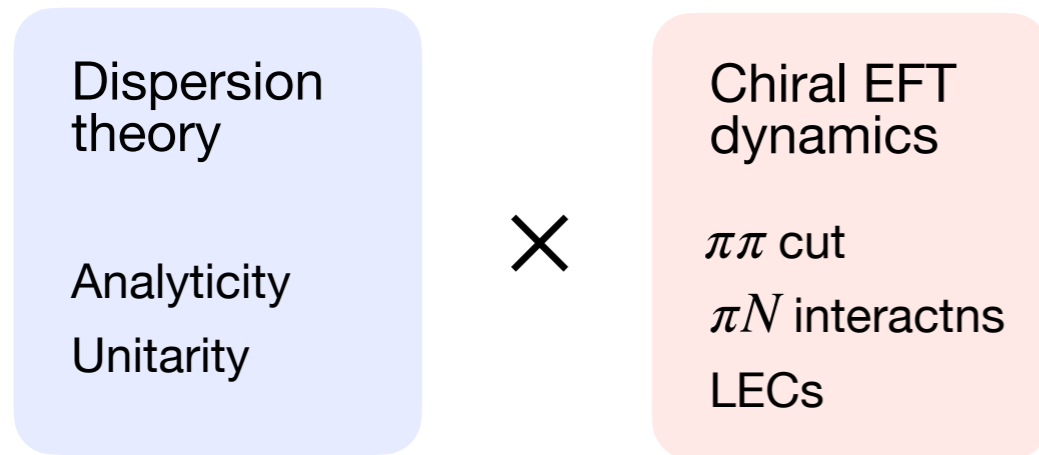


Low- Q^2 nucleon form factors from dispersion theory and chiral EFT

C. Weiss (JLab), NREC 2026 Workshop, Stony Brook U., 13-17 April 2026 [[Webpage](#)]



Analytic structure of form factors

Correlations $Q^2 = 0 \longleftrightarrow Q^2$ finite

Two-pion cut and spectral functions

DChEFT: Dispersion theory \times chiral EFT

Radius extraction

ep Mainz electric + magnetic

$\mu p/ep$ MUSE \leftarrow

Further applications

Peripheral charge/current densities

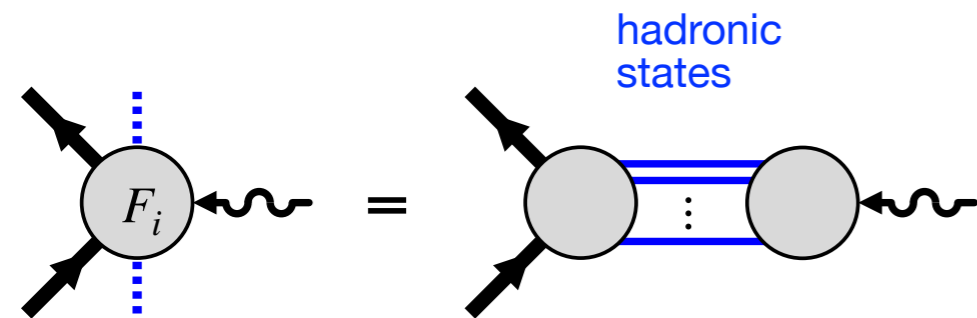
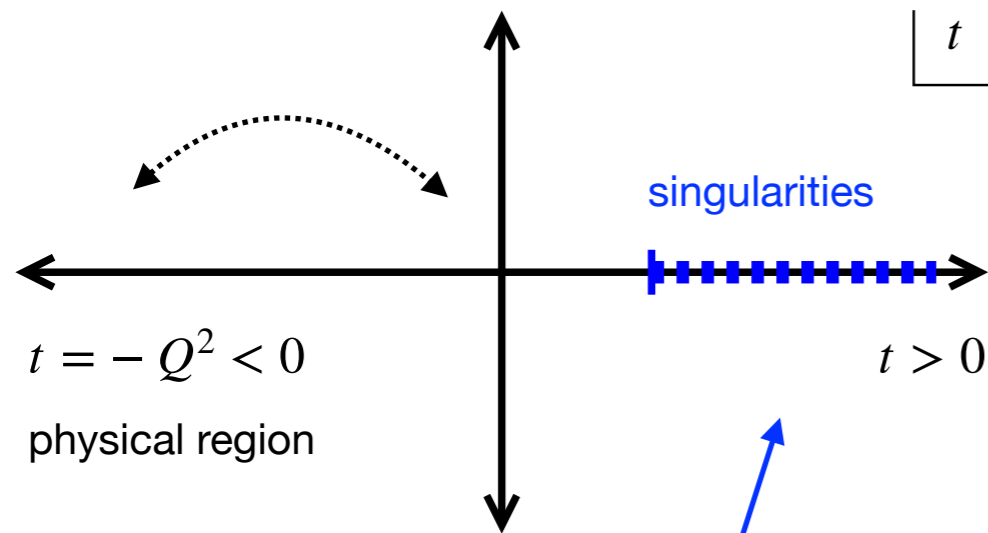
Gravitational form factors \leftarrow

F. Gil-Dominguez, J.M. Alarcon, C. Weiss PRD 108, 074026 (2023) [[INSPIRE](#)]

J.M. Alarcon, C. Weiss PRD 106, 054005 (2022) [[INSPIRE](#)]

J.M. Alarcon, D. Higinbotham, C. Weiss, PRC 102, 035203 (2020) [[INSPIRE](#)]

J.M. Alarcon, C. Weiss, PRC 96, 055206 (2017) [[INSPIRE](#)], PRC97, 055203 (2018) [[INSPIRE](#)], PLB784, 373 (2018) [[INSPIRE](#)]



FFs analytic functions of $t = -Q^2$

Singularities: Branch cuts at $t > 0$ from hadronic exchanges

Position of singularities: Hadron masses
Strength of singularities: Amplitudes \rightarrow Theory

Implications for radius extraction

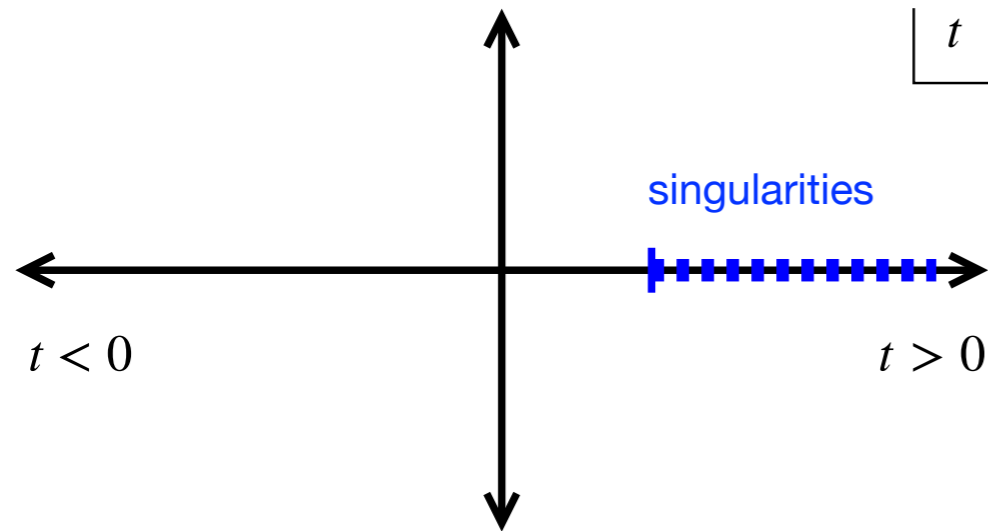
Correlates FF at $Q^2 > 0$ with derivatives at $Q^2 = 0$

Allows to use data at finite Q^2 for radius extraction, avoids “extrapolation to zero”

Necessary for magnetic radius extraction

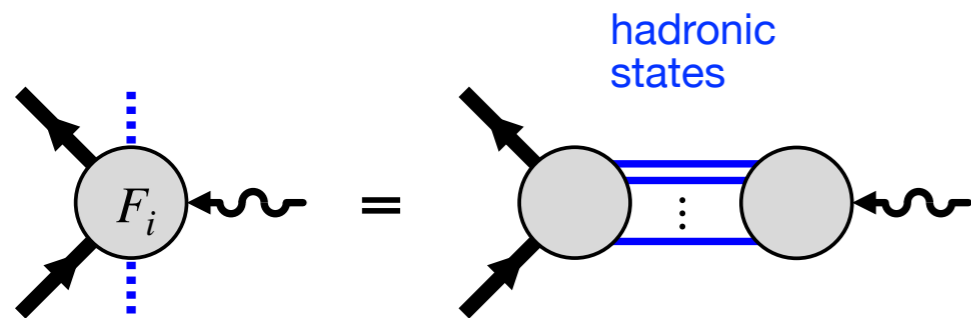
Predicts size and pattern of higher derivatives from singularities

Here: Discuss electromagnetic FFs; arguments can be extended to gravitational FFs



$$F_i(t) = \int_{t_{\text{thr}}}^{\infty} \frac{dt'}{\pi} \frac{\text{Im } F_i(t')}{t' - t - i0}$$

Expresses analytic structure of $F_i(t)$ on physical sheet



Spectral functions $\text{Im } F_i(t)$

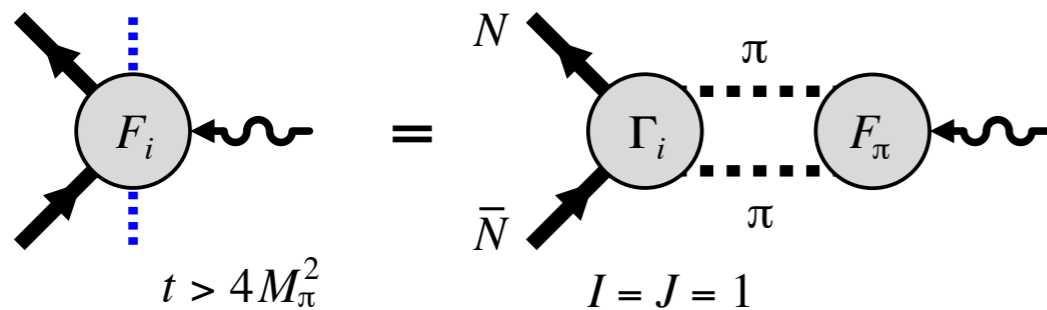
Transition amplitude
current \rightarrow hadronic states $\rightarrow N\bar{N}$

Processes in unphysical region $t < 4m_N^2$
below $N\bar{N}$ threshold

Isovector: $\pi\pi$ (incl. ρ), 4π , $K\bar{K}$, ...

Isoscalar: 3π (incl. ω), $K\bar{K}$ (incl. ϕ), ...

Needs to be calculated theoretically
Frazer, Fulco 1960; Höhler et al 1975+



Two-pion cut

Appears in isovector vector form factors

Lowest-mass state, dominates low- Q^2 spacelike form factors, peripheral densities

$\pi\pi$ system strongly interacting, ρ resonance

Spectral functions on two-pion cut

Analytic continuation of πN scattering data

Frazer, Fulco 1960; Höhler et al 1975+

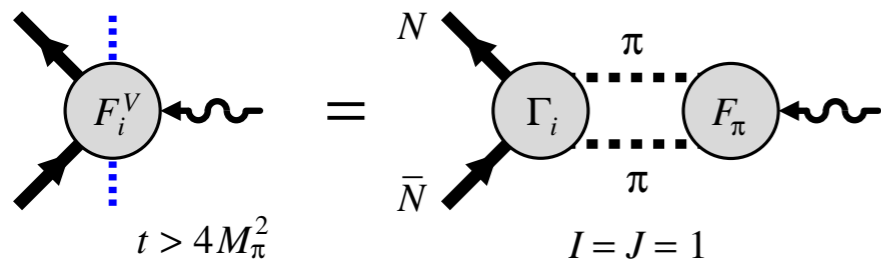
Roy-Steiner equations for πN scattering

Hoferichter, Kubis, Ruiz de Elvira, Hammer, Meissner 2016

Chiral EFT? Direct calculations poorly convergent because of strong $\pi\pi$ interactions

Gasser, Sainio, Svarc 1988; Becher, Leutwyler 1999; Kubis, Meissner 2001; Kaiser 2003; ...

