

JIMWLK on a quantum computer

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Based on: A. A. Agrawal, E. Budd, A. F. Kemper, A. Tarasov, V. V. Skokov, S.S.T : 2603.02516v1



JIMWLK: Jalilian-Marian, Iancu, McLerran, Weigert, Leonidov and Kovner

- JIMWLK evolution equation: small-x observables
- Important to understand high energy observables, gluon saturation, proton spin puzzle etc
- Current method of solving: Map to Langevin equation

*J.P. Blaizot, E. Iancu, H. Weigert
K. Rummukainen and H. Weigert*

JIMWLK: Jalilian-Marian, Iancu, McLerran, Weigert, Leonidov and Kovner

- JIMWLK evolution equation: small- x observables
- Important to understand gluon saturation, high energy collisions
- Current method: Map to Langevin equation

Problems

- Computationally expensive for better statistics
- Higher order JIMWLK lacks a Langevin formulation
- Not all observables can be evolved

Need a new method to solve JIMWLK!

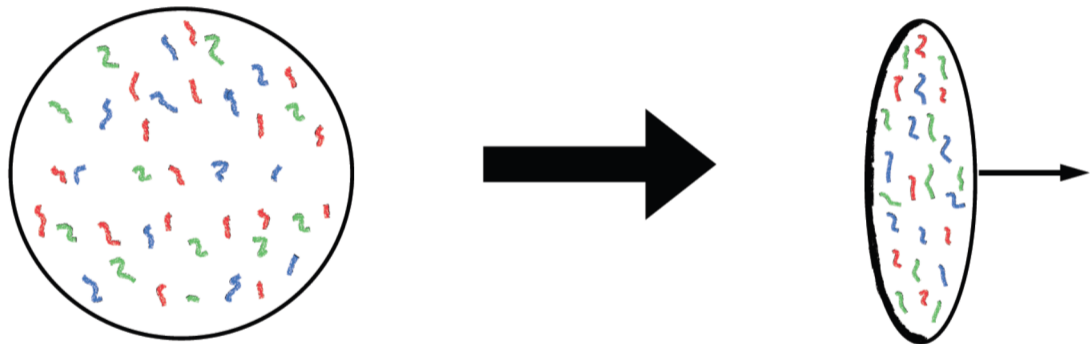
We propose a new algorithm to solve JIMWLK on quantum computers

- Evolve lindblad equation in a lattice gauge theory motivated basis
- No statistical noise
- Compute the evolution of a more general class of “off-diagonal” observables

A first study: Toy model of JIMWLK for $SU(2)$

Lin, Lin (2024)
N.Kleo, J.Stryker, M.Savage (2019)

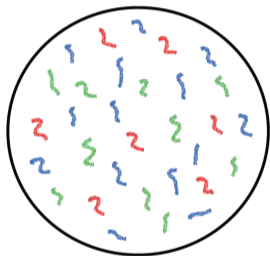
JIMWLK setup



Nucleus flattens due to lorentz contraction

Weight functional

QCD degrees of freedom separated into valence and soft based on rapidity



$W[j]$: Distribution of valence color charges

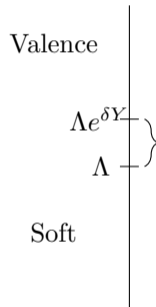
Valence
($k^+ > \Lambda$)

Λ

Soft
($k^+ < \Lambda$)

