

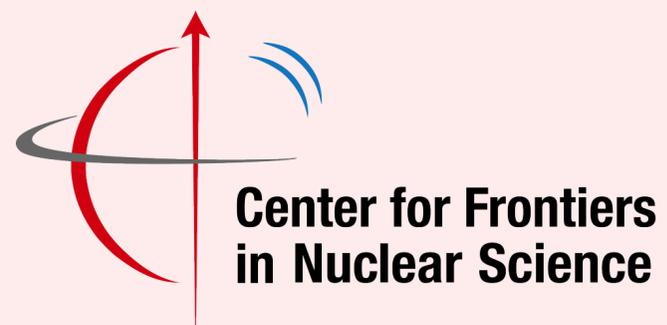
High precision A_1 and g_1 measurements at ePIC in the Early Running of EIC

Win Lin

Stony Brook University

CFNS Friday Meeting

03/06/2026



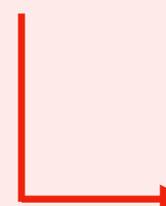
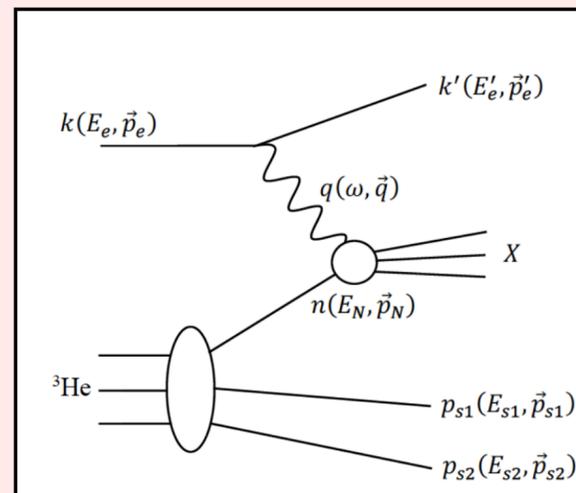
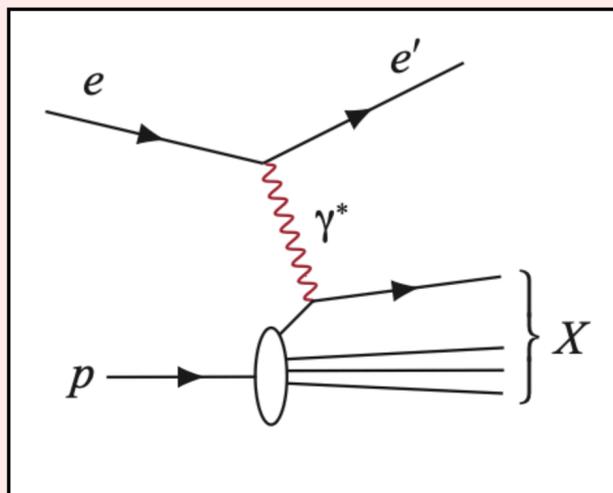
Current Early Science Matrix

| | Species | Energy (GeV) | Luminosity/year (fb ⁻¹) | Electron polarization | p/A polarization |
|---|---------------|----------------------|-------------------------------------|-----------------------|--------------------------|
| YEAR 1 | e+Ru or e+Cu | 10 x 115 | 0.9 | NO (Commissioning) | N/A |
| YEAR 2 | e+D e+p | 10 x 130 | 11.4 4.95 - 5.33 | LONG | NO TRANS |
| YEAR 3 | e+p | 10 x 130 | 4.95 - 5.33 | LONG | TRANS and/or LONG |
| YEAR 4 | e+Au e+p | 10 x 100 10 x 250 | 0.84 6.19 - 9.18 | LONG | N/A TRANS and/or LONG |
| YEAR 5 | e+Au e+3He | 10 x 100 10 x 166 | 0.84 8.65 | LONG | N/A TRANS and/or LONG |
| Note: the eA luminosity is per nucleon | | | | | |

Current Early Science Matrix

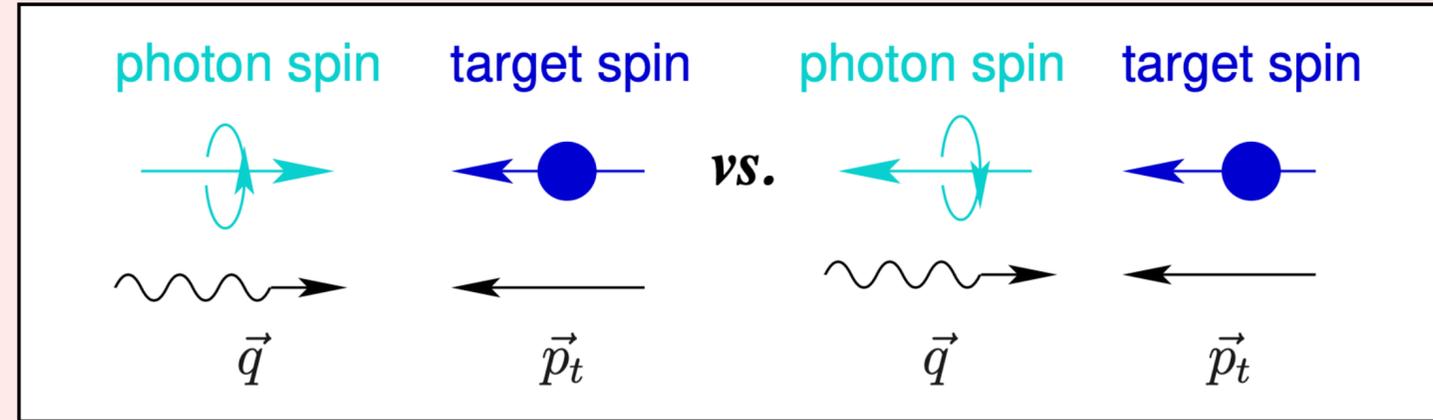
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Note: the eA luminosity is per nucleon



| | | |
|---------|---------|--------------------------------|
| A_1^p | A_1^n | Double spin asymmetry |
| g_1^p | g_1^n | Spin structure function |

Motivation: double spin asymmetry



$$A_1(x, Q^2) \equiv \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{A_{\parallel}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$$

p & γ^* spins anti-aligned

p & γ^* spins aligned

$$A_{\parallel} = \frac{\sigma_{\downarrow\uparrow} - \sigma_{\uparrow\uparrow}}{\sigma_{\downarrow\uparrow} + \sigma_{\uparrow\uparrow}}$$

p & e spins anti-aligned

p & e spins aligned

$$A_{\perp} = \frac{\sigma_{\downarrow\Rightarrow} - \sigma_{\uparrow\Rightarrow}}{\sigma_{\downarrow\Rightarrow} + \sigma_{\uparrow\Rightarrow}}$$

p & e spins perpendicular

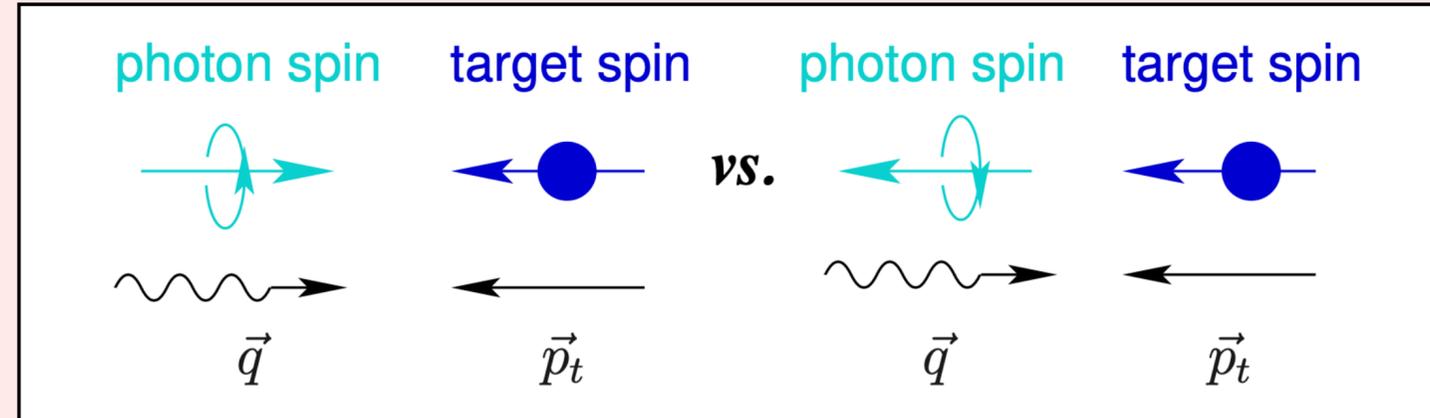
$$D = \frac{y(2 - y)(2 + \gamma^2 y)}{2(1 + \gamma^2)y^2 + (4(1 - y) - \gamma^2 y^2)(1 + R)}$$

$$R \equiv \frac{\sigma_L}{\sigma_T} = \frac{F_L}{F_2 - F_L}$$

$$d = \frac{D\sqrt{4(1 - y) - \gamma^2 y^2}}{2 - y}$$

$$\gamma^2 = \frac{4M^2 x^2}{Q^2} \quad \eta = \frac{4(1 - y) - \gamma^2 y^2}{(2 - y)(2 + \gamma^2 y)} \quad \xi = \frac{\gamma(2 - y)}{2 + \gamma^2 y}$$

Motivation: spin-dependent structure function



$$A_1(x, Q^2) \equiv \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{A_{\parallel}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$$

p & γ^* spins anti-aligned

p & γ^* spins aligned

$$g_1 = \frac{F_2}{2x(1 + R)}(A_1 + \gamma A_2)$$

Spin-dependent structure function

$$A_{\parallel} = \frac{\sigma_{\downarrow\uparrow} - \sigma_{\uparrow\uparrow}}{\sigma_{\downarrow\uparrow} + \sigma_{\uparrow\uparrow}}$$

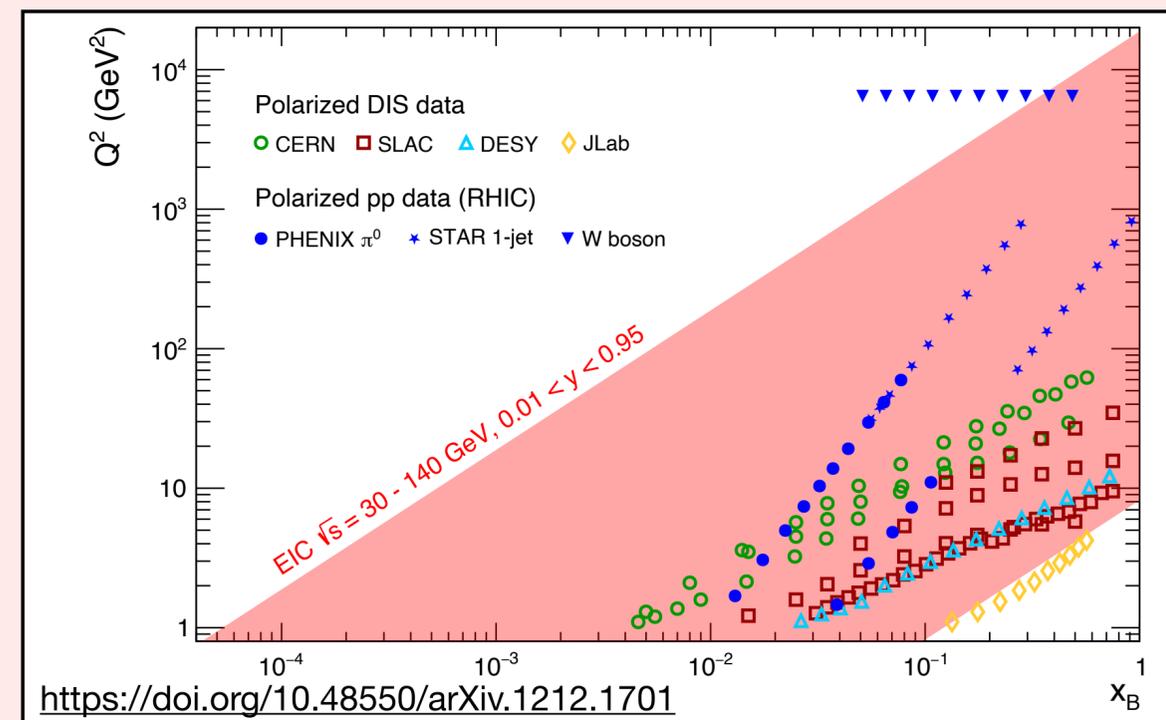
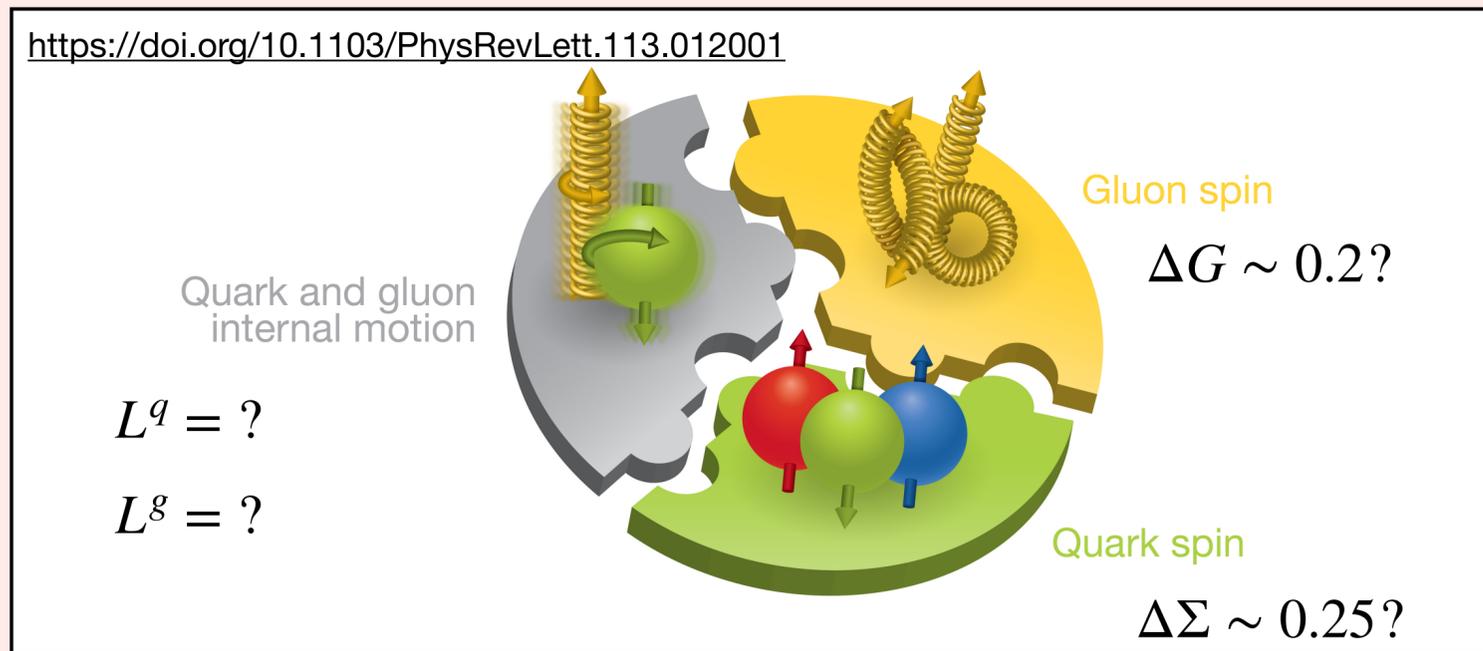
p & e spins anti-aligned

p & e spins aligned

$$A_{\perp} = \frac{\sigma_{\downarrow\Rightarrow} - \sigma_{\uparrow\Rightarrow}}{\sigma_{\downarrow\Rightarrow} + \sigma_{\uparrow\Rightarrow}}$$

p & e spins perpendicular

$$A_2(x, Q^2) \equiv \frac{2\sigma_{LT}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{\xi A_{\parallel}}{D(1 + \eta\xi)} - \frac{A_{\perp}}{d(1 + \eta\xi)}$$



Perturbative QCD:

$$g_1(x, t) = \frac{1}{2} \sum_{k=1}^{n_f} \frac{e_k^2}{n_f} \int_x^1 \frac{dy}{y} \left[C_q^S \left(\frac{x}{y}, \alpha_s(t) \right) \Delta \Sigma(y, t) + 2n_f C_g \left(\frac{x}{y}, \alpha_s(t) \right) \Delta G(y, t) + C_q^{\text{NS}} \left(\frac{x}{y}, \alpha_s(t) \right) \Delta q_{\text{NS}}(y, t) \right]$$

Bjorken sum:

$$\begin{aligned} \Gamma_1^{p-n}(\alpha_s) &= \Gamma_1^{p-n}(Q^2) = \int_0^1 (g_1^p(Q^2) - g_1^n(Q^2)) dx \\ &= \sum_{n>0} \frac{\mu_{2n}^{p-n}(\alpha_s)}{Q^{2n-2}} = \frac{g_A}{6} \left[1 - \frac{\alpha_s(Q^2)}{\pi} - 3.58 \left(\frac{\alpha_s(Q^2)}{\pi} \right)^2 - 20.21 \left(\frac{\alpha_s(Q^2)}{\pi} \right)^3 - 175.7 \left(\frac{\alpha_s(Q^2)}{\pi} \right)^4 - \mathcal{O}((\alpha_s)^5) \right] \end{aligned}$$

at finite Q^2

Why measuring double spin asymmetry

$$\frac{d^3\sigma(\beta)}{dQ^2 dx d\phi} = \frac{d^3\sigma_0}{dQ^2 dx d\phi} - \frac{d^3\Delta\sigma(\beta)}{dQ^2 dx d\phi}$$

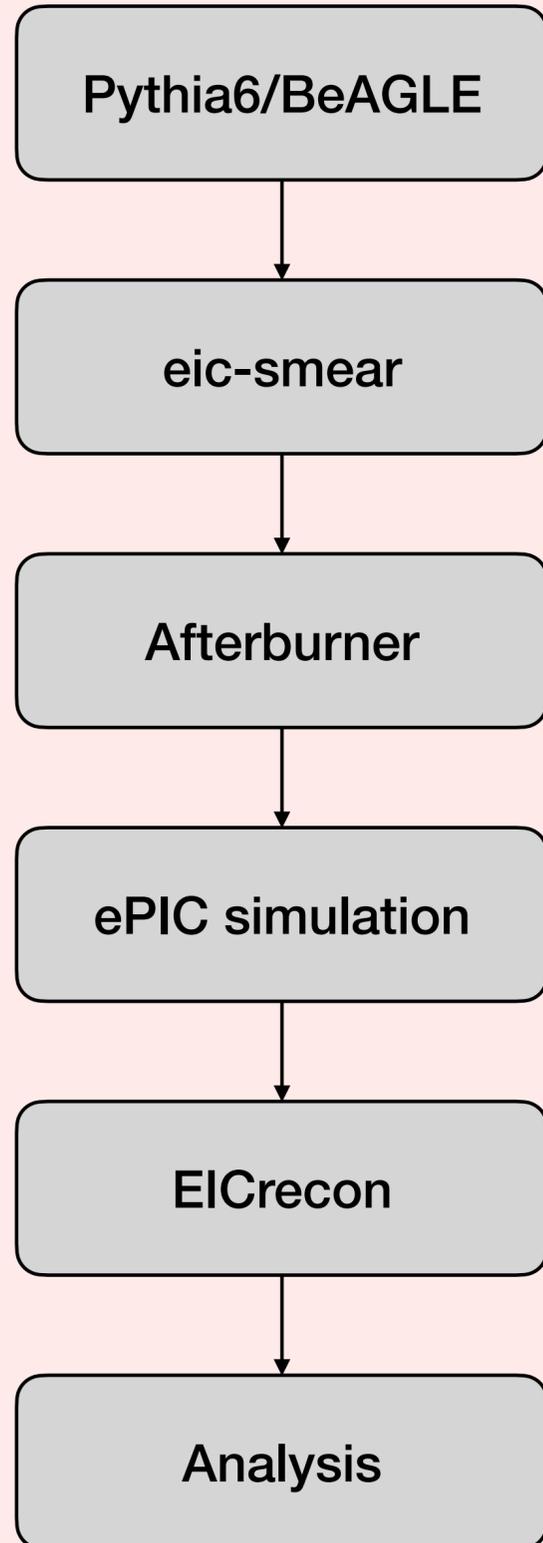
$$\frac{d^2\sigma}{dQ^2 dx} = \frac{4\pi\alpha^2}{Q^4} \frac{1}{x} \left[xy^2 F_1(x, Q^2) + \left(1 - y - \frac{Mxy}{2E} F_2(x, Q^2) \right) \right]$$

$$\frac{d^3\Delta\sigma(\beta)}{dQ^2 dx d\phi} = \frac{4\alpha^2}{Q^2} y \left\{ \cos\beta \left[\left(1 - \frac{y}{2} - \frac{\gamma^2 y^2}{4} \right) g_1(x, Q^2) - \frac{\gamma^2 y}{4} g_2(x, Q^2) \right] - \cos\phi \sin\beta \frac{\sqrt{Q^2}}{\nu} \left(1 - y - \frac{\gamma^2 y^2}{4} \right)^{\frac{1}{2}} \left[\frac{y}{2} g_1(x, Q^2) + g_2(x, Q^2) \right] \right\}$$

- Direct measurement is more challenging
- Easier to measure via cross section ratio
- Systematic is largely canceled out

NC Inclusive DIS: $e + p \rightarrow e' + X$

<https://arxiv.org/abs/2103.05419>



BeAGLE: eA event generator

A hybrid model consisting of DPMJet and PYTHIA with nPDF EPS09.

Nuclear geometry by DPMJet and nPDF provided by EPS09.