Need different approach!

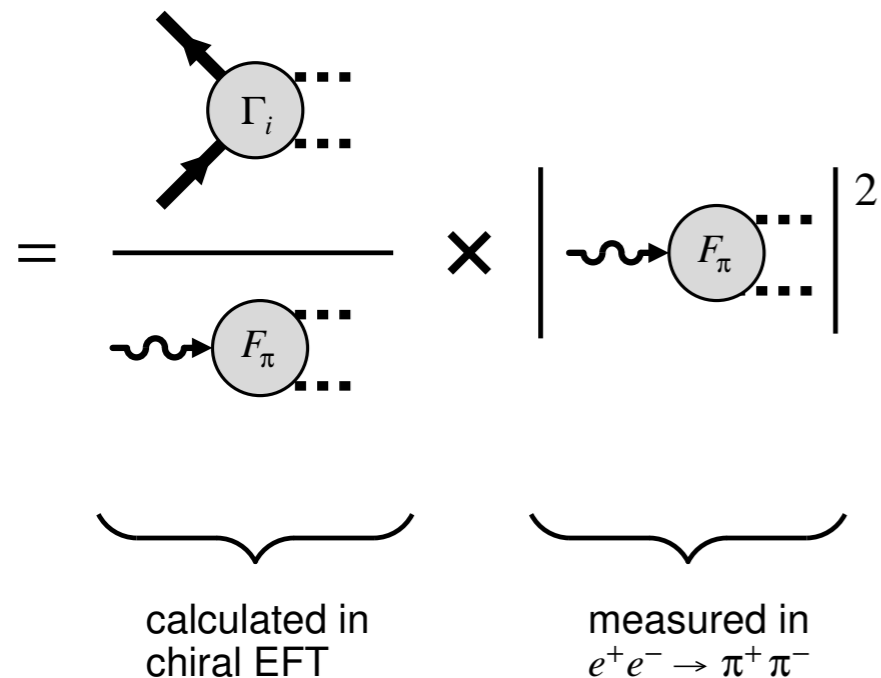


Elastic unitarity relation

$F_\pi(t)$ current $\rightarrow \pi\pi$ amplitude = pion timelike FF

$\Gamma_i(t)$ $\pi\pi \rightarrow N\bar{N}$ partial-wave amplitude

Amplitudes have same phase from $\pi\pi$ interactions:
Watson theorem



Factorize $\pi\pi$ interactions (N/D representation)

Γ_i/F_π free of $\pi\pi$ interactions

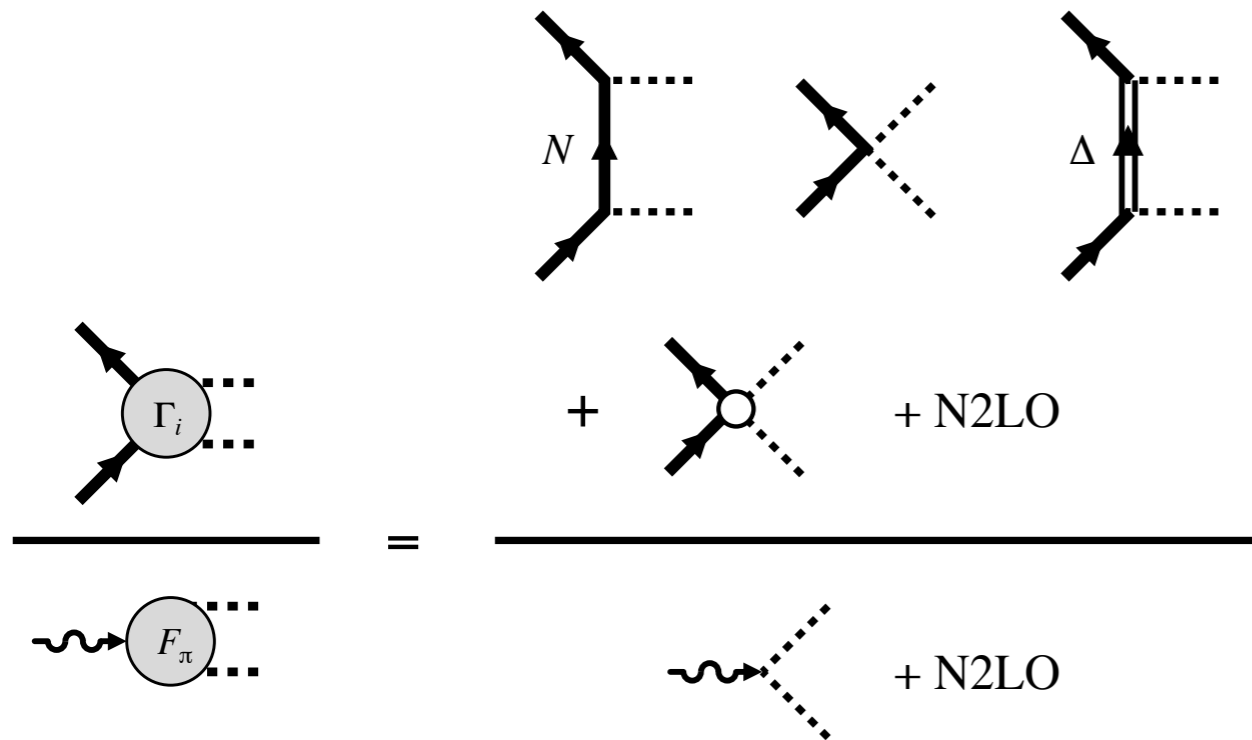
\rightarrow calculated in ChEFT with good convergence

$|F_\pi|^2$ contains $\pi\pi$ interactions

\rightarrow measured in e^+e^- annihilation

$$\begin{aligned} \text{Im } F_i(t) &= \frac{k_{\text{cm}}^3}{\sqrt{t}} \Gamma_i(t) F_\pi^*(t) \\ &= \frac{k_{\text{cm}}^3}{\sqrt{t}} \frac{\Gamma_i(t)}{F_\pi(t)} |F_\pi(t)|^2 \end{aligned}$$

Alarcon, Hiller Blin, Vicente Vacas, Weiss, NPA 96, 18 (2017)
Alarcon, Weiss, PLB 784 (2018) 373; PRC 97 (2018) 055203
Alt. formulation: Granados, Leupold, Perotti 2017



Relativistic ChEFT

Expansion in $\{M_\pi, k_\pi\}/\Lambda_{\text{chiral}}$

Include Δ isobar

ChEFT calculation of Γ_i/F_π

LO: Born terms + Weinberg-Tomozawa

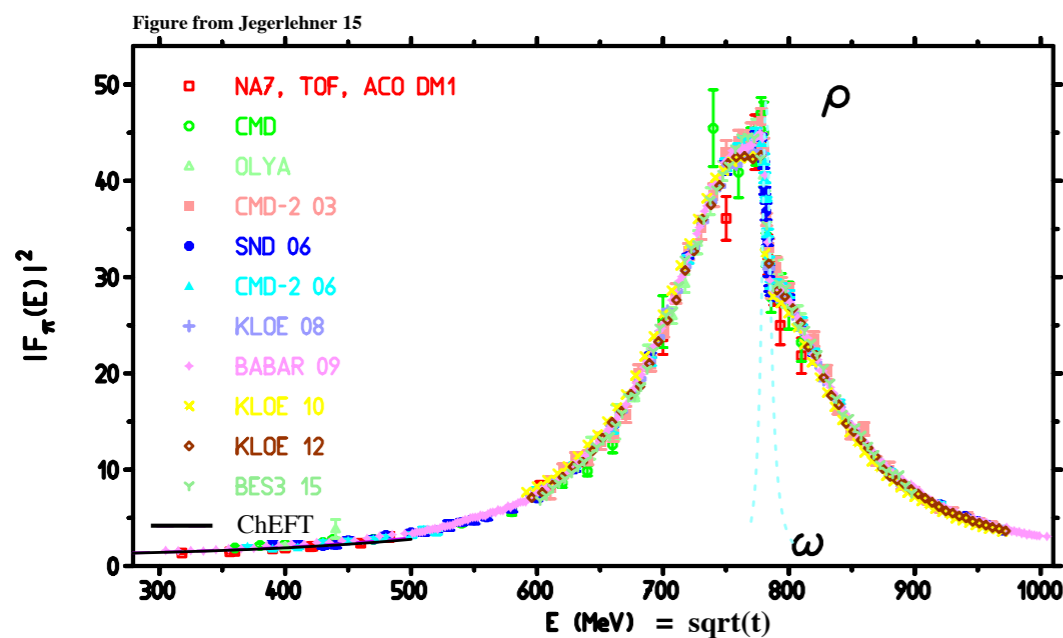
NLO: Contact term in Γ_i ($i = 2$)

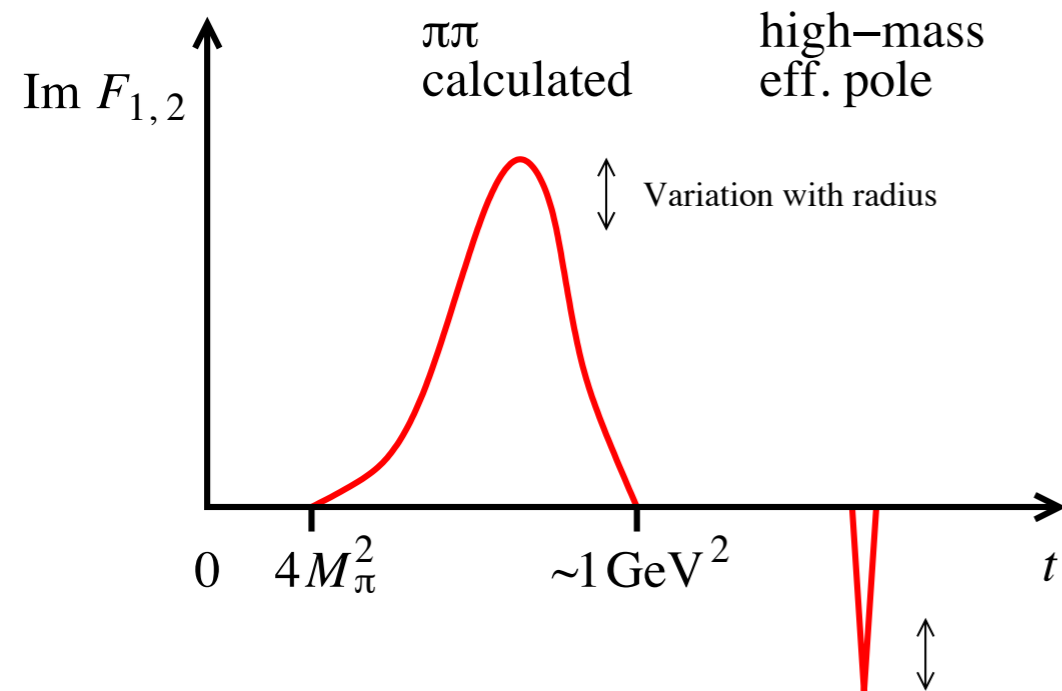
N2LO: Contact term and pion loops
Presently use partial result
Contains LEC, to be determined

Good convergence

Pion timelike form factor $|F_\pi|^2$

Measured accurately in $e^+e^- \rightarrow \pi^+\pi^-$





Spectral functions

$\pi\pi$ region calculated from unitarity + ChEFT

High-mass region parametrized by effective poles
Pole positions \rightarrow theoretical uncertainty

Sufficient for low- Q^2 form factors

Sum rules and parameters

Spectral functions constrained by sum rules
for $F(0)$, $F'(0)$ = charges, radii

Sum rules connect ChEFT LECs \leftrightarrow nucleon radii

Nucleon radii appear directly as parameters,
control finite- Q^2 behavior of form factors

