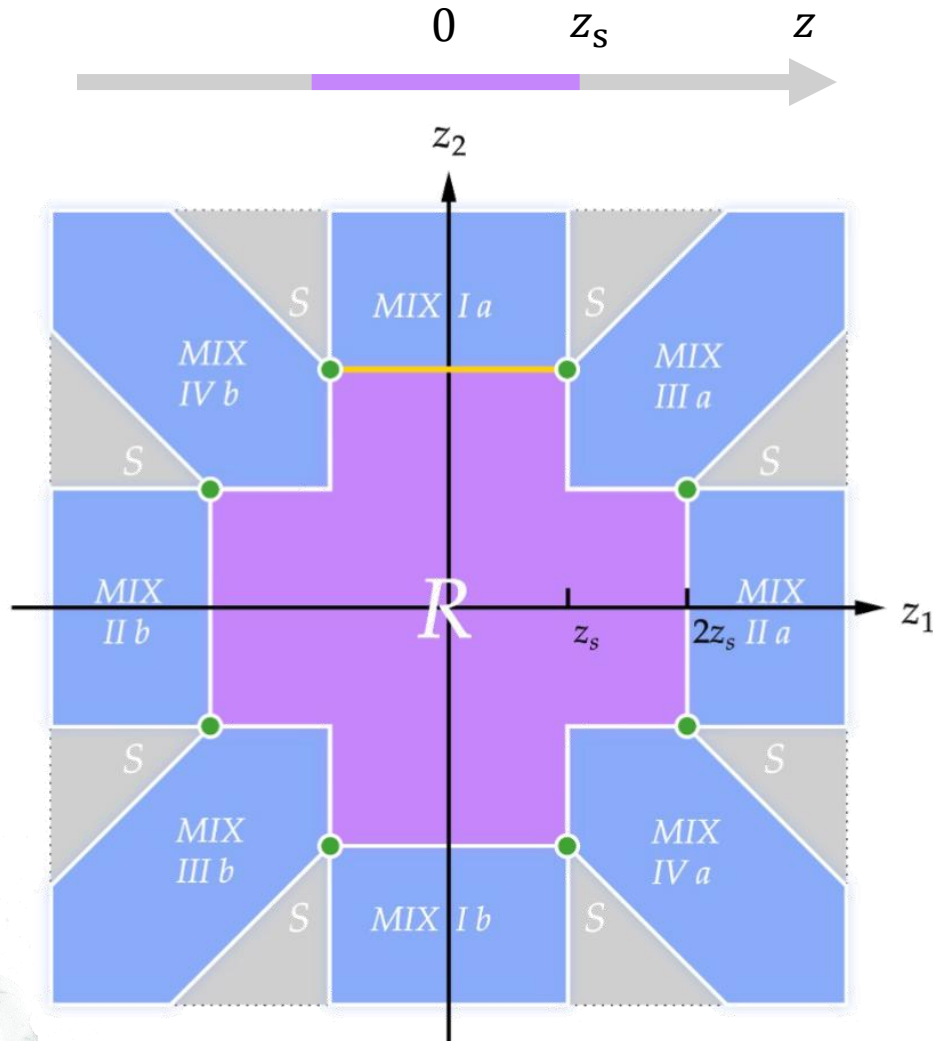
A large teal circle containing the text "Part 4".

Part 4



Hybrid scheme

Hybrid scheme: Use suitable schemes for different regions



- 1) Short-distance region

$$\frac{\hat{M}_{\overline{MS}}(z_1, z_2, 0, P^z, \mu)}{\hat{M}_{\overline{MS}}(z_1, z_2, 0, 0, \mu)} \cdot S_{\text{short}}(z_1, z_2)$$

- 2) Long-distance region

$$\frac{\hat{M}_{\overline{MS}}(z_1, z_2, 0, P^z, \mu)}{\hat{M}_{\overline{MS}}(\text{sign}(z_1)z_s, \text{sign}(z_2)2z_s, 0, 0, \mu)} \cdot S_{\text{long}}(z_1, z_2)$$