Functional fokker-planck JIMWLK: Evolution of the weight functional $W_Y[j]$

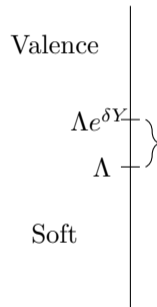
$$\partial_Y W_Y = H_{JIMWLK} W_Y$$



Scattering cross-sections averaged over different configurations of $W[j]$

Functional fokker-planck JIMWLK: Evolution of the weight functional $W_Y[j]$

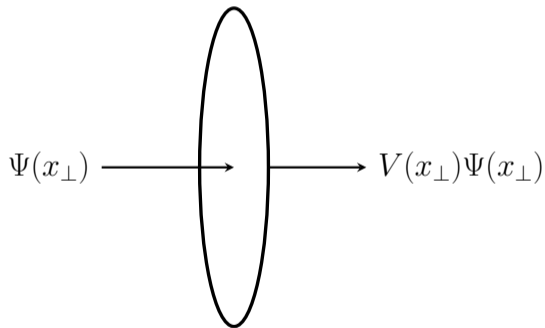
$$\partial_Y W_Y = H_{JIMWLK} W_Y$$



Too complex: rewrite as a langevin equation

Work with infinite wilson lines

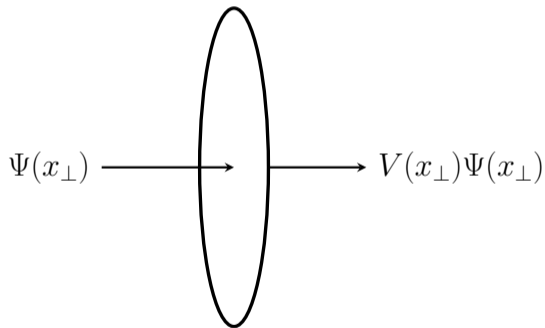
$$V_j(x_\perp) = \mathcal{P} \exp \left[i \int_{-\infty}^{\infty} dx^- \tau_j^a \alpha^a(x^-, x_\perp) \right]$$



Eikonal scattering of a color charge on a nucleus at high energy: Wilson lines

Work with infinite Wilson lines

$$V_j(x_\perp) = \mathcal{P} \exp \left[i \int_{-\infty}^{\infty} dx^- \tau_j^a \alpha^a(x^-, x_\perp) \right]$$



Write JIMWLK as a Langevin equation on Wilson lines

The procedure summarized: Evolving the wilson lines

- Generate an initial set of wilson lines from $W_Y[j]$
- Evolve the Wilson lines $V_{\alpha_n}^Y(x_\perp)$ for multiple realizations of gaussian noise ξ
- Compute expectation values by using Wilson lines at rapidity Y and average over noise and initial conditions

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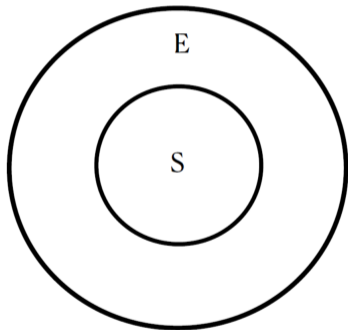
Problems with Langevin-JIMWLK

- Computationally expensive: requires statistical averaging over noise and initial conditions
- NLL JIMWLK and helicity JIMWLK do not admit langevin formulation

Need a new method: Lindblad formulation of JIMWLK

An alternative

Write JIMWLK as a Lindblad evolution for open quantum systems



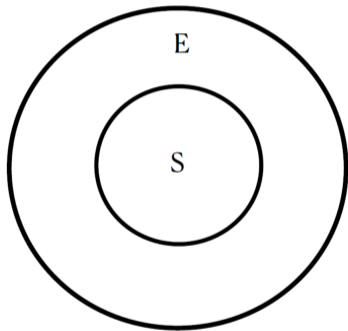
Lindblad evolution: Evolution of system density matrix ρ_S in the presence of environment

Interaction parameterized by **jump operators** Q_α

N.Armesto, F. Dominguez, A.Kovner, M.Lublinsky, V.V.Skokov (2019)

An alternative

Write JIMWLK as a Lindblad evolution for open quantum systems



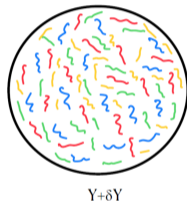
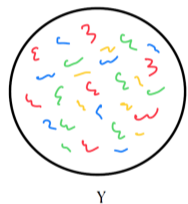
Lindblad evolution: Evolution of system density matrix ρ_S in the presence of environment