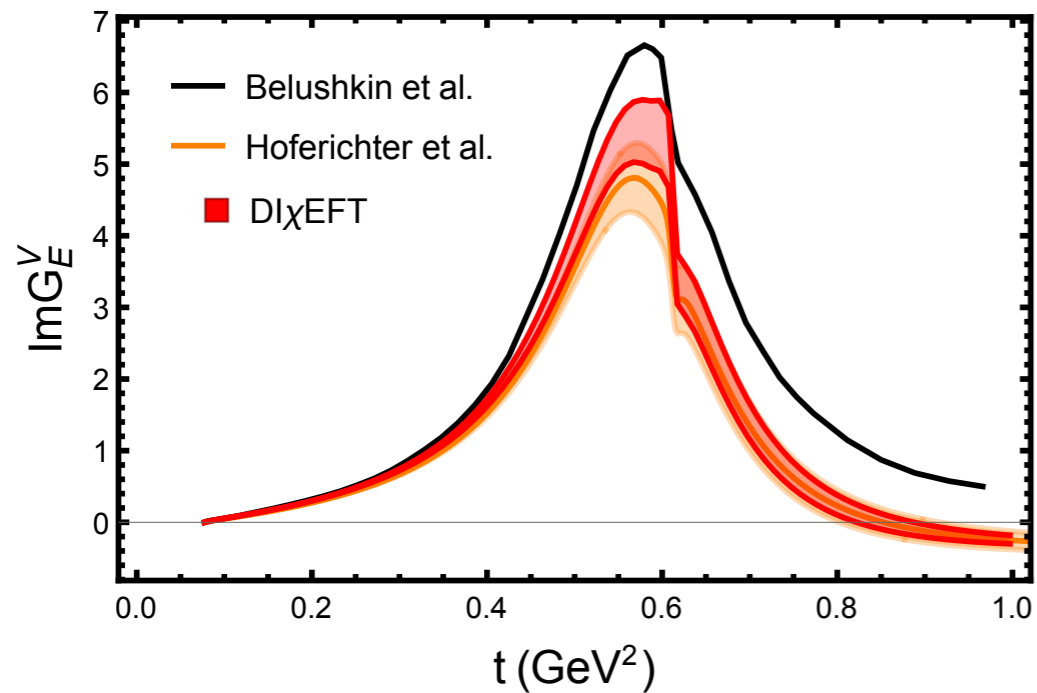
$$\frac{1}{\pi} \int_{t_{\text{thr}}}^{\infty} dt \frac{\text{Im } F_1(t)}{t} = Q$$

$$\frac{1}{\pi} \int_{t_{\text{thr}}}^{\infty} dt \frac{\text{Im } F_1(t)}{t^2} = \frac{1}{6} \langle r^2 \rangle_1$$

$$\frac{1}{\pi} \int_{t_{\text{thr}}}^{\infty} dt \frac{\text{Im } F_2(t)}{t} = \kappa$$

$$\frac{1}{\pi} \int_{t_{\text{thr}}}^{\infty} dt \frac{\text{Im } F_2(t)}{t^2} = \frac{1}{6} \kappa \langle r^2 \rangle_2$$

[+ asymptotic conditions]

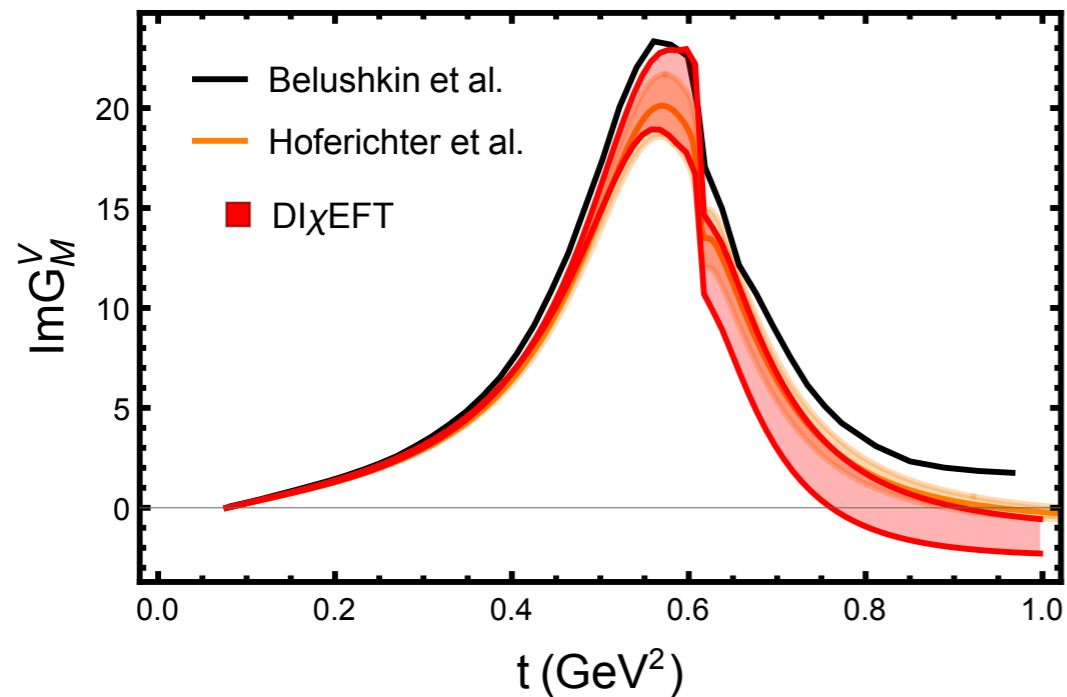


Spectral functions depend on radii as parameters

Bands show uncertainty from radii (PDG range)
Uncertainty from high-mass pole position → later

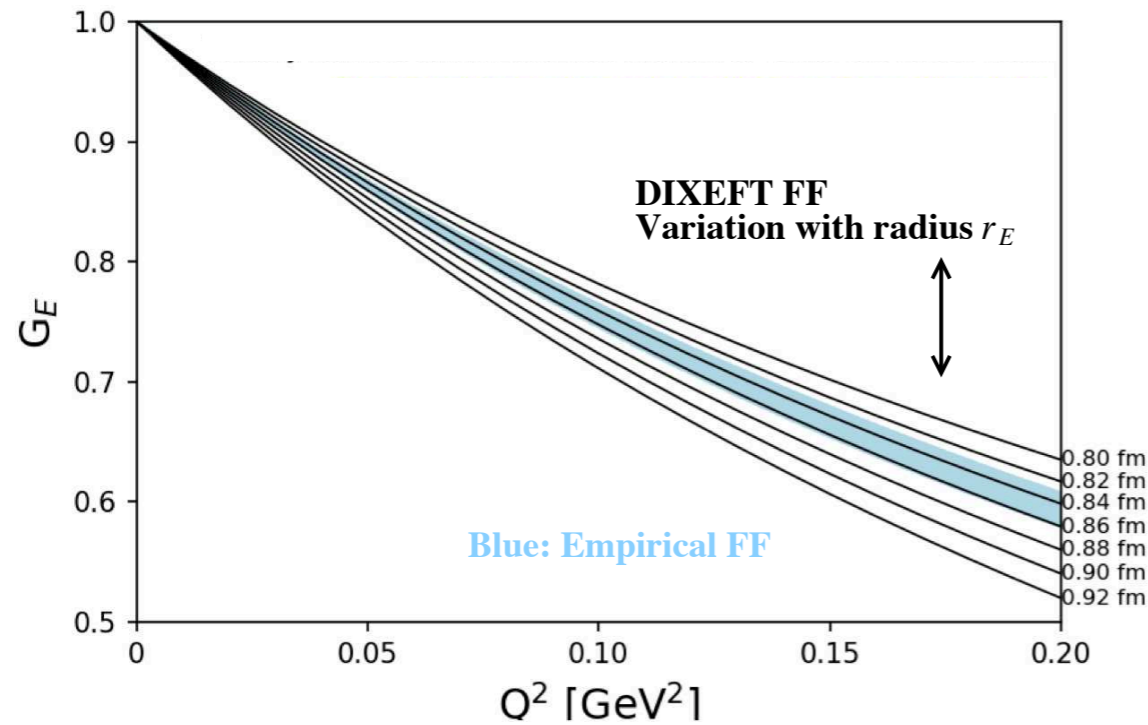
Good agreement with Roy-Steiner results

Hoferichter et al. 2017



Here: $G_E, G_M \leftrightarrow F_1, F_2$

Alarcon, Weiss, PLB784, 373 (2018) [INSPIRE]



$$G_i(t) = \int_{4M_\pi^2}^{\infty} \frac{dt'}{\pi} \frac{\text{Im } G_i(t')}{t' - t - i0}$$

Family of FF predictions depending on radii as parameters

Each member respects analyticity, sum rules

Each member has intrinsic theoretical uncertainty from high-mass states

Radius correlated with finite- Q^2 behavior!

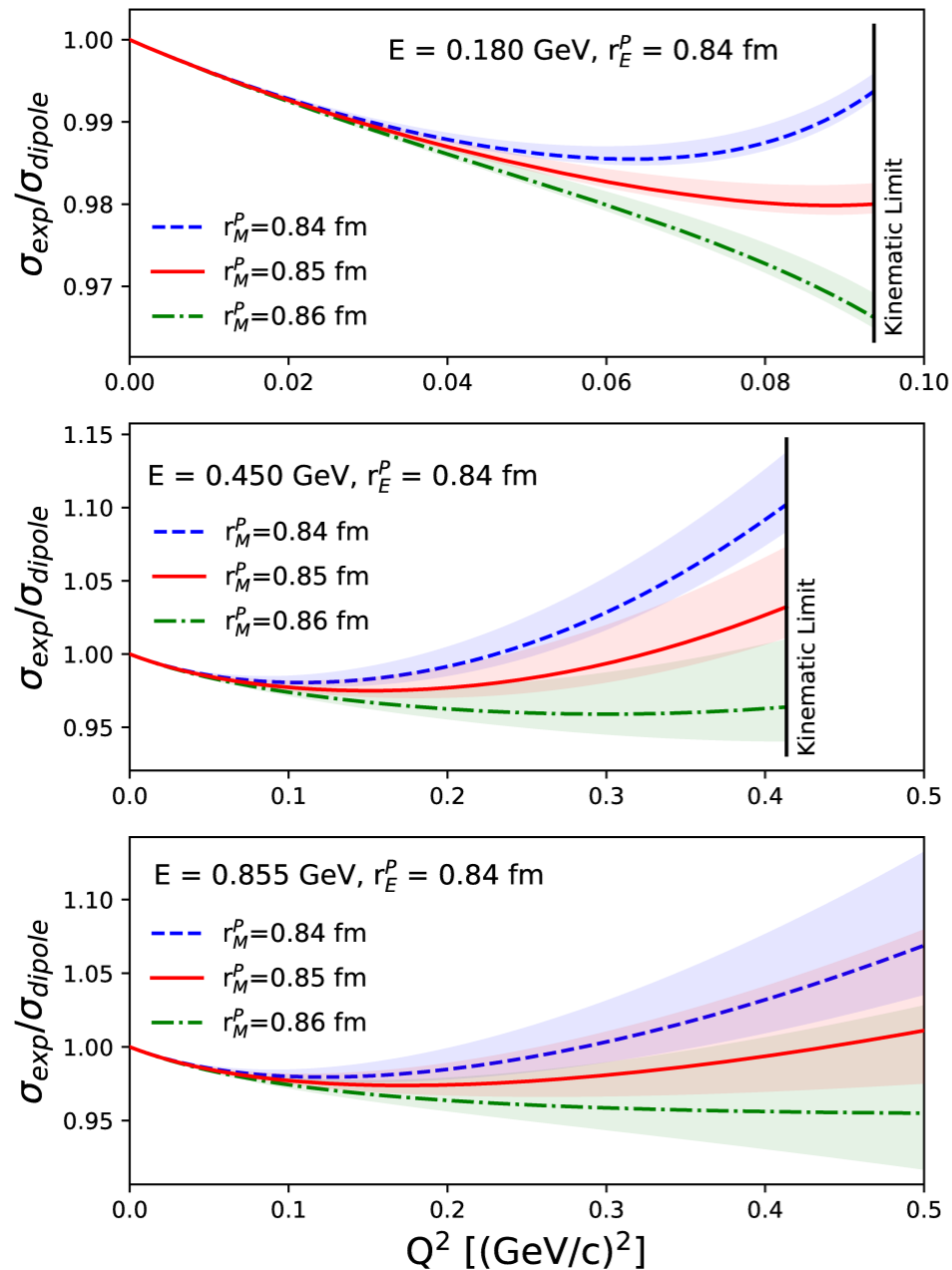
Radius extraction using DIChEFT

Compare DIChEFT FF predictions with data, for various values of radius parameter

Radius constrained by finite- Q^2 data

Optimal Q^2 range determined by interplay of radius sensitivity and exp+thy uncertainties