- 3) Mixing regions

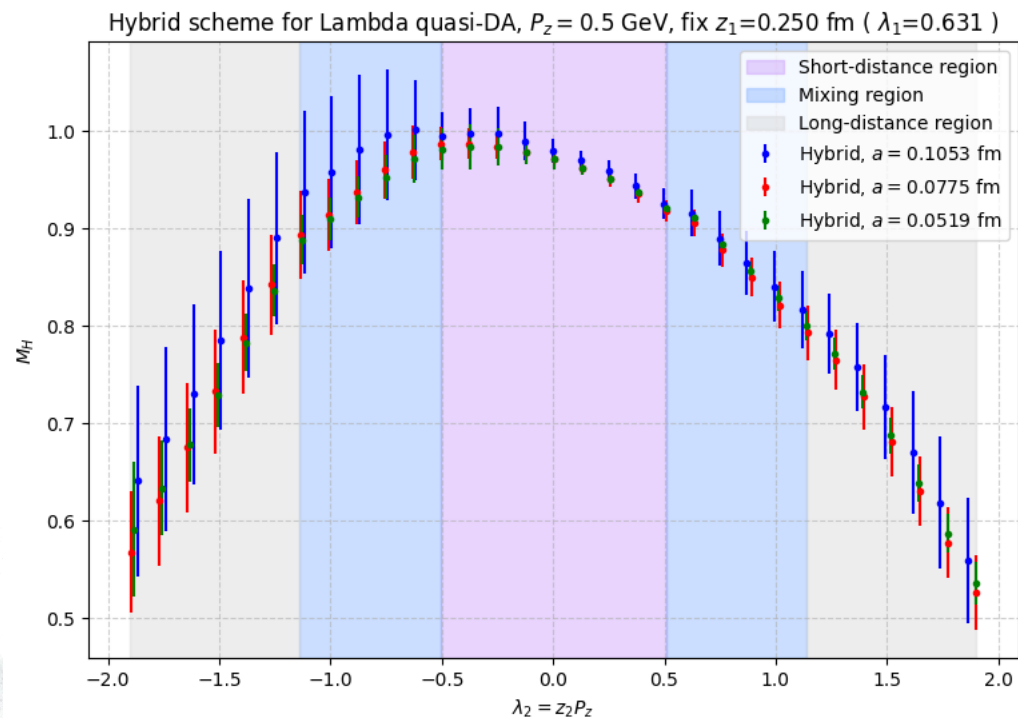
$$\frac{\hat{M}_{\overline{MS}}(z_1, z_2, 0, P^z, \mu)}{\hat{M}_{\overline{MS}}(z_1, \text{sign}(z_2)2z_s, 0, 0, \mu)} \cdot \theta(z_s - |z_1|)\theta(|z_2| - 2z_s)$$

Hybrid scheme quasi-DA $\tilde{\psi}^{\text{hybrid}}$

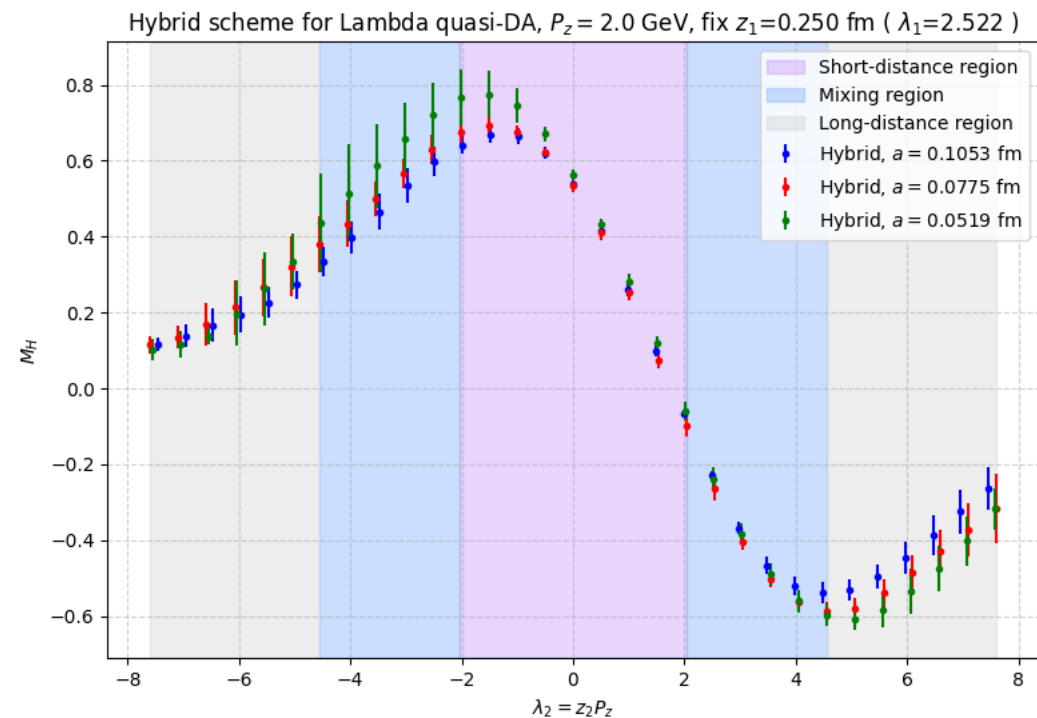
- z-dependence of quasi-DA $\tilde{\psi}^{\text{hybrid}}(z_1, z_2, P^z)$ on different lattice spacing