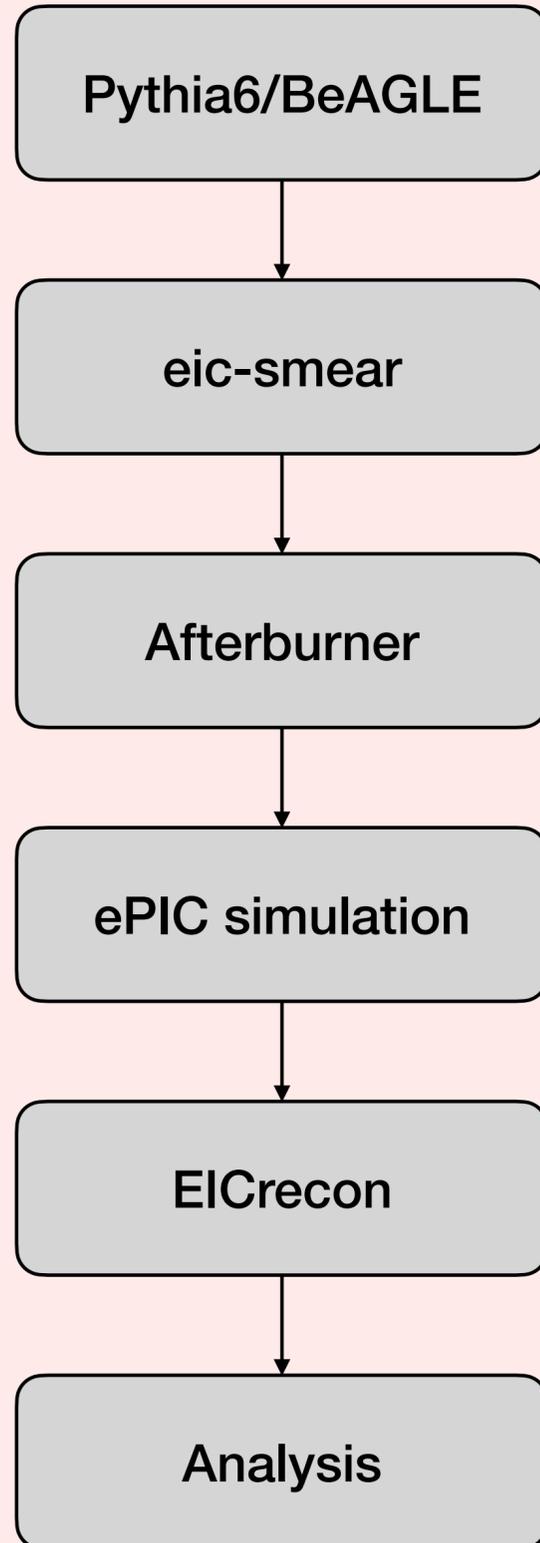
Parton level interaction and jet fragmentation completed in PYTHIA.

Nuclear evaporation (gamma deexcitation/nuclear fission/fermi break up) treated by DPMJet

Energy loss effect from routine by Salgado&Wiedemann to simulate the nuclear fragmentation effect in cold nuclear matter

- ☑ Parton distribution functions
- ☑ Short-range correlations
- ☑ Fermi motion
- ☑ Partonic (or “dipole”) MS
- ☑ Partonic gluon radiation
- ☑ Formation times
- ☑ Hadronic Cascade
- ☑ Nuclear evaporation, break up
- ☑ Photonic de-excitation of A*

Projection simulation & analysis flow



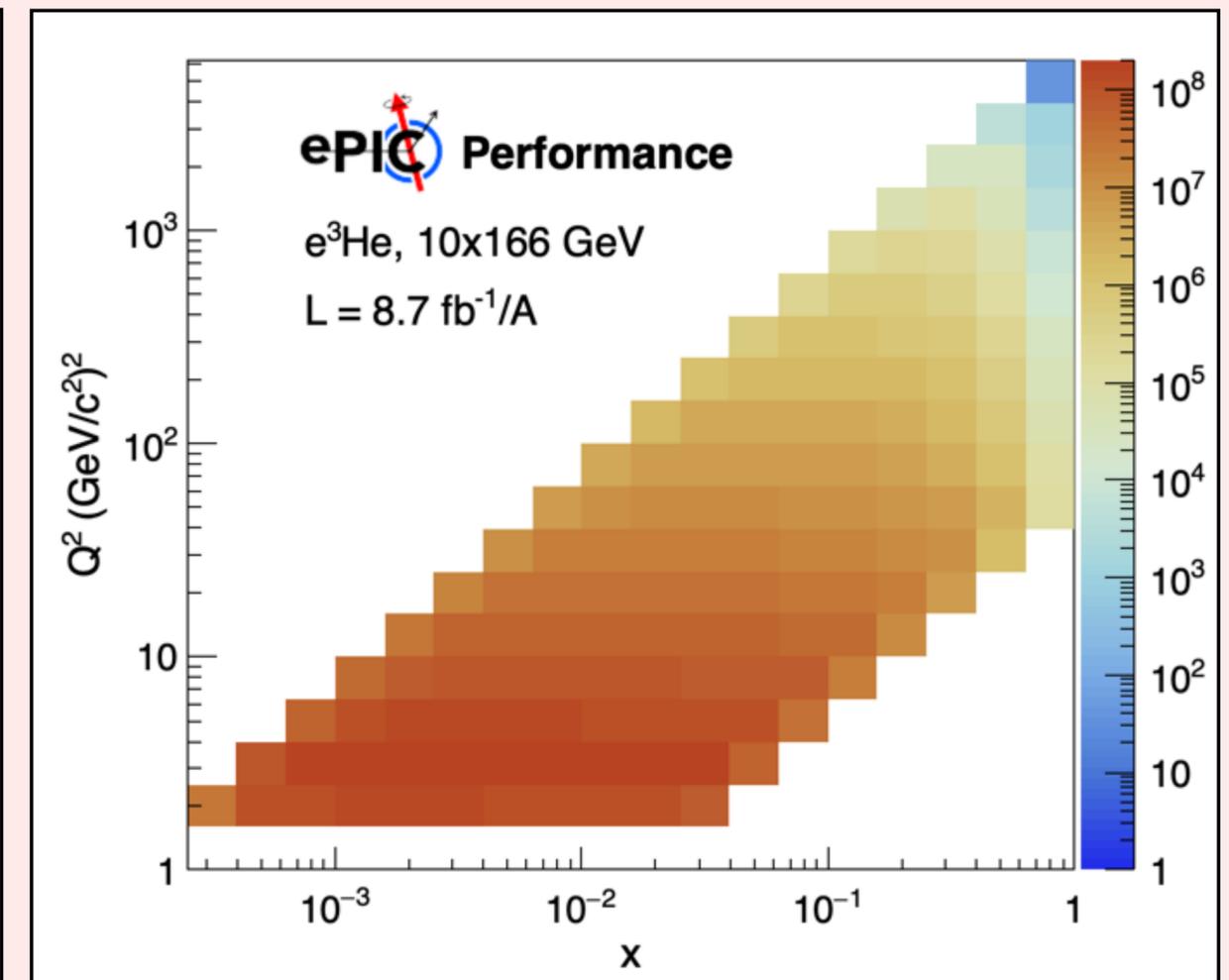
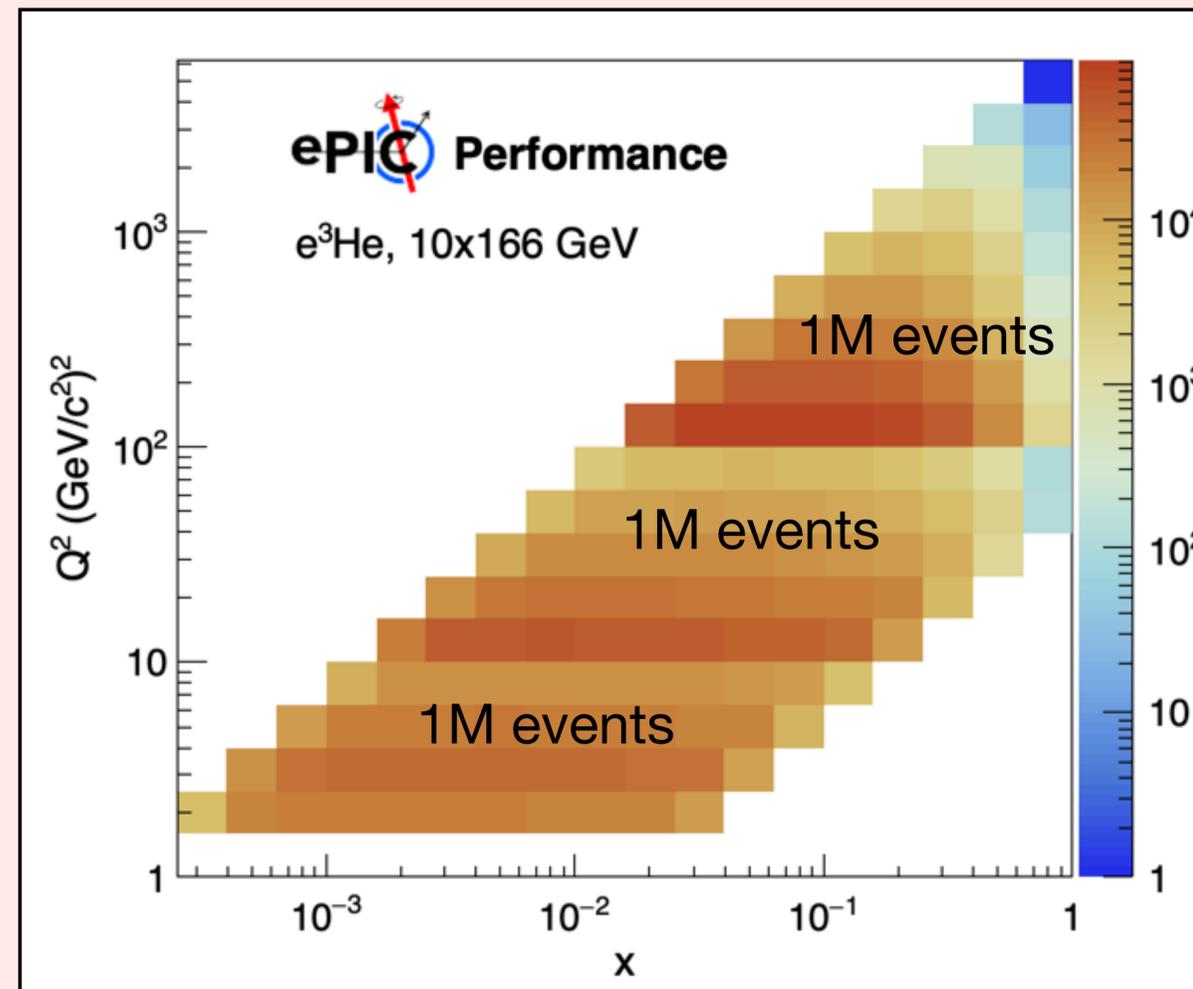
► Generated events are non-polarized, and have no QED radiative effect

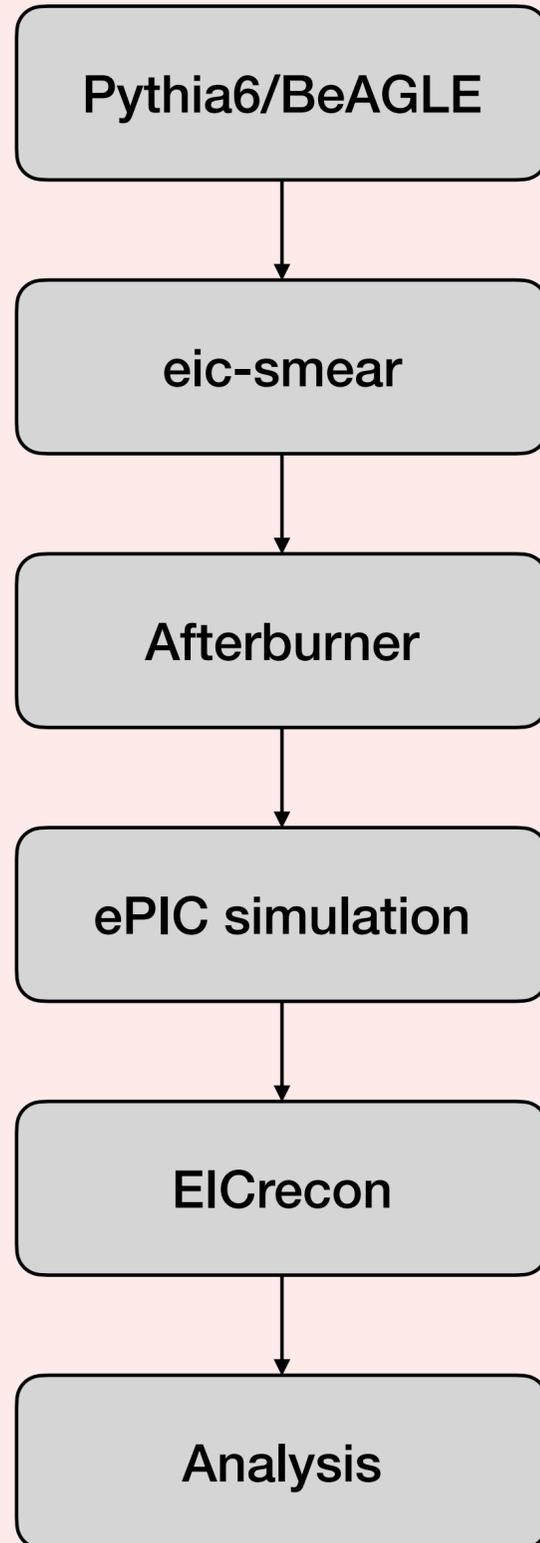
► Use generated event to estimate statistical uncertainty: $\delta A_{\parallel,\perp} = \frac{1}{\sqrt{N}P_e P_N}$

► Use parameterization for central values of A_1

generated events

scaled to planned luminosity

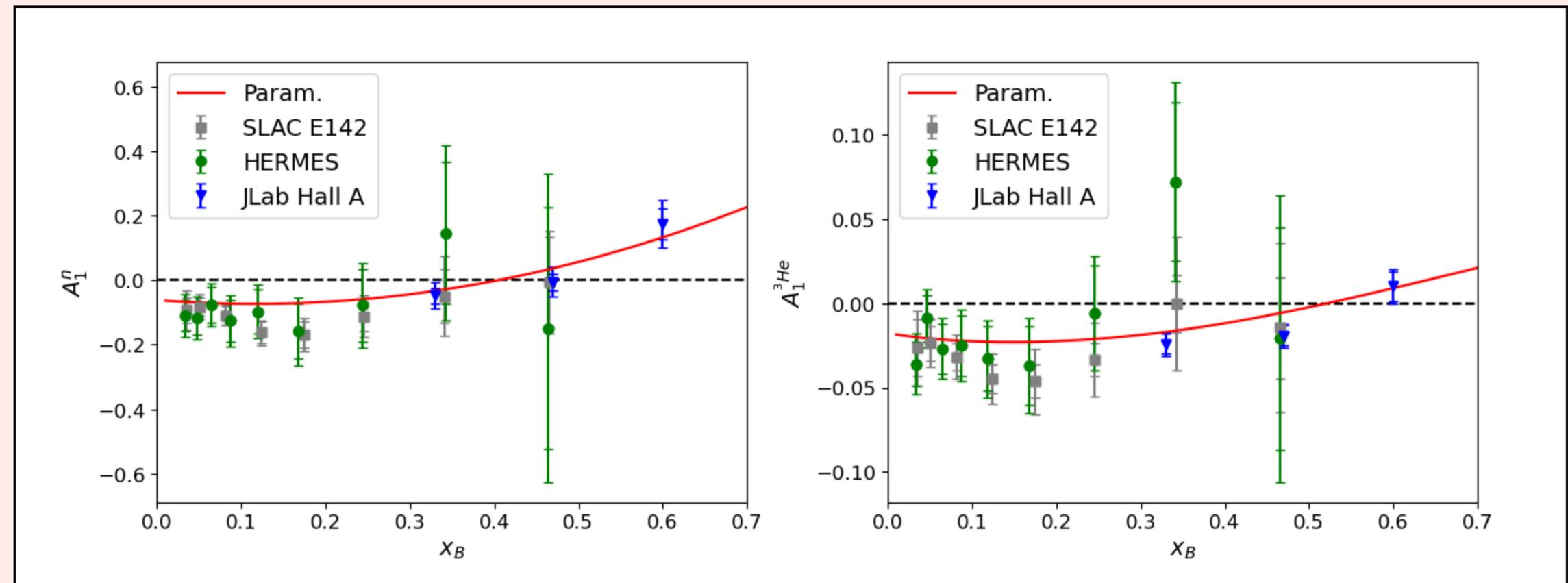




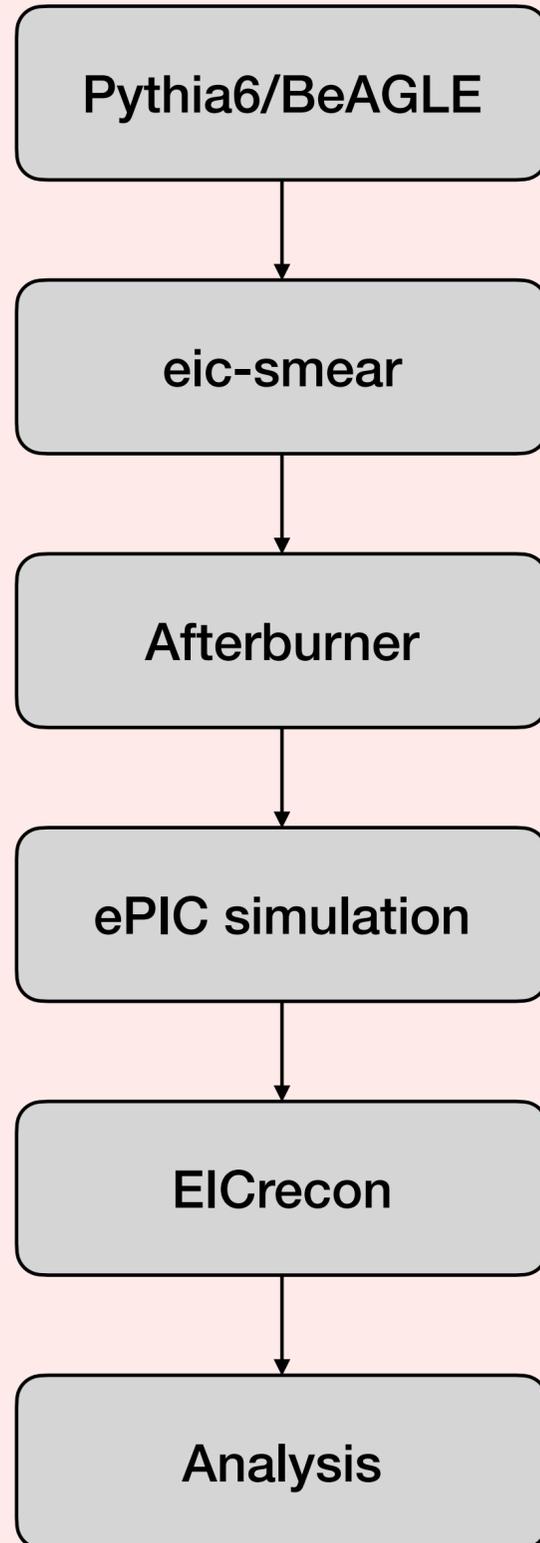
- ▶ A_1 calculated from: [Doi: 10.1103/PhysRevC.70.065207](https://doi.org/10.1103/PhysRevC.70.065207)

$$\text{▶ } A_1^{3\text{He}} = P_n \frac{F_2^n}{F_2^{3\text{He}}} A_1^n + 2P_p \frac{F_2^p}{F_2^{3\text{He}}} A_1^p \quad P_p = -0.028 \pm 0.004 \quad P_n = 0.86 \pm 0.02$$

- ▶ all F_2 's are taken from [JAM22](#)



- Parameterization at $Q^2 = 2.88 \text{ GeV}^2$
- Data points are at various Q^2 with majority $< 5 \text{ GeV}^2$



- Event reconstruction is done in ROOT codes: [GitHub link](#)

Scattered electron identification

Double spectator tagging

Kinematic reconstruction

Projecting physics results



▶ Start from “ReconstructedParticleCollection” ← Currently using MC to match clusters and tracks