Interactions parameterized by **jump operators** Q_α

$$\partial_t \rho_S(t) = -i[H, \rho_S(t)] + \sum_{\alpha} Q_{\alpha} \rho_S(t) Q_{\alpha}^{\dagger} - \frac{1}{2} \{Q_{\alpha}^{\dagger} Q_{\alpha}, \rho_S(t)\}$$

JIMWLK as Lindblad equation

Write JIMWLK as a Lindblad evolution for open quantum systems



time \rightarrow rapidity

system \rightarrow target(valence)

environment \rightarrow soft degrees of freedom

$$\frac{d}{dY}\rho_T(Y) = \int \frac{d^2 z_\perp}{2\pi} \left[Q_i^a[z_\perp], \left[Q_i^a[z_\perp], \rho_T(Y) \right] \right]$$

What do we want?

Ultimate goal: Simulate JIMWLK with an initial density matrix

- Difficult on a quantum computer
- Reduce 2 dimensions to one dimension
- Work with $SU(2)$ instead of $SU(3)$
- Need a discretization/truncation of field space

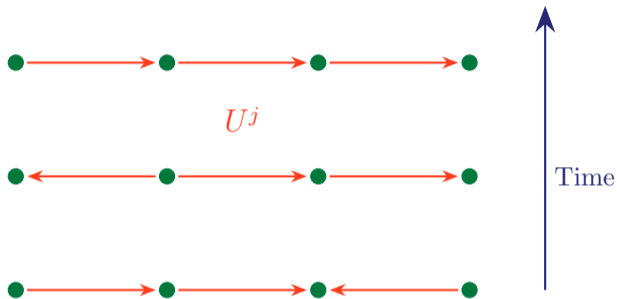
Reduced goal: Simulate reduced JIMWLK in 0 dimensions (2 points) with $SU(2)$

Reduced goal: Simulate reduced JIMWLK in 0 dimensions (2 points) with SU(2), starting with an initial density matrix.

$$\frac{d}{dY} \rho_T(Y) = a_{\perp}^2 \sum_{z_{\perp}} z_{\perp} \left[Q_i^a[z_{\perp}], \left[Q_i^a[z_{\perp}], \rho_T(Y) \right] \right]$$

Main problem: Choose a basis for truncated hilbert space of valence dof

Hamiltonian lattice gauge theory

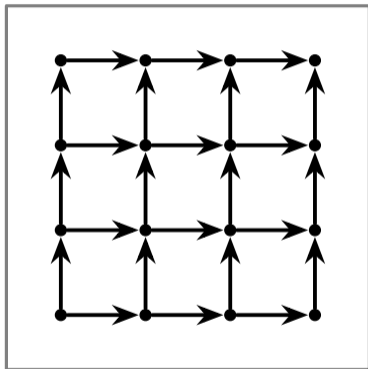


Space: Discretized, Time: Continuous

The setup

Wilson links are the degrees of freedom

$$U^j(\vec{r}, \mu) = \exp \left\{ ig a_s \tau_j^a A_\mu^a(\vec{r}) \right\}$$



Spatial lattice at a fixed time slice

The setup

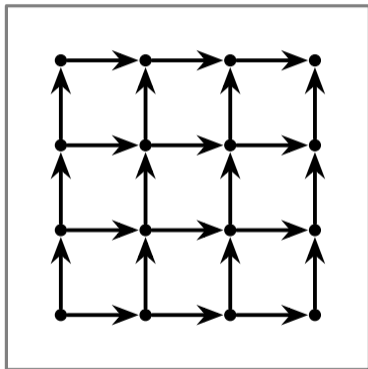
Conjugate to Wilson links: Electric fields