Alarcon, Higinbotham, Weiss, Ye PRC 99, 044303 (2019) [INSPIRE].
Fit to FF data up to $Q^2 \sim 0.5 \text{ GeV}^2$, uncertainties estimated



Extracted electric + magnetic radii from cross section fit

Used DIChEFT $G_{E,M}$ depending on r_E^p, r_M^p

Fitted cross sections with floating normalizations

Quantified fit and theoretical uncertainties

$$r_E^p = 0.842 \pm 0.002 \text{ (fit } 1\sigma) \begin{matrix} +0.005 \\ -0.002 \end{matrix} \text{ (theory full-range) fm}$$

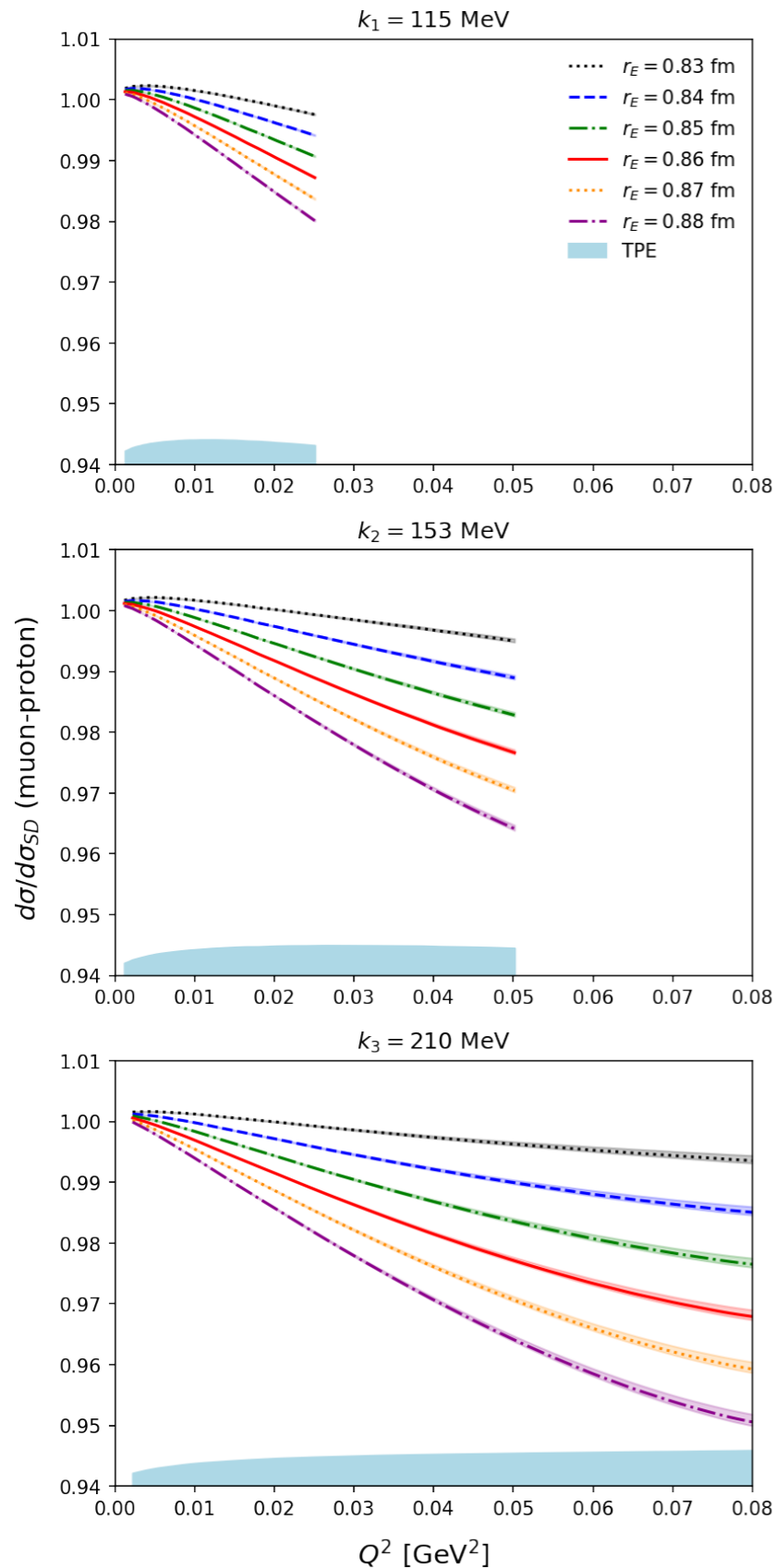
$$r_M^p = 0.850 \pm 0.001 \text{ (fit } 1\sigma) \begin{matrix} +0.009 \\ -0.004 \end{matrix} \text{ (theory full-range) fm}$$

Sensitivity to G_M only at $Q^2 > 0$, needs theory

DIChEFT enables accurate magnetic radius extraction

Conventional dispersion analysis: Lorenz, Hammer, Meissner 2012

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_{\text{Mott}} \frac{\epsilon [G_E^p]^2 + \tau [G_M^p]^2}{\epsilon(1 + \tau)}$$



First measurement of proton radius in $\mu p + ep$ scattering
 $k = 115 - 210$ MeV, $Q^2 = 0.001 - 0.08$ GeV² → [Talk Krahulik](#)

Studied radius extraction with DIChEFT

What kinematics has most impact on radius?
 What is the overall uncertainty including theory?

Tradeoff between:
 Sensitivity of cross section to radius
 Theoretical uncertainty
 Two-photon exchange corrections

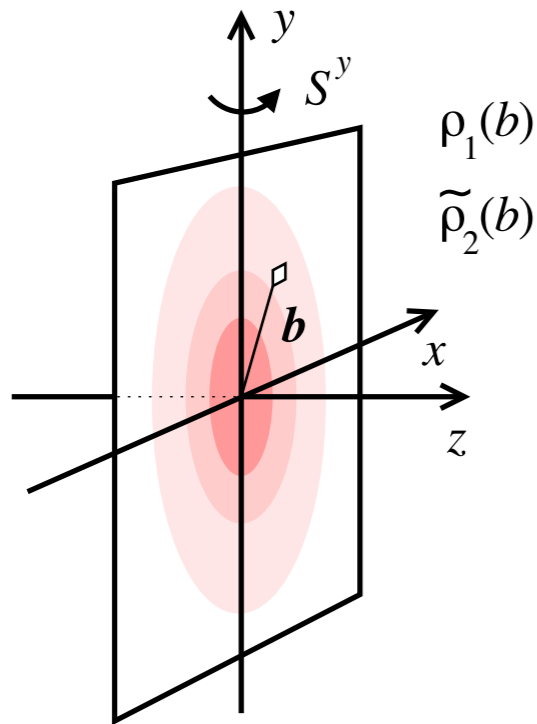
Findings

Influence of TPE on radius extraction diminished at higher Q^2

Optimal kinematics for radius extraction
 $k = 210$ GeV, $Q^2 = 0.05 - 0.08$ GeV²

F. Gil-Dominguez, J.M. Alarcon, C. Weiss PRD 108 074026 (2023) [[INSPIRE](#)]

Two-photon exchange: Tomalak, Vanderhaeghen 2016



$$F_{1,2}(t = -\Delta_T^2) = \int d^2b e^{i\Delta_T \cdot b} \rho_{1,2}(b)$$

Charge/magnetization densities at light-front time x^+ ,
Frame-independent, appropriate for relativistic systems

Soper 1976, Burkardt 2000, Miller 2007

Fourier transform of form factor data

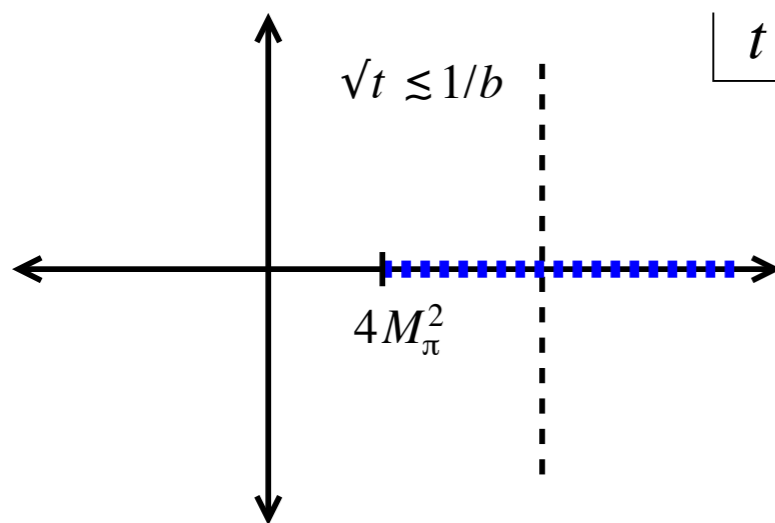
Miller 2007; Carlson Vanderhaeghen 2008; Venkat, Arrington, Miller, Zhan 2010