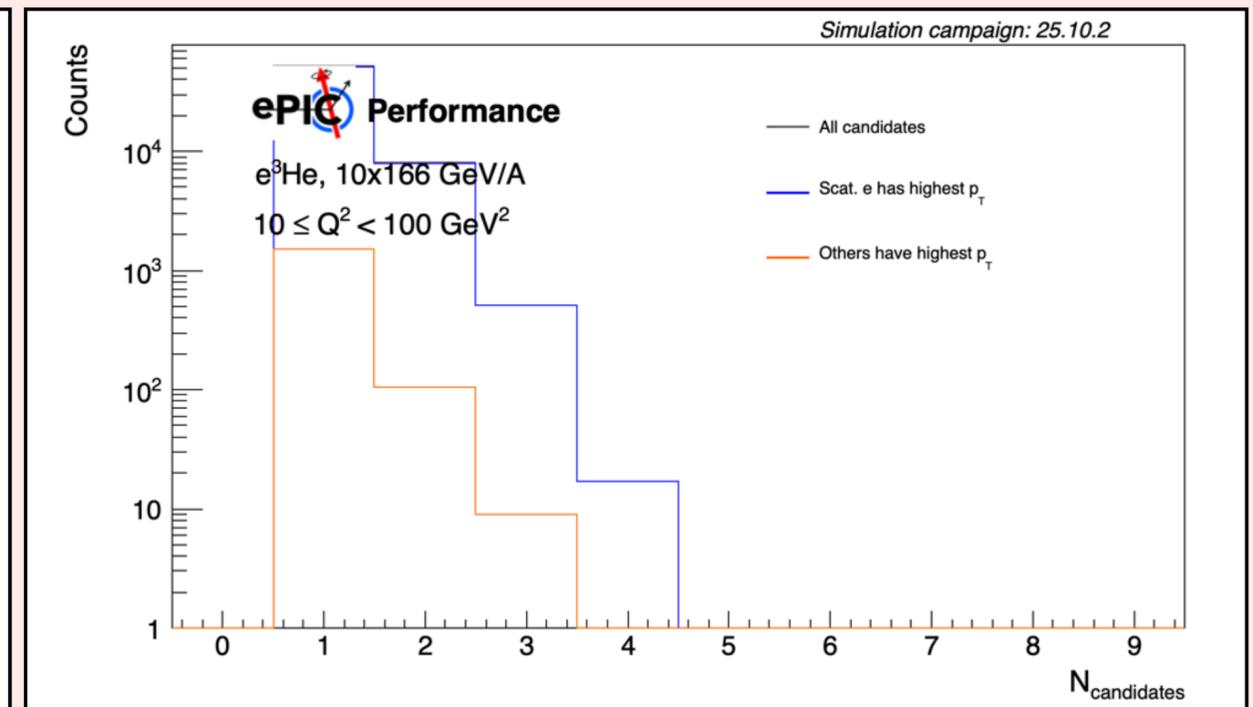
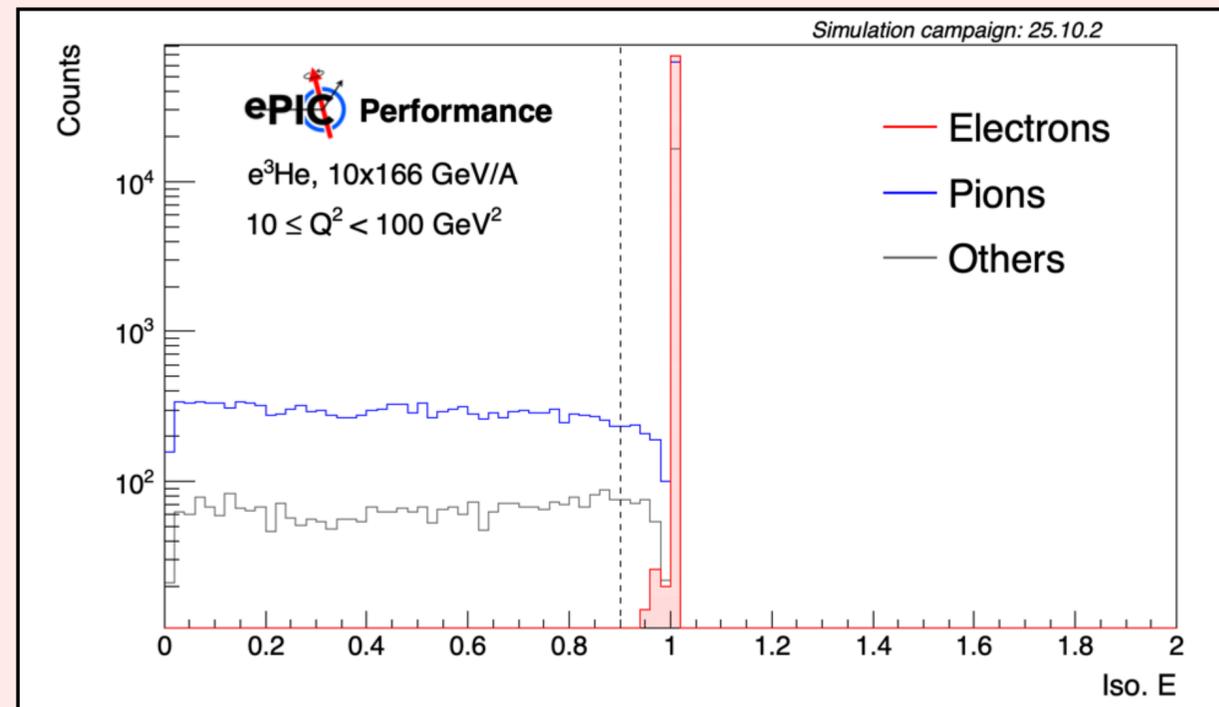
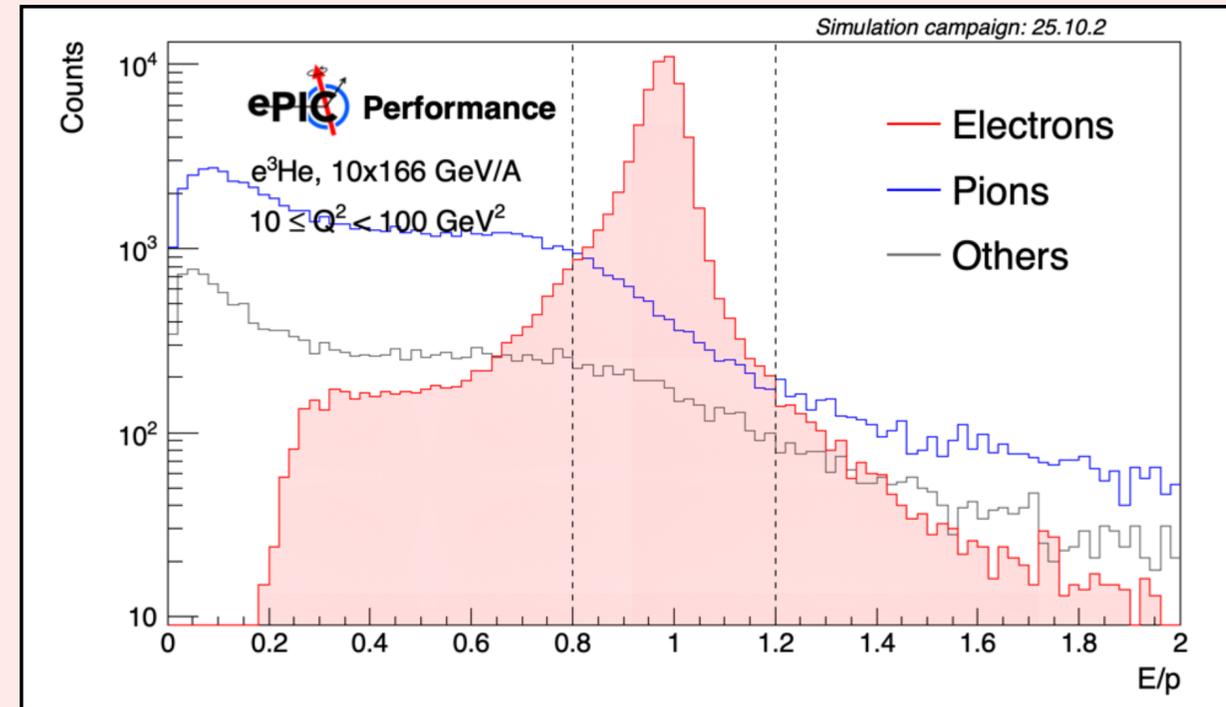
▶ Require clusters and negative tracks

▶ $0.8 < E/p < 1.2$

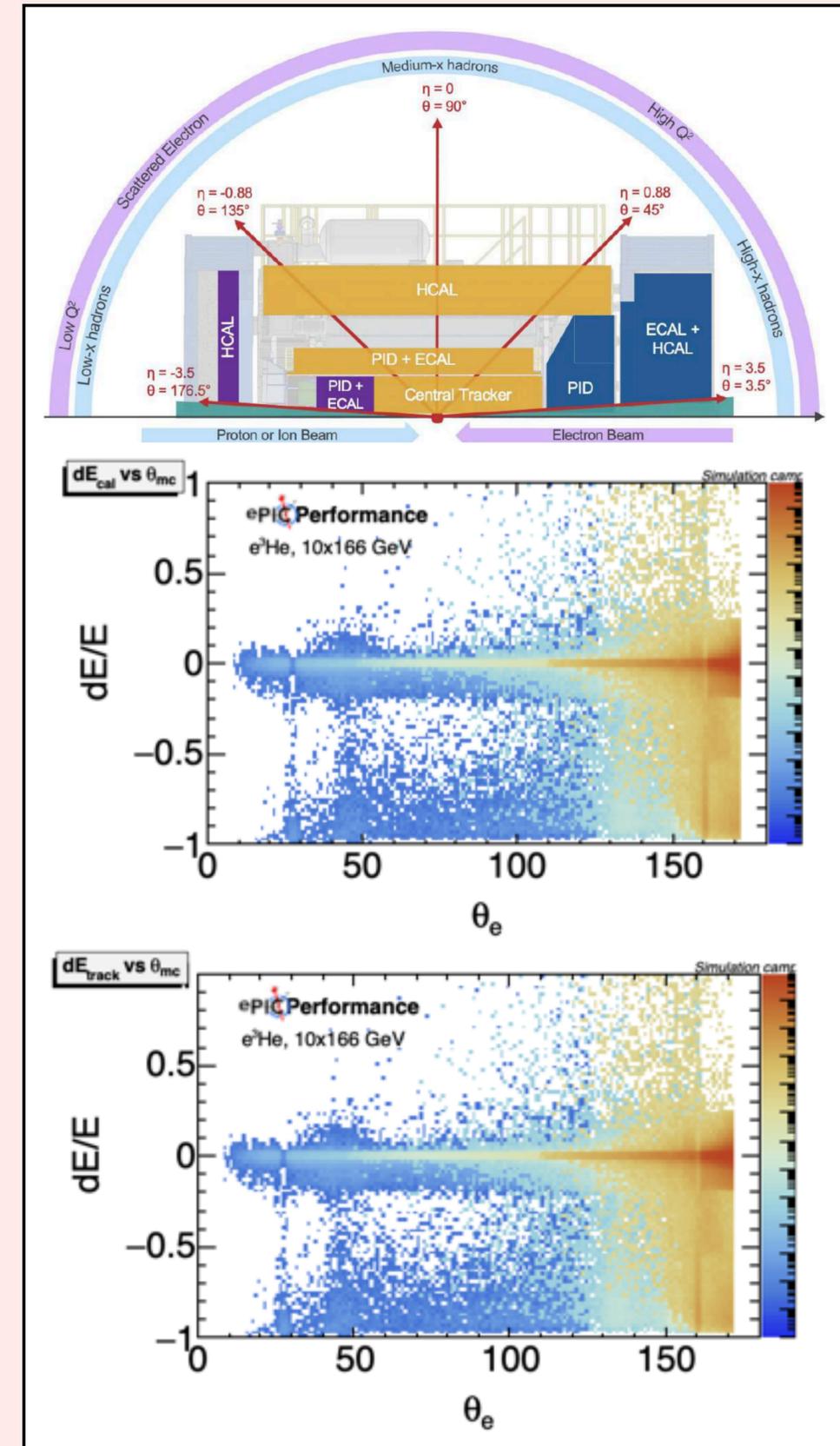
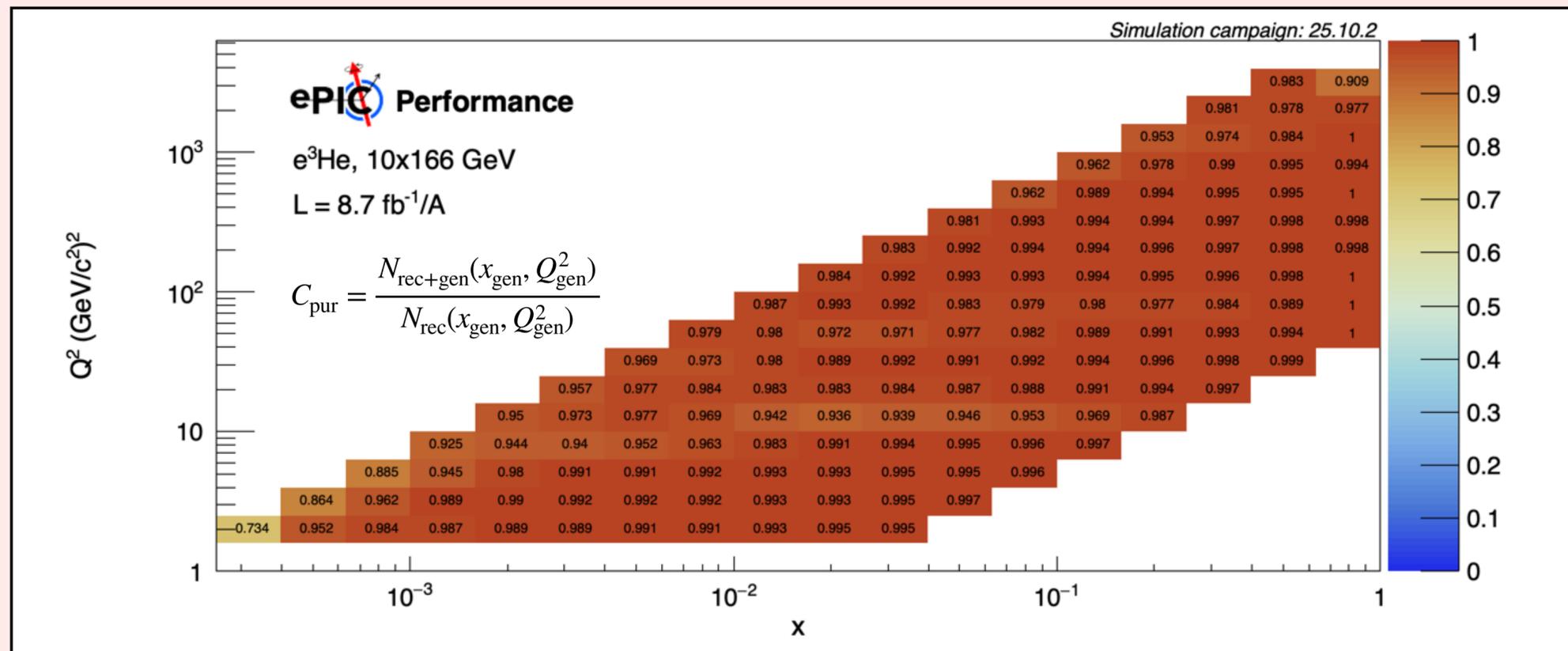
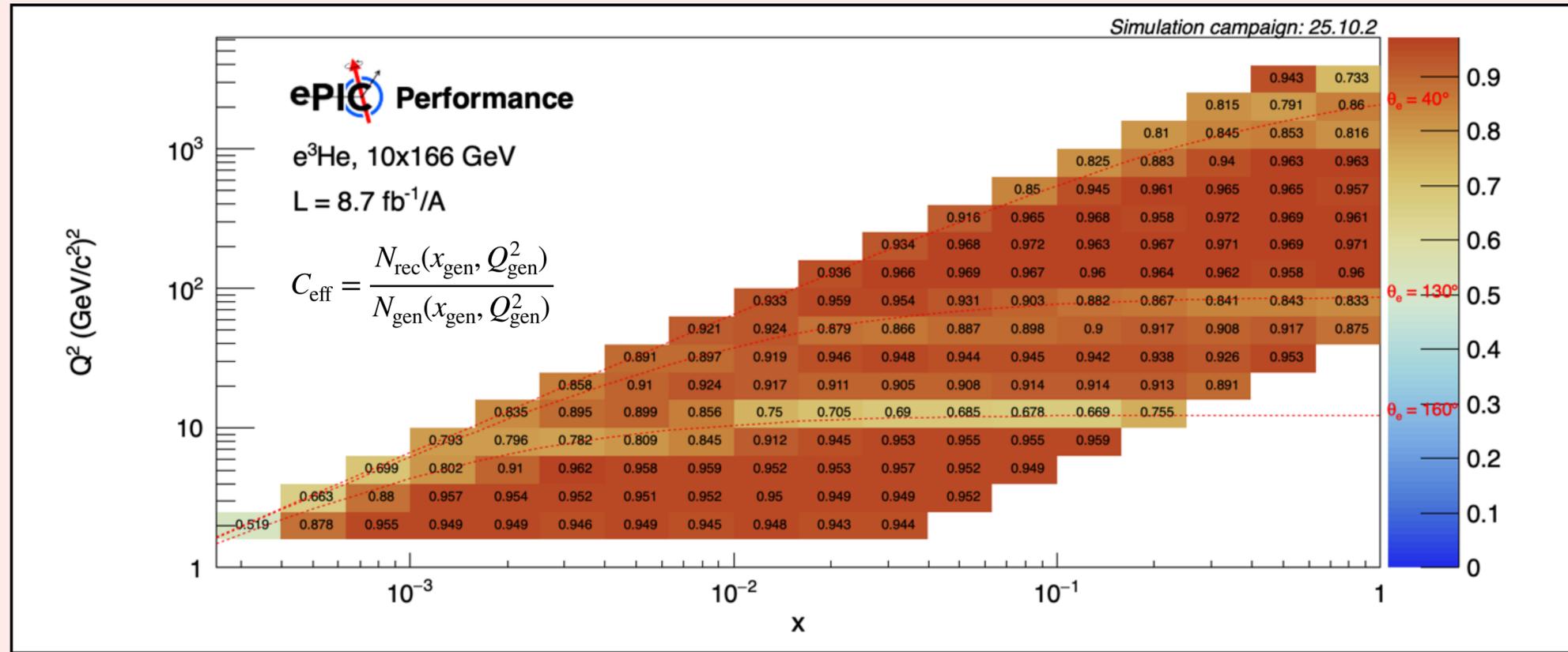
▶ Isolated cluster:

$$E_0 / \sum (E_{\Delta R < 0.4}) > 0.9, \Delta R = \sqrt{\Delta\eta^2 - \Delta\phi^2}$$

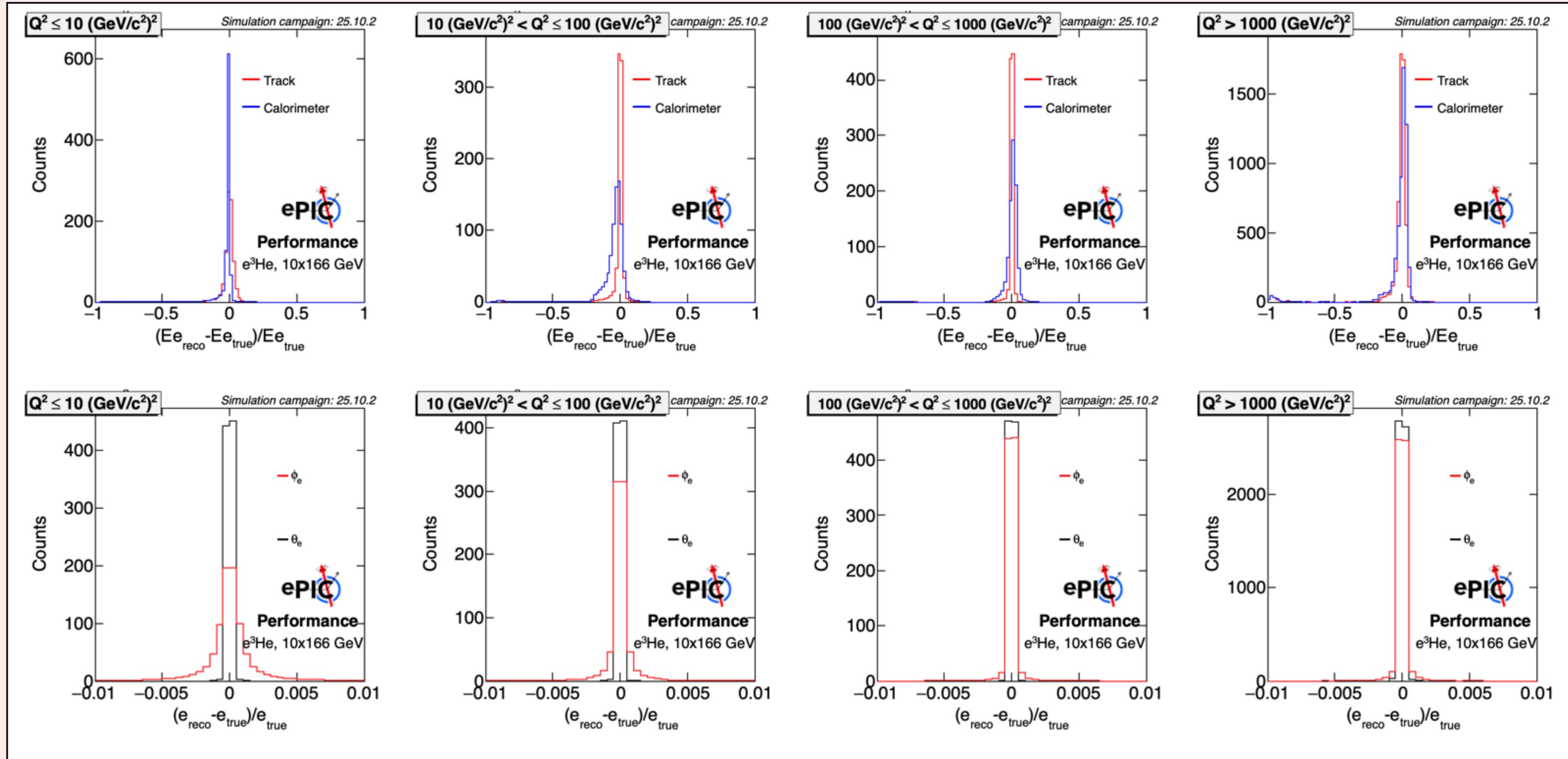
▶ Select the highest P_T particle



Scattered electron identification - efficiency & purity



Kinematic reconstruction - energy from track vs. cluster



Current available reconstruction in EICRecon

| Algorithm | Q^2 | Inelasticity y | Bjorken x |
|-----------------------|--|--|-------------------------|
| Electron (E) | $2E_0E_e(1 + \cos \theta_e)$ | $1 - \frac{E_e(1 - \cos \theta_e)}{2E_0}$ | $\frac{Q^2}{4E_0E_e y}$ |
| Jacquet-Blondel (JB) | $\frac{p_{t,h}^2}{1-y}$ | $\frac{\delta_h}{2E_0}$ | |
| Double-Angle (DA) | $\frac{4E_0^2}{\tan(\frac{\theta_e}{2})(\tan(\frac{\theta_e}{2}) + \delta_h/p_{t,h})}$ | $\frac{\delta_h/p_{t,h}}{\tan(\frac{\theta_e}{2}) + \delta_h/p_{t,h}}$ | |
| Sigma (Σ) | $\frac{E_e^2 \sin^2 \theta_e}{1-y}$ | $\frac{\delta_h}{\delta_h + E_e(1 - \cos \theta_e)}$ | x_Σ |
| E-Sigma ($e\Sigma$) | Q_E^2 | $\frac{Q_E^2}{4E_0E_e x_\Sigma}$ | |