Lambda A term fix $z_1=0.250$ fm

$P^z=0.5$ GeV



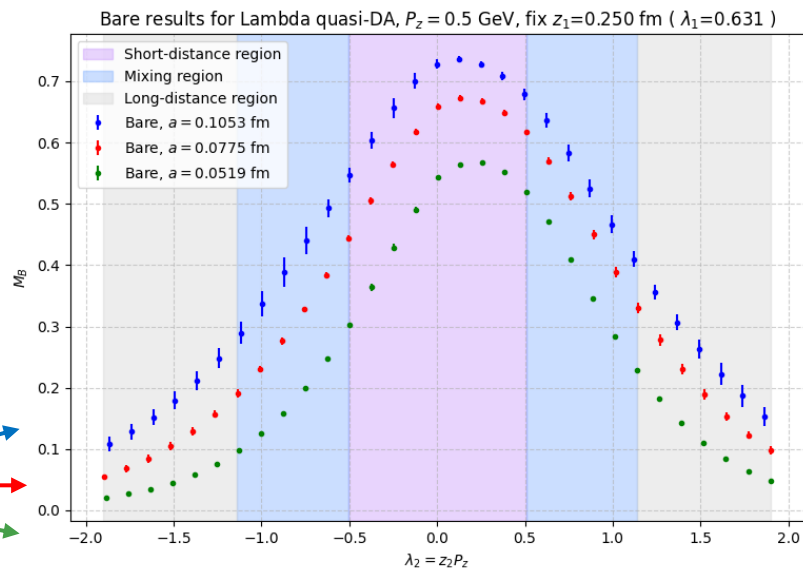
$P^z=2.0$ GeV



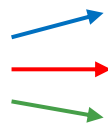
Schemes Comparison (Hybrid)