$$[\mathbf{E}^a, \mathbf{U}] = \tau^a \mathbf{E}^a,$$

$$[\mathbf{E}^a, \mathbf{E}^b] = i f^{abc} \mathbf{E}^c,$$

$$[\mathbf{F}^a, \mathbf{U}] = -\mathbf{F}^a \tau^a,$$

$$[\mathbf{F}^a, \mathbf{F}^b] = i f^{abc} \mathbf{F}^c.$$

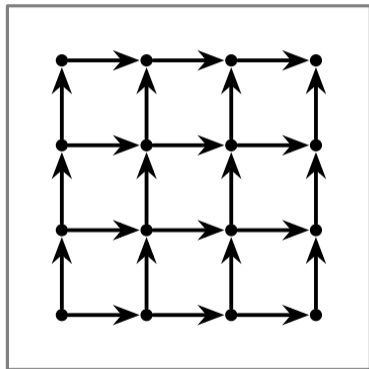


Spatial lattice at a fixed time slice

The setup

Two types of bases possible

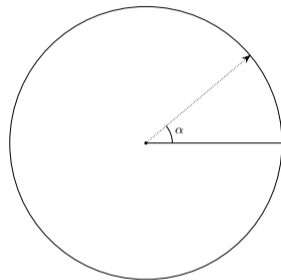
- **Group element basis:** Eigenstates of the U operator
- **Electric field basis:** Eigenstates of the E operator



Spatial lattice at a fixed time slice

$$U^j(\vec{r}, \mu)|\alpha\rangle = U^j(\vec{r}, \mu, \alpha)|\alpha\rangle$$

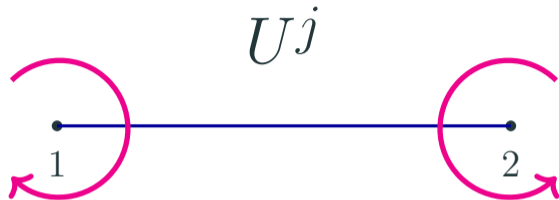
- α runs from 0 to 2π : is compact
- useful for formulating initial conditions



Hilbert space: Electric field basis

$$\mathbf{E}^a \mathbf{E}^a |j, m, n\rangle = j(j+1) |j, m, n\rangle$$

$$\mathbf{F}^a \mathbf{F}^a |j, m, n\rangle = j(j+1) |j, m, n\rangle$$

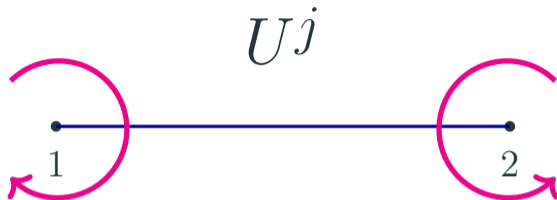


$$|j, m, n\rangle = \mathbf{U}_{mn}^j |0\rangle$$

T. Brynes, Y. Yamamoto (2006)

Hilbert space: Electric field basis

- j : Casimir
- m : Left rotation
- n : Right rotation

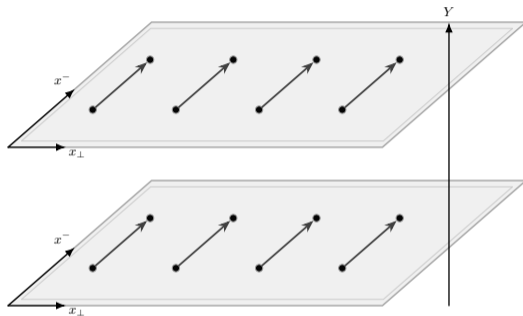


$$|j, m, n\rangle = \mathbf{U}_{mn}^j |0\rangle$$

$$U^j \rightarrow V(1) U^j V^{-1}(2)$$

Hilbert space of a rigid rotor

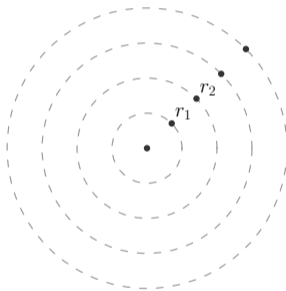
Mapping to JIMWLK setup



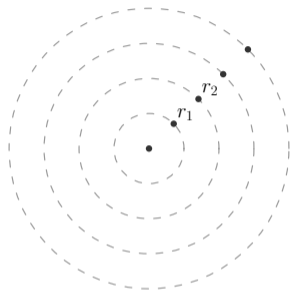
Discretize the infinite wilson link along x^- direction