Dispersive representation

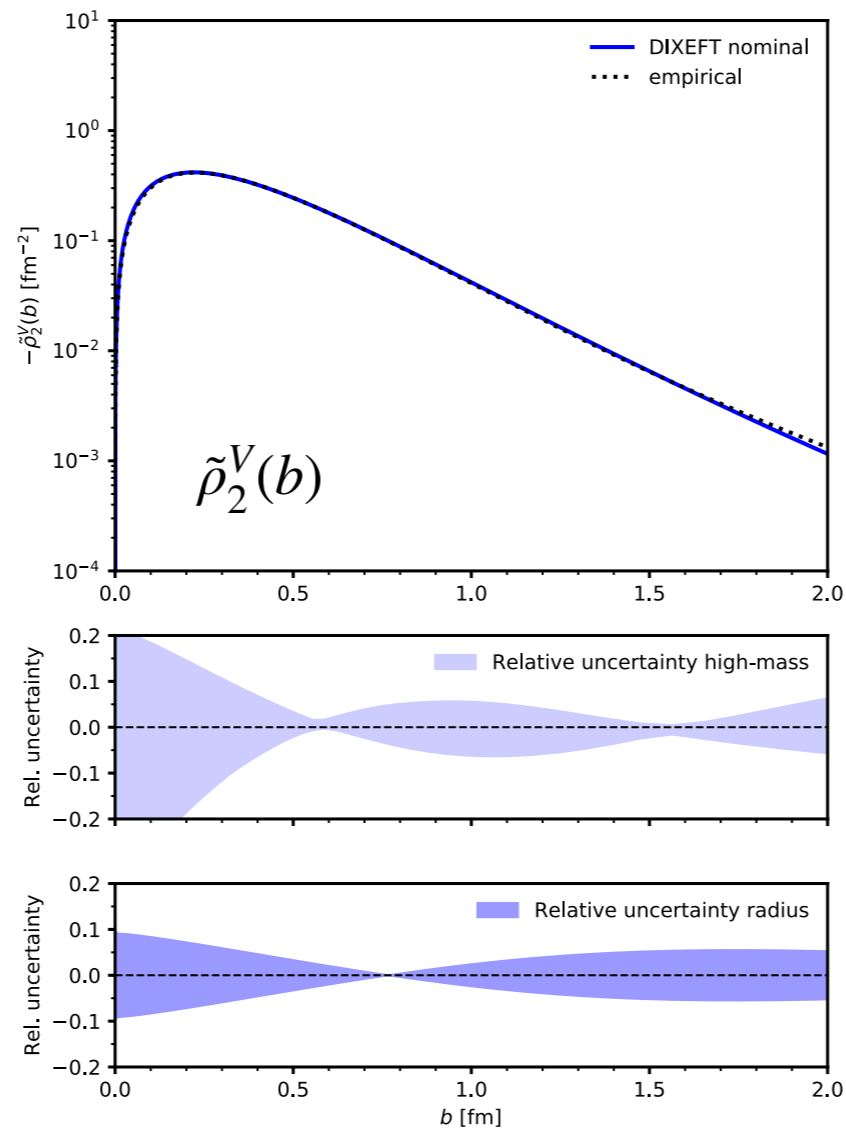
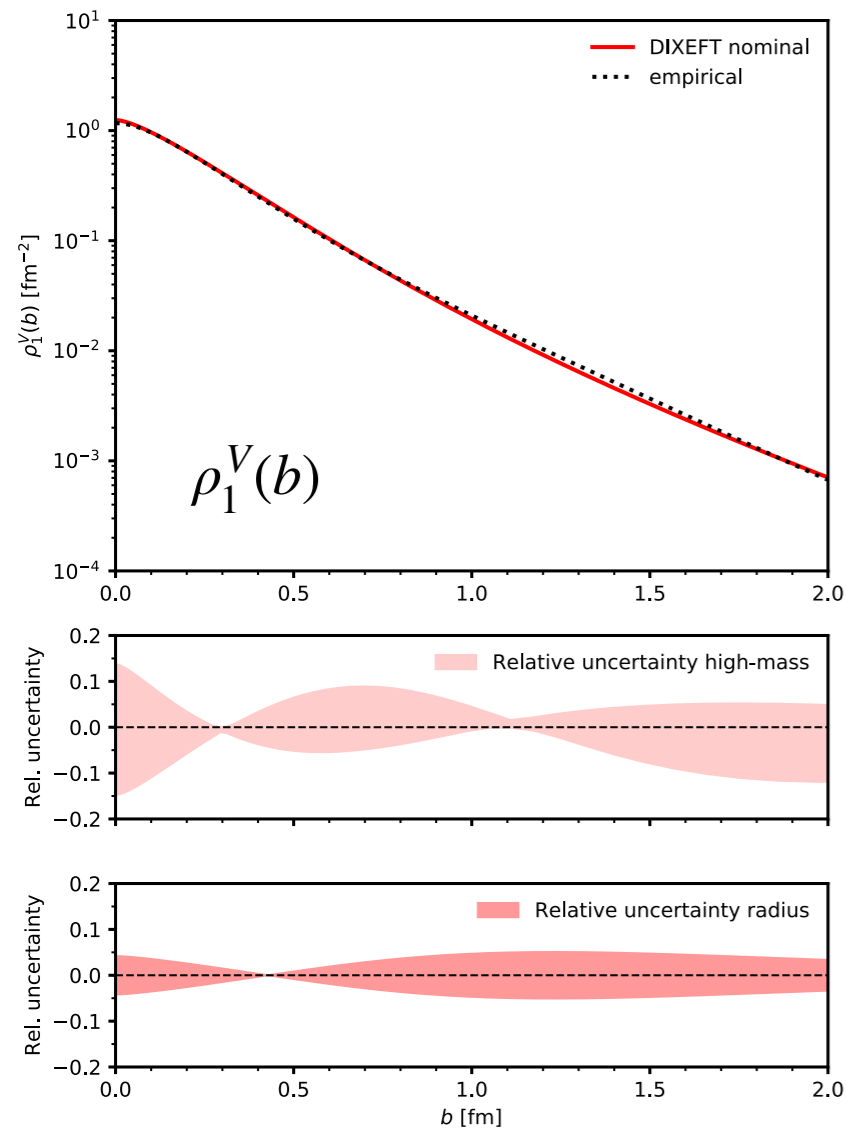
$$\rho(b) = \int_{4M_\pi^2}^{\infty} \frac{dt}{2\pi^2} K_0(\sqrt{t}b) \text{Im} F(t) \quad K_0 \sim e^{-b\sqrt{t}}$$

Exponentially convergent, acts as filter $\sqrt{t} \lesssim 1/b$
Large distances $b \leftrightarrow$ low masses \sqrt{t}

Peripheral densities }
Uncertainty quantification } with analyticity



Strikman, Weiss PRC 82, 042201 (2010) [INSPIRE]; Miller, Strikman, Weiss PRC 84, 045205 (2011) [INSPIRE]; Granados, Weiss JHEP 01, 092 (2014) [INSPIRE]



Transverse densities and uncertainties computed from DICH EFT spectral functions

Alarcon, Weiss PRD 106, 054005 (2022) [INSPIRE]

Large- b asymptotics governed by analyticity

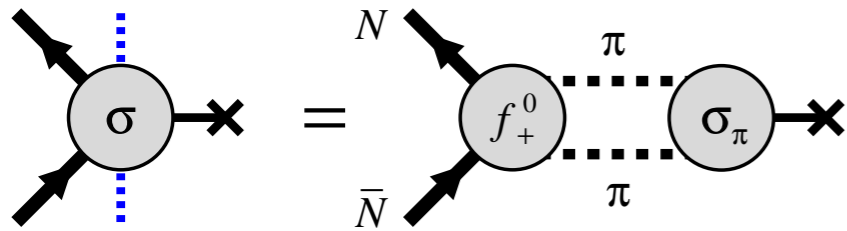
Difficult to obtain from Fourier transform. Requires proper analyticity of form factor!

Densities predicted with relative uncertainties $\lesssim 10\%$ at $b > 0.3$ fm

Excellent agreement with empirical densities

Uncertainty quantification using analyticity: Spectral functions \rightarrow densities

\rightarrow Supplement



Nucleon FFs of QCD energy-momentum tensor

Contain multipoles $J = 2, 1, 0$

Describe distributions of momentum, spin, forces, mass: Much interest → [Talk Schweitzer](#)

Can be analyzed using DIChEFT

Nucleon scalar form factor ($J = 0$)

$$\sigma(t) = \langle N' | m\bar{\psi}\psi | N \rangle \quad \text{sigma term}$$

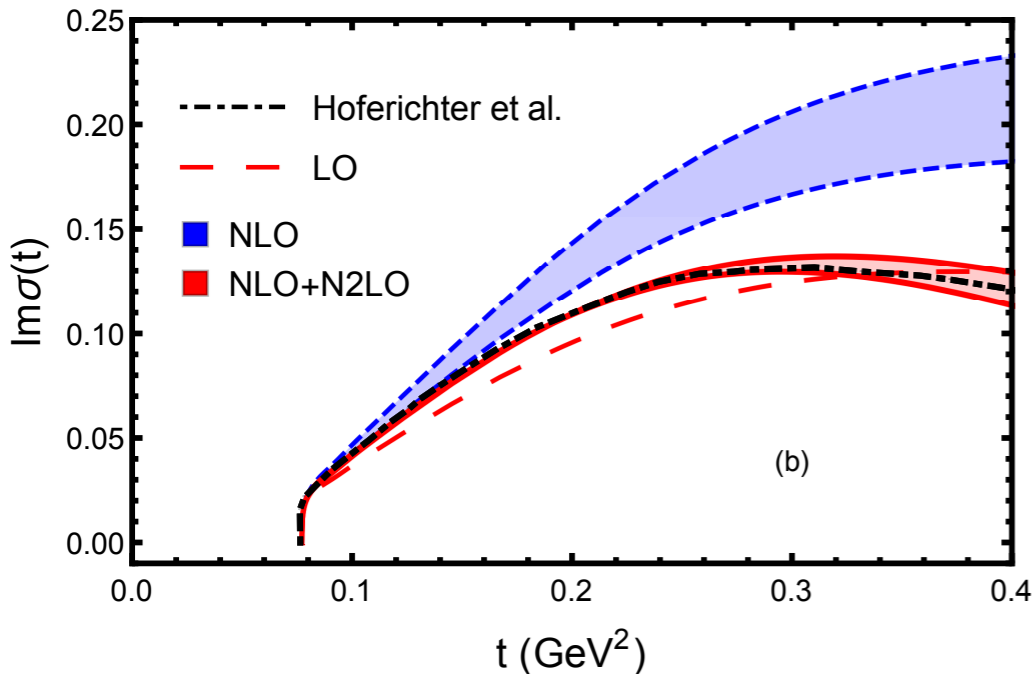
$\sigma_\pi(t)$ from empirical dispersion analysis
Colangelo et al 04; Celis, Cirigliano, Passemar 14

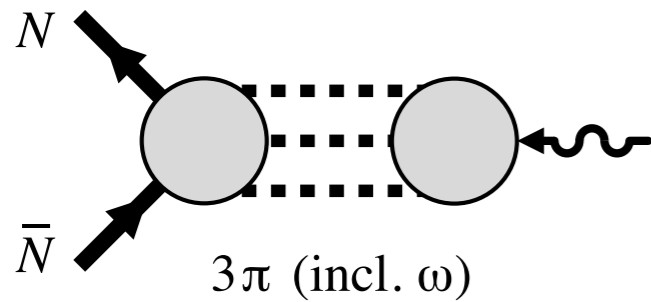
Used subtracted dispersion relation

$$\langle r^2 \rangle_\sigma = 1.03 - 1.13 \text{ fm}^2 \quad \text{nucleon scalar radius predicted}$$

Gravitational form factors ($J = 2$)

Possible with pion timelike FF as input: Constrained by LQCD, instanton vacuum
 χ QCD Collab B. Wang et al. 2024; Hackett et al. 2023. Liu, Shuryak, Weiss, Zahed, PRD 110, 054021 (2024)

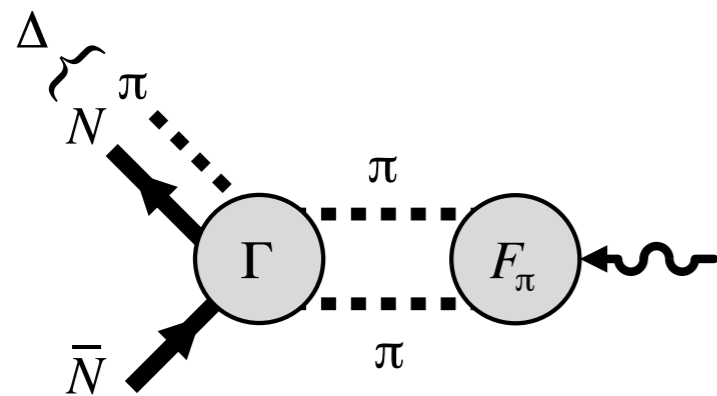




Nucleon FFs with 3π cut

Isoscalar vector current, isovector axial current

Use methods of 3-body unitarity
 Szczepaniak, Jackura, Piloni, Doering et al.