Part ④

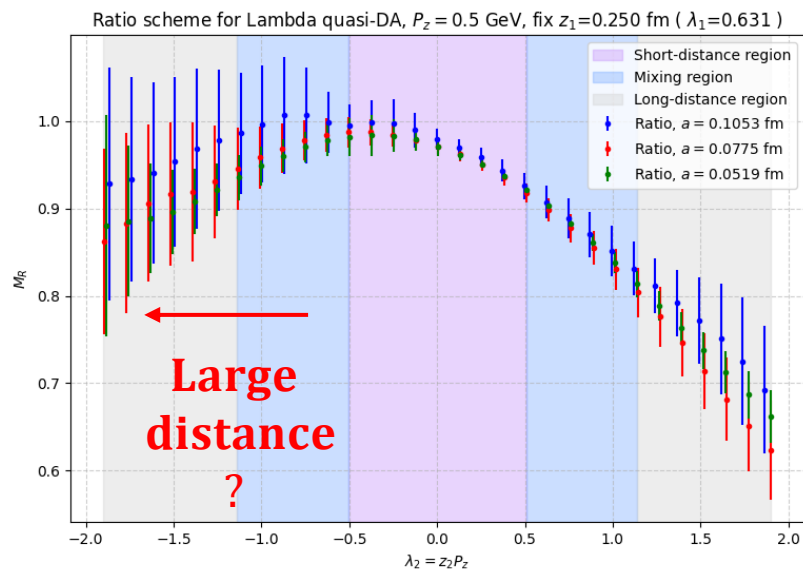
Bare Matrix Element



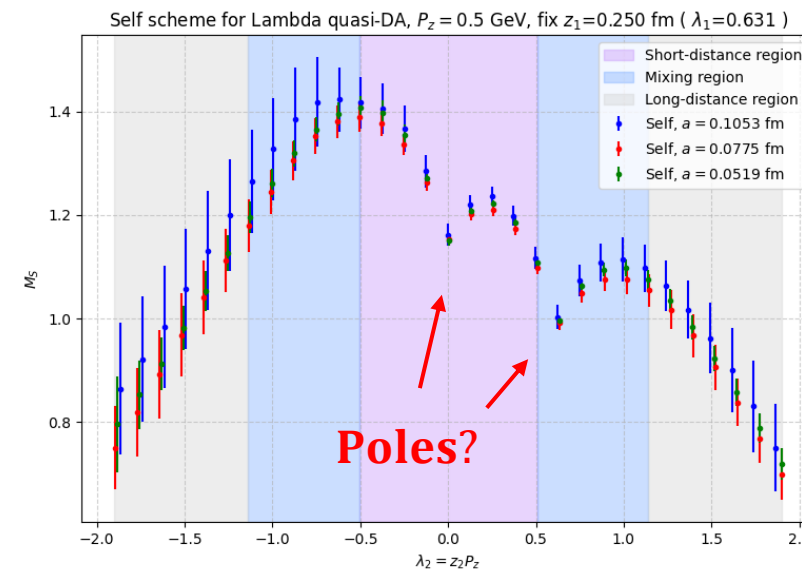
Divergence?



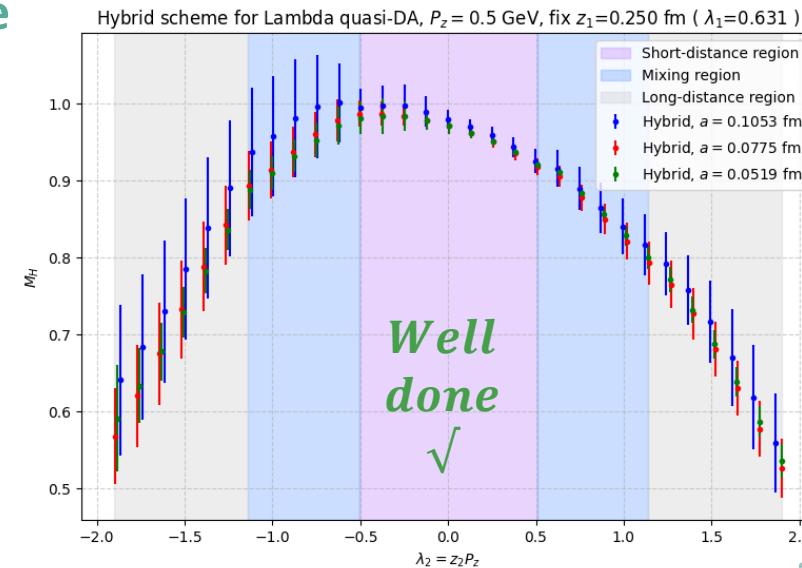
Ratio scheme



Self scheme



Hybrid scheme



- Hybrid renormalized large momentum baryon quasi-DA ($P^z = 2.0\text{GeV}$)