Space discretization

We approximate the JIMWLK jump operator with an azimuthally symmetric one

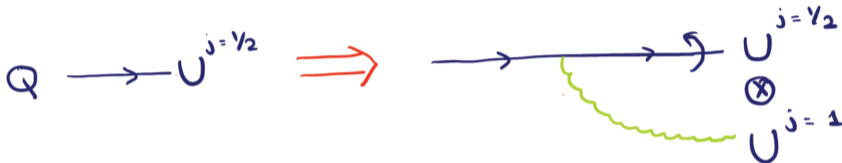


Consider two radial lattice points



- $j^{r_1}, j^{r_2} < j_{max}$
- $-j^r < n^r < j^r$
- For $j_{max} = 1/2$, 2 radial points: 25 states

Jump operators



$$U_{m_L m_R}^{1/2}(\alpha, x) U_{n_L n_R}^1(\alpha, x) = \sum_{J=1/2}^{3/2} C_{j, m_L j' n_L}^{J, M_L} C_{j m_R j, m_R}^{J, M_R} U_{M_L M_R}^J(\alpha, x)$$

Using this we can compute jump operators in the representation basis

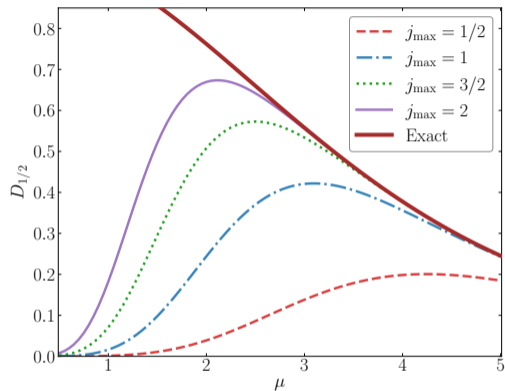
What do we want?

- Work with a maximum j : j_{max}
- Write initial conditions in group element basis
- Check convergence in representation basis
- Evolve using Lindblad-JIMWLK and check convergence of evolution

Work with a gaussian density matrix in group element basis and do a basis transformation

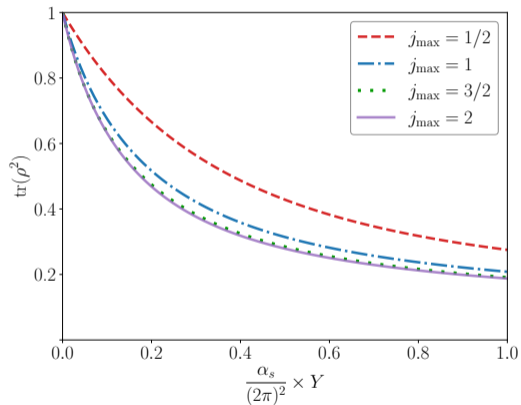
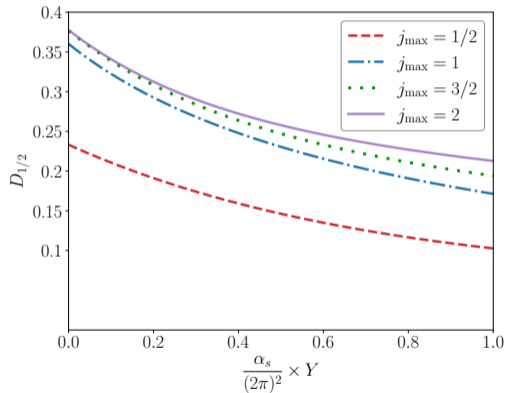
$$\langle \alpha | \rho | \alpha' \rangle \approx e^{-\frac{2(\alpha^2 + \alpha'^2)}{\mu^2}}$$