$$p_{t,h}^2 = \left(\sum_h p_{x,h} \right)^2 + \left(\sum_h p_{y,h} \right)^2$$

$$\delta_h = \sum_h E_h - p_{z,h}$$

Uses only scattered e^- info ←

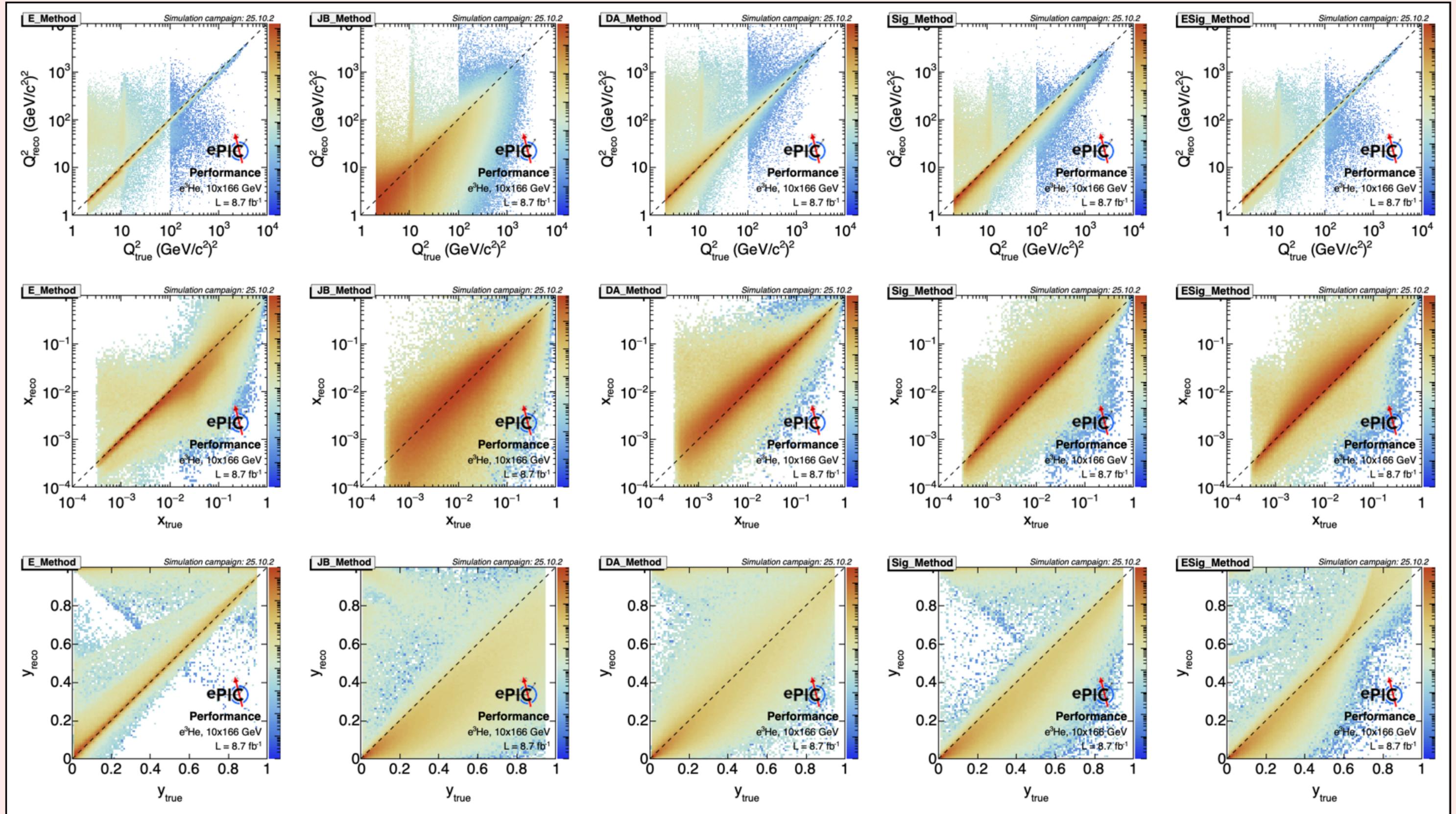
Uses everything else ←

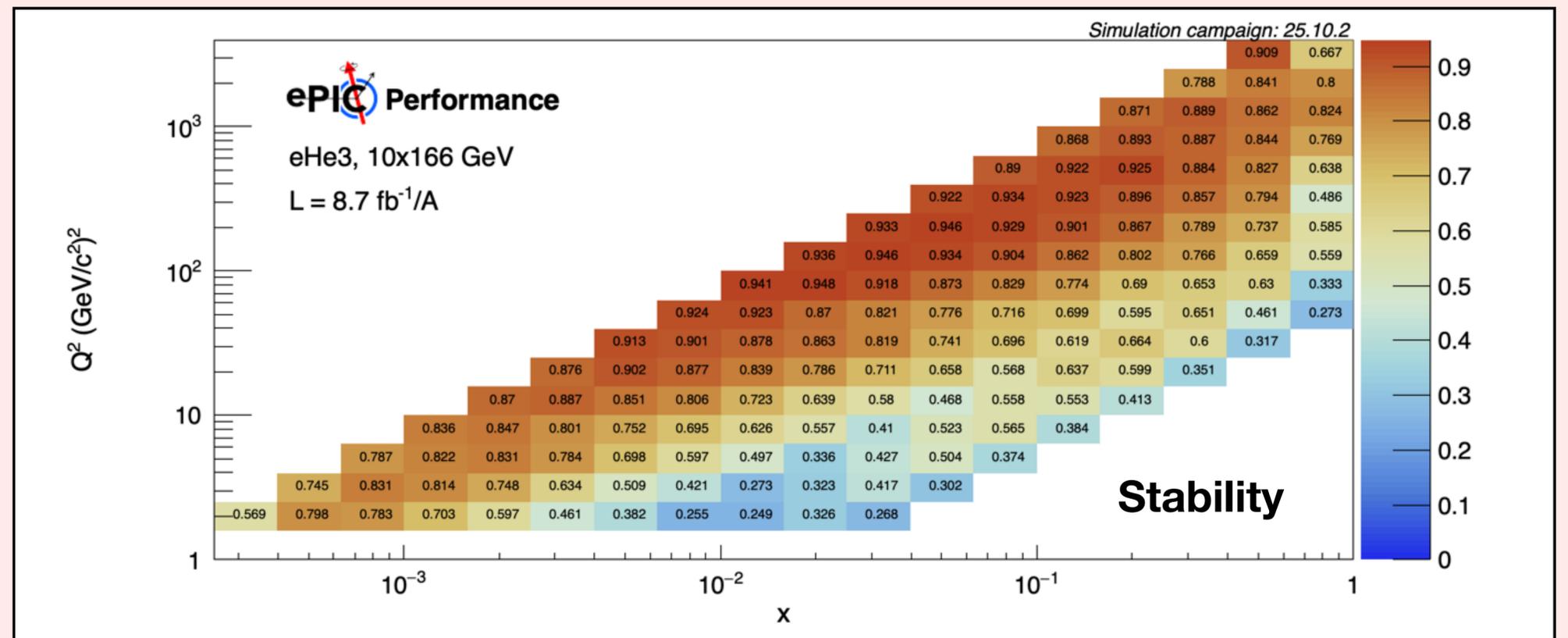
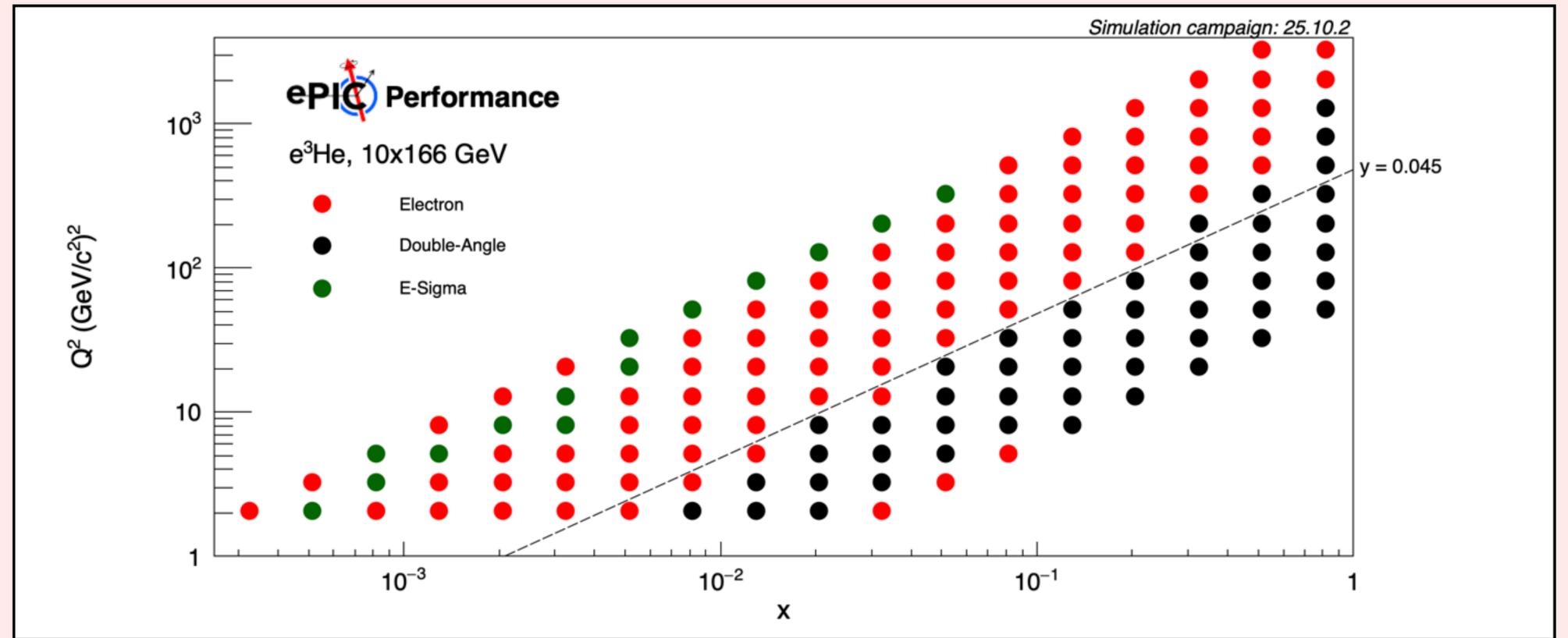
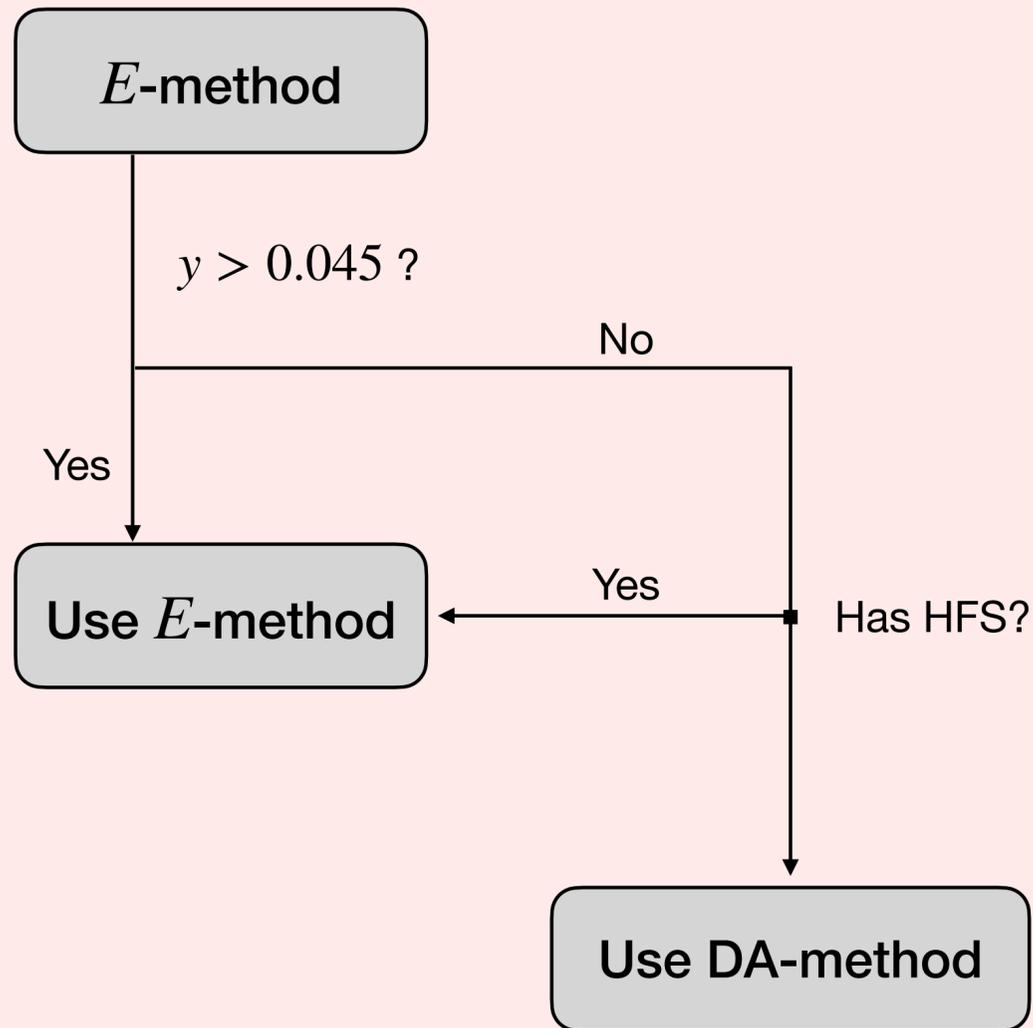
No use of scattered E_e ←

No use of scattered E_0 ←

Combine E and Σ method ←

Kinematic reconstruction - qualities





Projected A_1^P

$$\triangleright A_1(x, Q^2) \equiv \frac{\sigma_{1/2} - \sigma_{3/2}}{\sigma_{1/2} + \sigma_{3/2}} = \frac{A_{\parallel}}{D(1 + \eta\xi)} - \frac{\eta A_{\perp}}{d(1 + \eta\xi)}$$

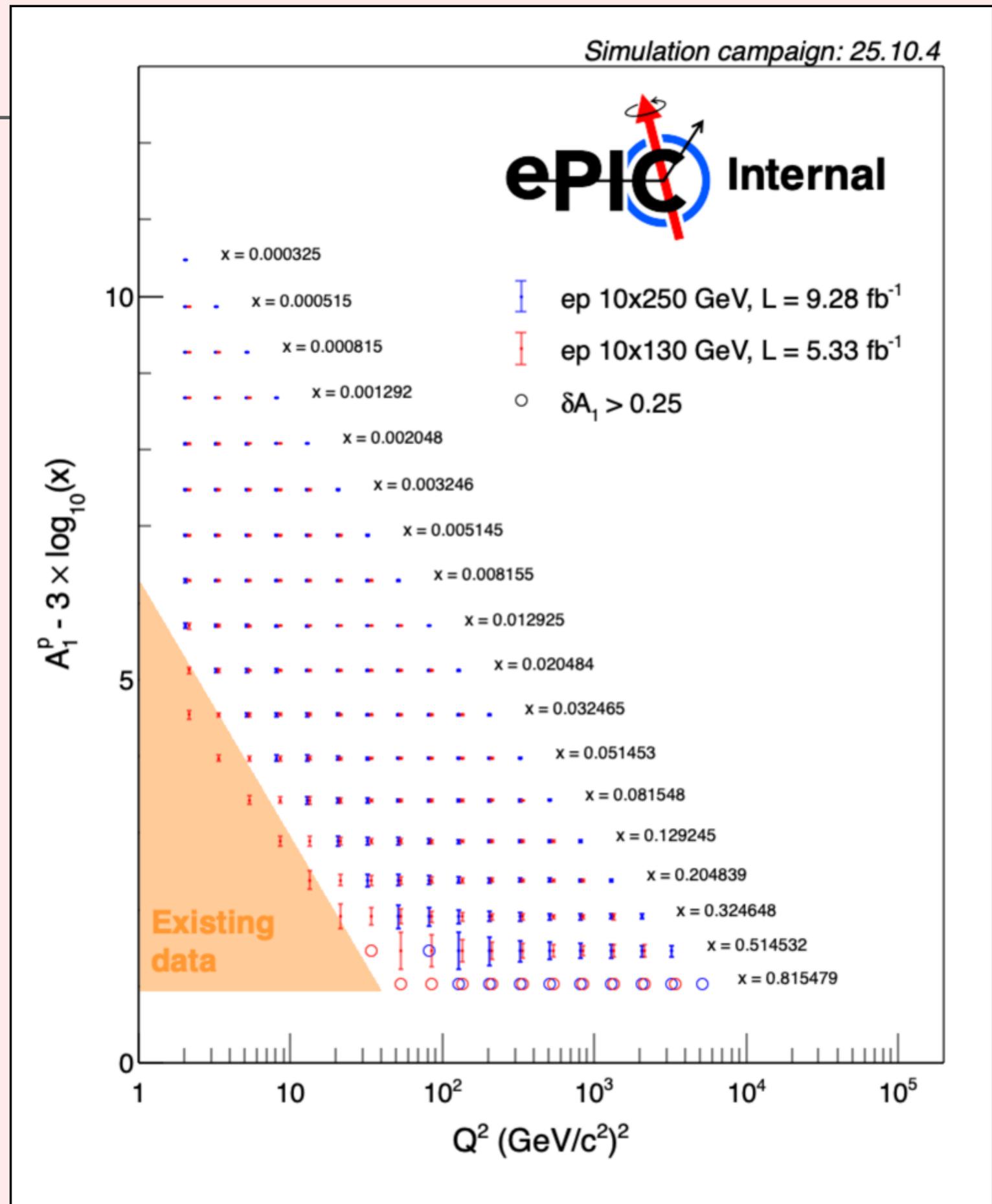
$$\triangleright \delta A_{\parallel, \perp} = \frac{1}{\sqrt{NP_e P_N}}$$

$$\triangleright P_e = P_p = 70\%$$

▶ Data split evenly between A_{\parallel} and A_{\perp}

▶ R calculated from [https://doi.org/10.1016/S0370-2693\(99\)00244-0](https://doi.org/10.1016/S0370-2693(99)00244-0)

▶ Statistical uncertainty only

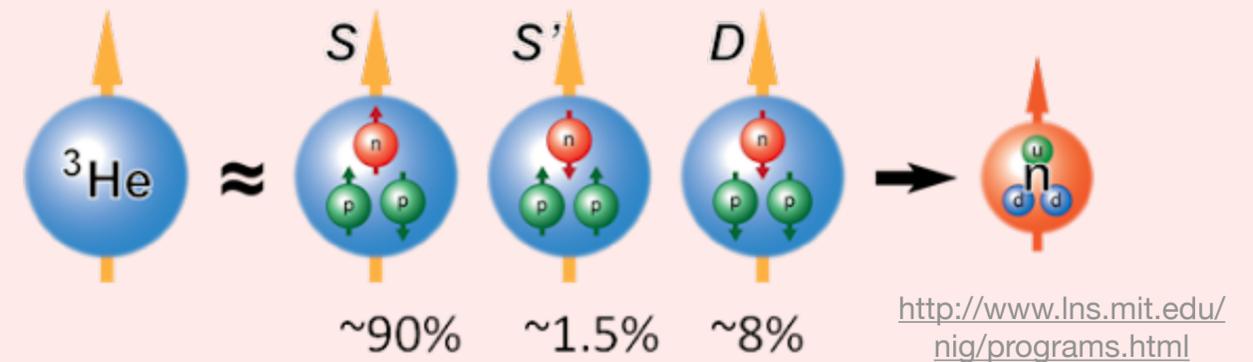


What about A_1^n ?

► Can be extracted from $A_1^{3\text{He}} = P_n \frac{F_2^n}{F_2^{3\text{He}}} A_1^n + 2P_p \frac{F_2^p}{F_2^{3\text{He}}} A_1^p$

$$P_n = 0.86 \pm 0.02 \quad P_p = -0.028 \pm 0.004$$

<https://doi.org/10.1103/PhysRevC.64.024004>

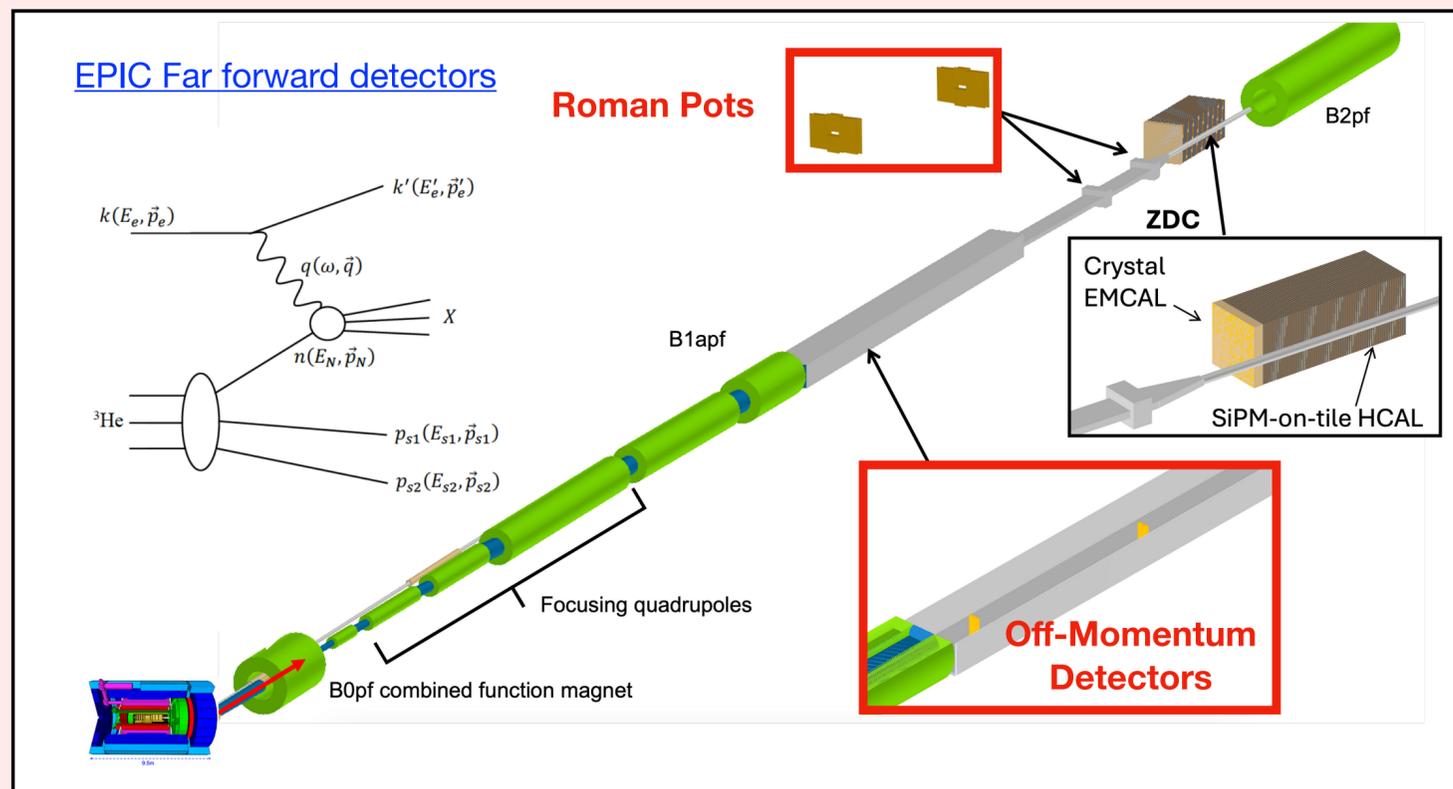
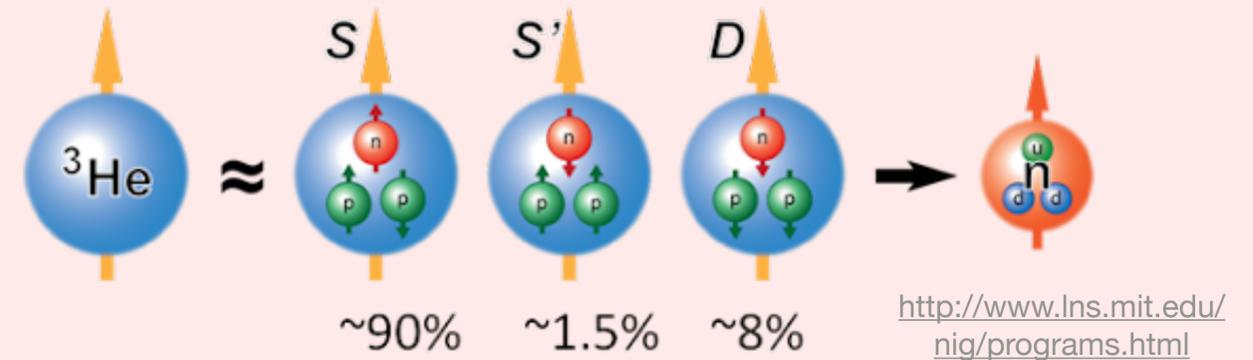


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- Or **directly** measured double spectator tagging:

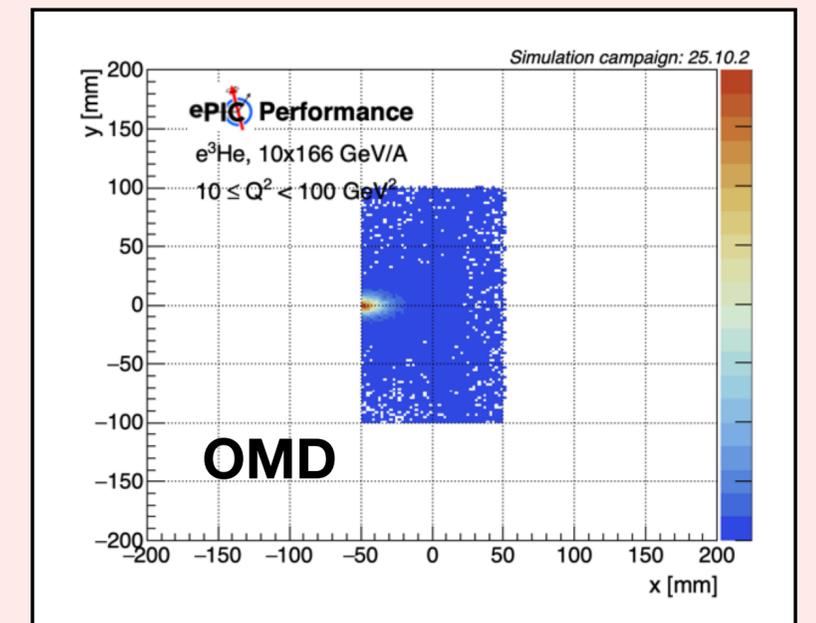
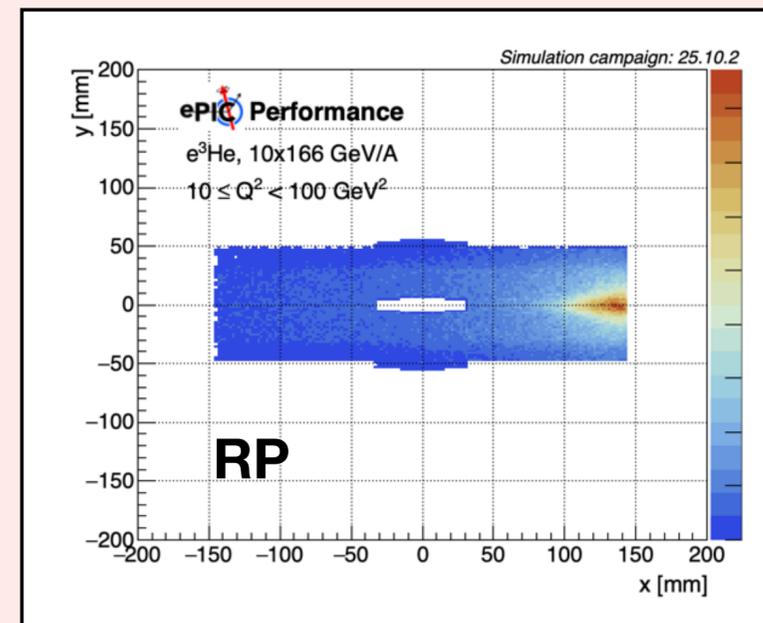
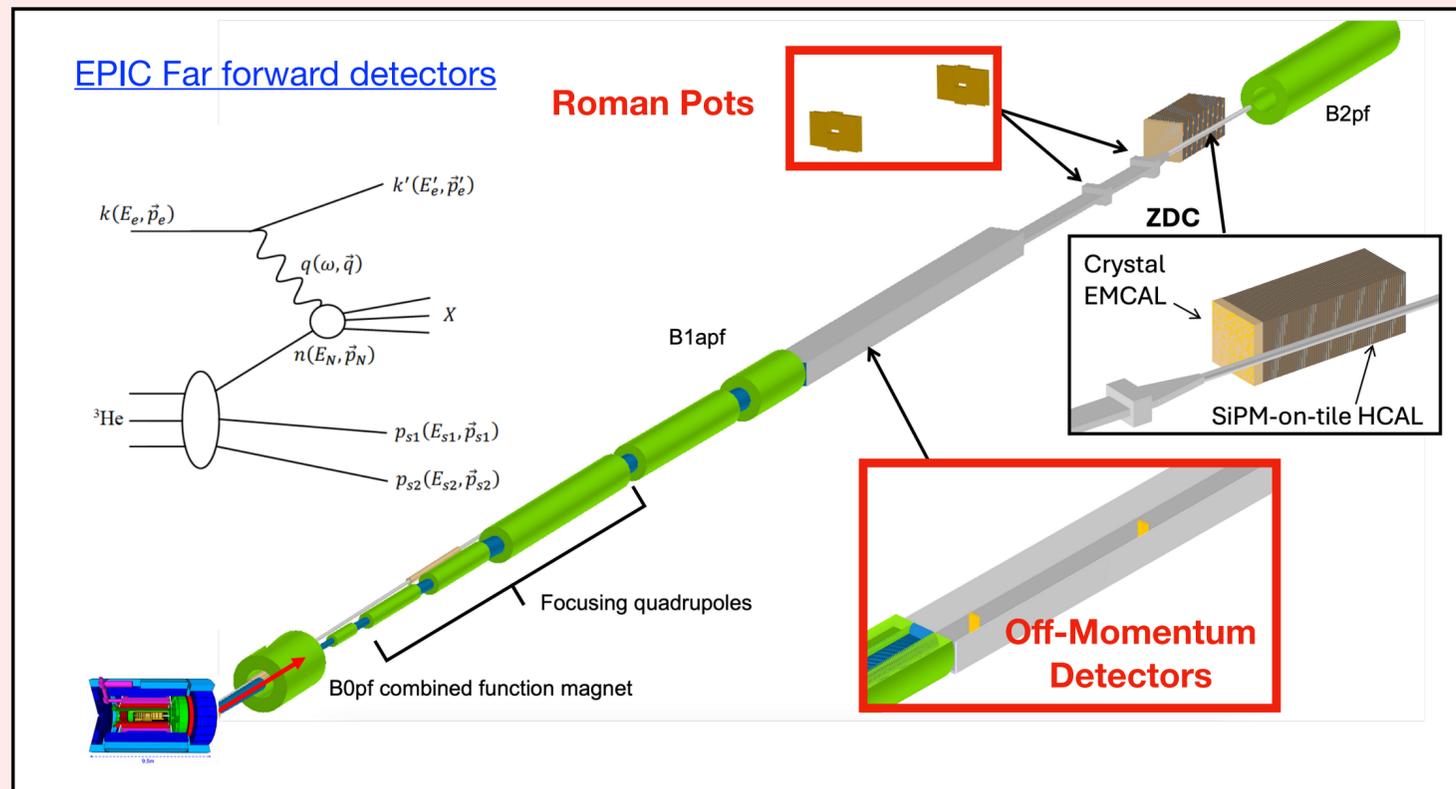
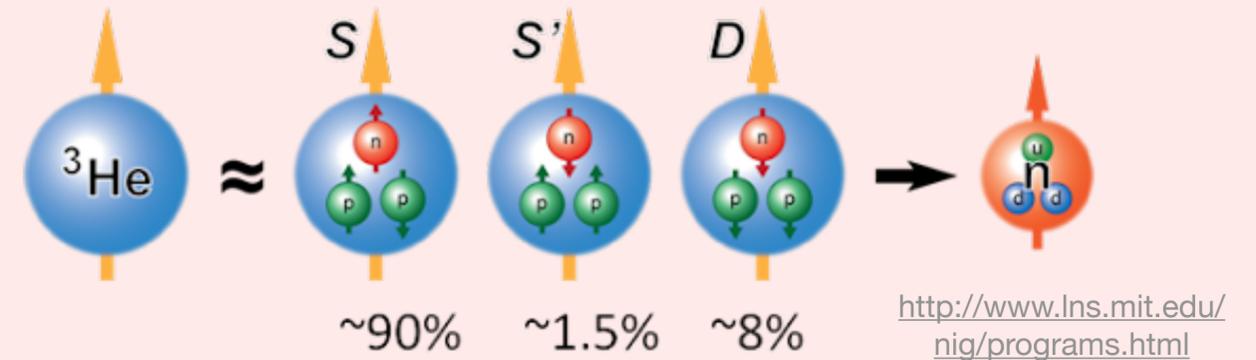


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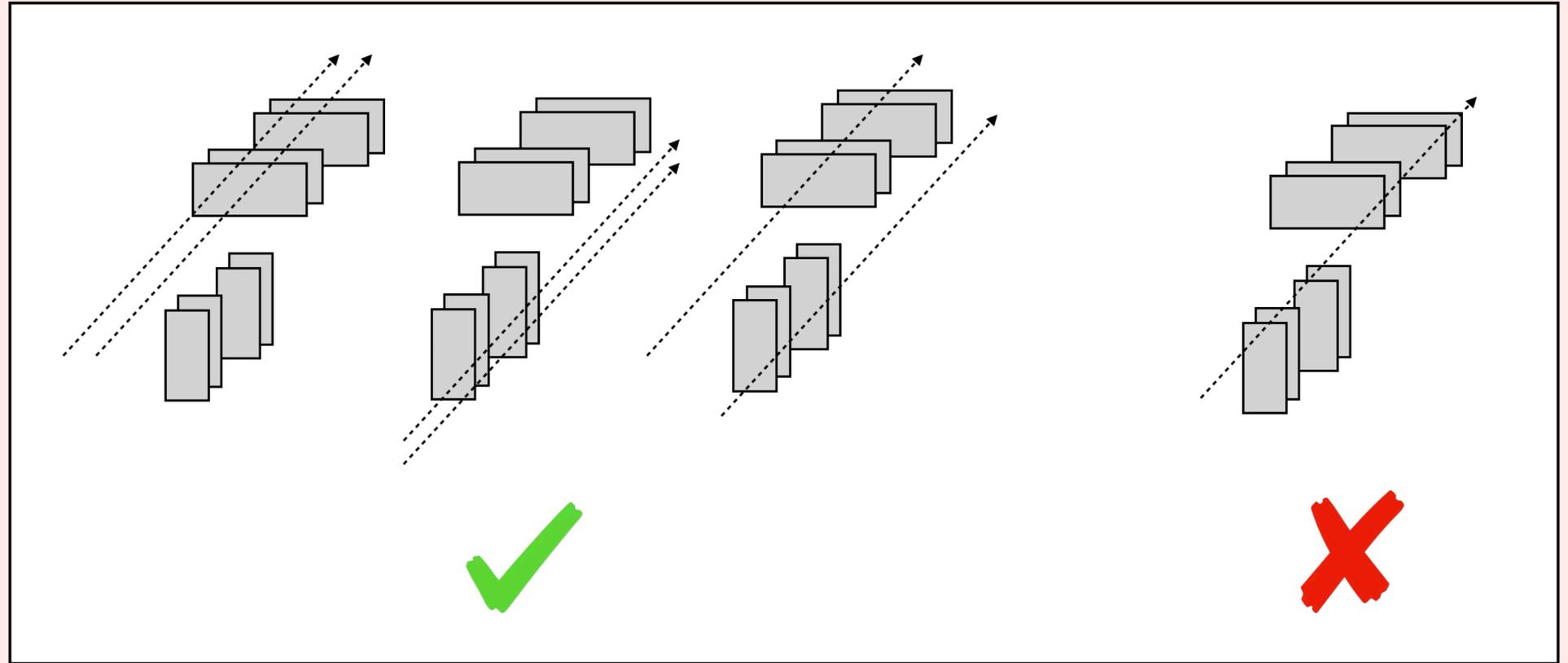
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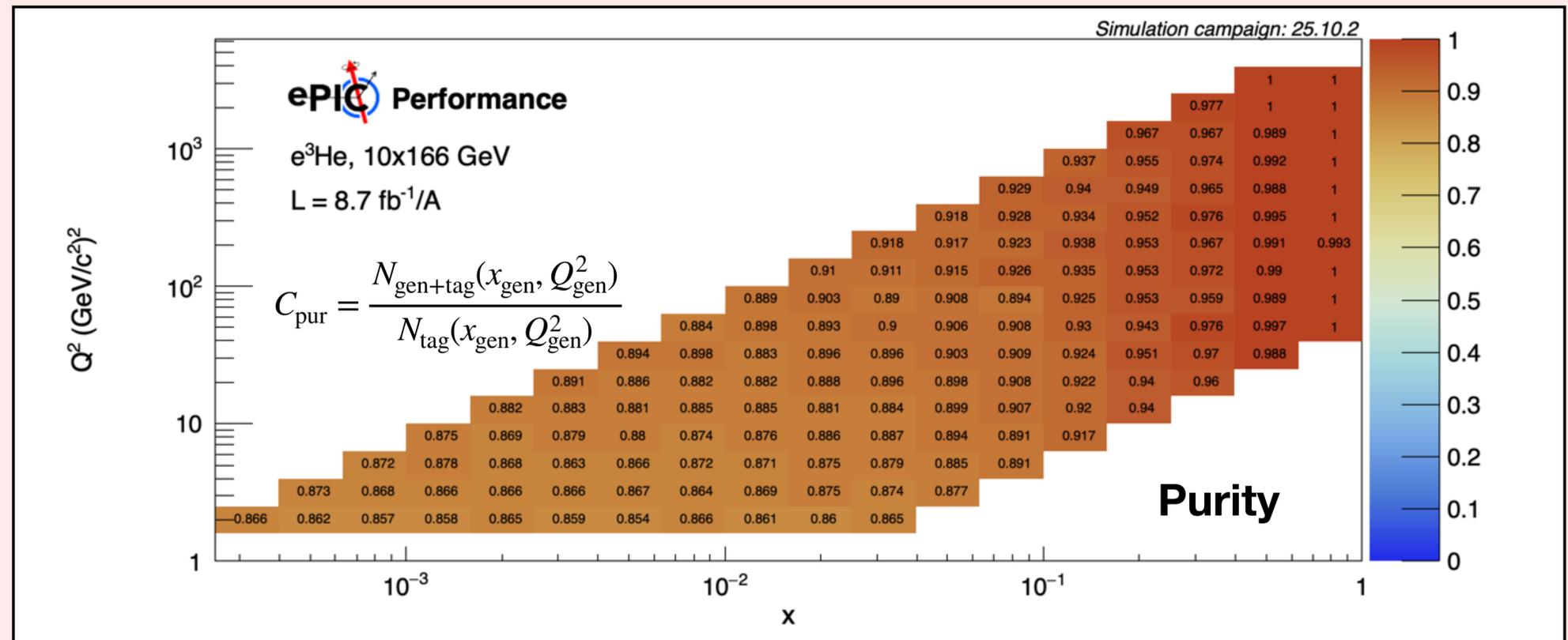
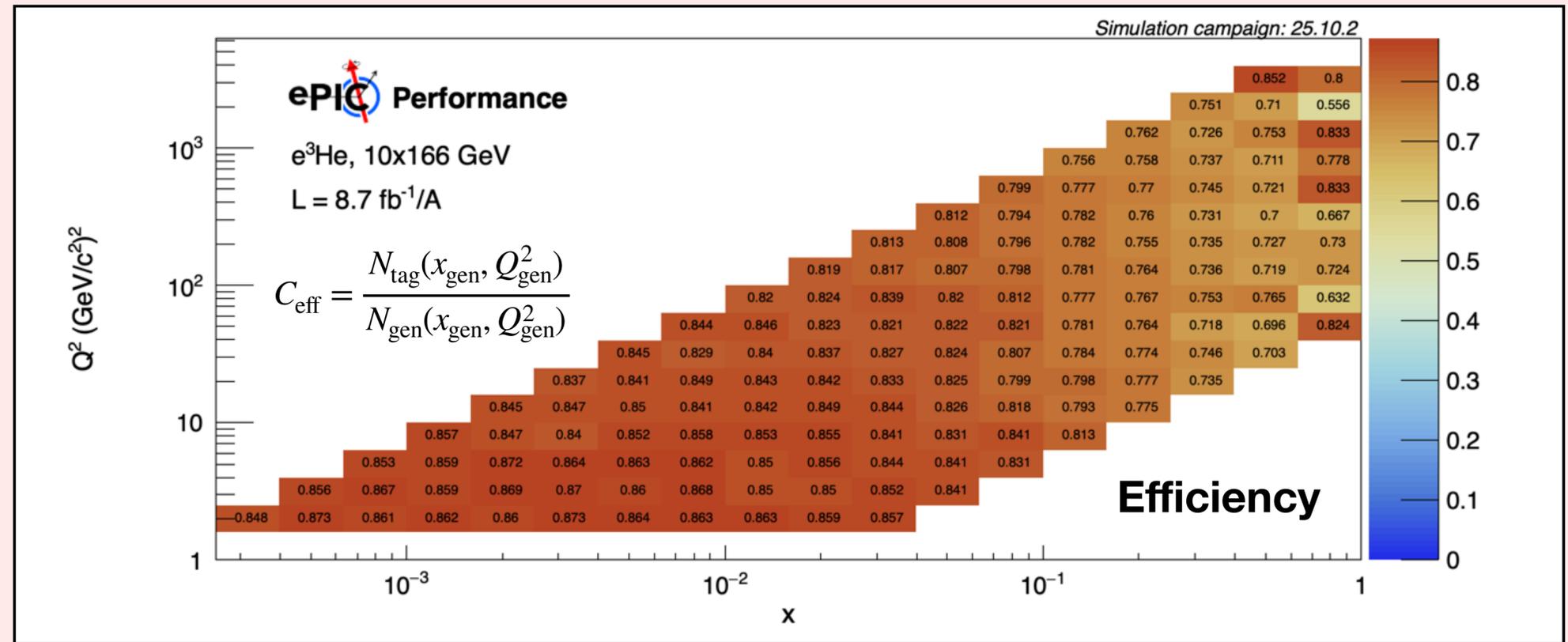
Double spectator tagging - method

- ▶ Hit-based tagging algorithm
- ▶ Number of proton tracks = min number of hits per plane per detector
- ▶ If number of proton tracks > 2 , then tag as double spectator



Double spectator tagging - efficiency & purity

- ▶ Hit-based tagging algorithm
- ▶ Number of proton tracks = min number of hits per plane per detector
- ▶ If number of proton tracks > 2, then tag as double spectator



Projected A_1^n from $e^3\text{He}$ DIS:

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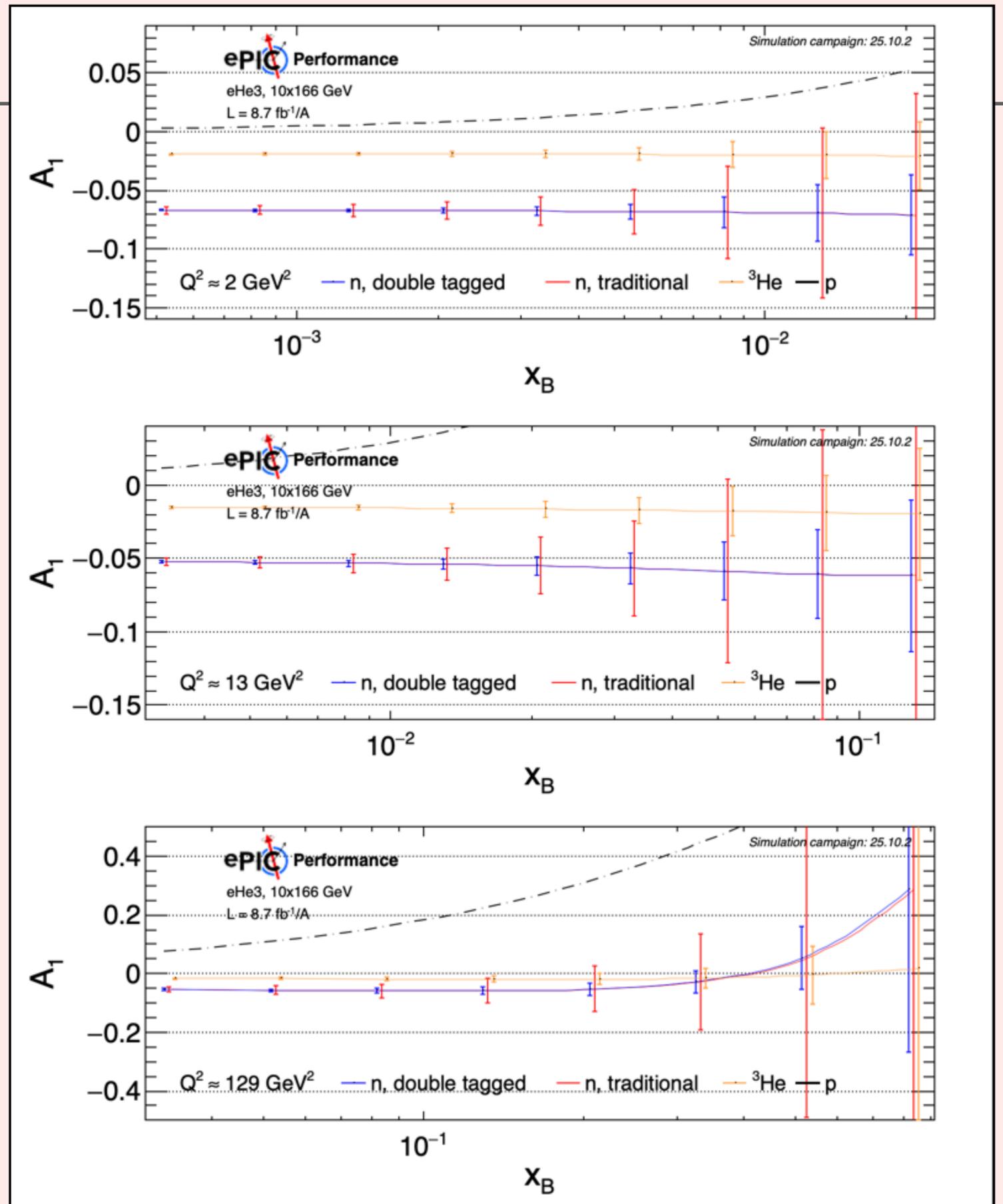
$$\triangleright P_e = P_n = 70\%$$

$$\triangleright 0.01 < y < 0.95, Q^2 \geq 2, W^2 > 4$$

\triangleright R calculated from [https://doi.org/10.1016/](https://doi.org/10.1016/S0370-2693(99)00244-0)

[S0370-2693\(99\)00244-0](https://doi.org/10.1016/S0370-2693(99)00244-0)

\triangleright Included statistical uncertainty and model uncertainties



Projected A_1^n :