Λ quasi-DA

Proton quasi-DA

Re

Im

Re

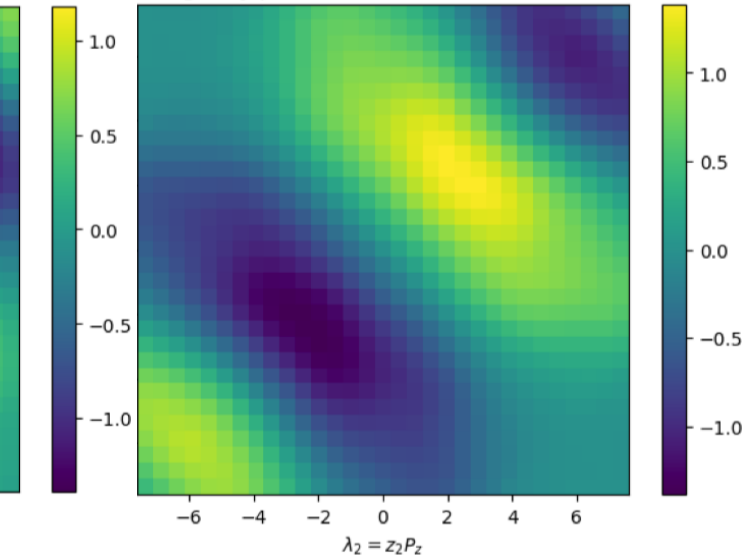
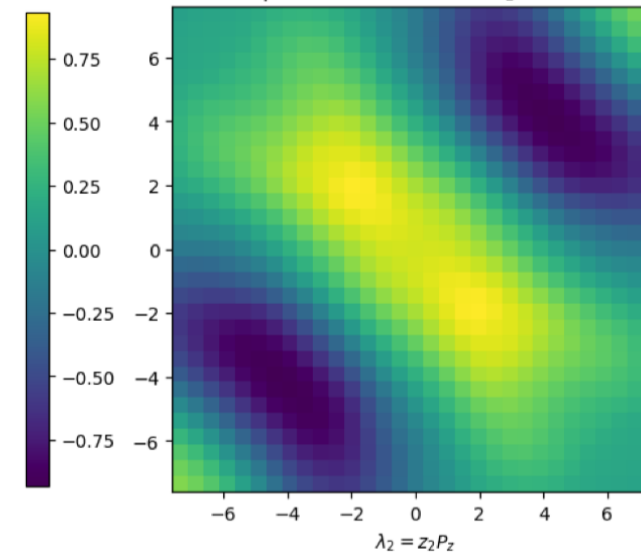
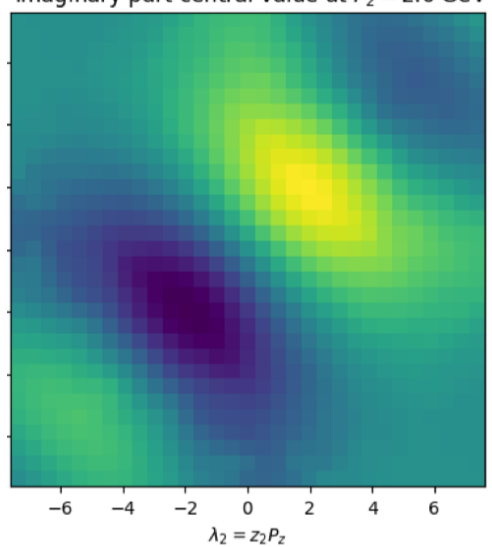
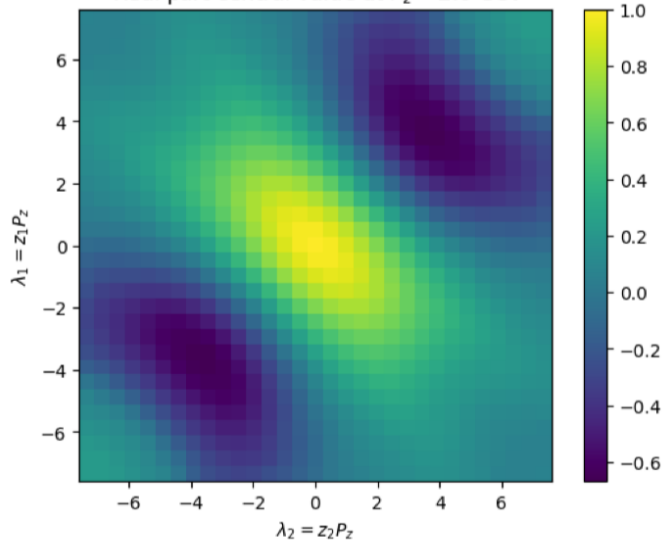
Im

Lambda hybrid renormalized quasi-DA
Real part central value at $P_z = 2.0\text{ GeV}$