$N \rightarrow \Delta$ transition form factors

Compute transition matrix element $\langle N\pi | J^\mu | N \rangle$
 Continue to pole in $s_{\pi N} = m_\Delta^2$, extract residue

ChEFT calculations
 Ledwig et al. 2010

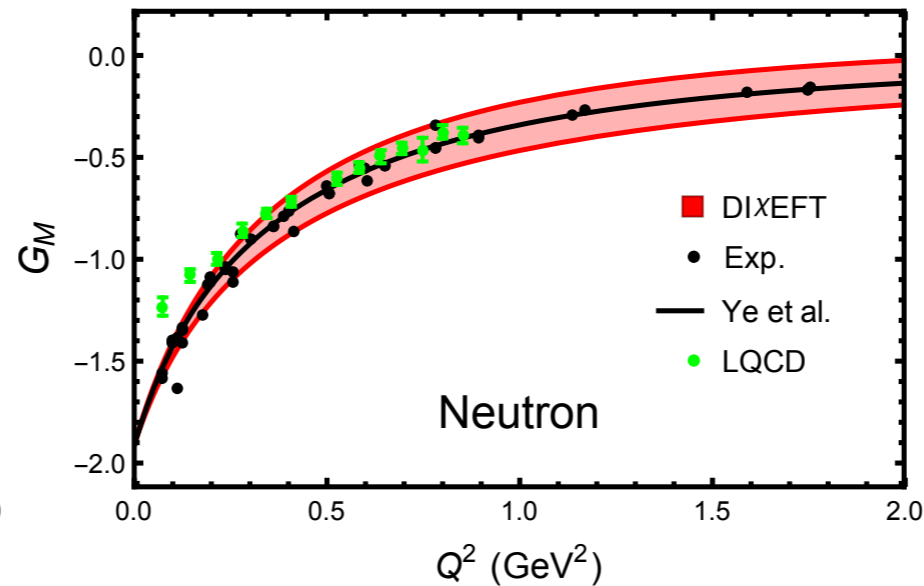
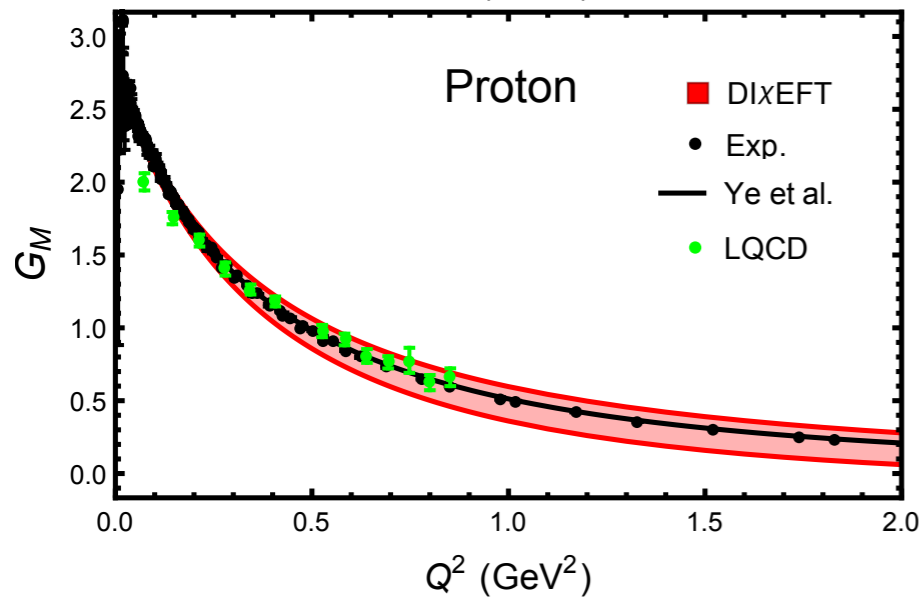
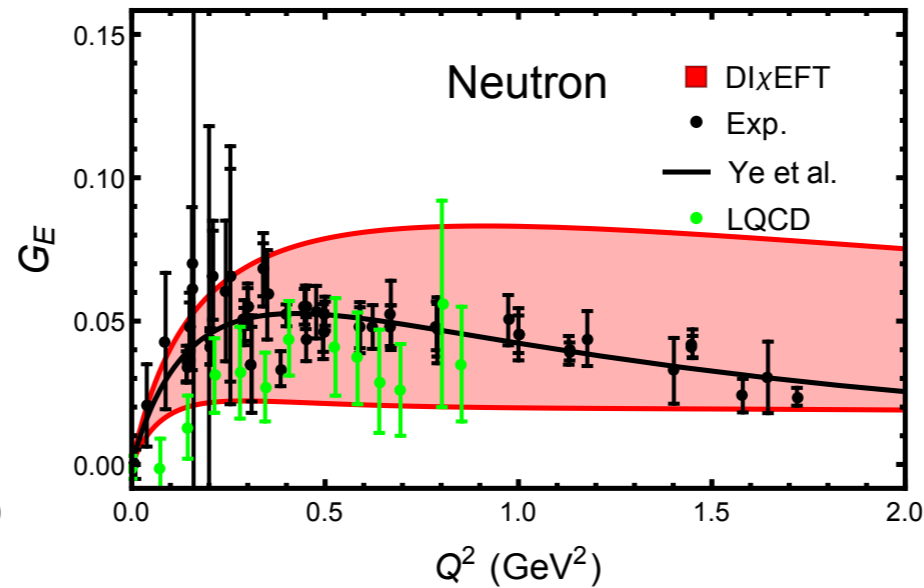
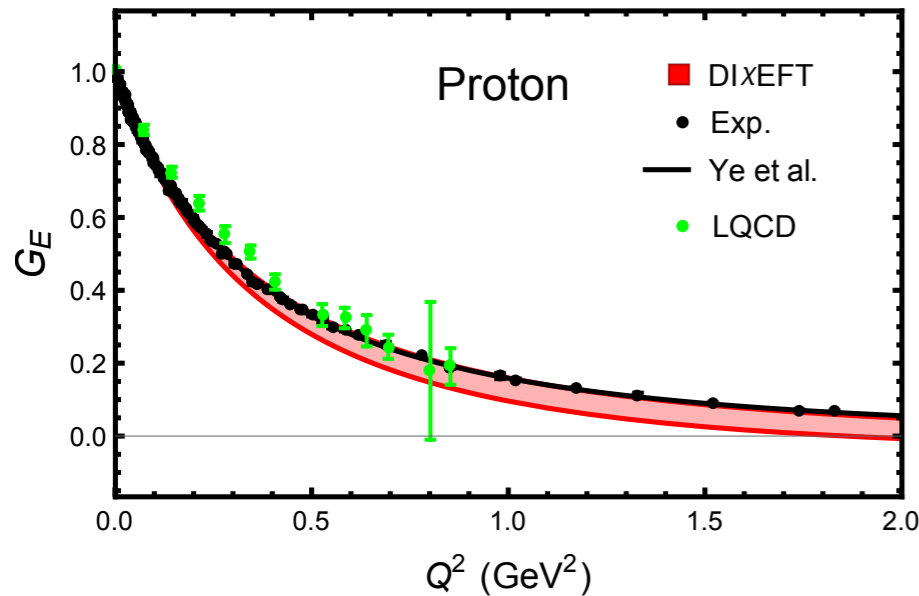
LQCD results
 Alexandrou et al. 08; Aubin, Orginos, Pascalutsa, Vanderhaeghen 08

Large- N_c spin-flavor symmetry connects $N \rightarrow N$ and $N \rightarrow \Delta$

→ [Talk Paolone](#)

- Analyticity correlates FF at $Q^2 = 0$ and finite Q^2 , plays essential role in radius extraction
- DIChEFT: Combines dispersion theory (analyticity, unitarity) with ChEFT (long-range dynamics), permits first-principles calculations of $\pi\pi$ spectral functions and low- Q^2 form factors
- DIChEFT-based radius extraction implements analyticity and resulting "information flow"
- Highest impact on radius from finite Q^2 data, no need for "extrapolation to zero". Assessments depend on actual exp + thy uncertainties, can be updated
- DIChEFT determines peripheral densities with quantified uncertainties, analyticity essential
- DIChEFT can be extended to gravitational form factors: Pion input needed

Supplemental material



$$G_i(t) = \int_{4M_\pi^2}^{\infty} \frac{dt'}{\pi} \frac{\text{Im } G_i(t')}{t' - t - i0}$$

Isovector: DIChEFT
 Isoscalar: ω +high-mass

Alarcon, Weiss,
 PLB784, 373 (2018)

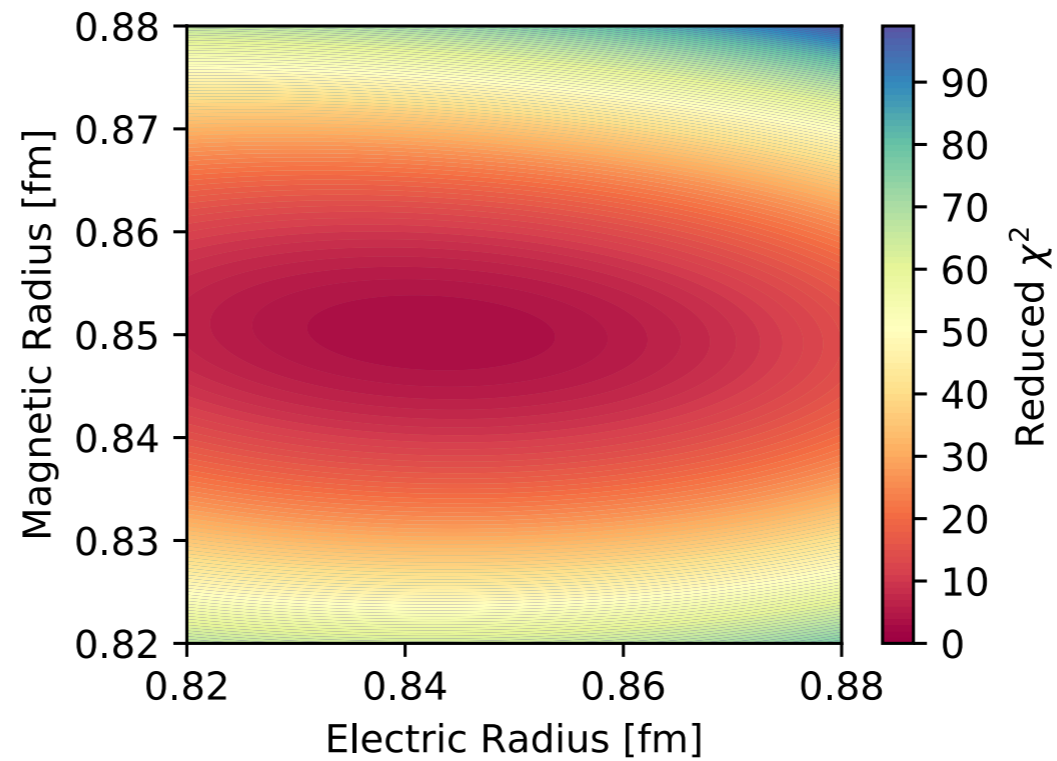
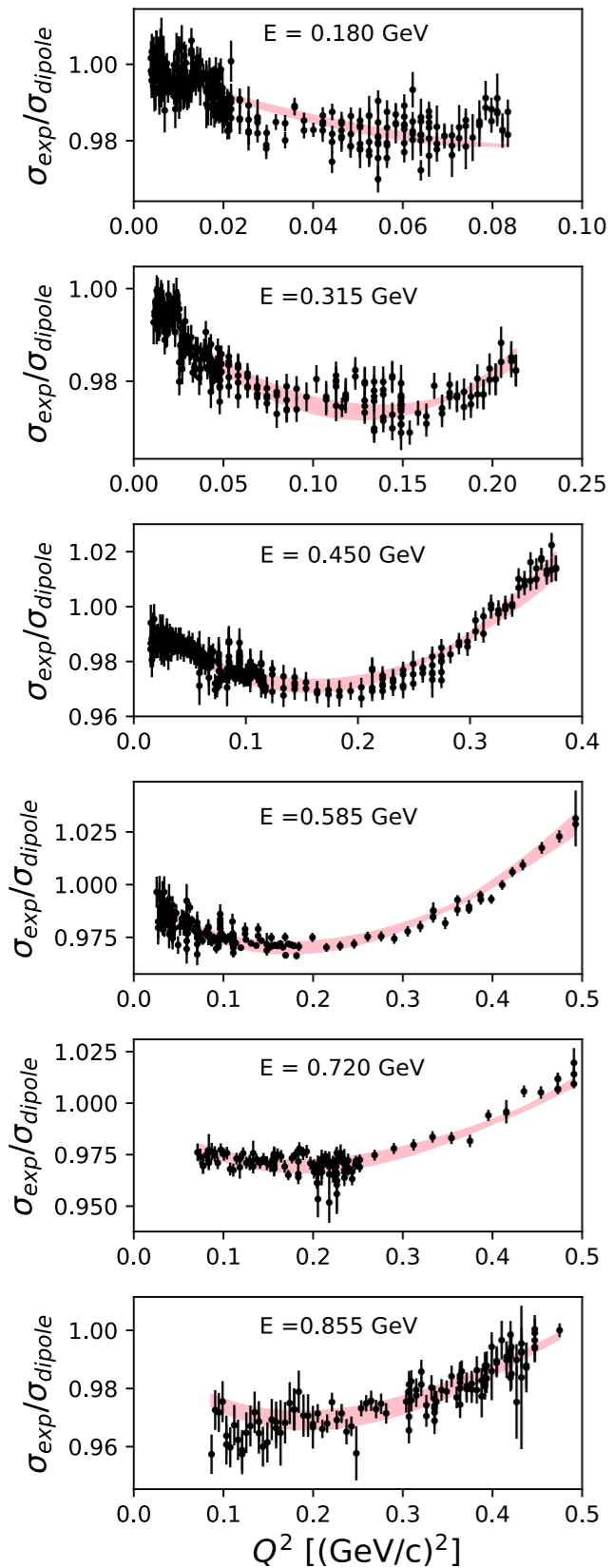
LQCD: Alexandrou et al. 2017