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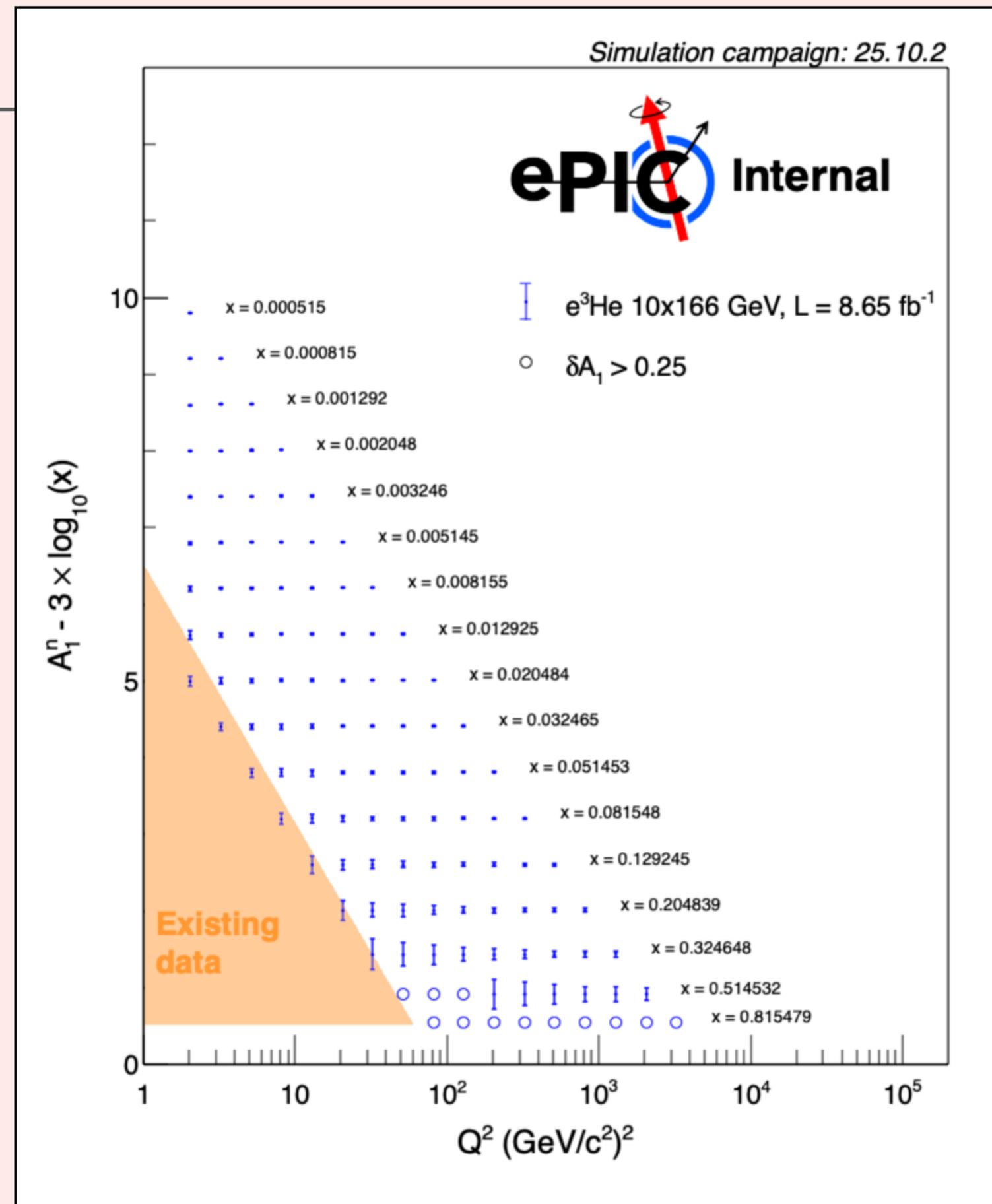
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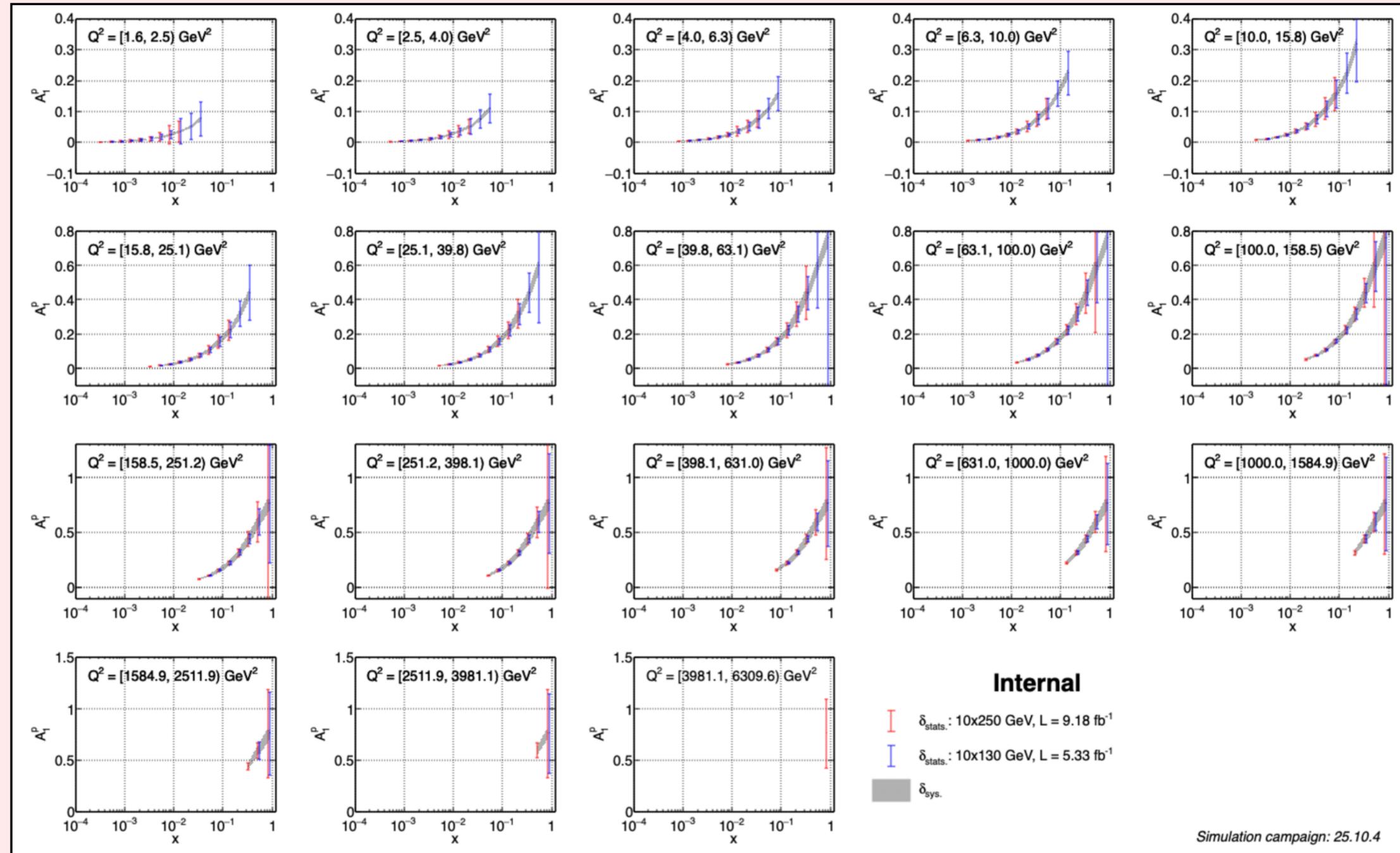


▶ estimation for ATHENA for A_{\parallel}

- ▶ 1.5% point-by-point uncorrelated
- ▶ 3% pion contamination
- ▶ 4.7% normalization
 - ▶ 3% δP_e , 3% δP_N
 - ▶ 1-2% detector effects

▶ Estimated using $A_1 \approx \frac{A_{\parallel}}{D}$

- ▶ Systematics to be studied in more details

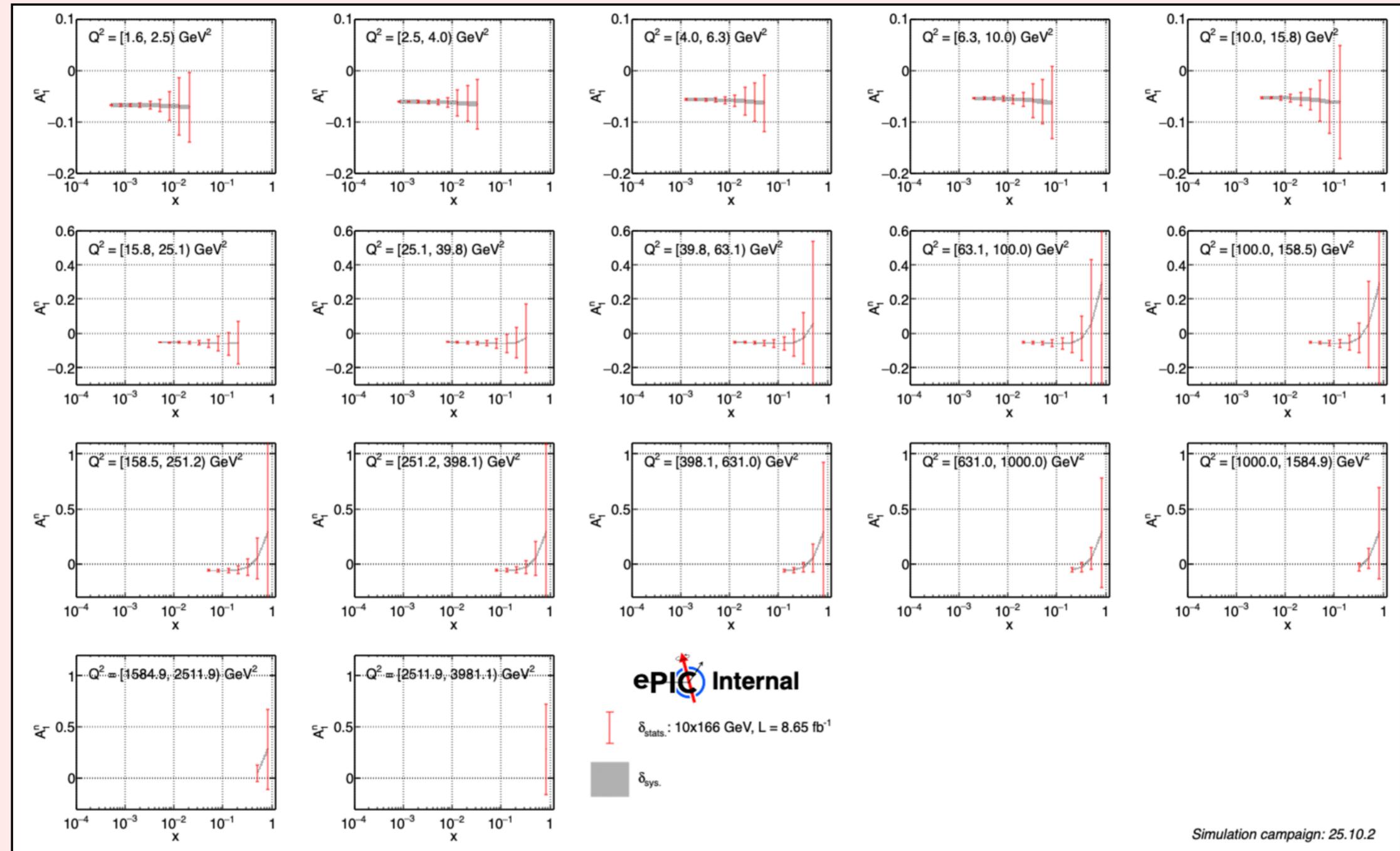


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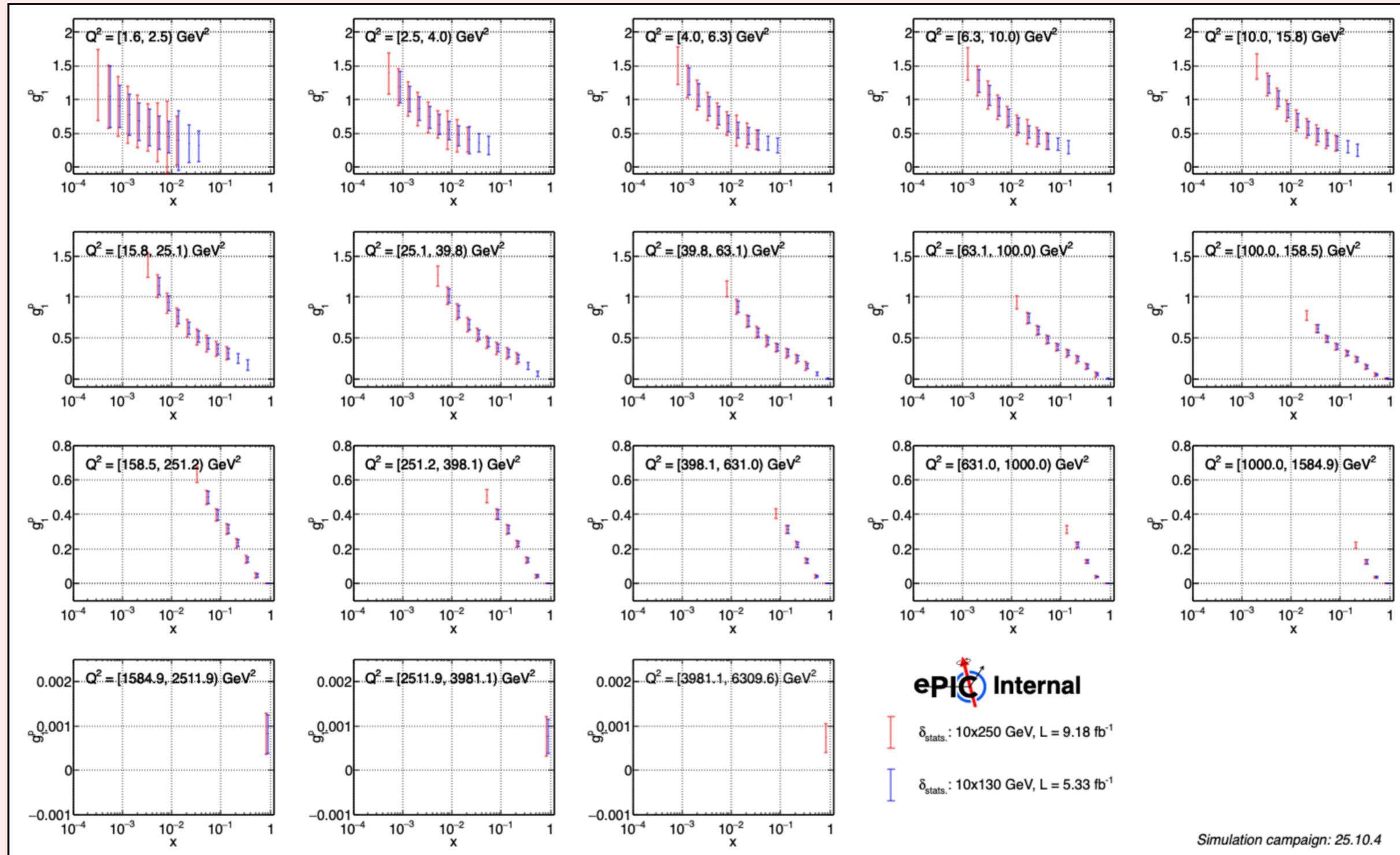
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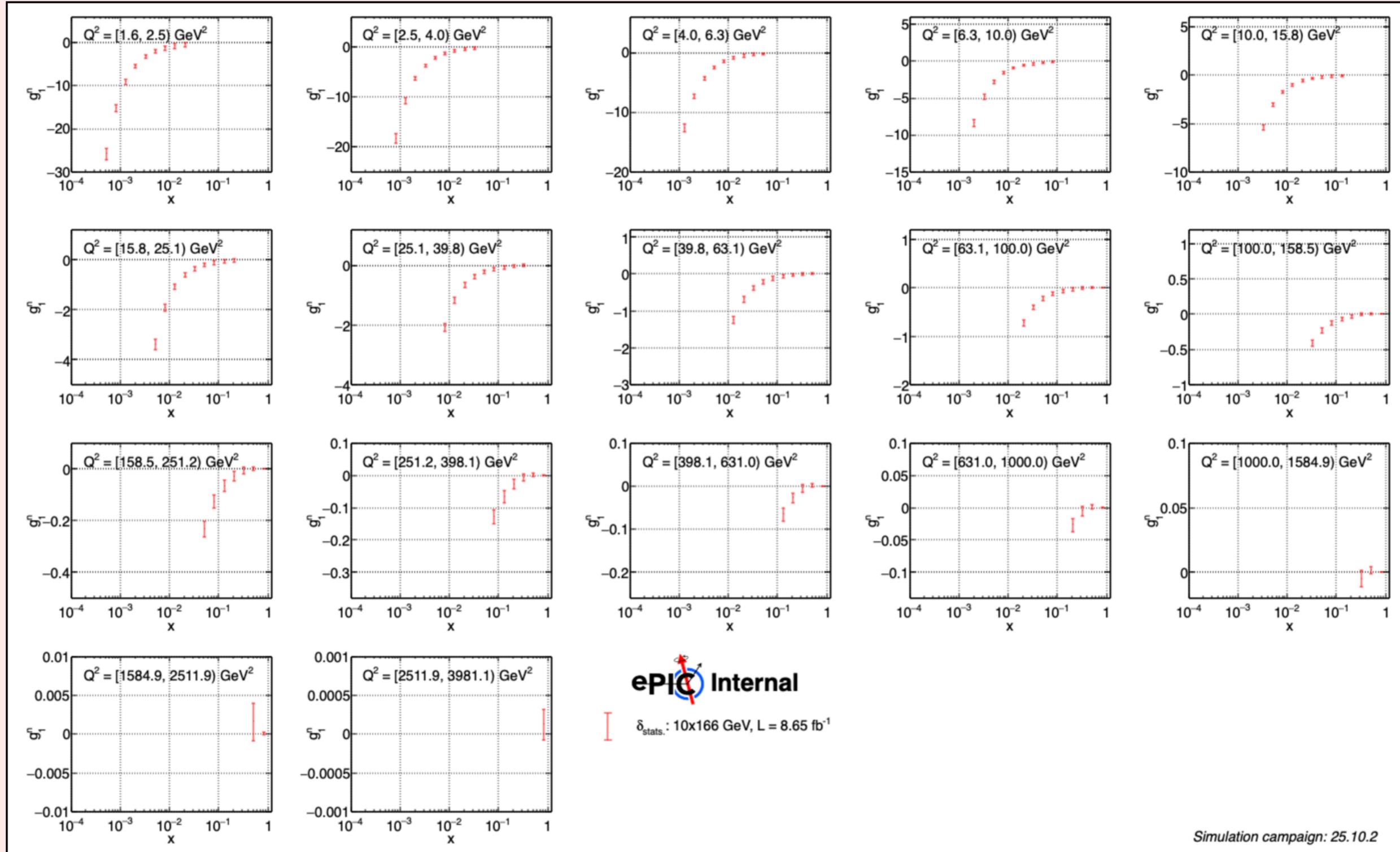
▶ Systematics to be studied in more details



- ▶ $A_1 \approx g_1/F_1$ with F_1 calculated from JAM22



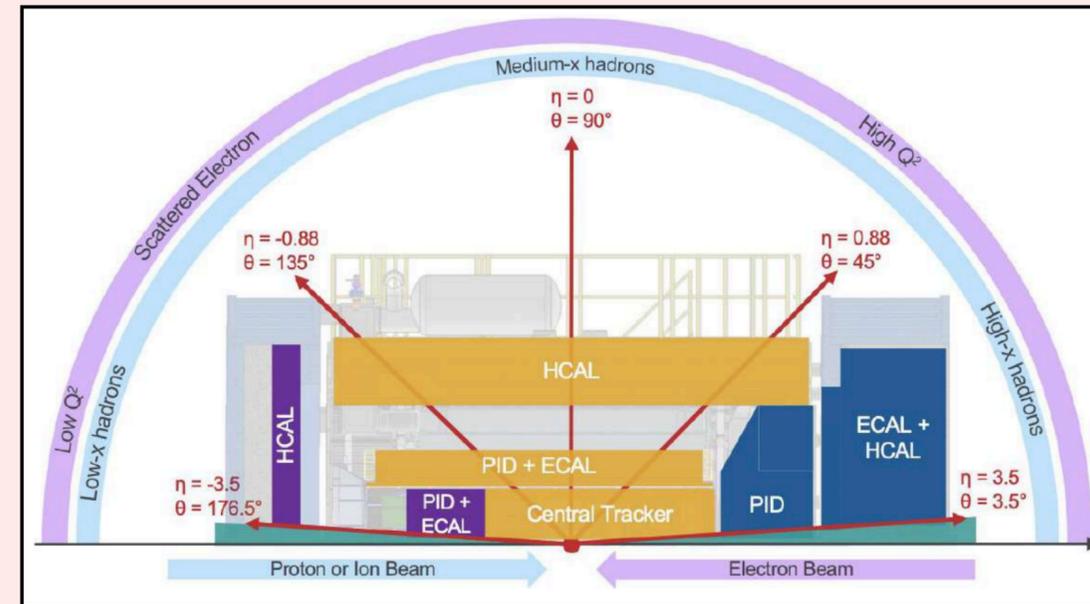
- ▶ $A_1 \approx g_1/F_1$ with F_1 calculated from JAM22