Lambda hybrid renormalized quasi-DA
Imaginary part central value at $P_z = 2.0\text{ GeV}$

Proton hybrid renormalized quasi-DA
Real part central value at $P_z = 2.0\text{ GeV}$

Proton hybrid renormalized quasi-DA
Imaginary part central value at $P_z = 2.0\text{ GeV}$



well-defined and smooth in all regions

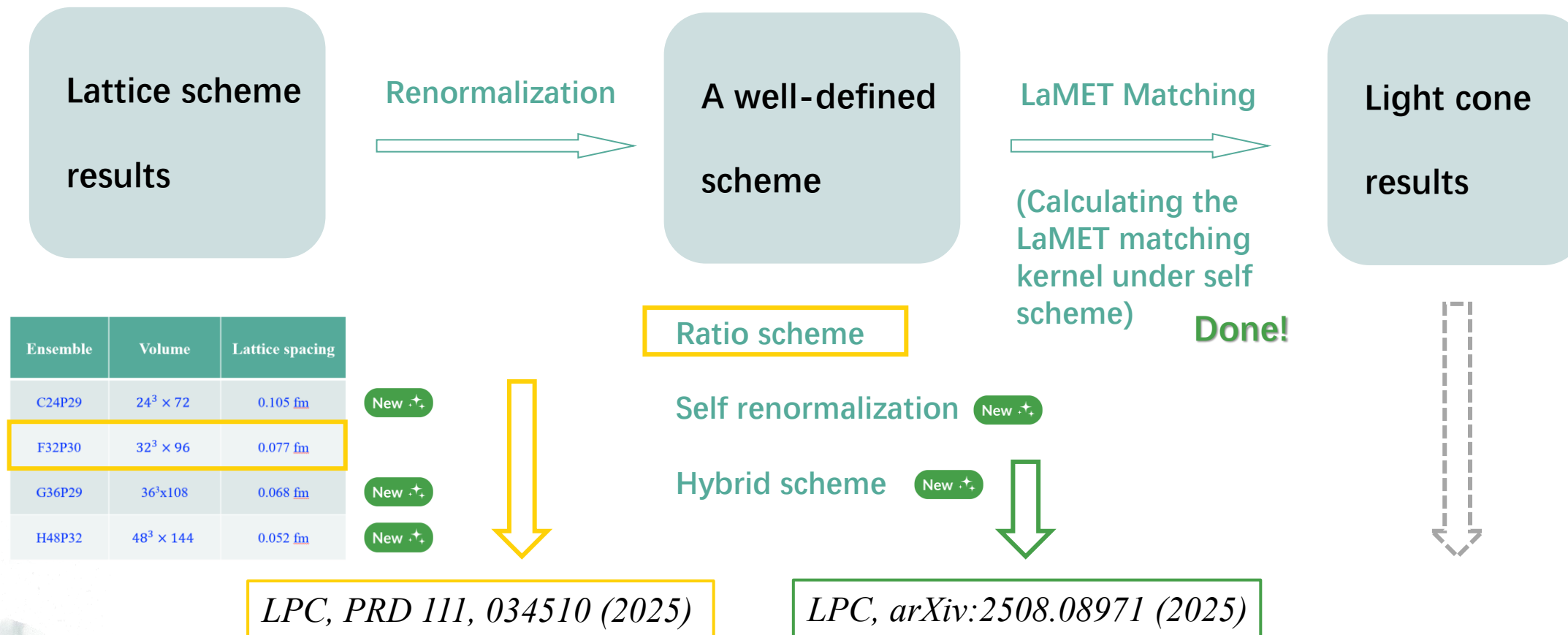


Part 5



Summary

Towards Reliable Light Baryon LCDA on Lattice:



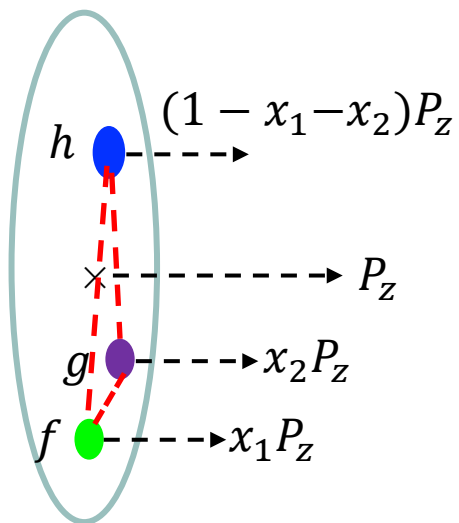
Thanks For Your Attention!