Form factors

Dispersion integral evaluated with spectral functions (including $\pi\pi$ and high-mass part)

Band shows uncertainty from radii (uncertainty from high-mass pole position \rightarrow later)

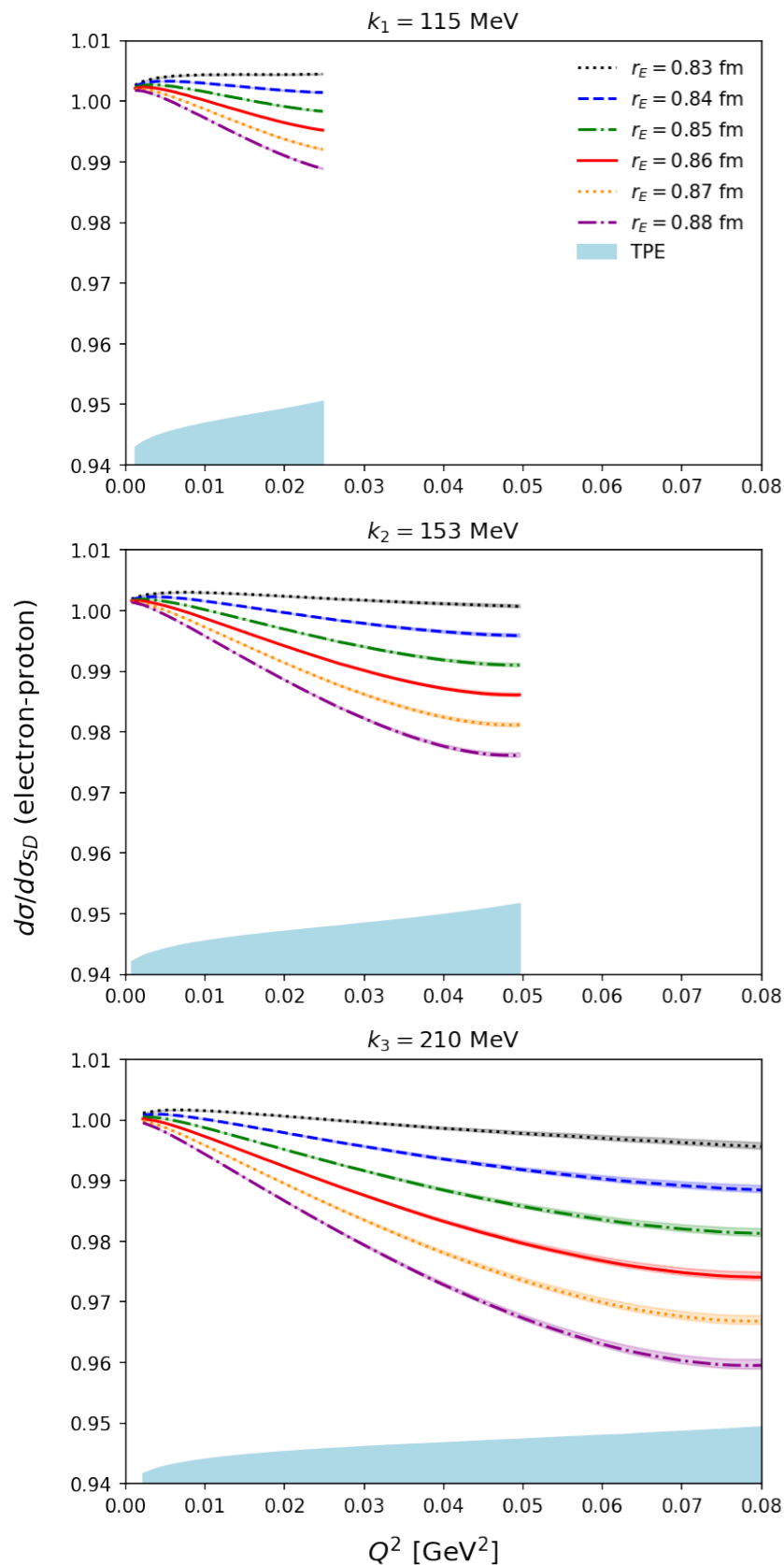
Excellent agreement with data. No fit, but prediction based on dynamics



χ^2 profile in electric and magnetic radius

Mainz A1 data and $\text{D}\chi\text{EFT}$ fit
 Bands: Fit uncertainty

Alarcon, Higinbotham, Weiss, PRC 102, 035203 (2020) [\[INSPIRE\]](#)



Here: MUSE ep scattering (for μp see above)

What kinematics has most impact on radius?
What is the overall uncertainty including theory?

Tradeoff between:
Sensitivity of cross section to radius
Theoretical uncertainty
Two-photon exchange corrections

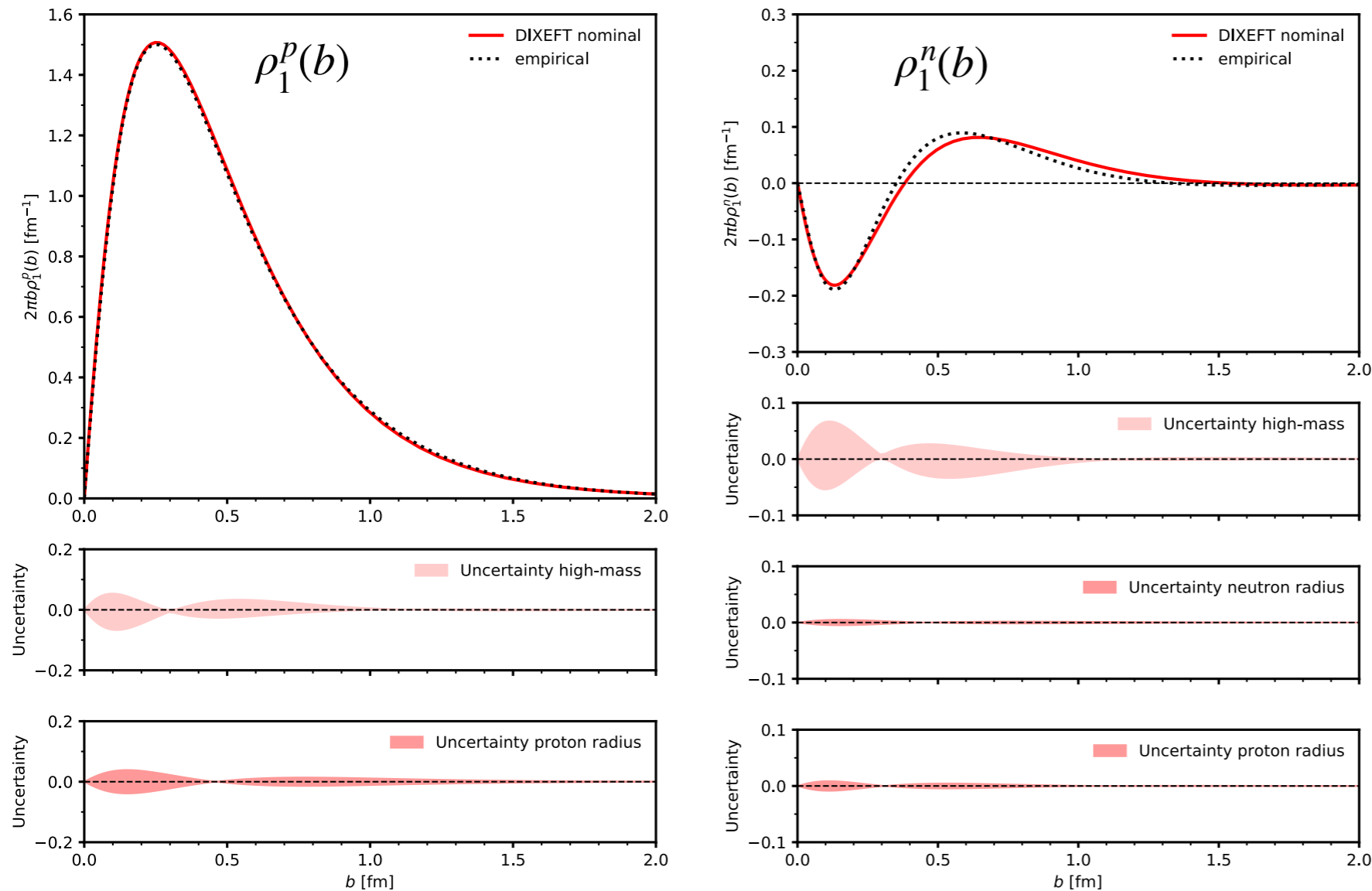
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F. Gil-Dominguez, J.M. Alarcon, C. Weiss PRD 108 074026 (2023) [[INSPIRE](#)]