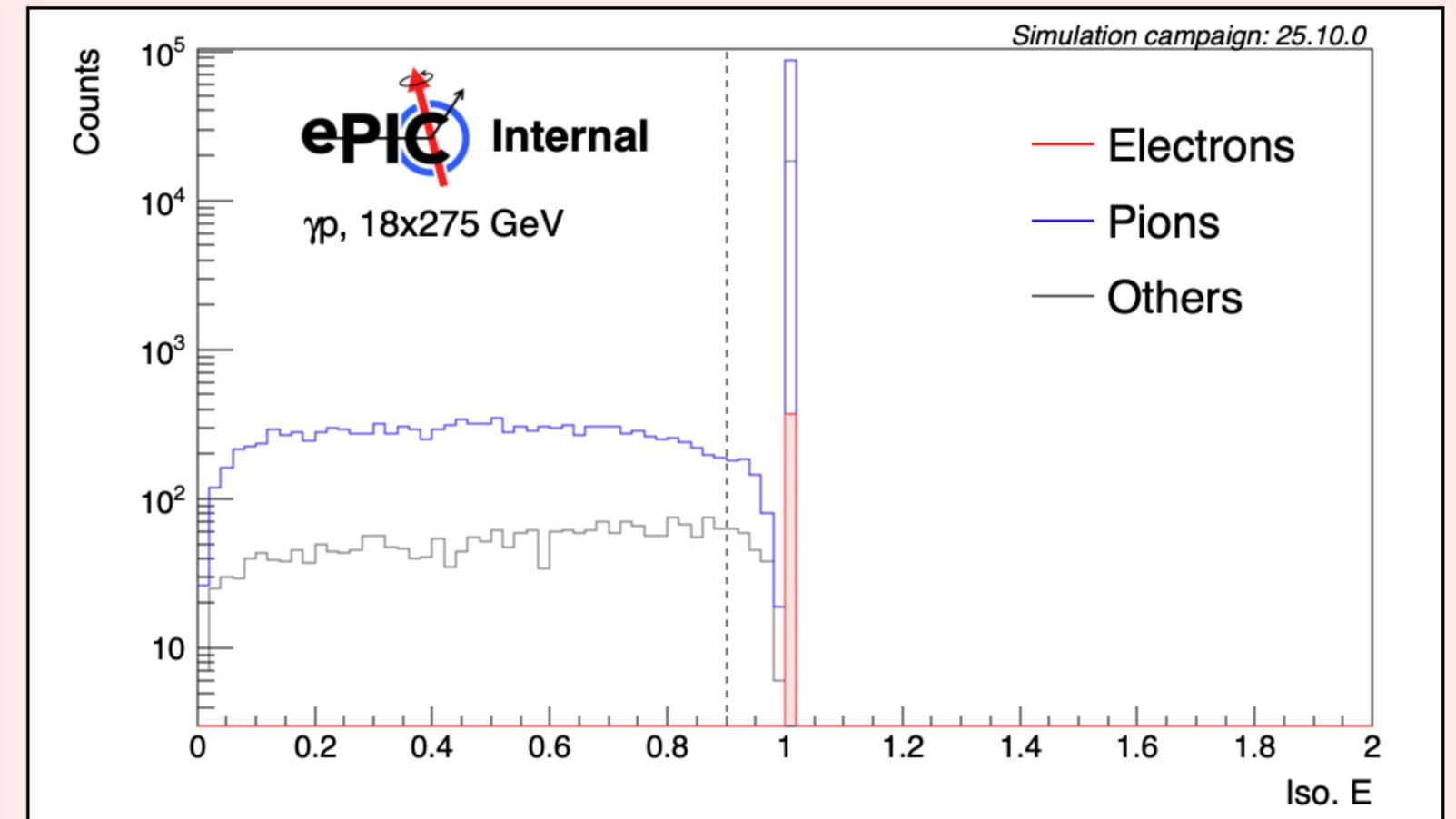
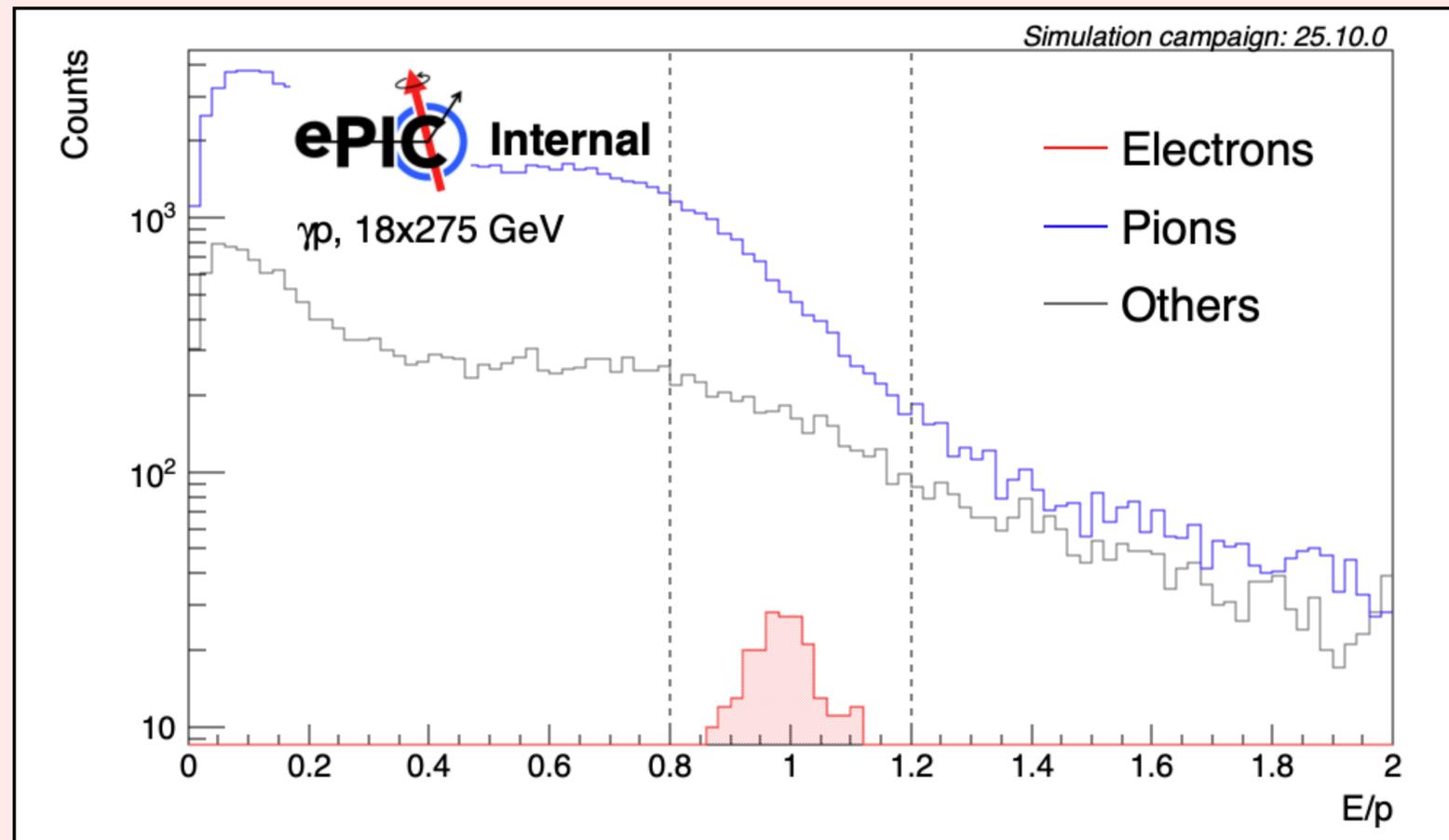
- ▶ Pion contamination?
- ▶ Beam background?
- ▶ Radiative corrections?
- ▶ ...

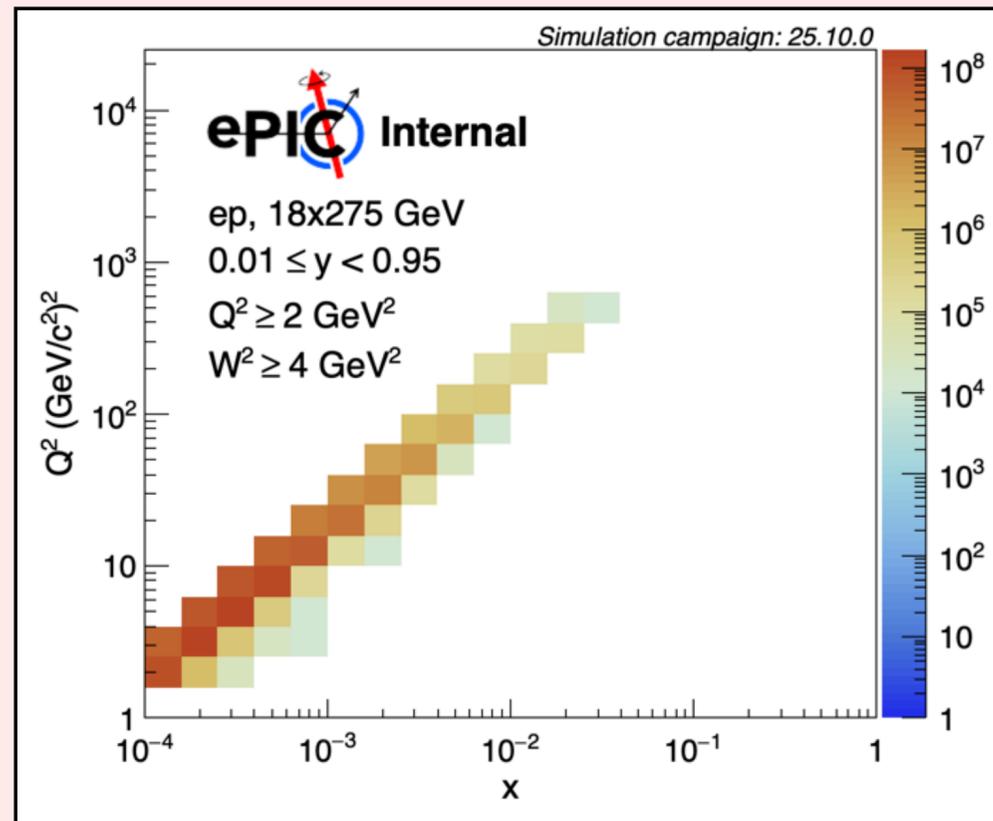
Photoproduction background

π background lower the purity at very high y at low x and low Q^2

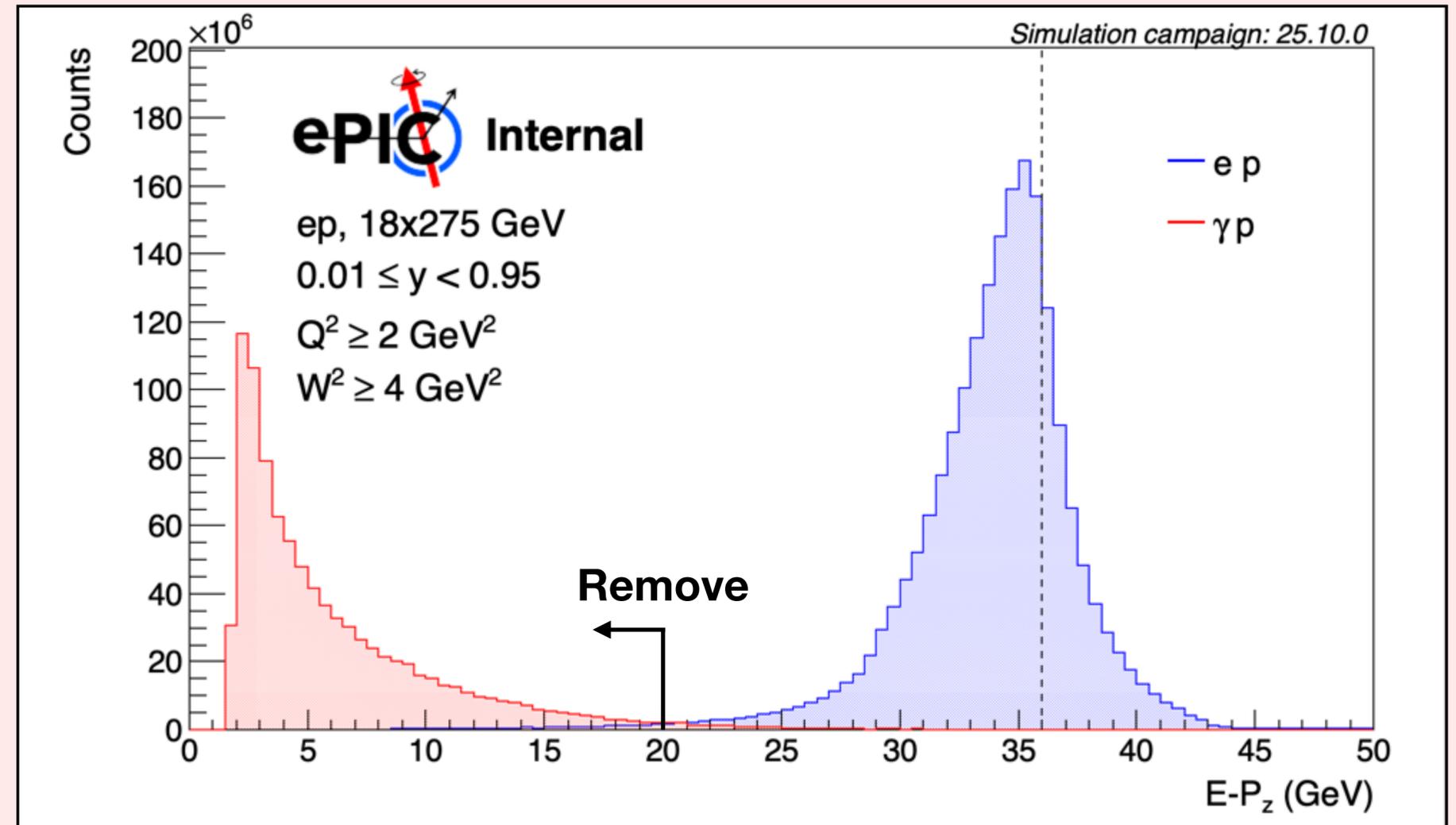
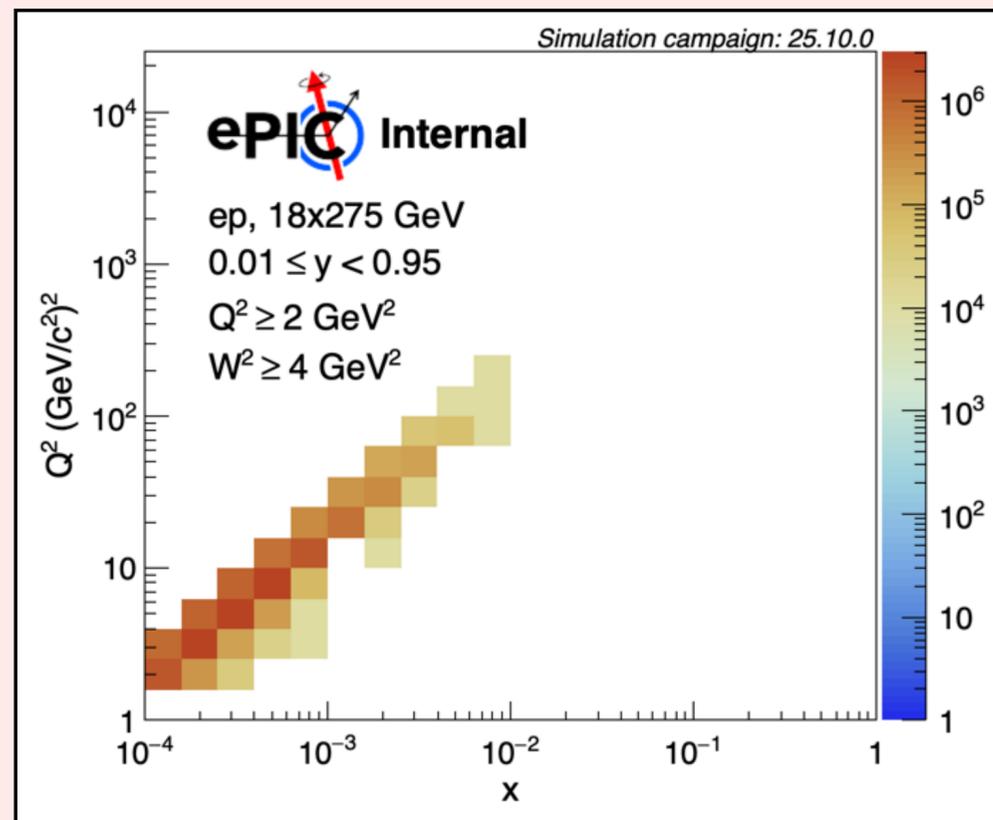