Part 6



Backup



- Definition of baryon LCDA:**

$$M_L(\xi_1, \xi_2; P) = \langle 0 | \epsilon^{ijk} f_\alpha^{i'}(\xi_1 n) W_{i'i}(\xi_1, \xi_0) g_\beta^{j'}(\xi_2 n) W_{j'j}(\xi_2, \xi_0) h_\gamma^k(\xi_0 n) | B(P, \lambda) \rangle$$

$$\Phi(x_1, x_2, \mu) = \int \frac{dP^+\xi_1}{2\pi} \frac{dP^+\xi_2}{2\pi} e^{ix_1P^+\xi_1 + ix_2P^+\xi_2} \frac{M_L(\xi_1, \xi_2; P, \mu)}{M_L(0, 0; P, \mu)}$$

C.Han et.al. JHEP 07019 (2024);

V.L.C & I.R.Z NPB 24652(1984); G.R.Farrar et.al. NPB 311585(1989)

- Leading twist octet baryon LCDA:**

$$\langle 0 | f_\alpha(z_1 n) g_\beta(z_2 n) h_\gamma(z_3 n) | B(P_B, \lambda) \rangle$$

$$= \frac{1}{4} f_V \left[(P_B C)_{\alpha\beta} (\gamma_5 u_B)_\gamma V^B(z_i n \cdot P_B) + (P_B \gamma_5 C)_{\alpha\beta} (u_B)_\gamma A^B(z_i n \cdot P_B) \right]$$

$$+ \frac{1}{4} f_T (i\sigma_{\mu\nu} P_B^\nu C)_{\alpha\beta} (\gamma^\mu \gamma_5 u_B)_\gamma T^B(z_i n \cdot P_B),$$

| Octet | n | p | Λ |
|-------|-----|-----|-----------|
| fgh | ddu | uud | uds |

➤ Definition of the Operators on Lattice

Matrix element

$$C_2(z_a, z_2; t, \vec{P}) = \int d^3x e^{-i\vec{P}\cdot\vec{x}} \langle 0 | \mathcal{O}_{\text{Sink}}(\vec{x}, t; z_1, z_2) \bar{\mathcal{O}}_{\text{Src}}(0, 0; 0, 0) T | 0 \rangle$$

Source-side Operator

$$\begin{aligned} \bar{O}_{\text{mod}}^P &= (u^T C \gamma_5 \gamma^t d) u, \\ \bar{O}_{\text{mod}}^\Lambda &= \frac{2(u^T C \gamma_5 \gamma^t d) s + (u^T C \gamma_5 \gamma^t s) d + (s^T C \gamma_5 \gamma^t d) u}{\sqrt{6}}. \end{aligned}$$

Projection Operator

$$T = \gamma_t + \gamma_z$$

Sink-side Operator

$$\begin{aligned} O_P^V &= [u^T(z_1 n_z) (C \gamma^t) u(z_2 n_z)] \gamma_5 d(z_3 n_z) & O_\Lambda^V &= [u^T(z_1 n_z) (C \gamma^t) d(z_2 n_z)] \gamma_5 s(z_3 n_z) \\ O_P^A &= [u^T(z_1 n_z) (C \gamma_5 \gamma^t) u(z_2 n_z)] d(z_3 n_z) & O_\Lambda^A &= [u^T(z_1 n_z) (C \gamma_5 \gamma^t) d(z_2 n_z)] s(z_3 n_z) \\ O_P^T &= [u^T(z_1 n_z) \frac{1}{2} C [\gamma^t, \gamma^\mu] u(z_2 n_z)] \gamma_5 \gamma_\mu d(z_3 n_z) & O_\Lambda^T &= [u^T(z_1 n_z) \frac{1}{2} C [\gamma^t, \gamma^\mu] d(z_2 n_z)] \gamma_5 \gamma_\mu s(z_3 n_z) \end{aligned}$$