Two-photon exchange: Tomalak, Vanderhaeghen 2016

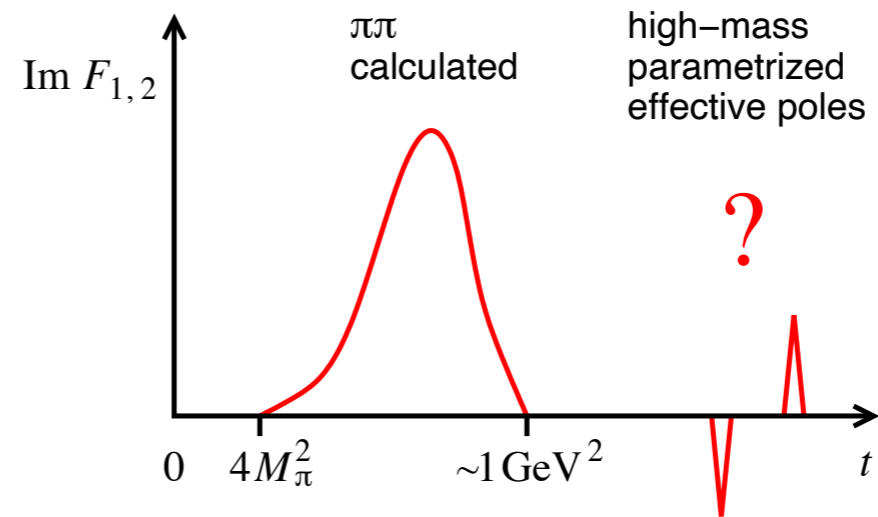
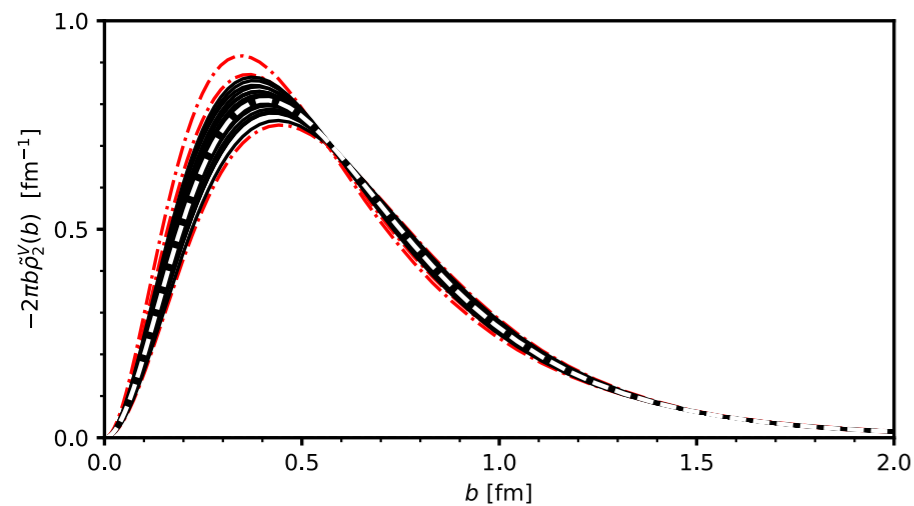
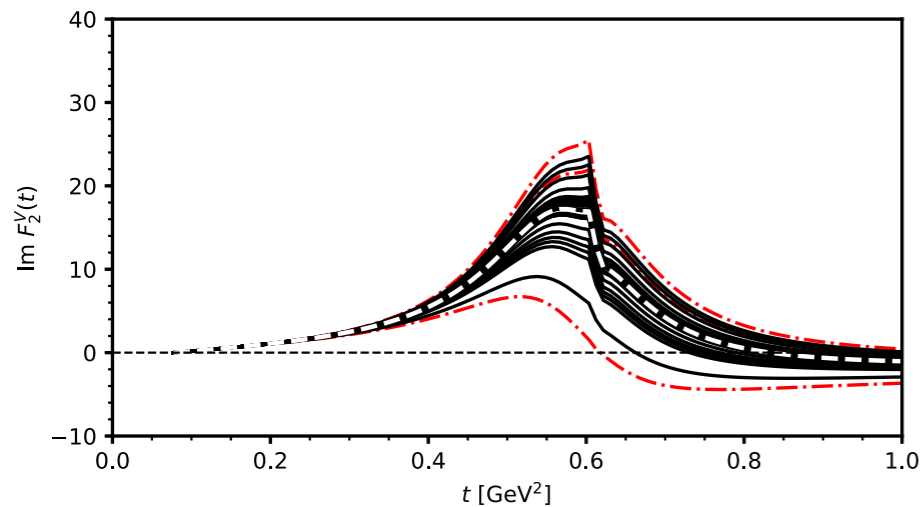
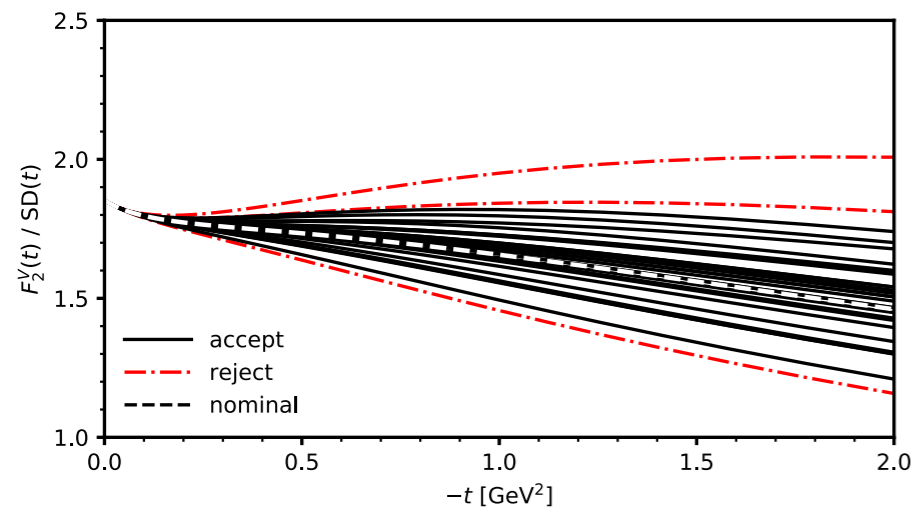


Alarcon, Weiss PRD
106, 054005 (2022)

Radial charge densities $2\pi b \rho_1^{p,n}$ of proton and neutron

Obtained realistic nucleon densities with controlled uncertainties

Reproduced positive charge density in neutron at intermediate $b \sim 0.5 - 1$ fm
Miller 2007



Shape of high-mass spectral function unknown
 → treat as theoretical uncertainty

Parametrize through effective poles
 $\text{Im } F_1[\text{high-mass}] = a_0 \delta(t - t_0) + a_1 \delta'(t - t_1)$

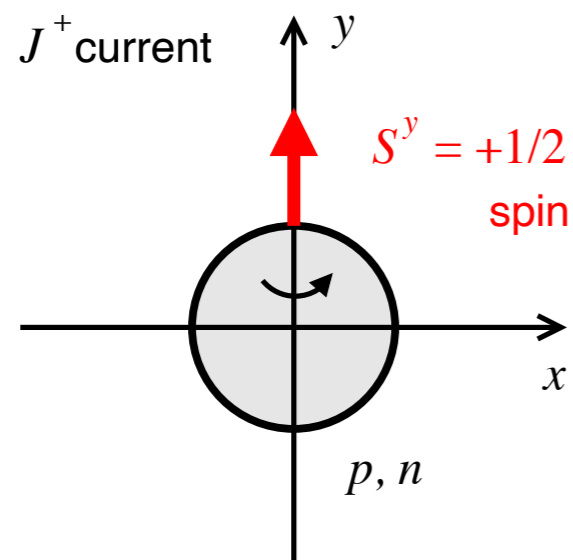
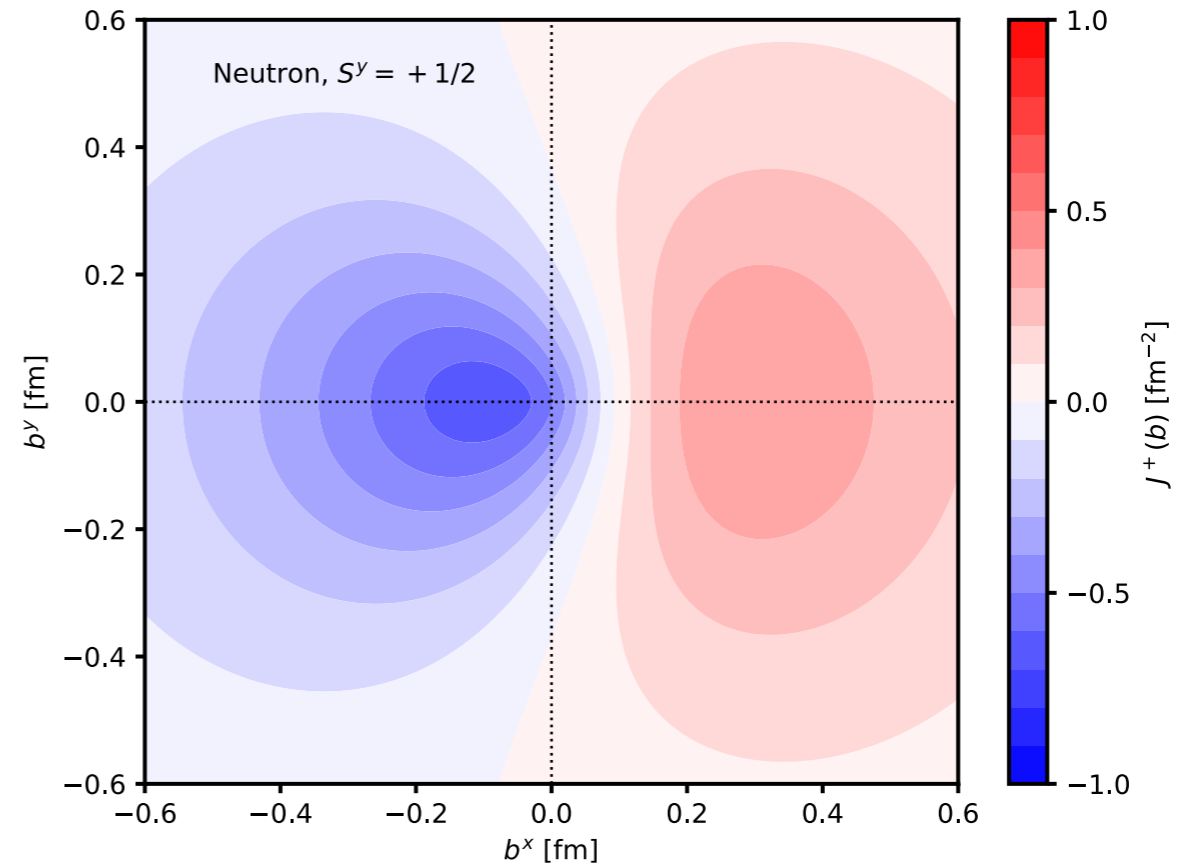
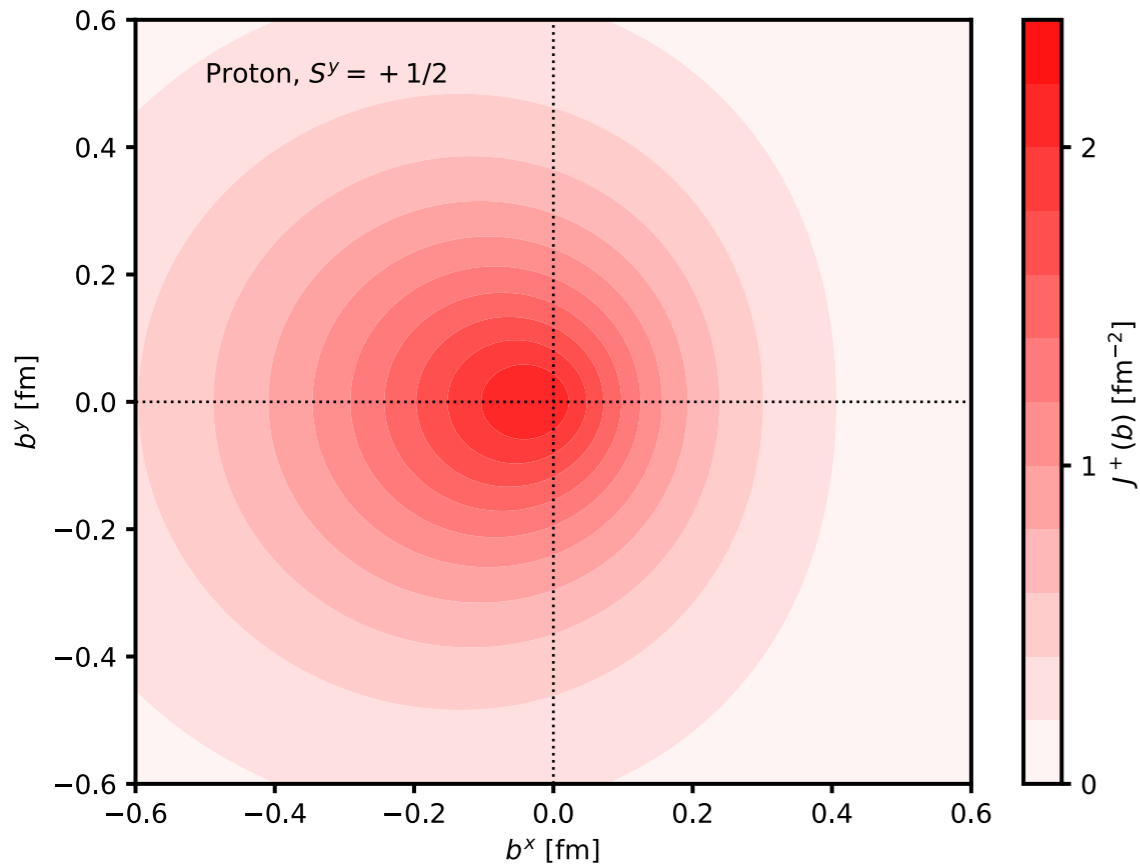
Pole positions considered unknown (in reasonable range)
 Pole coefficients fixed by sum rules

Generate MC ensemble of spectral functions
 Propagate variation into form factors, densities, etc.

Uncertainty quantification consistent with analyticity!

Peripheral densities not sensitive to unknown
 high-mass spectral function - robust predictions!

Transverse densities: J^+ current in polarized nucleon 23



Plus current density in transversely polarized nucleon localized at $x = y = 0$

$$\langle J^+(\mathbf{b}) \rangle = \rho_1(b) + (2S^y) \cos \phi \tilde{\rho}_2(b)$$

This is how an electromagnetic probe coupling to J^+ “sees” the nucleon in transverse space

Computed from DIChEFT results