

# Quasi-Static Lineshape Theory for Rydberg Excitations in Various Bose-Einstein Condensates



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## Introduction

- Rydberg atom: an atom with a very excited electron ( $n > 25$ ). In such atoms, the core resembles a hydrogen nucleus.
- Lineshape: a distribution of relative frequencies of detunings for electron excitations.
- Detuning: the difference in energetic shift between an undisturbed excitation and an excitation in the presence of perturbors, which are atoms within the electron's orbit.

## Theoretical Approach

The detuning of an excitation is the effect of perturbors on initial and final energies during photon absorption:

$$\omega = (E_f - E_i) + (\Delta E_f - \Delta E_i) = \omega_0 + \Delta\omega \quad (1)$$

The interaction between the Rydberg electron and perturber has a 1<sup>st</sup> and 2<sup>nd</sup> order component associated with s/p wave scattering respectively. The pseudopotential in (2) shows a dirac function (a semiclassical analog) instead of the wave function squared, which reflects a quantum perspective [3]:

$$V_{e-p} = 2\pi a_s(k) \delta^{(3)}(\mathbf{r} - \mathbf{R}) + 6\pi (a_p(k))^3 \delta^{(3)}(\mathbf{r} - \mathbf{R}) \nabla^2 \quad (2)$$

The neutral perturbors will also polarize and interact with the ionic Rydberg core as follows [3]:

$$V_{c-p} = \frac{-\alpha}{2r^4} \quad (3)$$

The perturber distributions are calculated using statistical mechanics. The Rydberg atom is positioned at the center of an harmonic trap in a gaseous Bose-Einstein condensate. The density distribution can be represented as the sum of BEC and Thermal atom densities [3, 4]:

$$N_{BEC}(\rho) = \begin{cases} \frac{m^2 \omega^2}{8\pi a_{bb}} (R_{TF}^2 - \rho^2) & 0 \leq \rho \leq R_{TF} \\ 0 & \rho > 0 \end{cases} \quad (4)$$

$$N_{th}(\rho) = \frac{1}{\lambda_T^3} Li_{3/2} \left( e^{\frac{k_B T}{\hbar^2} \left( \frac{1}{2} m \omega^2 \rho^2 - R_{TF}^2 + \frac{8\pi a_{bb}}{m} N_{th}(\rho) \right)} \right) \quad (5)$$

## Bose-Einstein Condensate

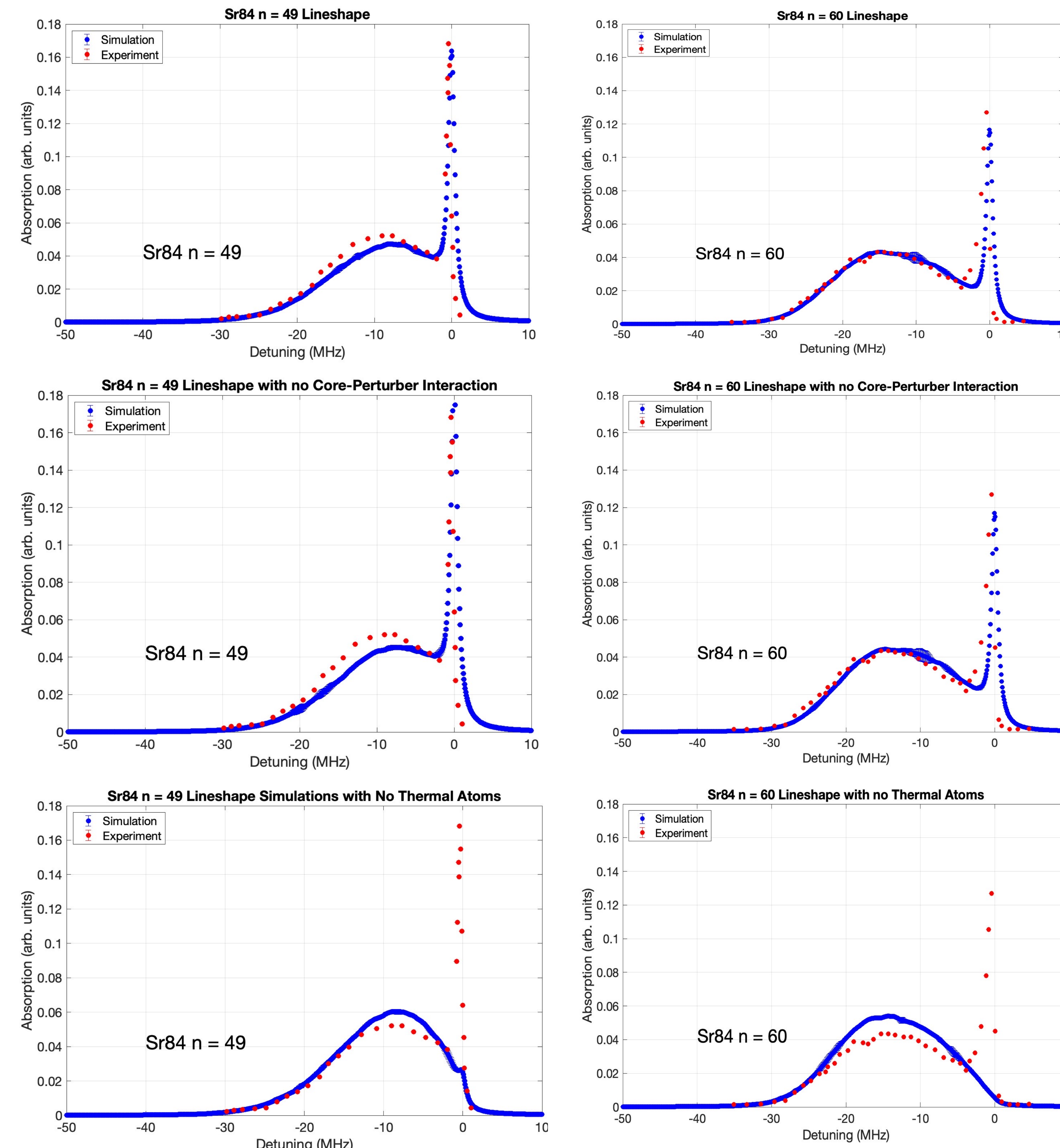
Indian physicist Satyendra Bose postulated that at low temperatures, the wave-particle duality of matter would result in wave packets with a macroscopic thermal de Broglie wavelength [1]:

$$\lambda = \frac{h}{p} = \sqrt{\frac{h^2}{mk_B T}} \quad (6)$$

At sufficiently low temperature, every particle's associated wavelength would become long enough to overlap with the others, leading every particle to occupy the same ground state and act as a unified quantum system. The BEC arises as a novel state of matter that allows physicists to observe and study quantum mechanics at a macroscopic scale.

## Purpose

Develop a general computational method for simulating lineshapes of Rydberg atoms in a Bose-Einstein Condensate (BEC) and analyze the model's dependence on BEC density using Sr84 and Rb87 as test cases. Theoretical lineshape simulation is useful because lineshapes can illuminate certain hard-to-measure Rydberg properties, such as electron-neutral interactions or the density of the media.



## Future Works

- Threshold densities: the current BEC densities ensure a single electron-perturber interaction and a relatively low dipole-dipole interaction potential
- Simulate higher densities to find transition to multiple perturbors and significant dipole-dipole interaction

## References

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- E. Fermi, "Sopra lo spostamento per pressione delle righe elevate delle serie spettrali," Il Nuovo Cimento (1924-1942), vol. 11, no. 3, p. 157, 2008.
- J. Perez-Rios, "Introduction to Cold and Ultracold Chemistry", Springer, 2020.
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## Ultracold Regime

- At temperatures below 100 nK, The particles' de Broglie wavelengths are macroscopic and they each occupy the same quantum state (BEC) [3].
- The Rydberg atom is placed within an harmonic trap that contains a distribution of mostly BEC gas with a smaller number of thermal atoms.
- Rydberg electrons can bind to perturbors to form a diatom known as a Rydberg polaron, which has a very large dipole moment despite being homonuclear [3].

## Considerations

- Quasi-static theory – Calculations assume negligible particle motion during excitation
- Perturbation Order – 1<sup>st</sup>/2<sup>nd</sup> (s/p)
- Thermal atoms
  - Internally excited beyond ground-state
- Multiple perturbors:  $1 < M < 5$
- Induced dipole-dipole interactions between perturbors

## Methodology

- Mathematica – original coding language for simulations
- Matlab – simulation notebooks converted from Mathematica

The lineshape simulation process is as follows:

- The harmonic trap probability distributions are scaled by radial distance from the center squared:  $P(\rho) \sim N(\rho) \rho^2$  (7)
- Choose number of nearby perturber atoms by designating the threshold radius as  $2.5(n^*)^2$ , and run different different numbers of perturbors according to a Poisson distribution.
- Choose the distance of each perturber from the Rydberg core. The vicinal density is constant since the max perturber distance is much less than the trap radius. Thus,  $P(r) \sim r^2$ .
- Total effect of perturbors is summed up using interaction potentials to get detunings:

$$V_{c-p} + \sum_1^M V_{e-p} + \sum_{i=1}^{N-1} \sum_{j=i+1}^N V_{p_i - p_j} \quad (8)$$

- This set of detunings is converted to a histogram to generate an approximate lineshape, which is then converted to a proper lineshape by accounting for the light bandwidth [4]:

$$I(\nu) = \sum_i \frac{1}{(\nu - \nu_i^2) + (\Gamma/2)^2} \quad (9)$$