Convergence of the initial condition

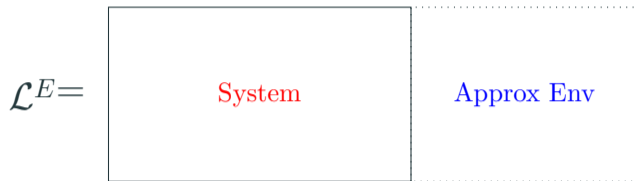


$$\hat{D}_{1/2} = \text{tr}_c(\mathbf{U}^{\dagger j}(r_1) \mathbf{U}^{1/2}(2))$$

Convergence of evolution



Implementation on quantum computers

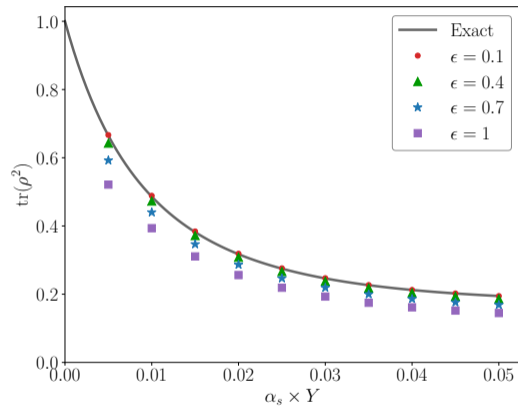
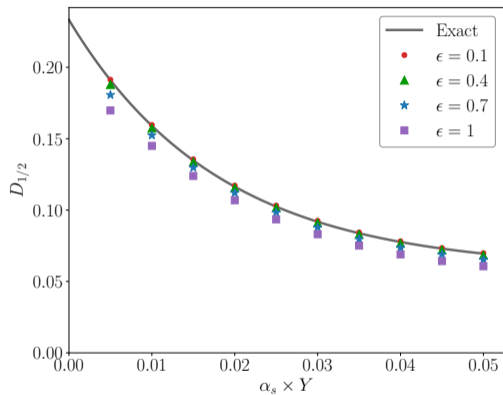


- Enlarge the hilbert space to make the lindbladian operator unitary

$$\frac{d|\rho\rangle^E}{dY} = \mathcal{L}^E |\rho\rangle^E$$

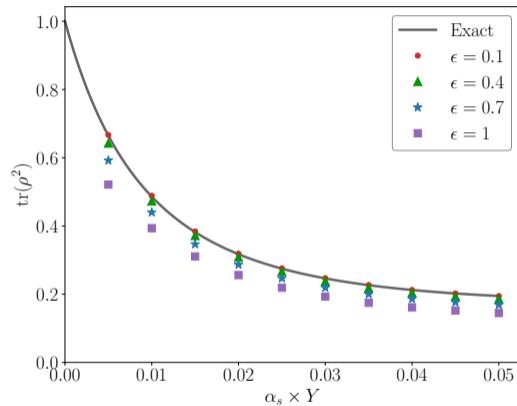
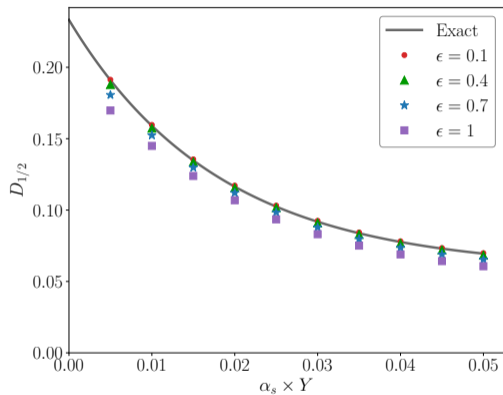
Schlimgen et al (2021, 2022)

Implementation on quantum computers



ϵ : Error parameter of the enlargement scheme

Implementation in quantum computers



$\epsilon \rightarrow 0$: Exact but noisy on quantum computers

JIMWLK in 2-dimensions (SU(3))

Reduce dimensions \downarrow 2 points

JIMWLK in 0-dimensions (SU(3))

Reduce generators \downarrow 3 colors

JIMWLK in 0-dimensions (SU(2))

Truncate field space \downarrow j_{max}

QM JIMWLK in 0-dimensions (SU(2))

Relax the approximations

- Generalize to infinite wilson lines
- Work with the $SU(3)$ group
- Intelligent truncations to include more transverse points
- Compute more interesting (off-diagonal) observables