Low Q^2 event samples:
SIDIS/pythia6-eic/
1.0.0/18x275/q2_0to1/

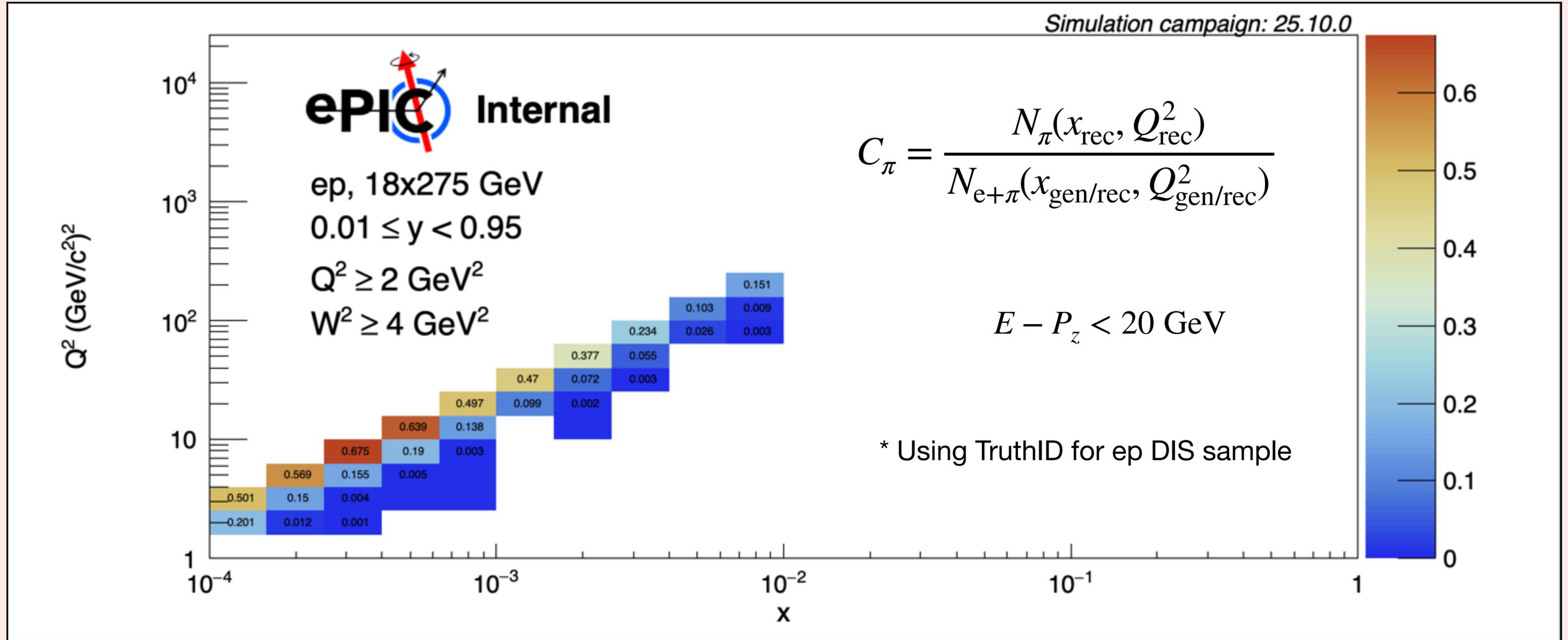


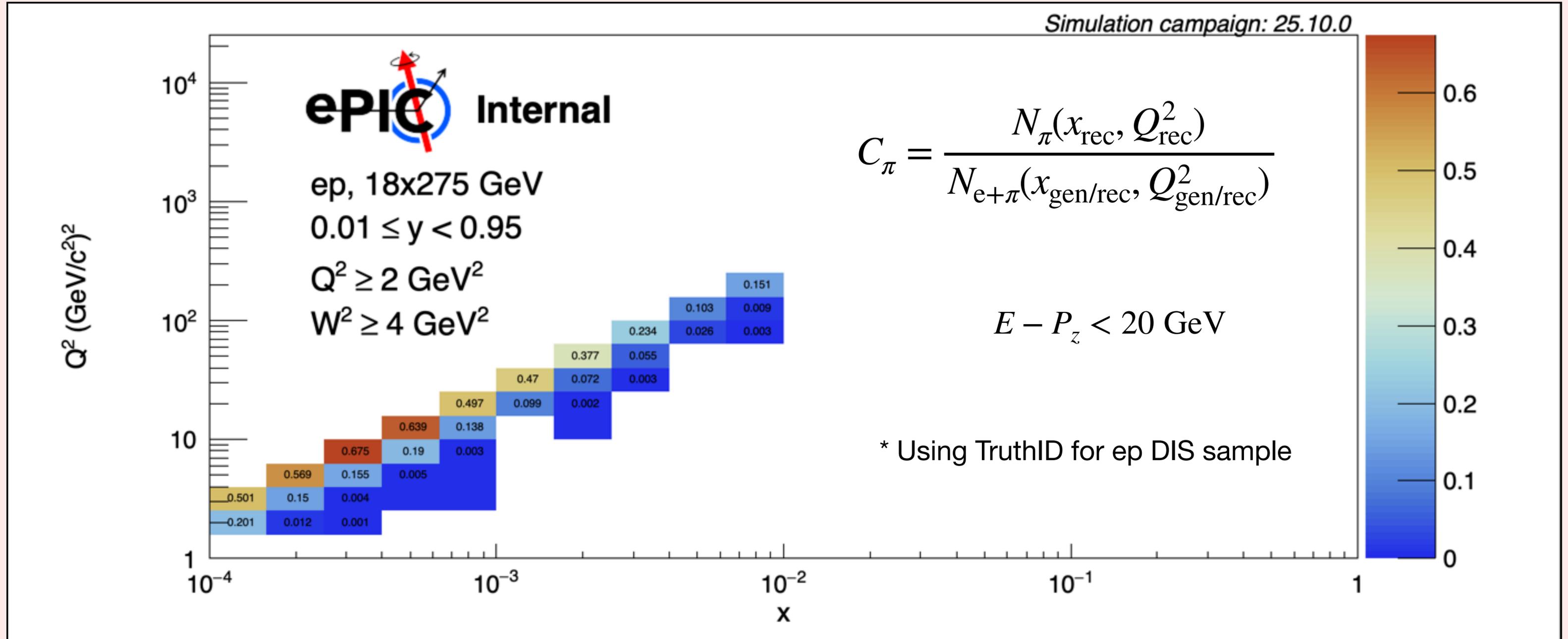


Reduced by
~97.95%



Both samples are scaled to $L = 10 \text{ fb}^{-1}$



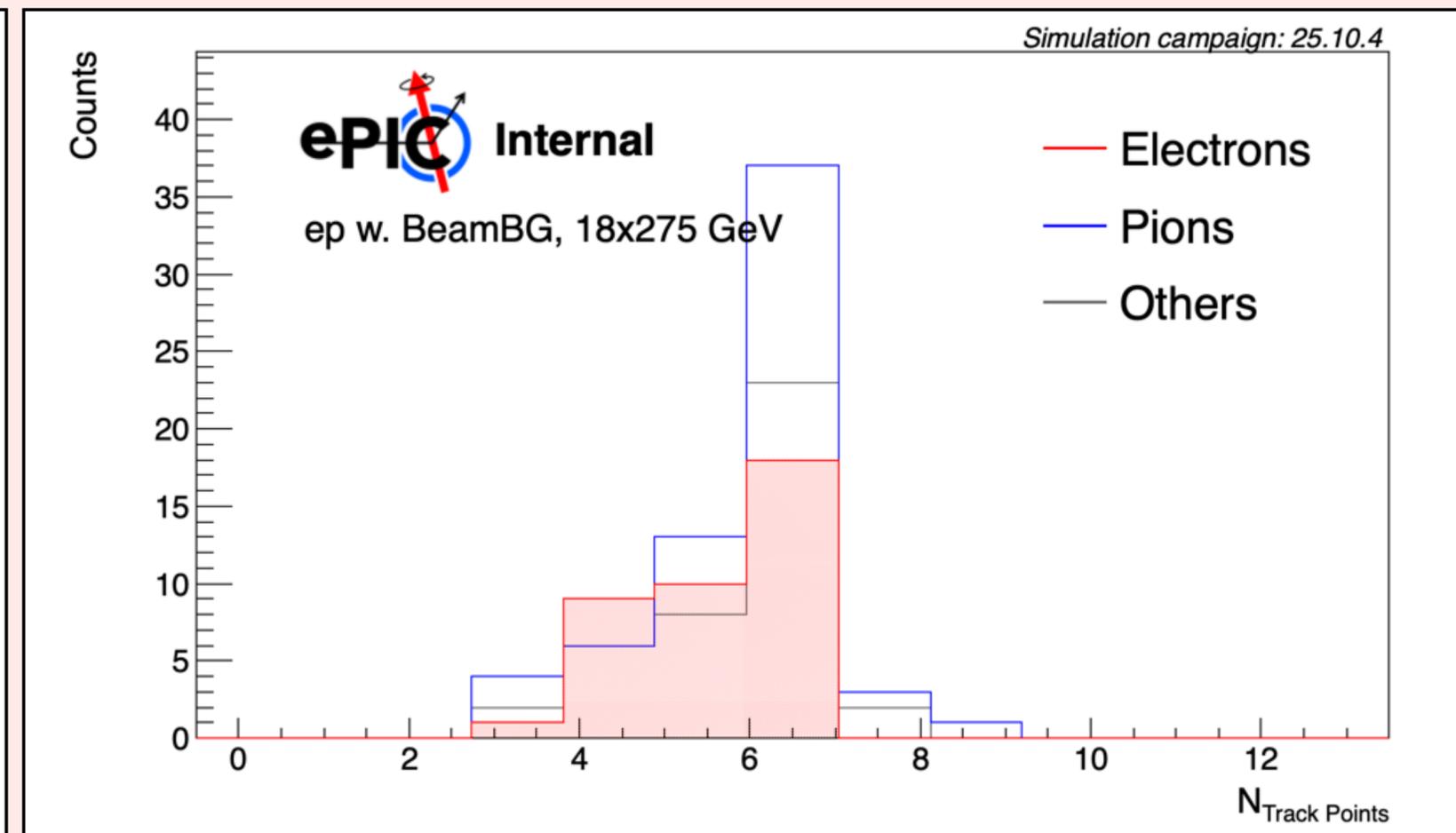
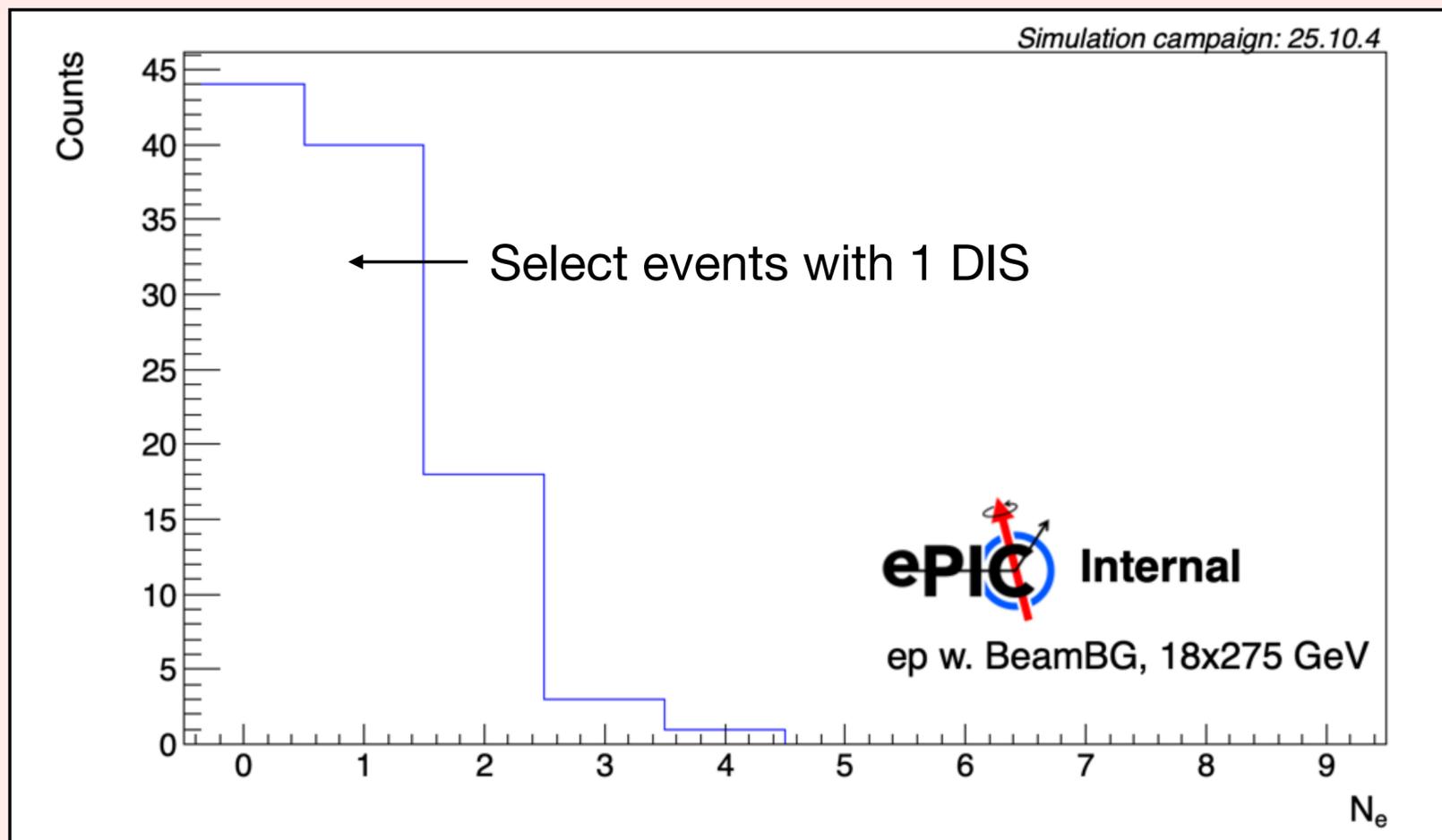


Need to use the same event generator for background and signal!

Machine background sample at 18x275 GeV:

Number of scattered electrons

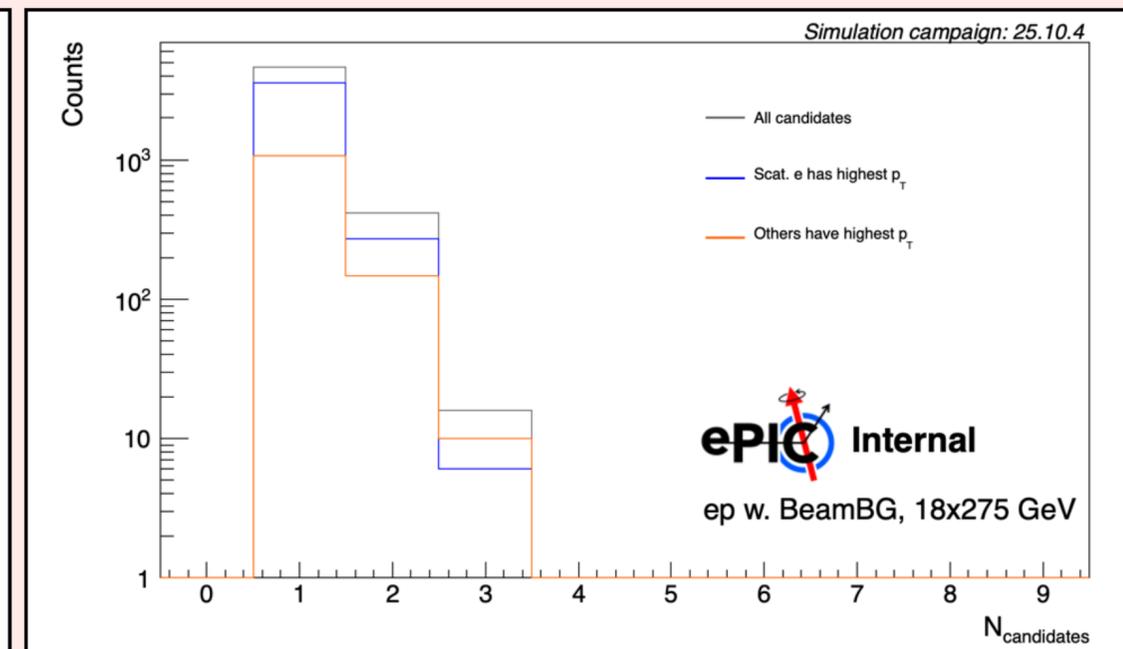
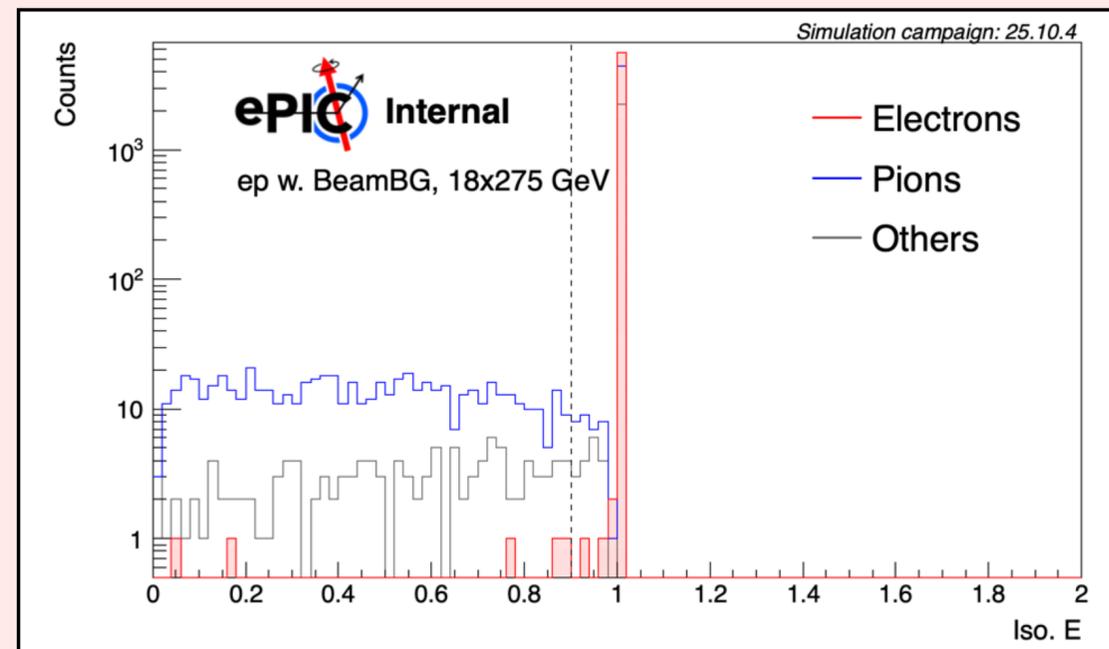
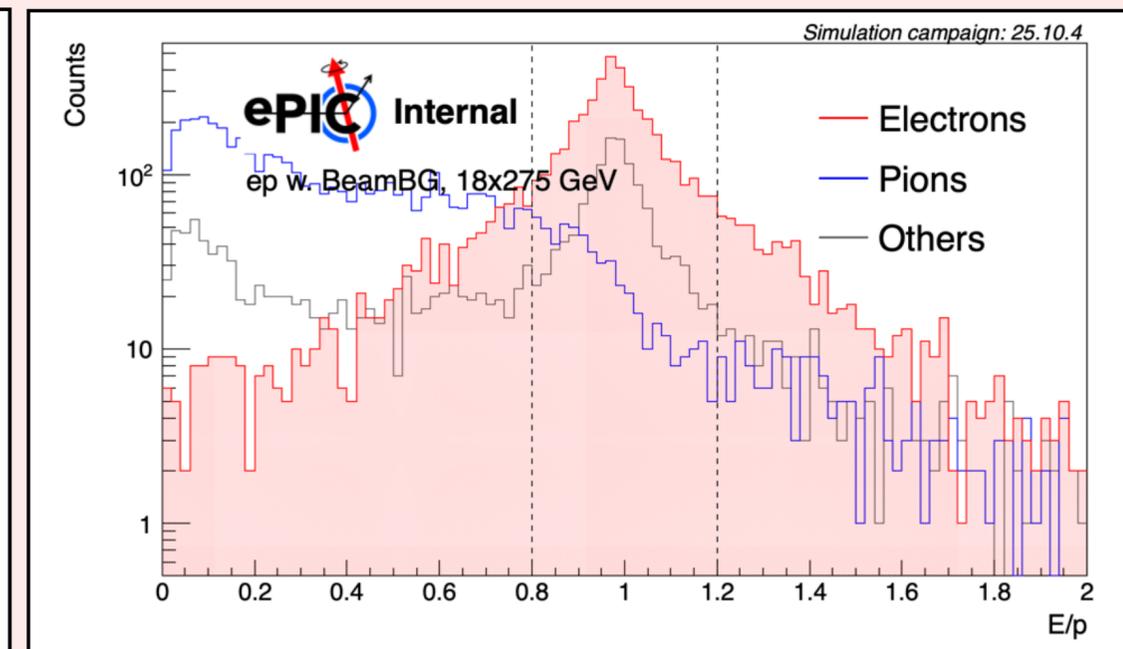
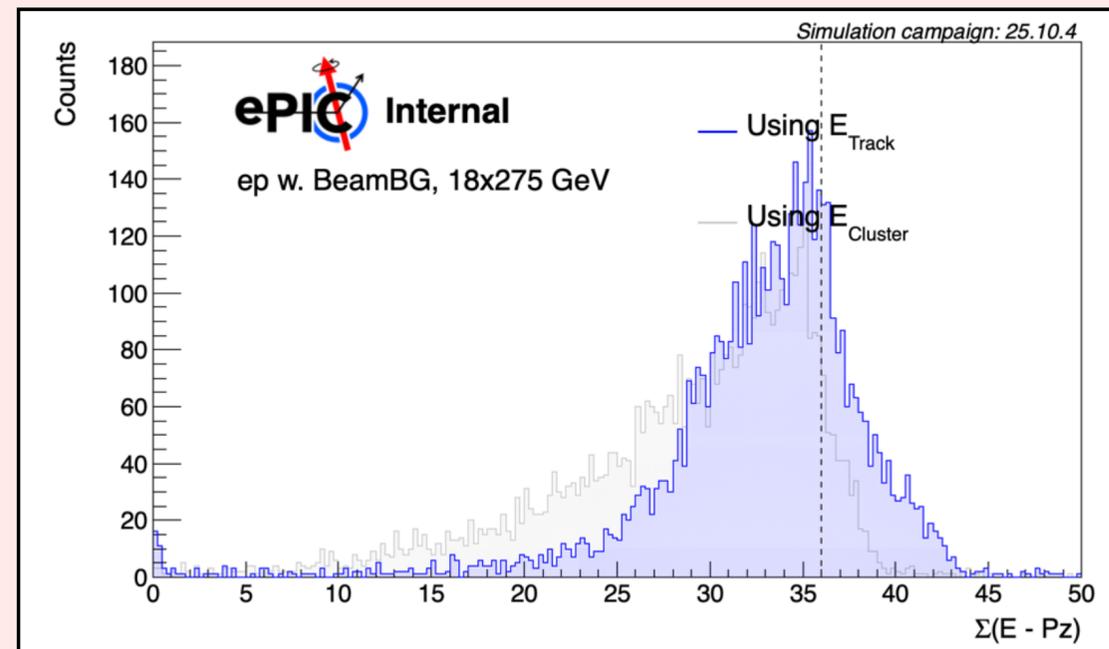
Number of points used for tracks



- ▶ Effect on eID at 18x275 GeV is small; need to understand E and p resolution better (low resolution might be due to other problem in current simulation and analysis)

- ▶ Need to look at impact on kinematic reconstruction
- ▶ 10 GeV beam electron is expected to have larger impact. Will show study soon.

Plots shown at collaboration meeting Jan 2026



Projection and analysis status:

- Realistic physics and detector simulation
- Electron identification algorithm
- Kinematic reconstruction algorithm
- Primitive tagging algorithm

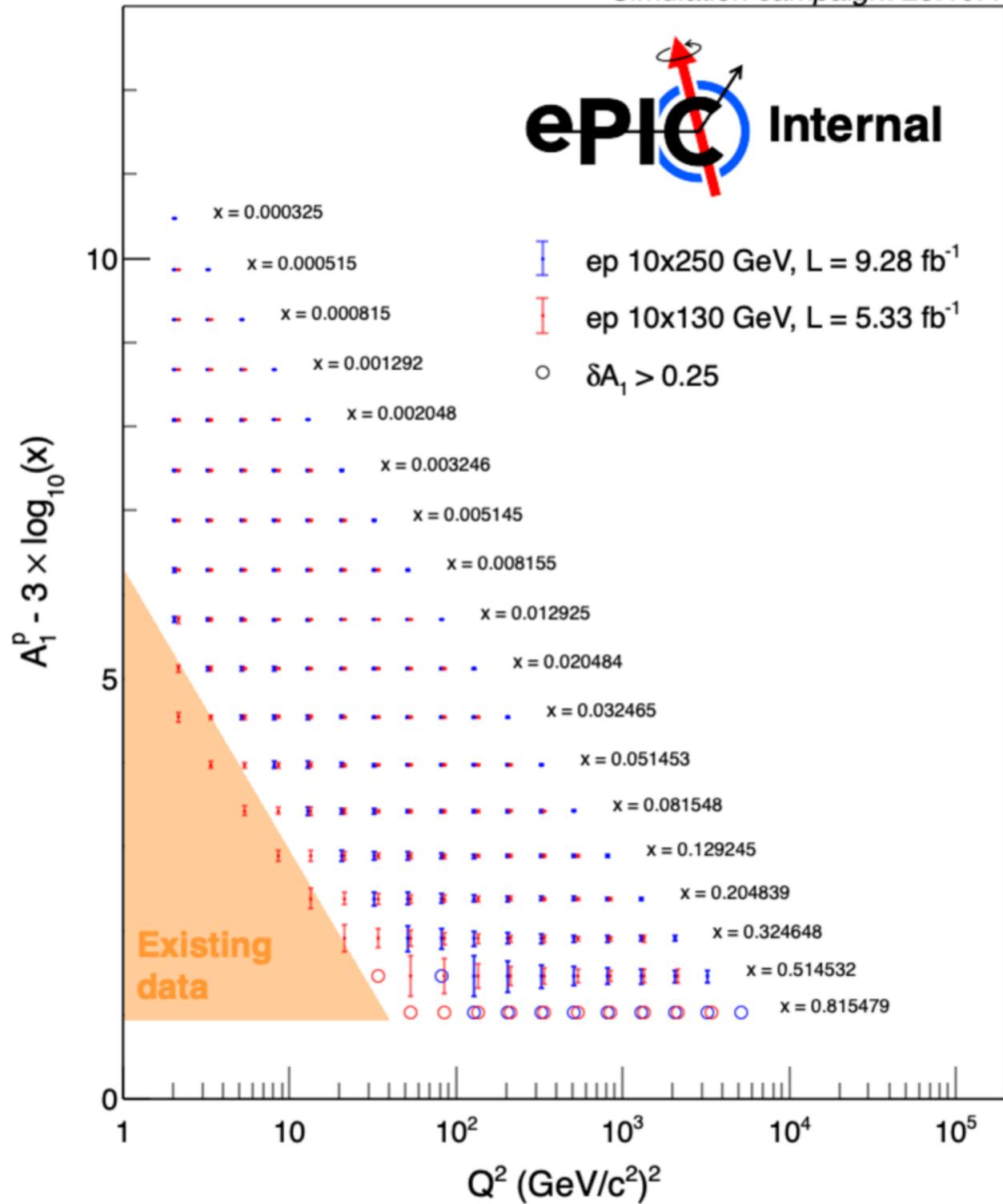
- Projected A_1^p, A_1^n with uncertainty estimate
- Projected g_1^p, g_1^n with uncertainty estimate

On-going study:

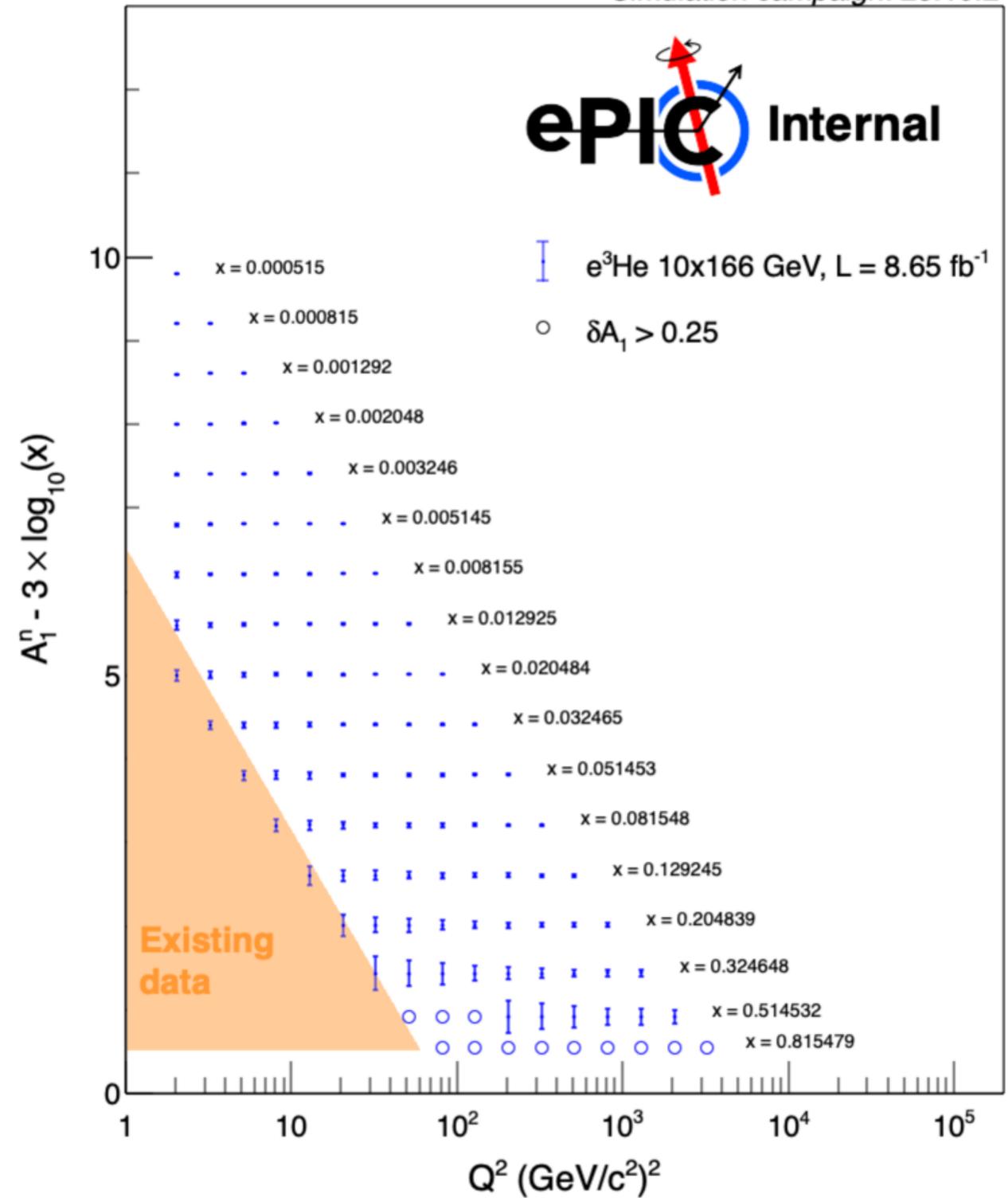
- Pion contamination
- Beam background

- Better tagging
- Use ZDC?

Simulation campaign: 25.10.4



Simulation campaign: 25.10.2



Thank you :)